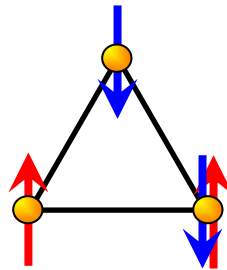




Neutron Scattering in Magnetism

- *focus on dynamics*



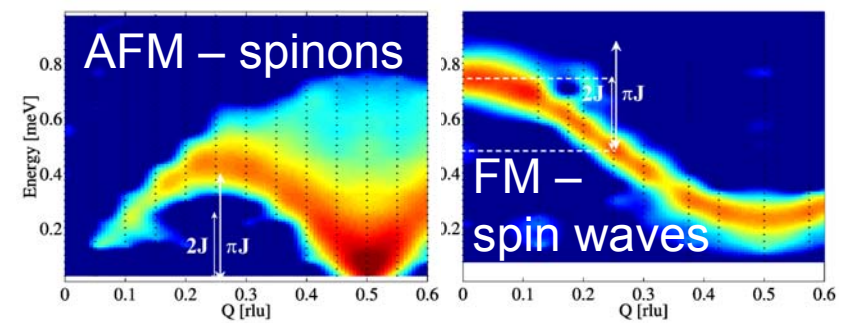
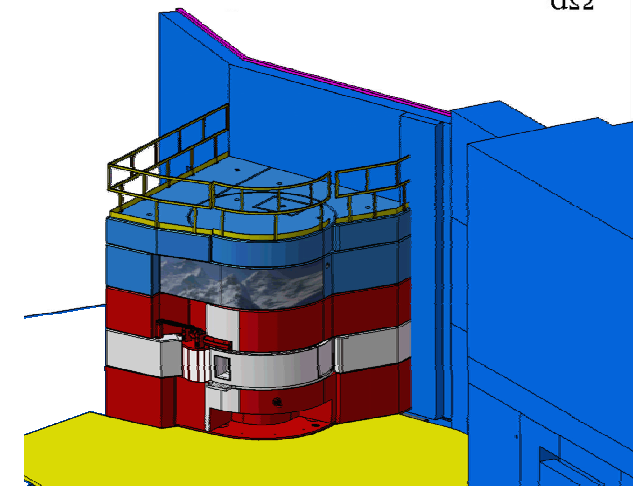
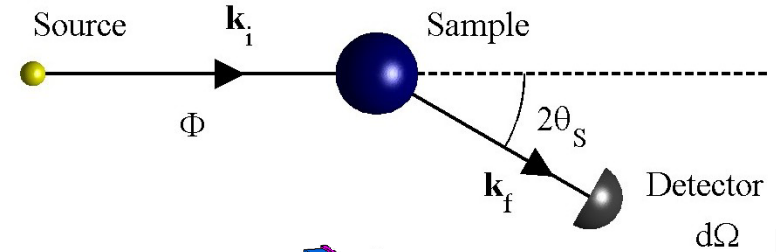
Winter School on Magnetism, Stuttgart 2008

Henrik Moodysson Rønnow
Laboratory for Quantum Magnetism
EPFL Switzerland

- 1) Theory
 - what can we measure
the cross-section

- 2) Experiment
 - how do we measure
instruments and methods

- 3) Physics
 - what do we measure
examples





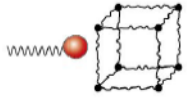
“Theory” – what do we measure



- Plane waves: diffraction and magnetic structures
- Conservation: momentum, energy, spin
- Cross-section
- Magnetic scattering
- Correlation function

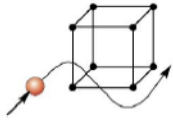


Neutrons



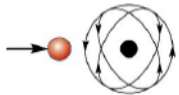
The **ENERGIES** of thermal neutrons are similar to the energies of elementary excitations in solids. Both have similar

- molecular vibrations,
- lattice modes, and
- dynamics of atomic motion.



The **WAVELENGTHS** of neutrons are similar to atomic spacings. They can determine

- structural sensitivity,
- structural information from 10^{-13} to 10^{-4} cm, and
- crystal structures and atomic spacings.



Neutrons “see” **NUCLEI**. They

- are sensitive to light atoms,
- can exploit isotopic substitution, and
- can use contrast variation to differentiate complex molecular structures.

Neutrons are **NEUTRAL** particles. They

- are highly penetrating,
- can be used as nondestructive probes, and
- can be used to study samples in severe environments.

Neutrons have a **MAGNETIC** moment. They can be used to

- study microscopic magnetic structure,
- study magnetic fluctuations, and
- develop magnetic materials.



Neutrons have **SPIN**. They can be

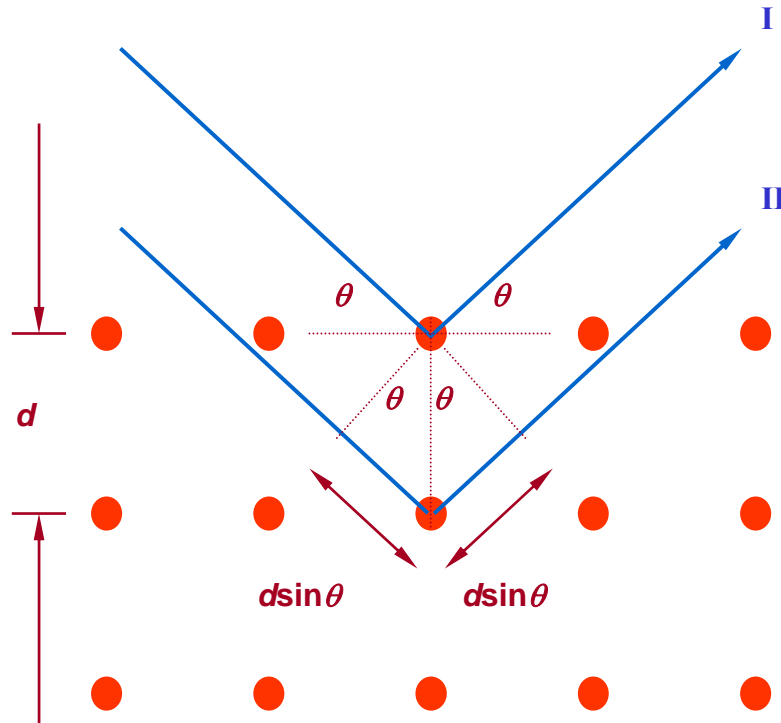
- formed into polarized neutron beams,
- used to study nuclear (atomic) orientation, and
- used for coherent and incoherent scattering.



Neutron Plane Waves



- Plane waves and Bragg's law



$$n\lambda = 2d \sin\theta$$

- Bragg's law:

