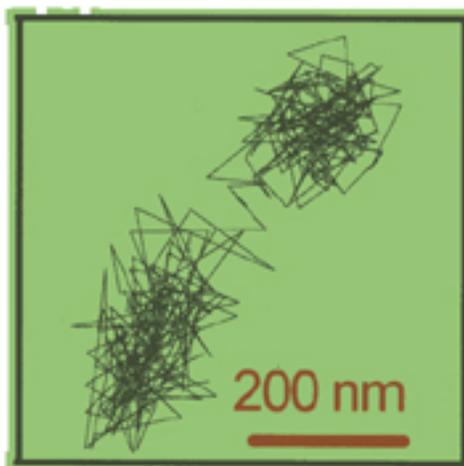
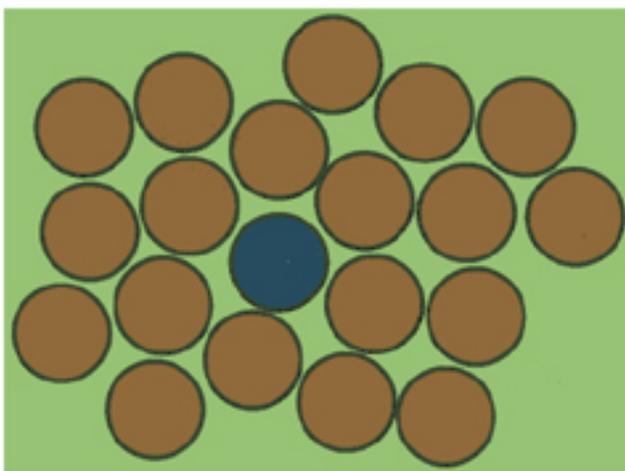


Cage trapping:



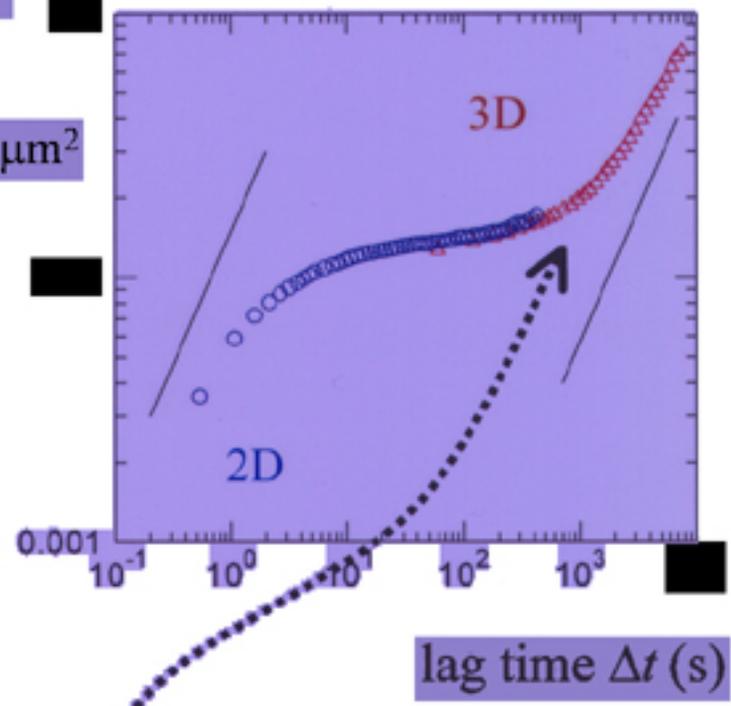
$\phi=0.56$, 100 min
(supercooled fluid)

- Short times: particles stuck in “cages”
- Long times: cages rearrange

Intermittency

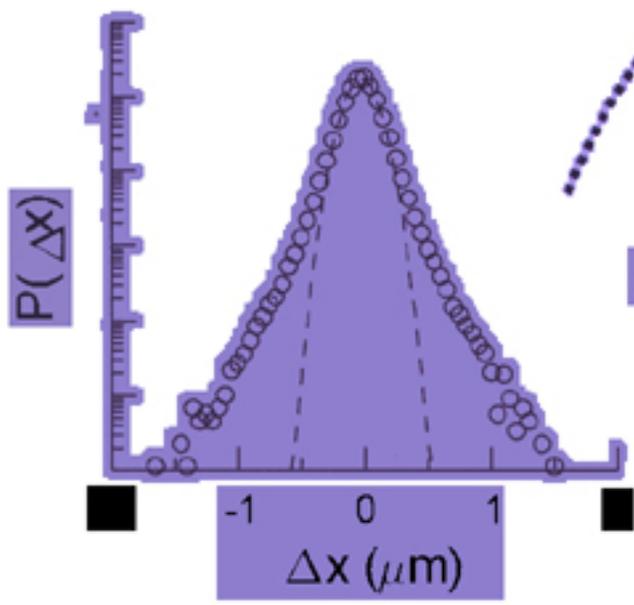
Volume fraction $\phi=0.5$
"supercooled fluid"

$\langle x^2 \rangle \mu\text{m}^2$



lag time Δt (s)