$$n\lambda = 2d \sin\Theta$$

- Reciprocal lattice

$$\tau = ha^* + kb^* + lc^*$$

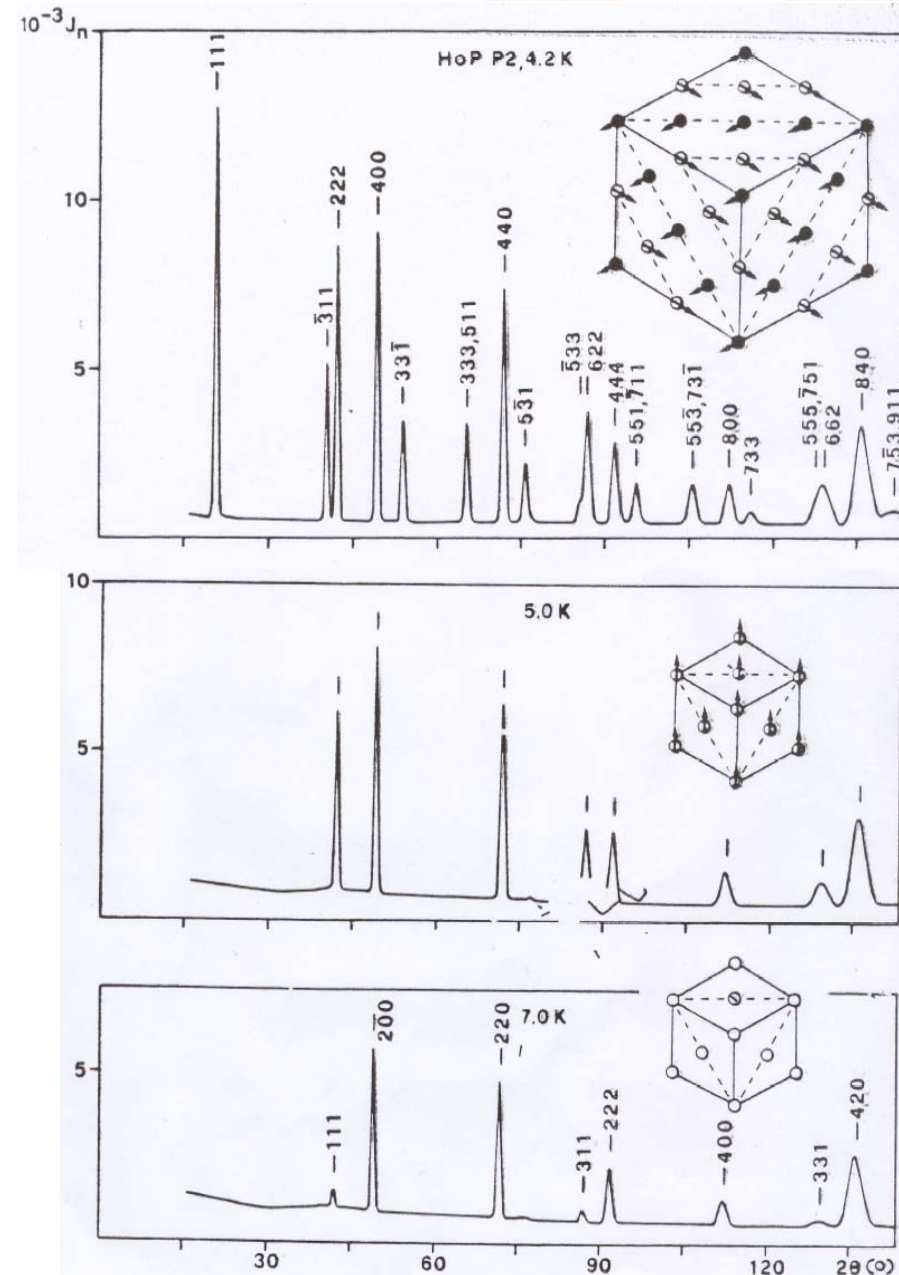
$$\tau = 2k \sin\Theta$$

- Powder diffraction – example: HoP

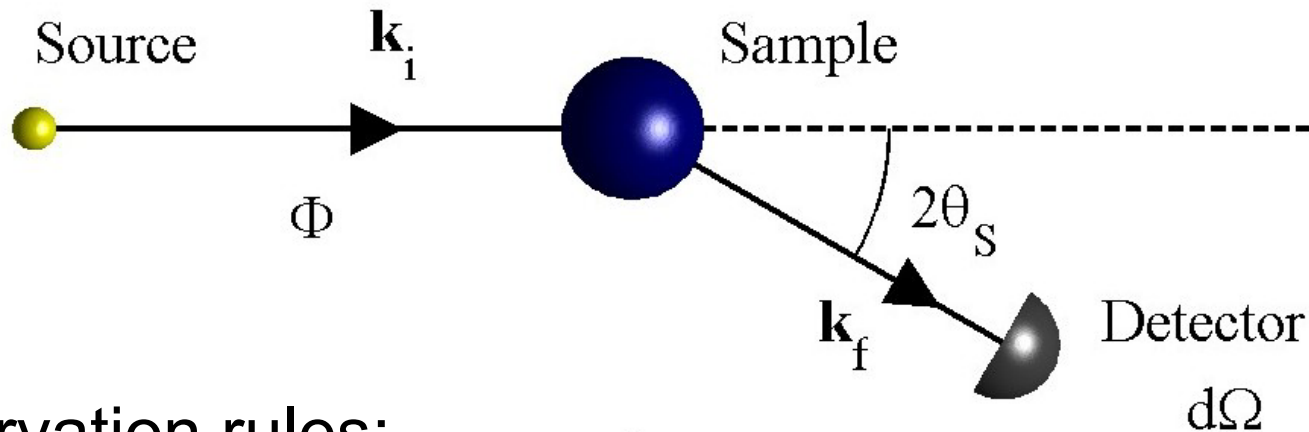
- Structural peaks at 7K
- Ferromagnet at 5K
- Non-collinear at 4.2K

- Structure refinement:

- Standard 'black box' methods
- Absorption, multiple scattering, texture etc.
- Positions, displacements, moments, directions



Consider scattering neutrons on a sample into a detector.



Conservation rules:

	Sample	=	Neutron	
Momentum	$\hbar\mathbf{Q}$	=	$\hbar\mathbf{k}_i - \hbar\mathbf{k}_f$	
Energy	$\hbar\omega$	=	$E_i - E_f$	$= \hbar(k_i^2 - k_f^2)/2m_n$
Spin	ΔS	=	$\sigma_i - \sigma_f$	

⇒ We can control and measure these quantities !



Cross-section



Intensity in angular $d\Omega$ and energy dE_f elements:

$$I_{d\Omega dE_f} = \left(\frac{d^2\sigma}{d\Omega dE_f} \right) \Phi d\Omega dE_f$$

From initial state i to final state f of neutron \mathbf{k} and sample λ

$$\left(\frac{d^2\sigma}{d\Omega dE_f} \right)_{\lambda_i \rightarrow \lambda_f} = \frac{k_f}{k_i} \left(\frac{m_n}{2\pi\hbar^2} \right)^2 |\langle \mathbf{k}_f \lambda_f | V | \mathbf{k}_i \lambda_i \rangle|^2 \delta(E_{\lambda_i} - E_{\lambda_f} + \hbar\omega)$$

Neutrons treated as plane waves:

$$|\mathbf{k}\mathbf{s}_n\rangle = V^{-1/2} \exp(i\mathbf{k} \cdot \mathbf{r}_n) |\mathbf{s}_n\rangle$$

Energy conservation \Rightarrow integral rep.:

$$\delta(\hbar\omega + E_i - E_f) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} e^{i(\hbar\omega + E_i - E_f)t/\hbar} dt$$

Fourier transform in
- space/momentum
- time/energy

$$|\langle \mathbf{k}_f \lambda_f | V | \mathbf{k}_i \lambda_i \rangle|^2$$

Dipole interaction – electron spin and orbit

$$V_{\text{mag}}(\mathbf{r}) = -\frac{\mu_0}{4\pi} 2\gamma\mu_N\mu_B \boldsymbol{\sigma}_n \cdot \left(\nabla \times \left(\frac{\mathbf{s} \times \hat{\mathbf{R}}}{|\mathbf{R}|^2} \right) + \frac{1}{\hbar} \frac{\mathbf{p} \times \hat{\mathbf{R}}}{|\mathbf{R}|^2} \right)$$

Non-polarised neutron average over spin states:

$$\left(\frac{d^2\sigma}{d\Omega dE_f} \right)_{\text{mag}} = \frac{(\gamma r_0)^2 k_f}{2\pi\hbar k_i} \sum_{\alpha\beta} (\delta_{\alpha\beta} - \hat{Q}_\alpha \hat{Q}_\beta) \sum_{l'd'ld} \frac{1}{4} g_{d'} g_d F_{d'}^*(\mathbf{Q}) F_d(\mathbf{Q}) \int_{-\infty}^{\infty} dt \exp(-i\omega t) \langle \exp\{-i\mathbf{Q} \cdot \mathbf{R}_{l'd'}(0)\} \exp\{i\mathbf{Q} \cdot \mathbf{R}_{ld}(t)\} \rangle \langle S_{l'd'}^\alpha(0) S_{ld}^\beta(t) \rangle$$

fudge factor dipole factor magnetic form factor spin-spin correlation function
Fourier transform

Polarised neutrons:

$$\left(\frac{d^2\sigma}{d\Omega dE_f} \right)_{s_i \rightarrow s_f} = \frac{k_f}{k_i} \sum_{if} P_{\lambda_i} |\langle \lambda_f | \sum_l \exp(i\mathbf{Q} \cdot \mathbf{r}_l) U_l^{s_i s_f} | \lambda_i \rangle|^2 \delta(E_{\lambda_i} - E_{\lambda_f} + \hbar\omega)$$

$$U_l^{s_i s_f} = \langle s_f | b_l - p_l \boldsymbol{\sigma}_n \cdot \mathbf{S}_l \perp + B_l \boldsymbol{\sigma}_n \cdot \mathbf{I}_l | s_i \rangle$$



Spin-spin correlation function

$$S^{\alpha\beta}(\mathbf{Q}, \omega) = \frac{1}{2\pi} \sum_{l'd'ld} \exp(-i\mathbf{Q} \cdot \{l' + d' - l - d\}) \int_{-\infty}^{\infty} dt \exp(-i\omega t) \langle S_{l'd'}^{\alpha}(0) S_{ld}^{\beta}(t) \rangle$$

Dynamic structure factor

$$\left(\frac{d^2\sigma}{d\Omega dE_f} \right)_{\text{mag}} = \frac{1}{\hbar} \frac{k_f}{k_i} p^2 \exp(-2W) \sum_{\alpha\beta} (\delta_{\alpha\beta} - \hat{Q}_{\alpha} \hat{Q}_{\beta}) S^{\alpha\beta}(\mathbf{Q}, \omega)$$

Fluctuation dissipation theorem \Rightarrow gen. susceptibility

$$S(\mathbf{Q}, \omega) = [n(\omega) + 1] \chi''(\mathbf{Q}, \omega) = \frac{\chi''(\mathbf{Q}, \omega)}{1 - \exp(-\hbar\omega/k_B T)}$$

intrinsic dynamics



response to perturbation

Theory !

- Dynamic structure factor: inelastic

$$S(\mathbf{Q}, \omega) \propto \int_{-\infty}^{\infty} dt e^{-i\omega(t-t')} \langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(t') \rangle$$

– periodic: $\sin(\omega_0 t) \Rightarrow$ peak: $\delta(\omega_0 - \omega)$

- Static structure factor: elastic

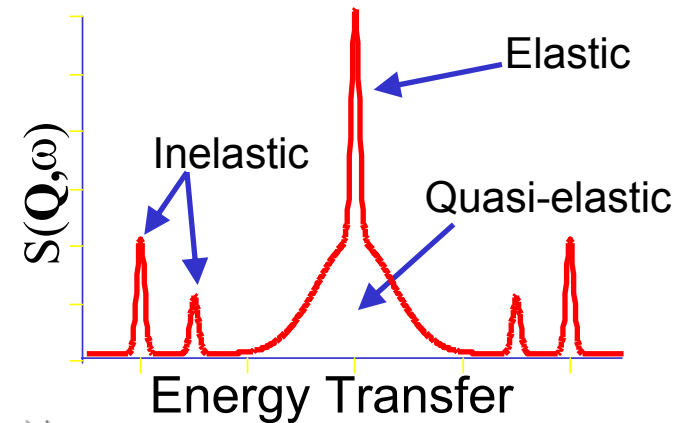
$$S(\mathbf{Q}, \omega = 0) \propto \int_{-\infty}^{\infty} dt \langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(t') \rangle \simeq \langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(\infty) \rangle$$

– Bragg peaks at $\omega = 0$

- Instantaneous structure factor

$$S(\mathbf{Q}) = \int d\omega S(\mathbf{Q}, \omega) \propto \int_{-\infty}^{\infty} dt \delta(t - t') \langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(t') \rangle = \langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(t) \rangle$$

– Exponential decay: $\exp(-t/\tau) \Rightarrow$ Lorentzian: $1/(1+\omega^2\tau^2)$