$\Delta t = 1000$ s



Nongaussian Parameter α_2

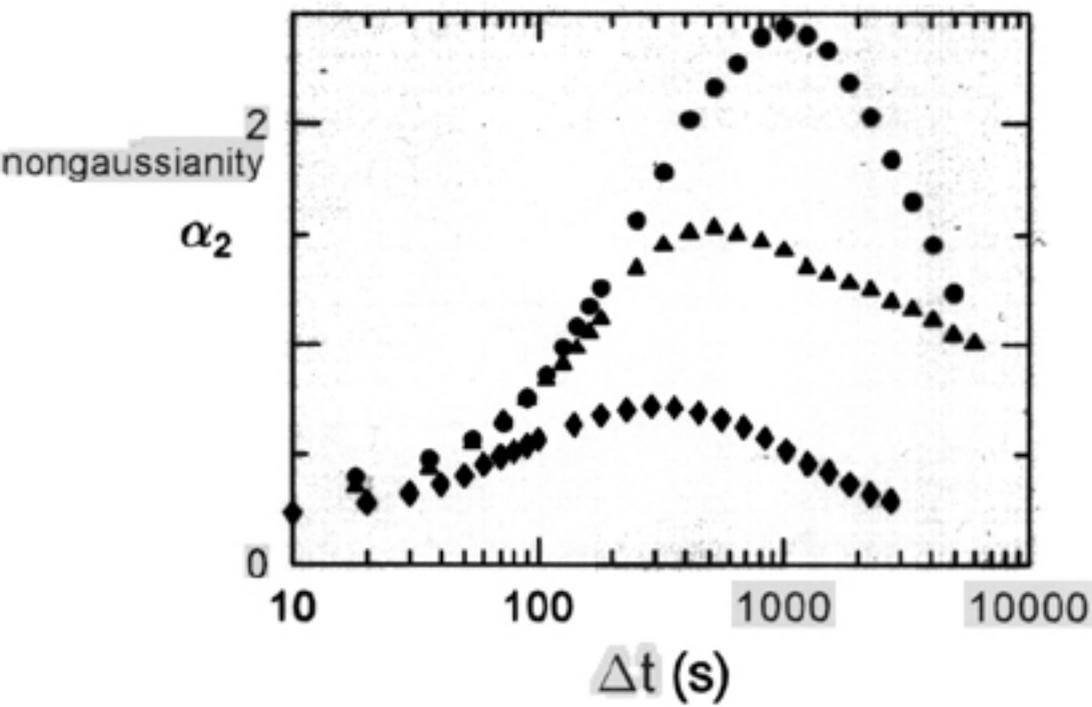
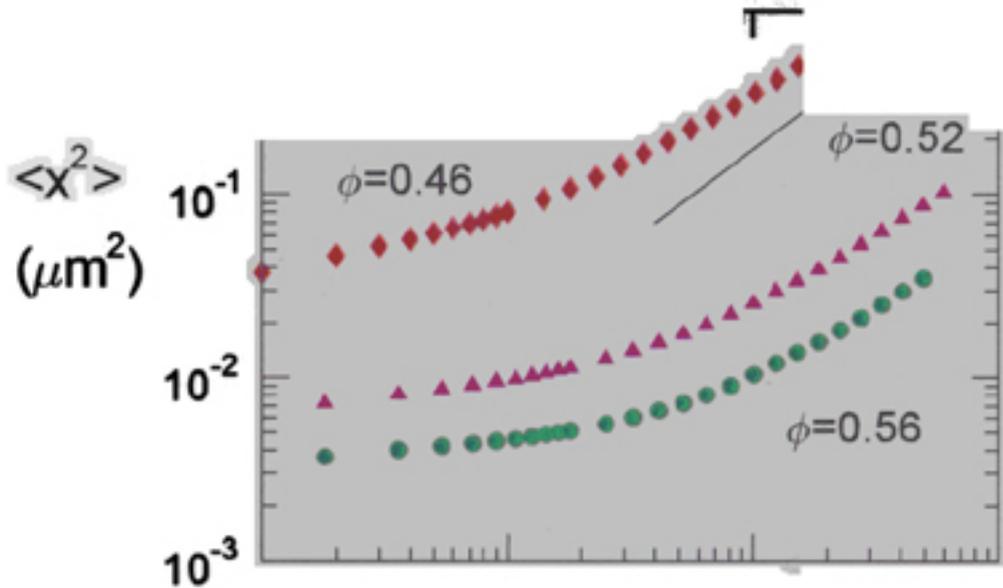
Rahman, 1964

$$\alpha_2 = \frac{\langle x^4 \rangle}{3\langle x^2 \rangle^2} - 1$$

Gaussian: $\alpha_2 = 0$

Exponential: $\alpha_2 = 1$





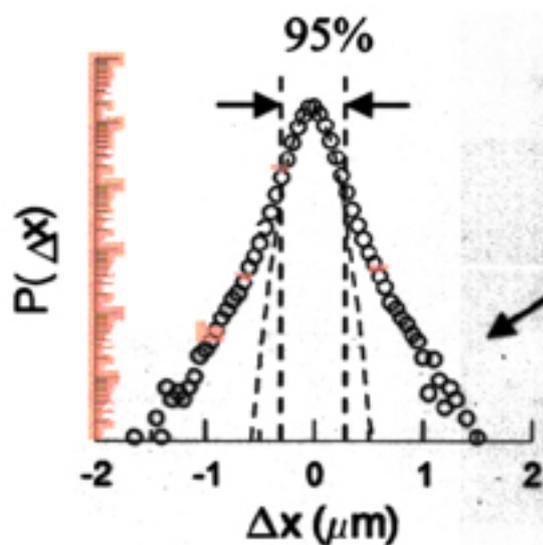
Time scale:

Δt^* when nongaussian parameter α_2 largest

Length scale:

Δr^* on average, 5% of particles have $\Delta r(\Delta t^*) > \Delta r^*$

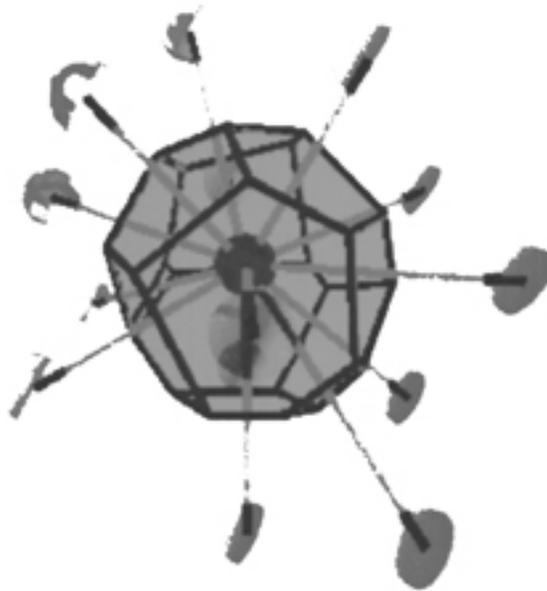
\approx cage rearrangements



top 5% = tails
of Δx distribution

($\phi=0.53$, supercooled fluid)