Recap: what do we measure



Conservation rules \Rightarrow momentum, energy, spin
dependence

Intensity \Rightarrow Cross-section, magnetic dipole interaction

Correlation function \Rightarrow dynamic structure factor:

$$\text{Intensity} \propto \left(\frac{d^2\sigma}{d\Omega dE_f} \right)_{\text{mag}} = \frac{1}{\hbar} \frac{k_f}{k_i} p^2 \exp(-2W) \sum_{\alpha\beta} (\delta_{\alpha\beta} - \hat{Q}_\alpha \hat{Q}_\beta) S^{\alpha\beta}(\mathbf{Q}, \omega)$$



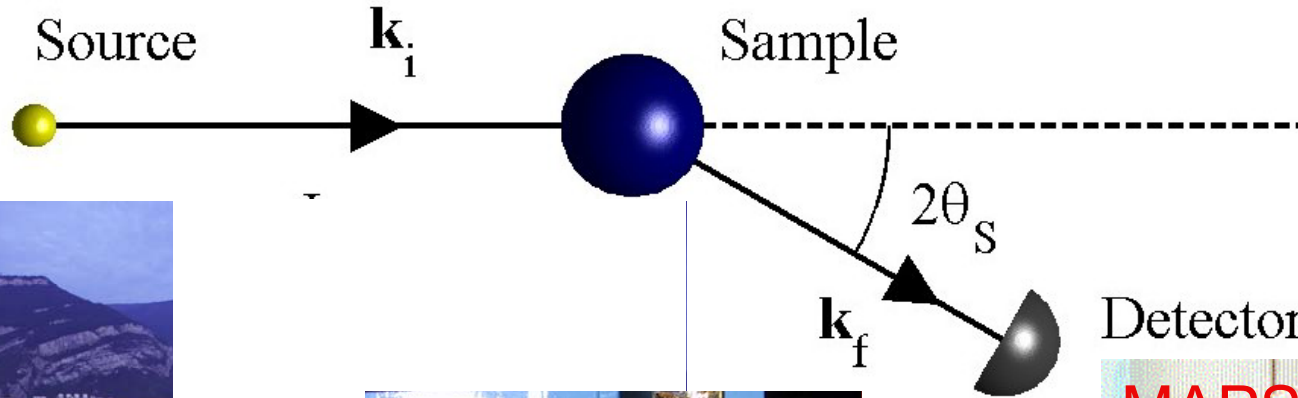
Experiment



Theory



Experiment – how do we measure *instruments and methods*



ILL



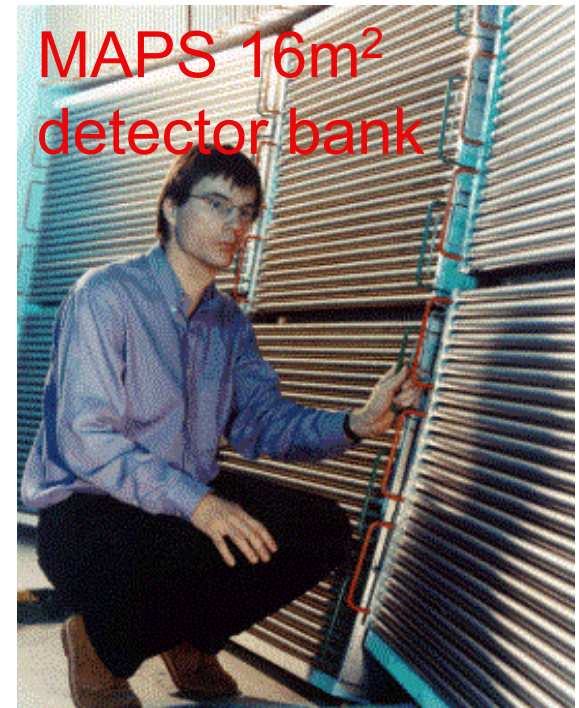
SNS



CuSO4 \cdot 5D2O

Christian Rüegg

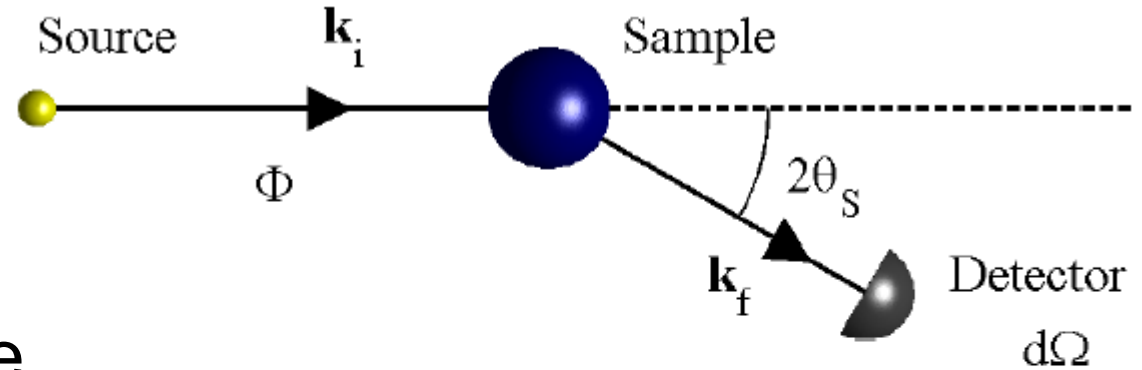
Detector



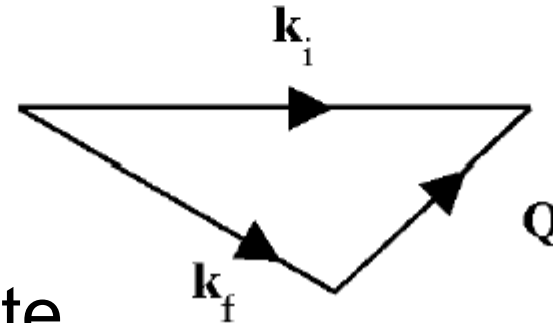
MAPS 16m²
detector bank

Chris Frost

- Need to define \mathbf{k}_i , \mathbf{k}_f and angle between them relative to sample



- Define $E_i \propto |\mathbf{k}_i|^2 \Rightarrow$ monochromate
- Know angle 2θ between k_i and k_f relative to sample orientation
- Determine $E_f \propto |\mathbf{k}_f|^2 \Rightarrow$ analyse

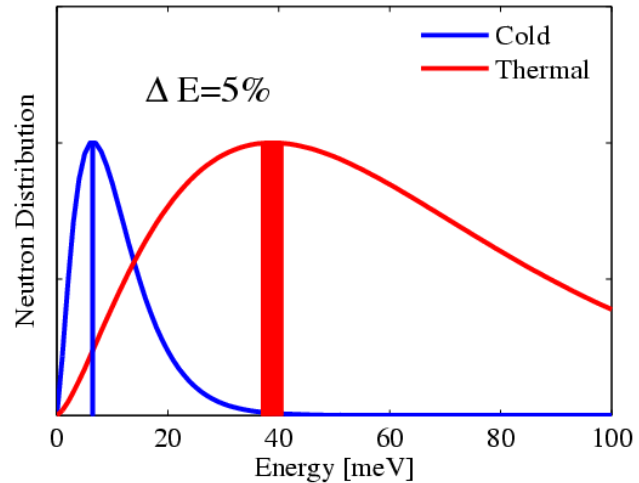




The neutron energy



- We cannot (controlled) change the neutron energy



Must select neutrons with wanted energy from a (Maxwellian) distribution

(Liouville & the phase space transformer)

- We have no energy sensitive neutron detector... yet!
 $n + {}^3\text{He} \rightarrow {}^3\text{H} + {}^1\text{H} + 0.764 \text{ MeV}$ – want **meV** need 10^{-9} precision!
- Make another selection (analyse) to know energy
- Or, if we know starting time we can use time-of-flight



Useful numbers



- Energy units:

eU 1 meV =

$k_B T$ 11.6 K =

$h\nu$ 0.24 THz =

hc/λ 8.06 cm⁻¹ =

$\mu_B H$ 17.3 T =

E 1.6x10⁻²² J

- Neutron Energies:

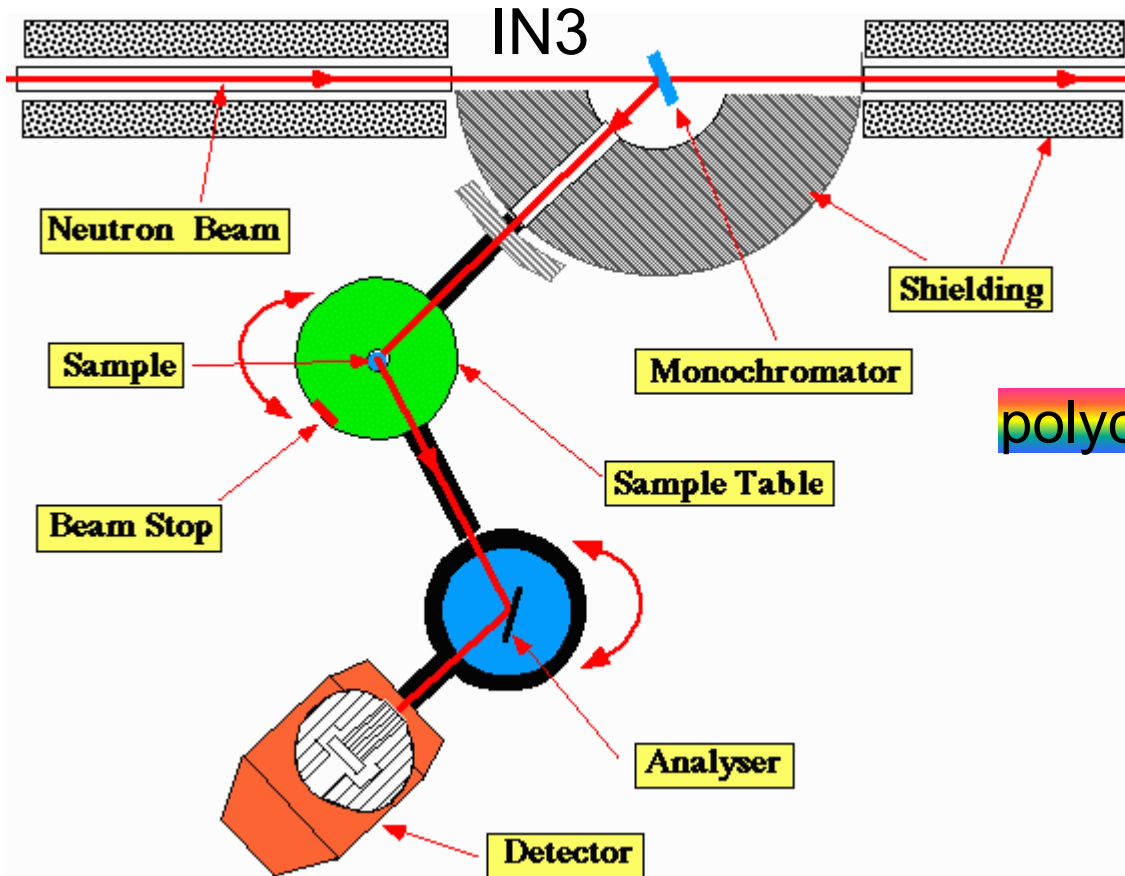
E [meV] =

2.072 k² [meV Å²] =

81.8 / λ^2 [meV Å²]