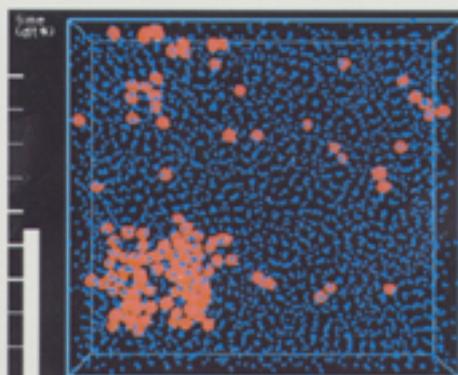
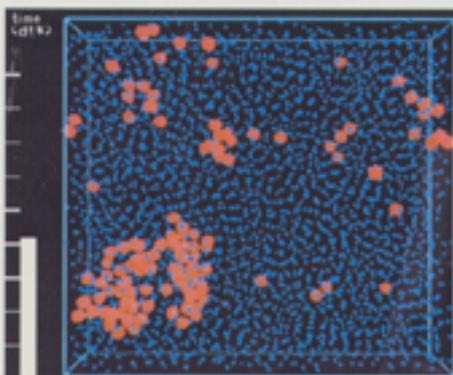
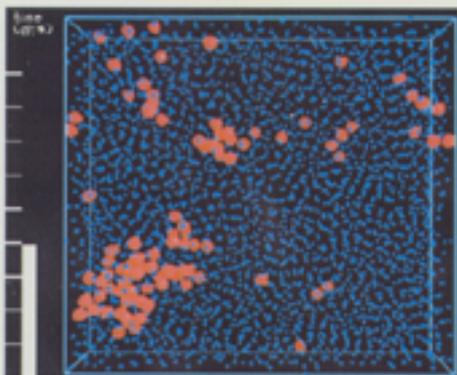
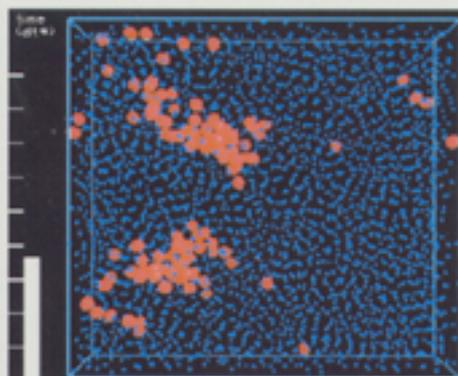
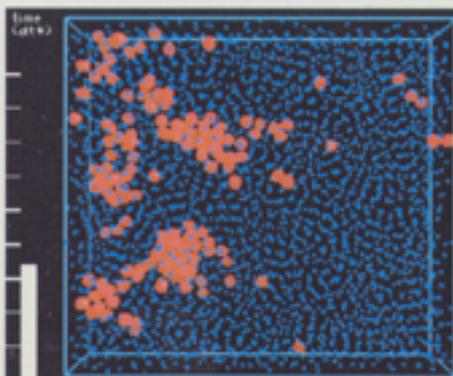
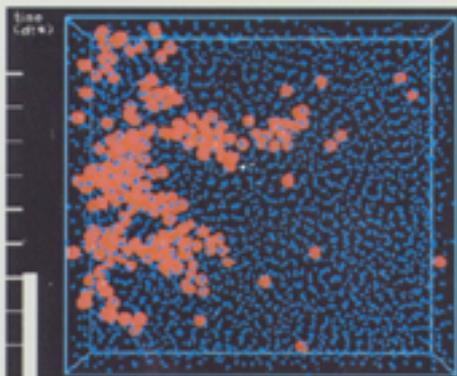
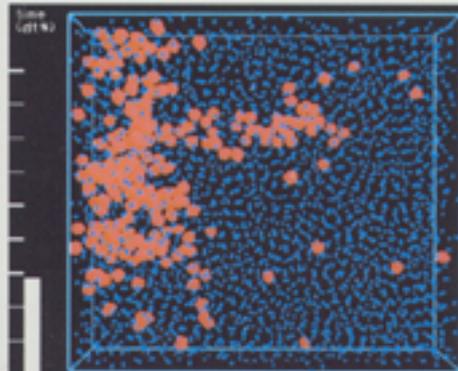
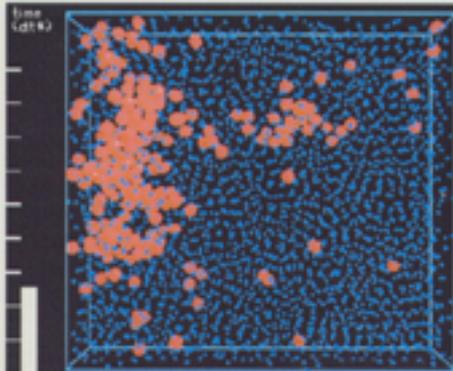
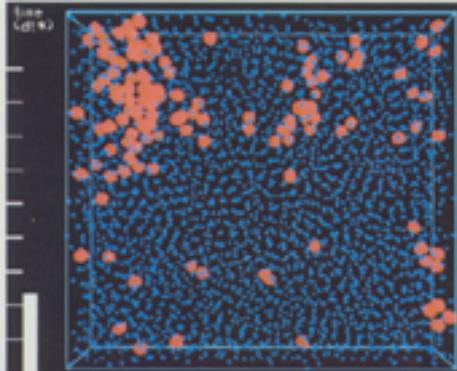
Voronoi polyhedra -- Delaunay triangulation