TAS: three axis spectrometer



Bragg's law

$$n\lambda = 2d \sin\theta$$

$$n\tau = 2k \sin\theta$$



$k, 2k, \dots$

Higher order ($2k$) filters:

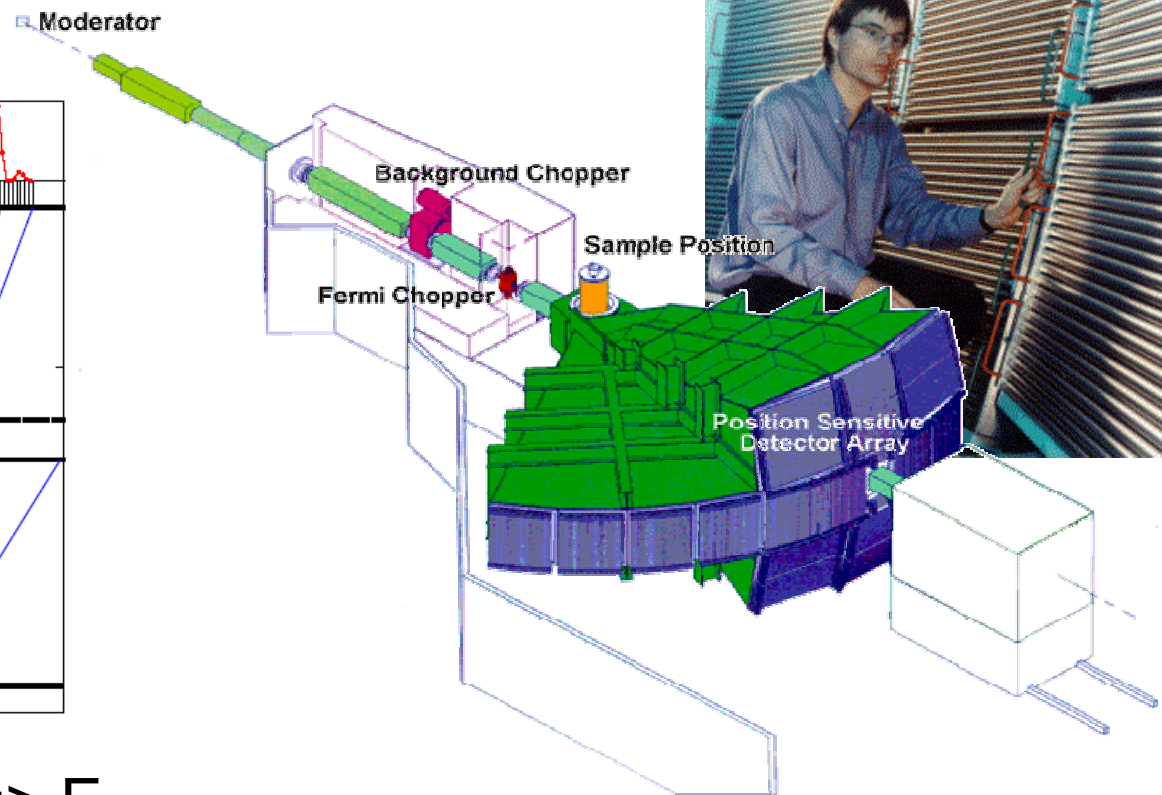
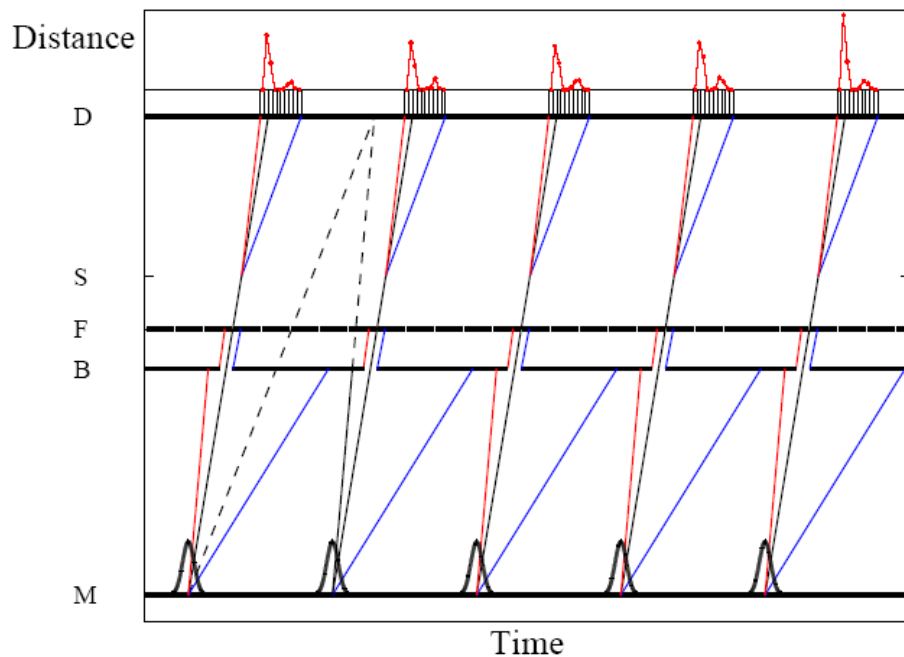
Be: <5 meV

BeO: <3.8 meV

PG: 14.7 & 35 meV

velocity selector: 'any'

Pulsed beam
known start time



Monochromatising chopper => E_i

Time-of-flight => E_f

Advantage: nothing between sample and detector => easy multiplexing



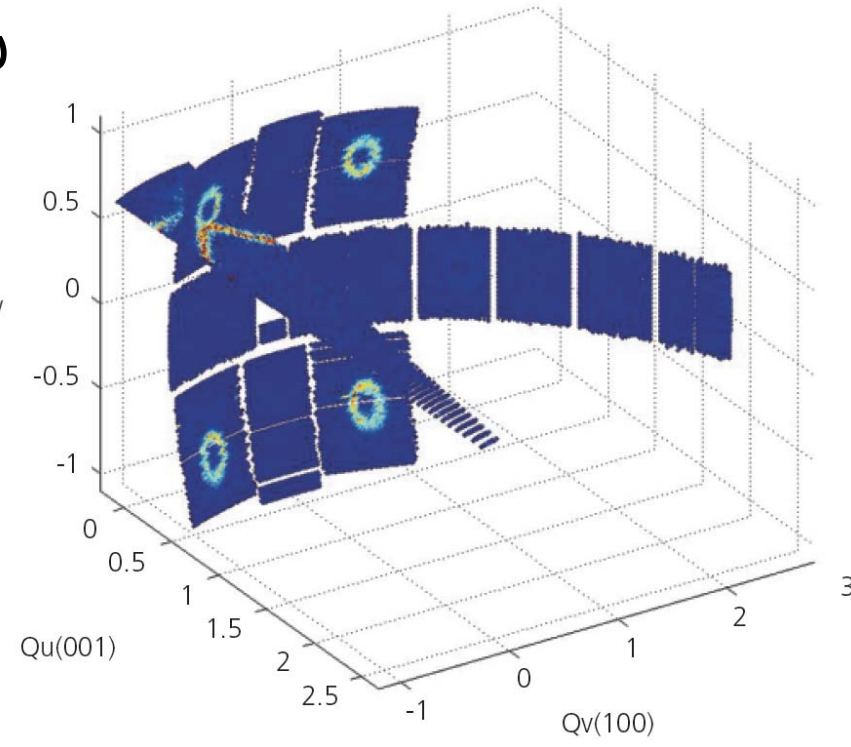
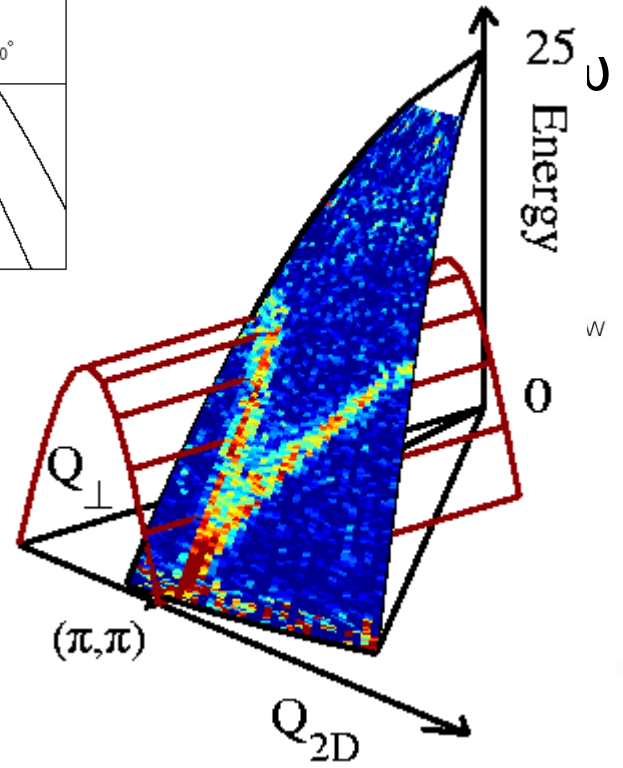
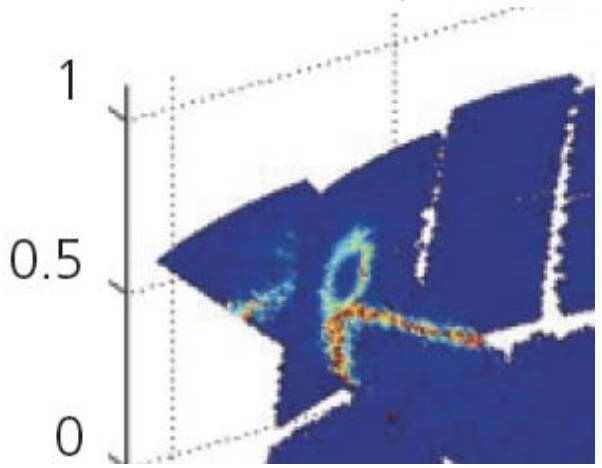
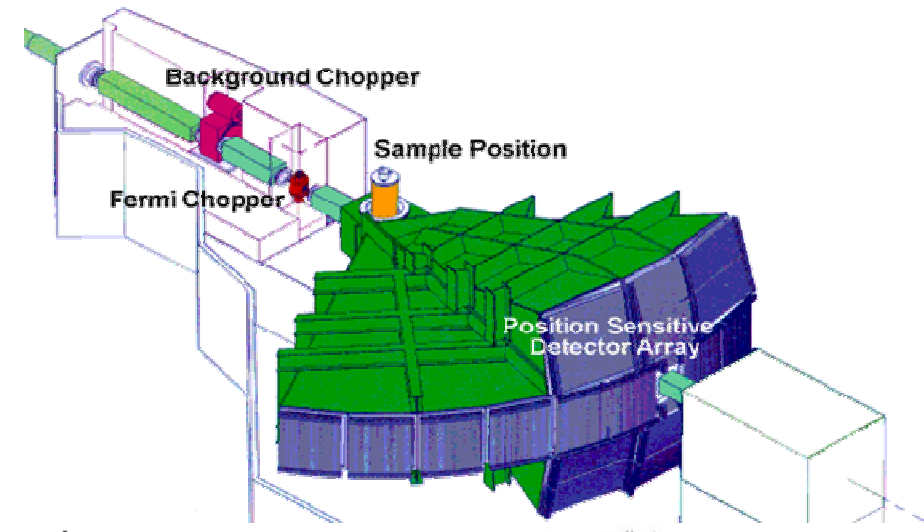
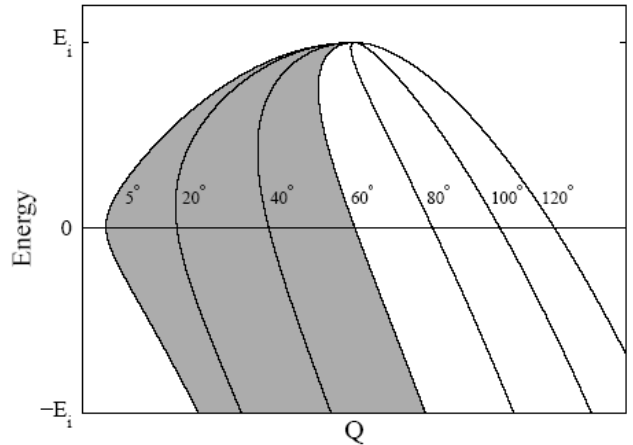
TOF: Q-E parabolas and maps



- Q for fixed E_i and θ :

$$Q^2 = k_i^2 + k_f^2 - 2|k_i||k_f|\cos 2\theta$$

$$= 2E_i - \hbar\omega - 2(E_i^2 - E_i\hbar\omega)^{1/2} \cos 2\theta$$



- Pulsed sources:
MAPS, Merlin, LET...
 $E_i = 20\text{meV}$ to 2eV

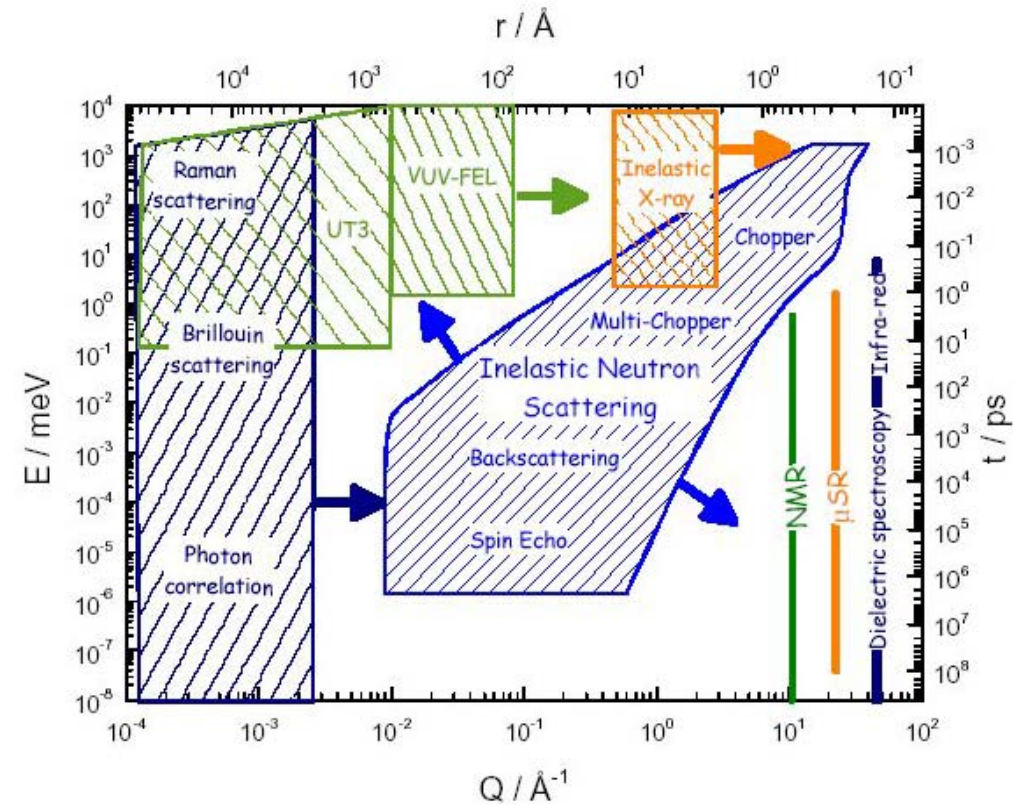
- Continuous sources:
chopper define pulse
 $\Delta E \sim \mu\text{eV}$ to meV

- Sample: 5-50g



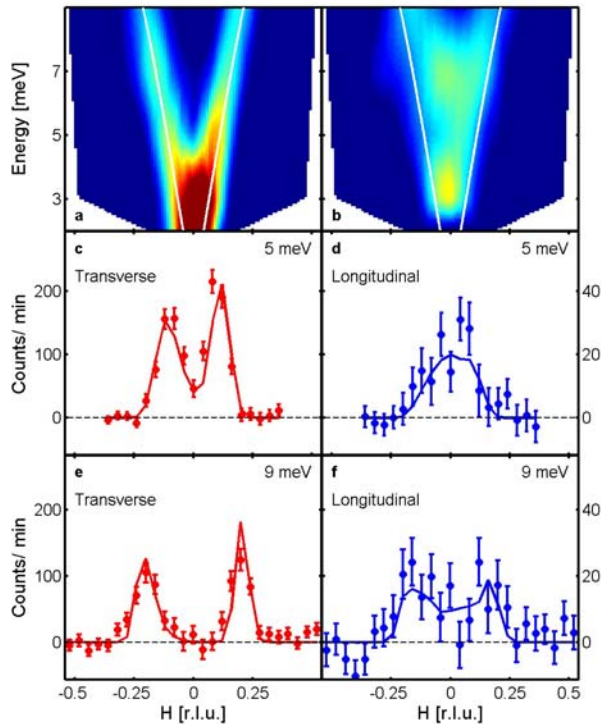
- Counting: few days!
 - Need $\times 10$ - 100 for T-dependence
 - Need $\times 100$ - 1000 for P-dep.

Time-Distance Scales Probed



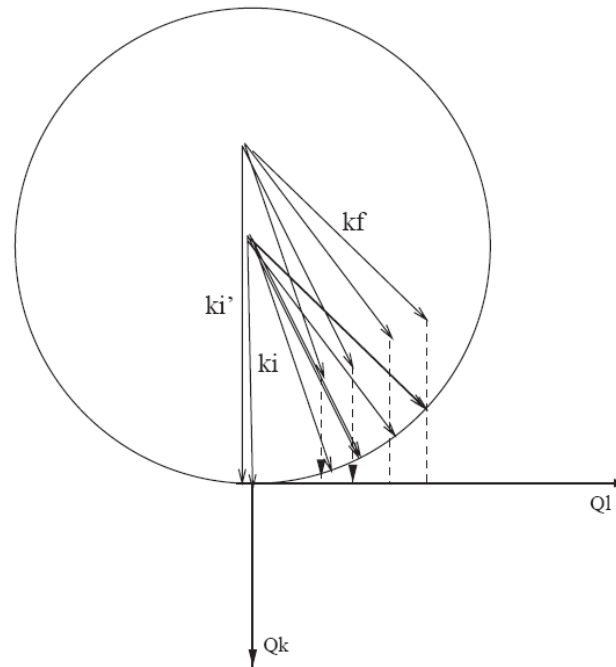
TAS

- Focus on one Point
- Flexible
- Polarisation analysis
- Already “optimal”