(“Wigner-Seitz cell”)

polyhedron face: *set of all points
equidistant from two particles*

defines nearest neighbor particles

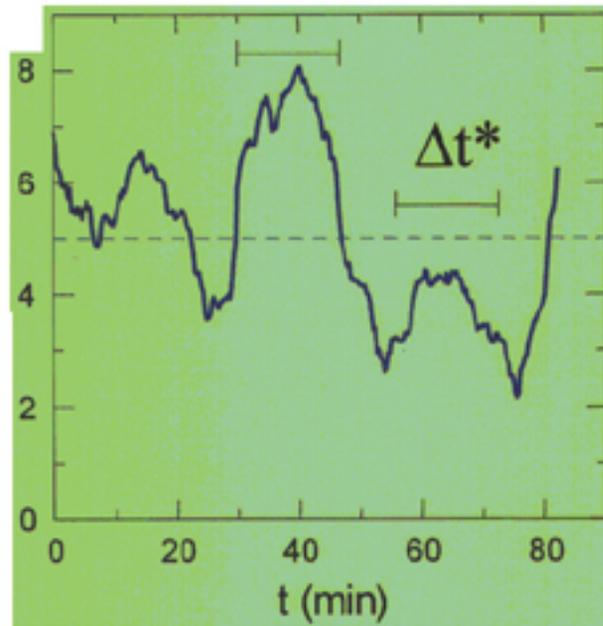


F

f

Supercooled fluid

$$\phi = 0.56$$

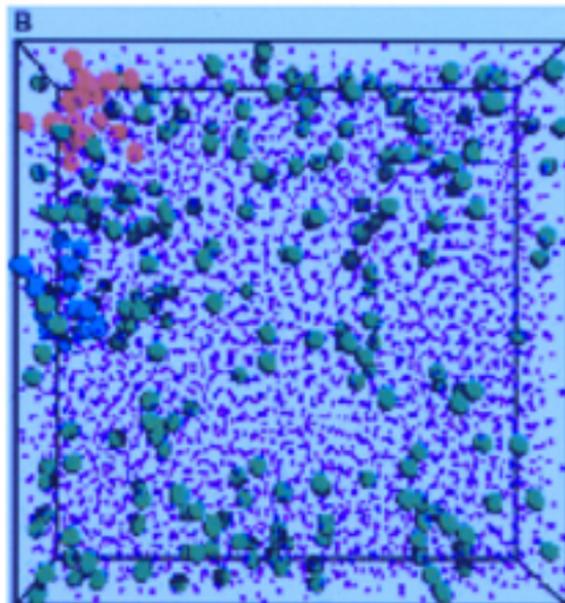
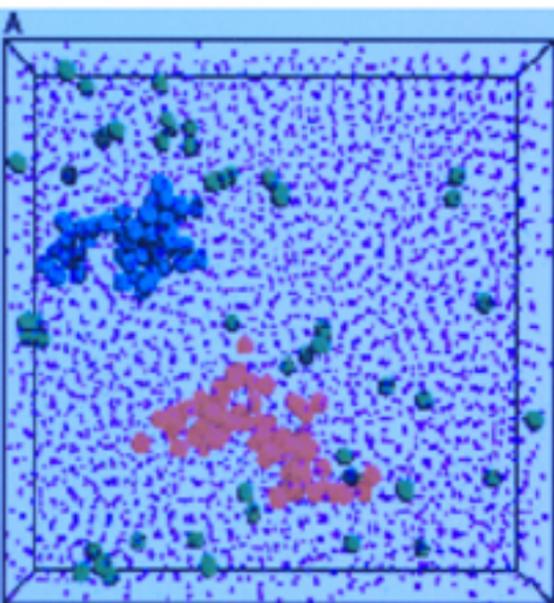