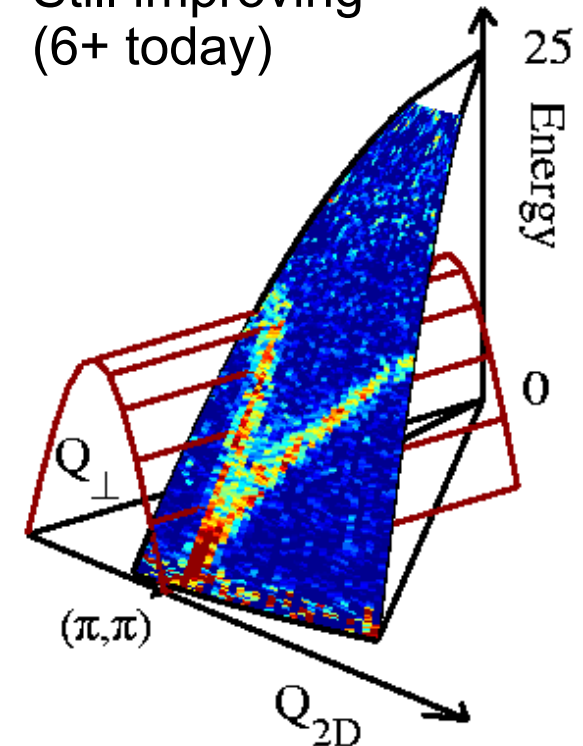
Multi-TAS

- a line in momentum-energy-space
- More Neutrons recorded than TAS
- More flexible than TOF



TOF

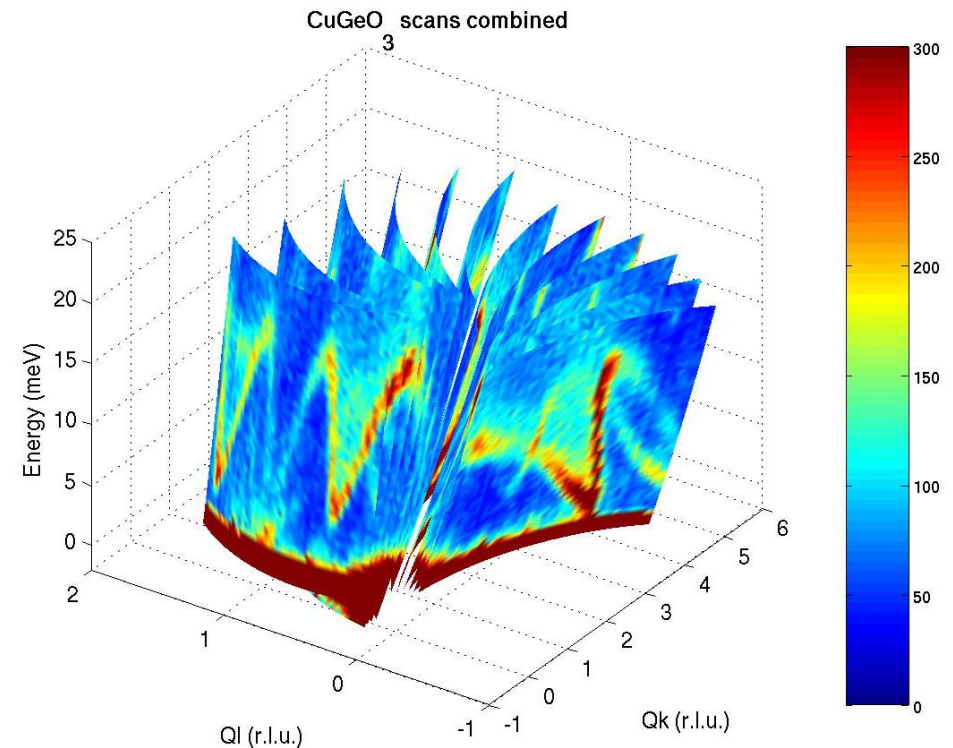
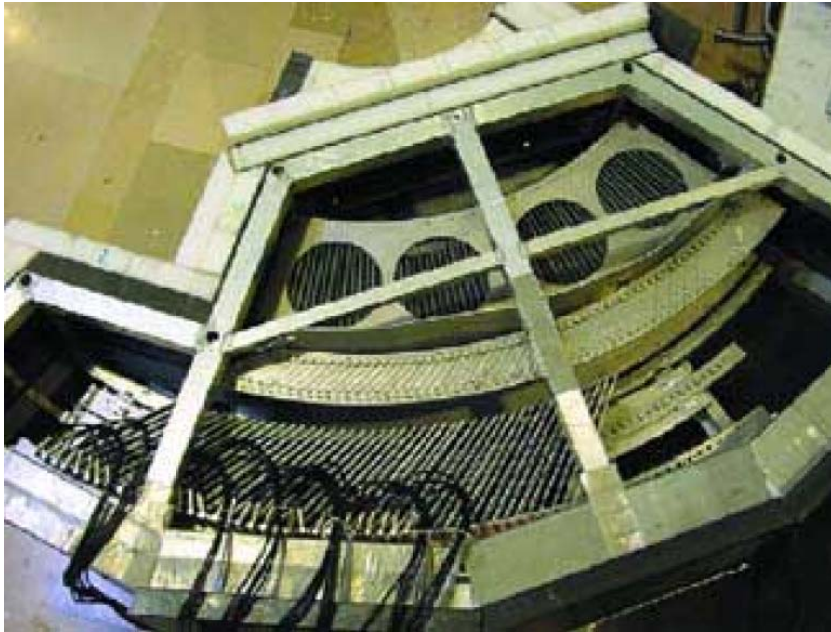
- 2-3D manifold
- Overview – sees “everything”
- Less flexible
- Still improving (6+ today)



Multi (47) analyser-detector

- Analysers Cu(200) $E_f=31$ meV
- Detectors ^3He , 0.3° each 1°
- IN8, 7 samples in 1 week

(Jimenez-Ruiz, Demmel & 7 students)





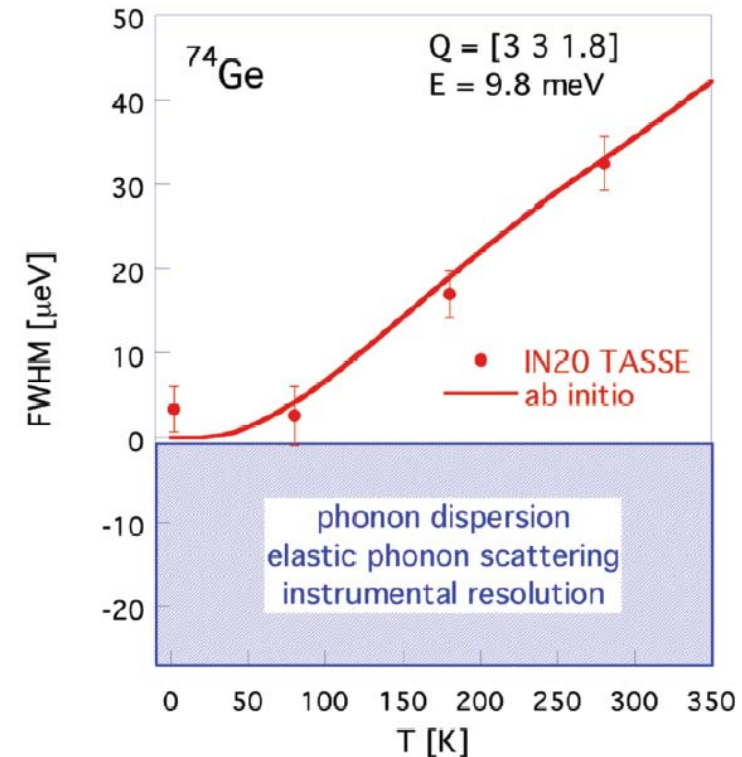
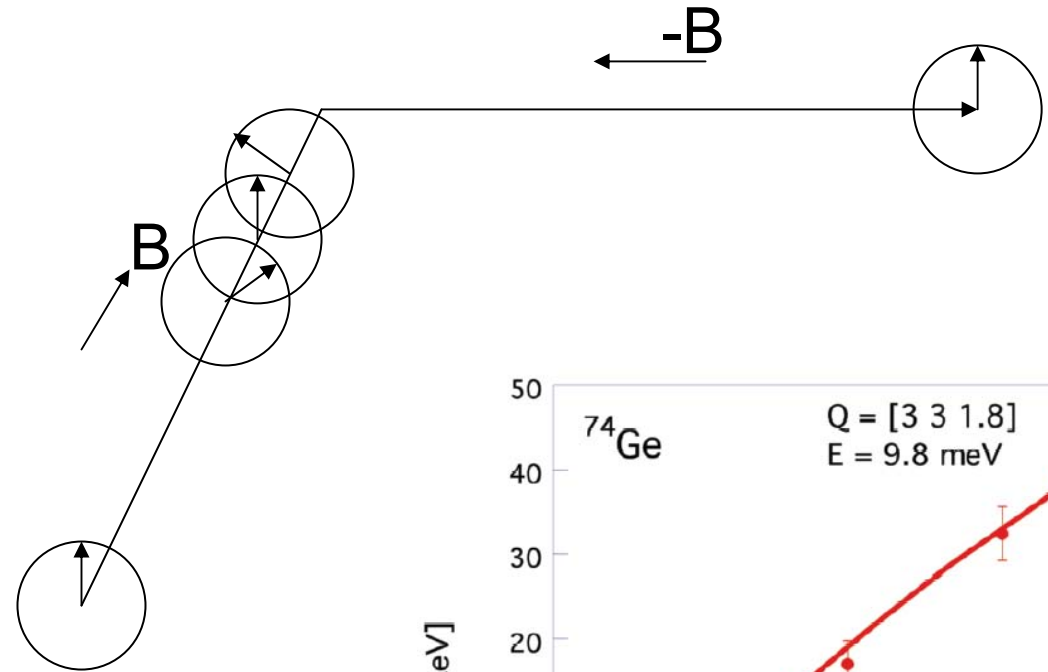
Recap.: How do we measure



- Control k_i and k_f
- Three axis spectrometers (TAS)
- Time-of-flight spectrometers (TOF)
- Energies from μeV to eV
- New sources in recent years!
- New spectrometers emerging!
- Many other neutron techniques in magnetism



- Indirect TOF
- Backscattering:
 - better resolution
- Spin-Echo:
 - ‘count’ precessions
 - combined with TAS
 - μeV resolution



- Powder & single crystal diffraction
- Reflectivity (surfaces, films and multilayers)
- Small angle scattering (large objects: domains, magnetic vortices in superconductors etc.)
- ...