top 5%

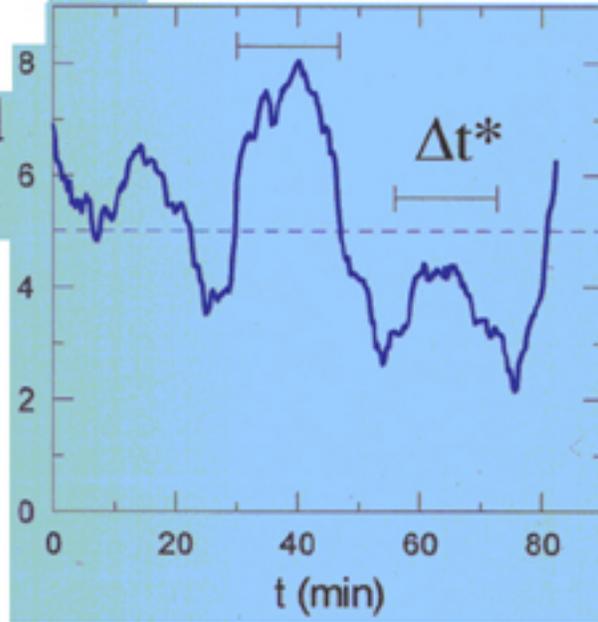
Very large fluctuations

- Sample is not ergodic

Glass Transition

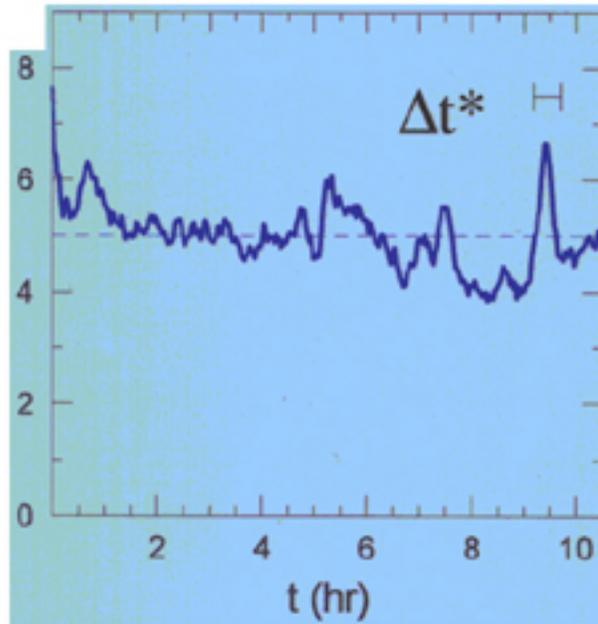


supercooled
fluid $\phi=0.56$



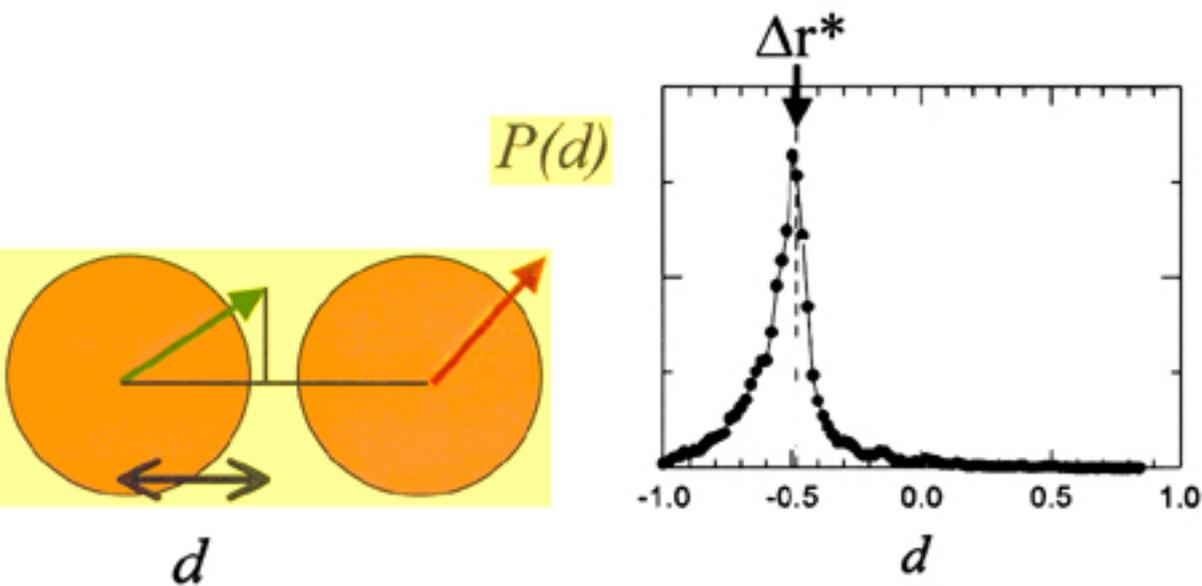
top 5%

glass
 $\phi=0.61$

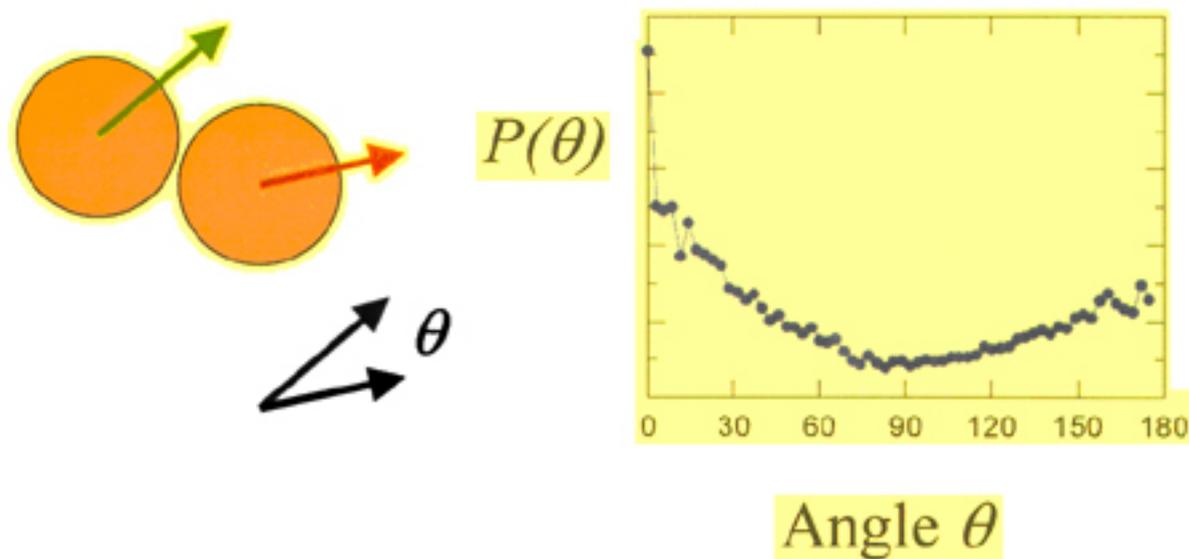


top 5%

Particles move towards neighbors

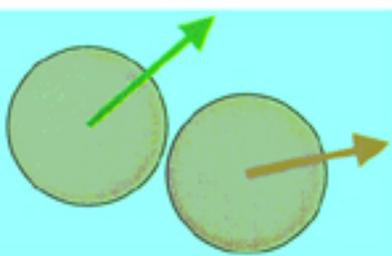


Particles move in same direction as neighbors

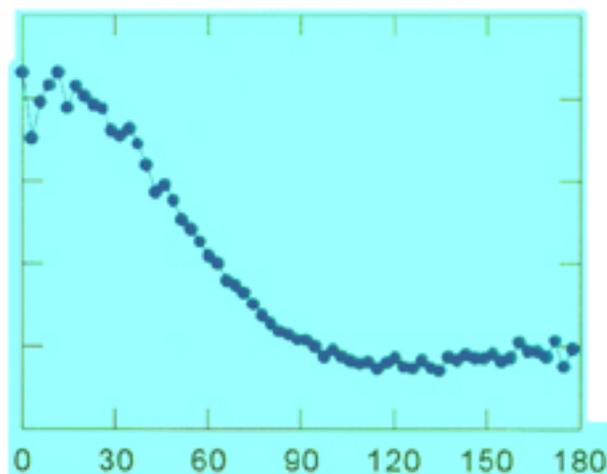


($P(\theta)$ corrected for spherical projection)