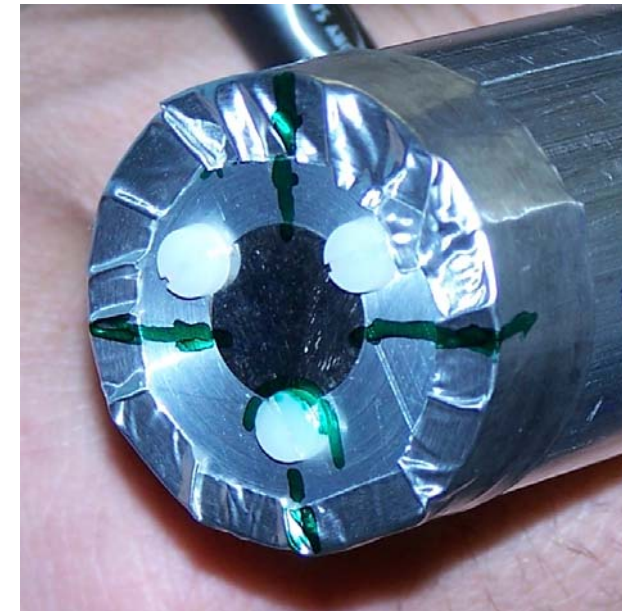


Other techniques: Imaging



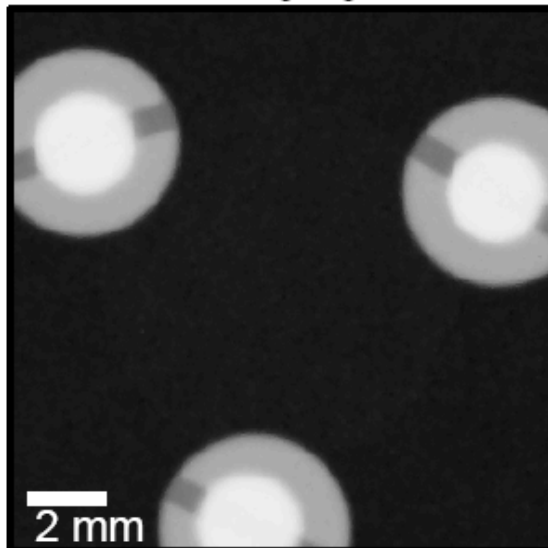
Coherent neutron phase imaging

- FM domains in FeSi_{3%} disk
- Promise of 3D domain imaging!



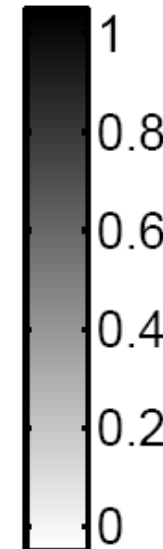
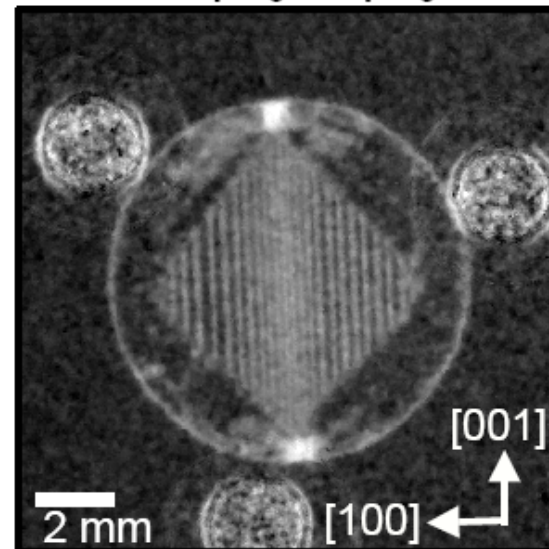
Absorption

$$a_0^s / a_0^r$$

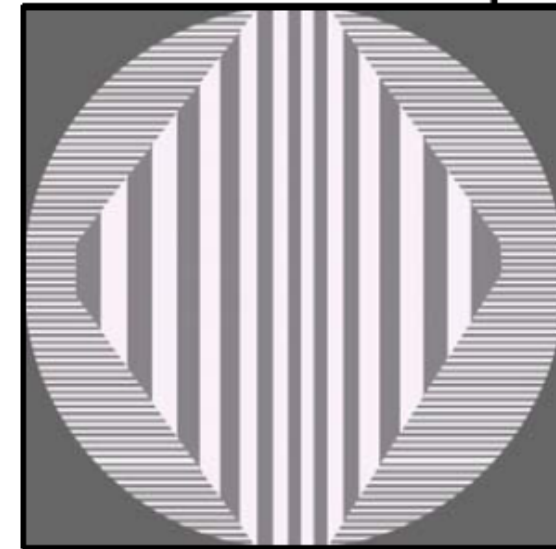


Phase contrast

$$a_1^s \cdot a_0^r / a_1^r \cdot a_0^s$$



schematic map





Physics – what do we measure

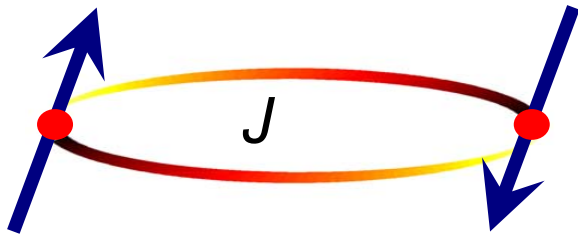
examples



~~• Spin-waves, Holstein-Primakof transformation ...~~

• A pair of spins:

$$\vec{H} = J \sum \mathbf{S}_i \cdot \mathbf{S}_j$$

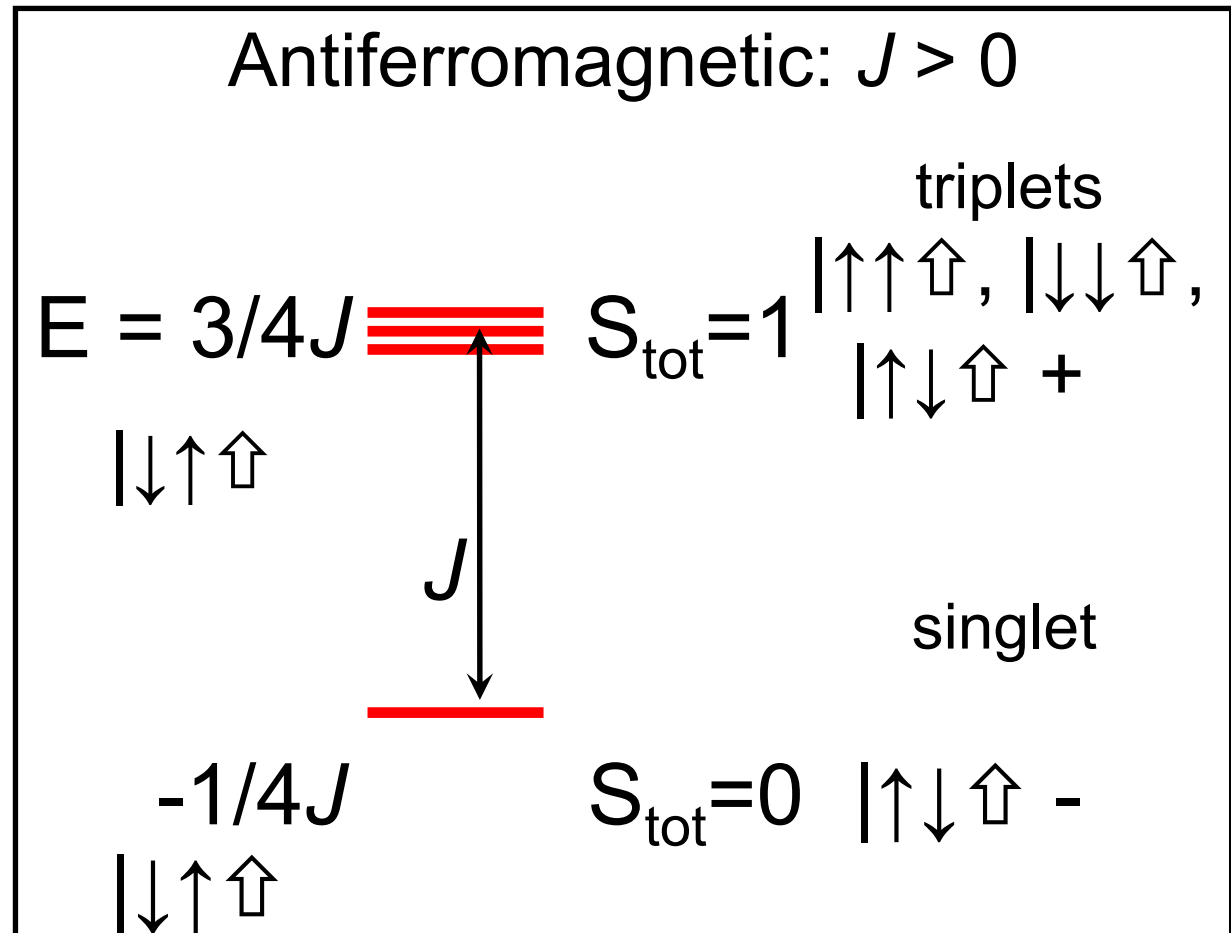


Ferromagnet: $J < 0$

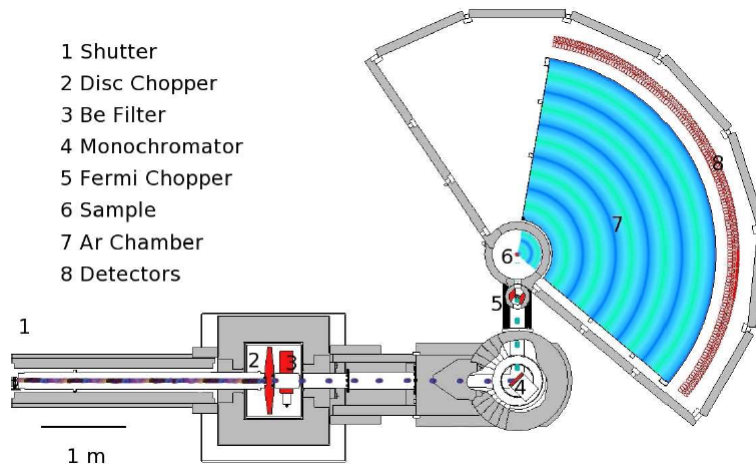
|GS \uparrow = | $\uparrow\uparrow\uparrow\uparrow$ or