Particles follow neighbors

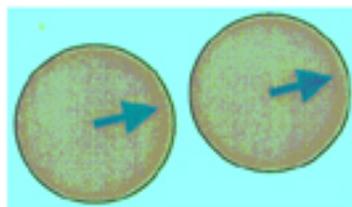


$P(\theta)$



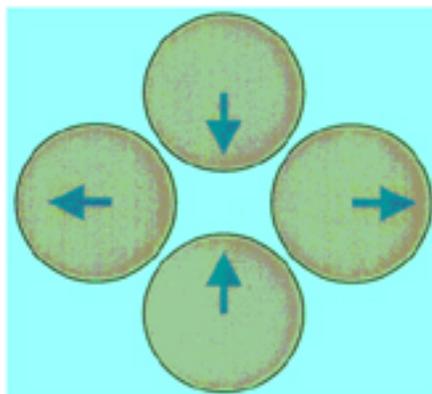
Angle θ

($P(\theta)$ corrected for spherical projection)



$\theta \approx 0^\circ$

common



$\theta \approx 180^\circ$

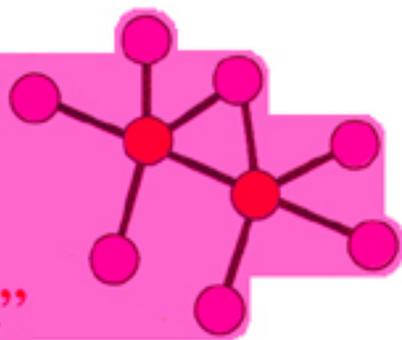
rare

Look for local orientational order:

- $q_{6m}(i)$ forms a 13-dimensional vector

- find unit vectors $q_{6m}(i)$

- if $q_{6m}(i) \cdot q_{6m}(j) > 0.5$, then bond (ij) is “crystal-like”

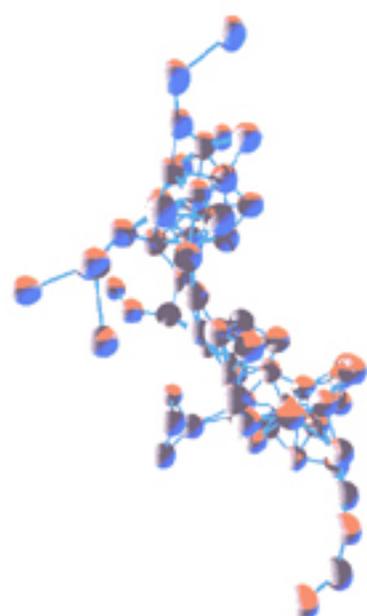
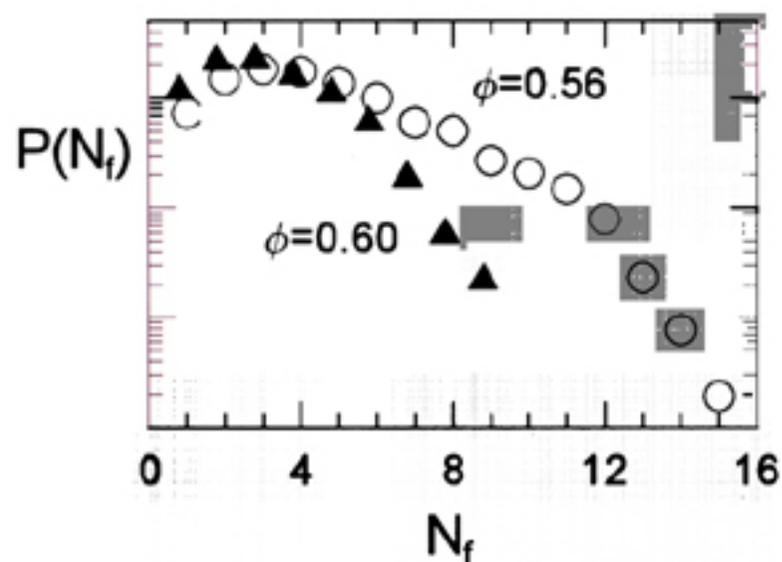


- if particle has ≥ 8 “crystal-like” bonds, it is a **crystal-like particle**

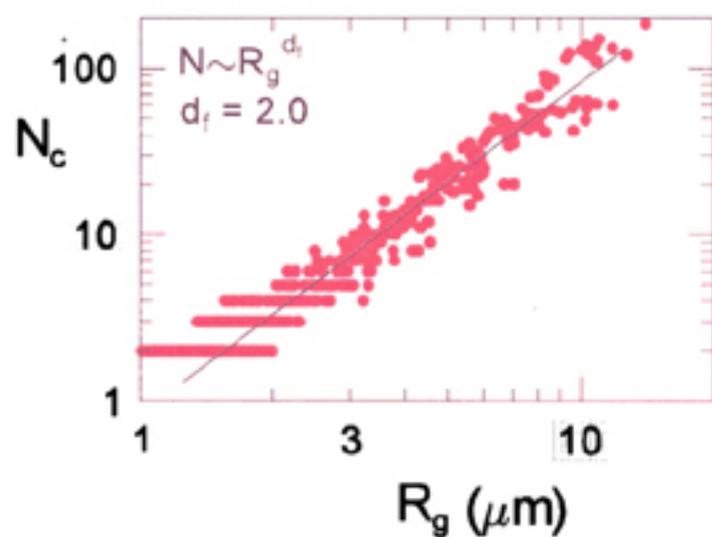


rotationally invariant, insensitive to type of crystal

Number N_f of fast neighbors to a fast particle:

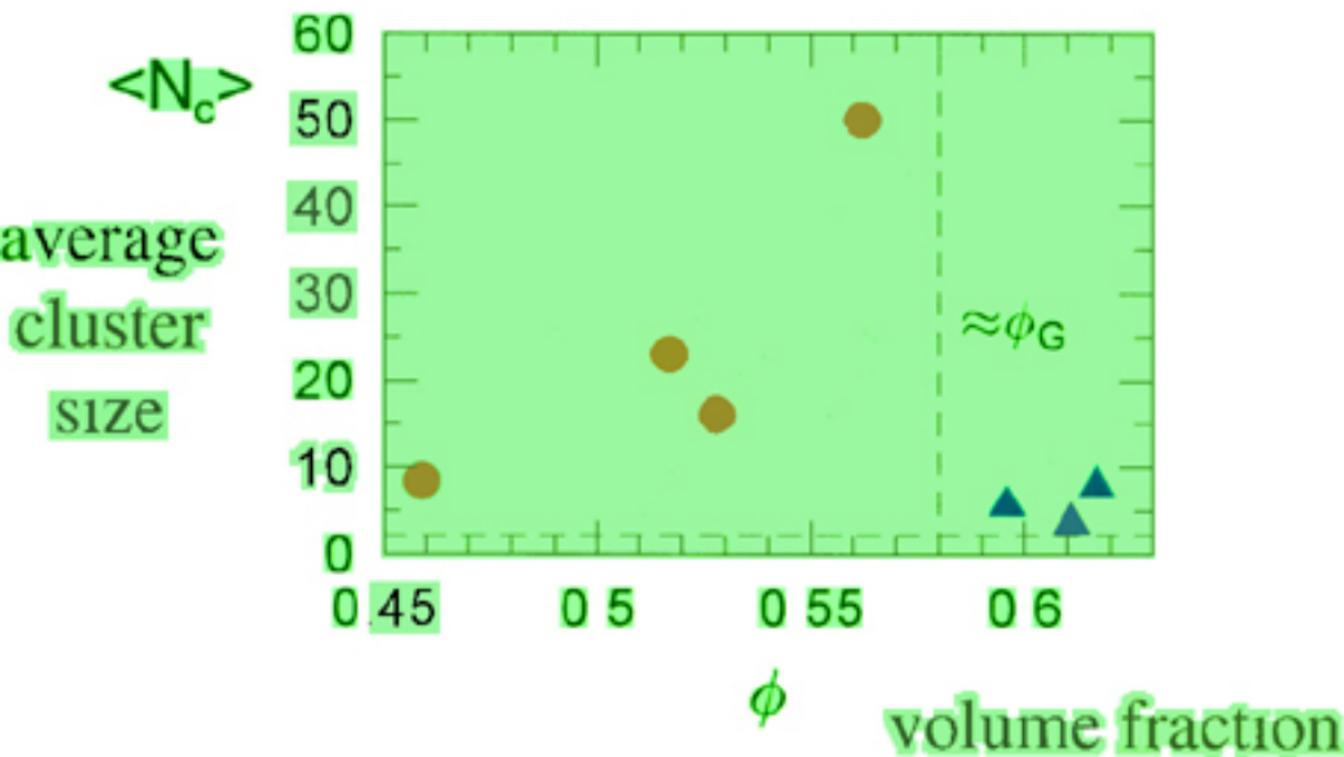


Fractal dimension:

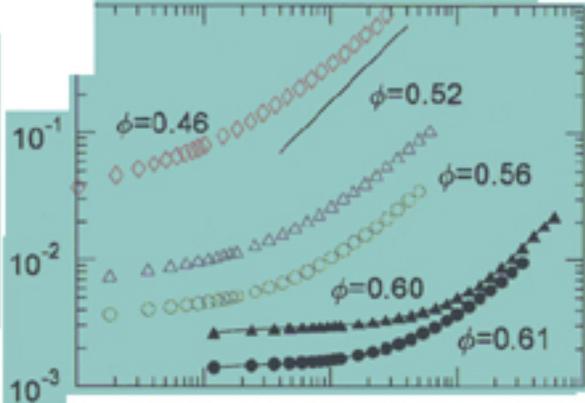


$\phi = 0.56$
supercooled fluid

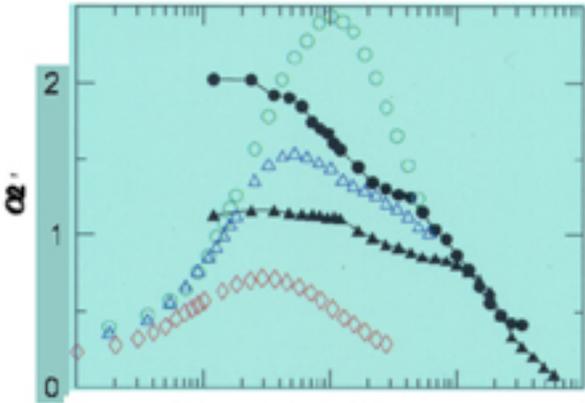
Cluster size grows as glass transition is approached



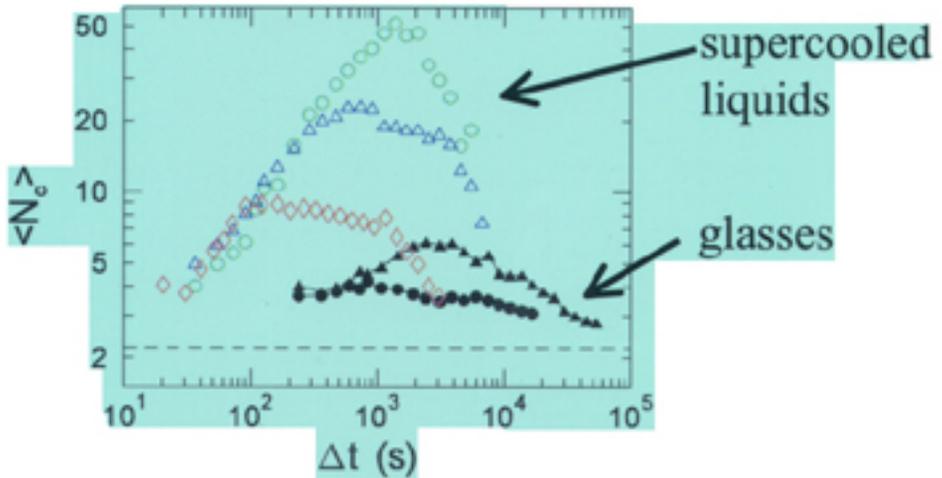
mean square displacement
 $\langle x^2 \rangle$ (μm^2)



nongaussian
parameter



average
cluster
size



Relaxation modes of glasses & supercooled fluids

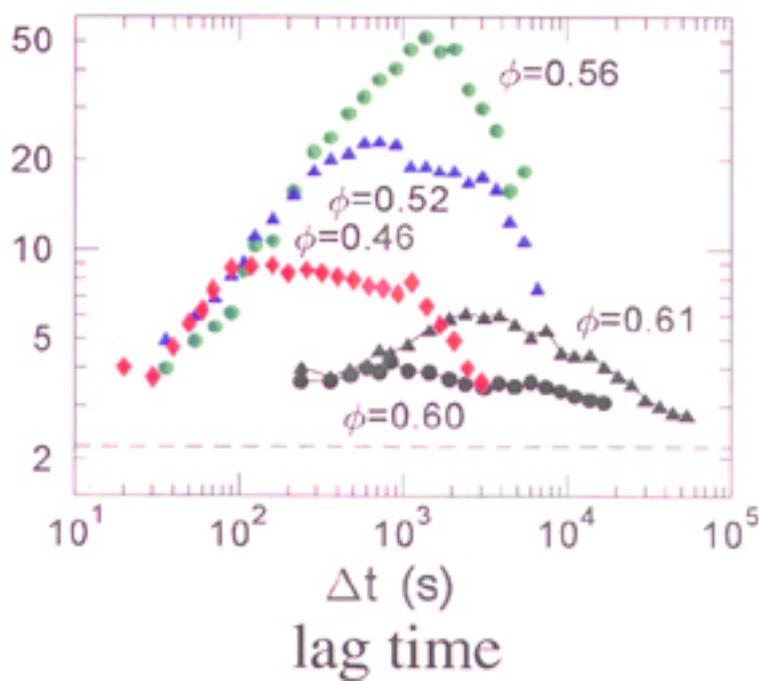
“ α ” and “ β ” relaxation processes:
cage breakdown for fluids
short time scale relaxations,
both fluids & glasses

Our results: cluster structure similar for

- Supercooled fluids, small Δt (“ β ”)
- glasses, all Δt (“ β ”)

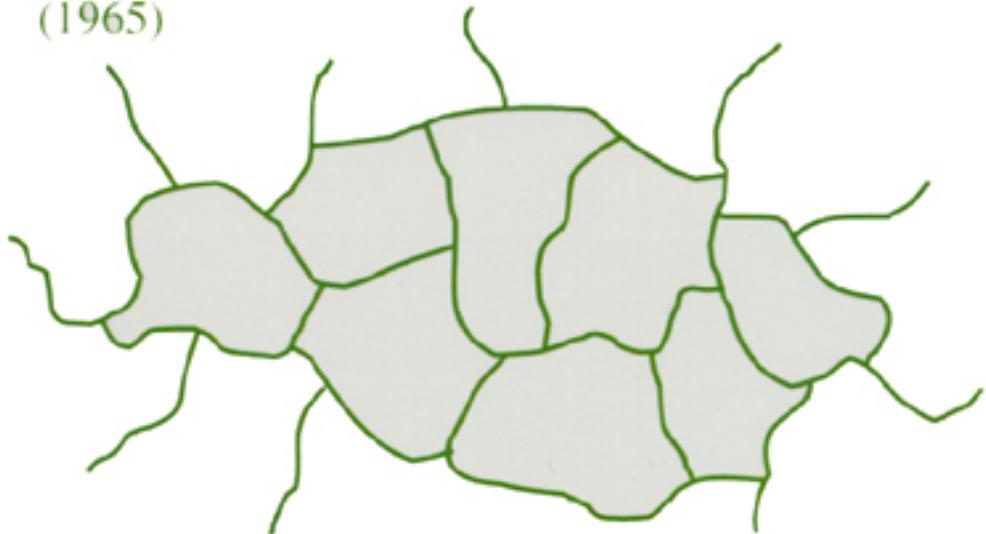
average
cluster
size

$\frac{\Lambda_c}{N_V}$



Dynamical Heterogeneity: possible *dynamic* length scale

Adam & Gibbs: "cooperatively rearranging regions"
(1965)



NMR experiments:

- Schmidt-Röhr & Spiess (1991, polymers)
- Tang, Johnson, et al (1998, metallic glasses)
- Sillecu et al (1992, o-terphenyl)

Photobleaching:

- Cicerone & Ediger (1995, o-terphenyl)

Simulations:

- Glotzer, Kob, Donati, et al (1997, Lennard-Jones)

Molecular Dynamics Simulations:

W. Kob, S. C. Glotzer, C. Donati, P. H. Poole, ...
1997 - now

- Binary mixture of Lennard-Jones particles
- Equilibrium supercooled liquids
- 8000 particles, 3D