| $\downarrow\downarrow\downarrow\downarrow$

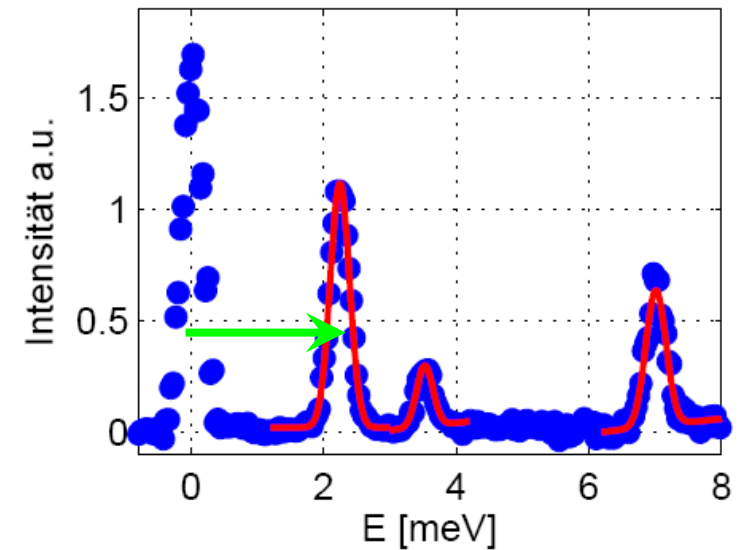
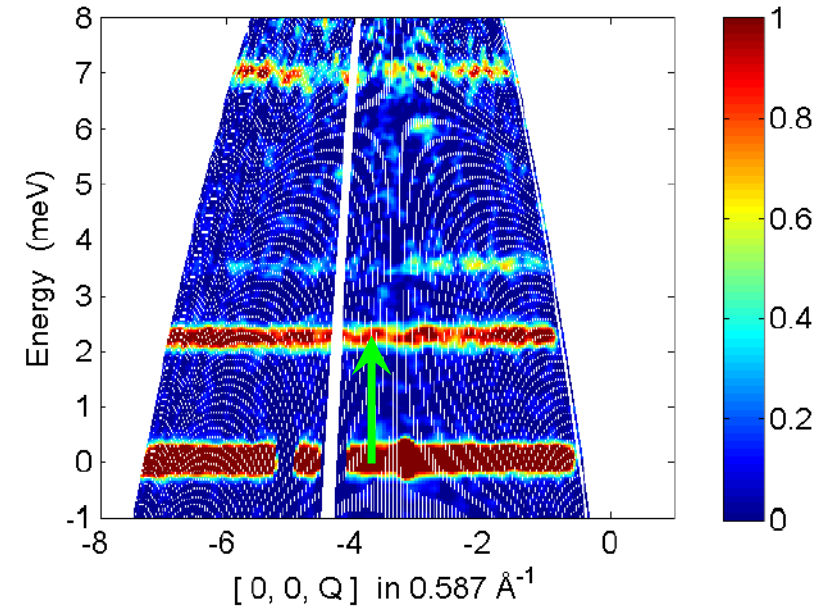
“Classical”



- Transition from ground to excited state



- TOF powder spectrum
(Crystal fields in LiErF_4)



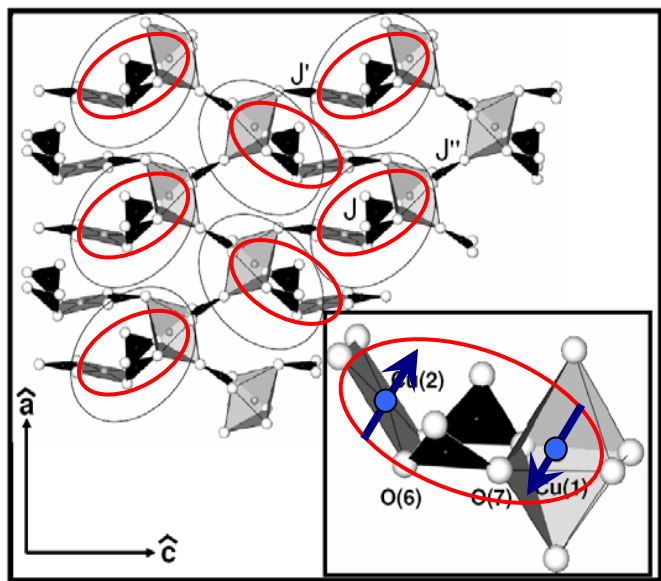


Q-dependence – Structure factor

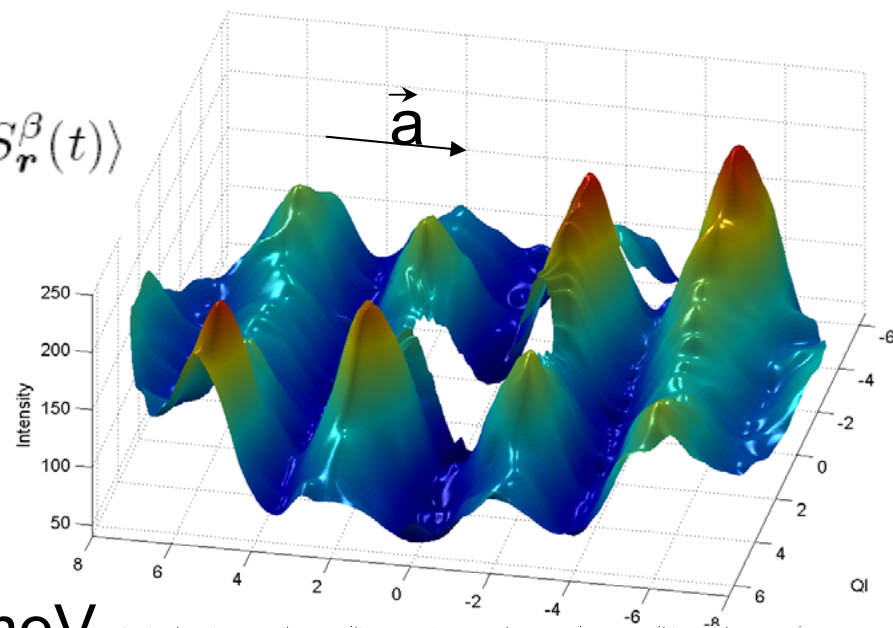


$$S^{\alpha\beta}(\mathbf{Q}, \omega) = \frac{1}{2\pi} \sum_{\mathbf{r}' \neq \mathbf{r}} e^{-i\mathbf{Q} \cdot (\mathbf{r}' - \mathbf{r})} \int_{-\infty}^{\infty} dt e^{-i\omega t} \langle S_{\mathbf{r}'}^{\alpha} S_{\mathbf{r}}^{\beta}(t) \rangle$$

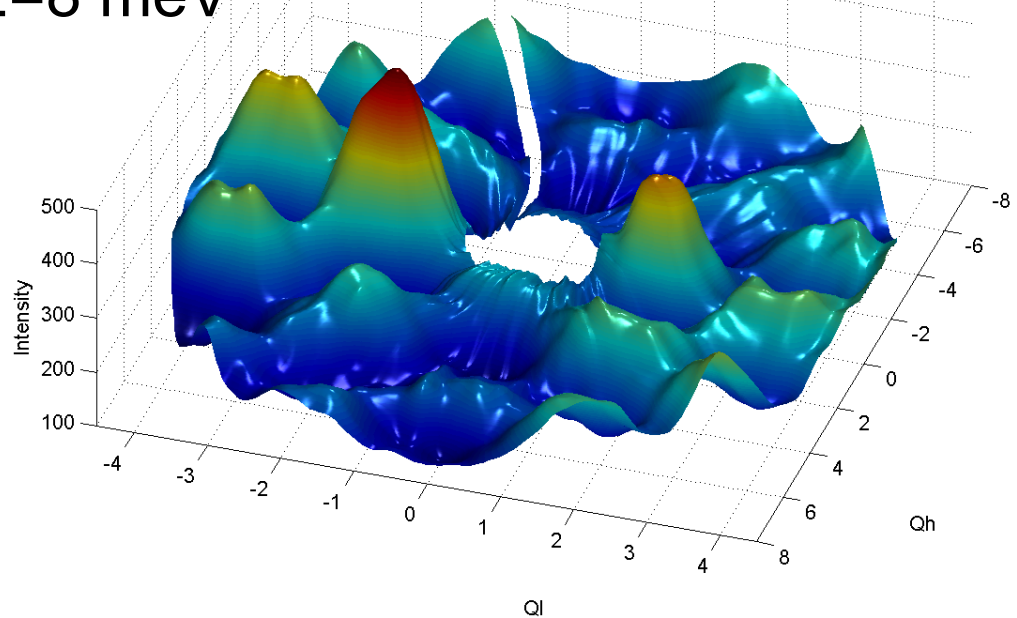
$$S(\vec{q}) = \frac{1}{2} (1 - \cos(\vec{q} \cdot \vec{a}))$$



Ba₂Cu(BO₃)₂
Spectrometer IN8+MAD



E = 8 meV

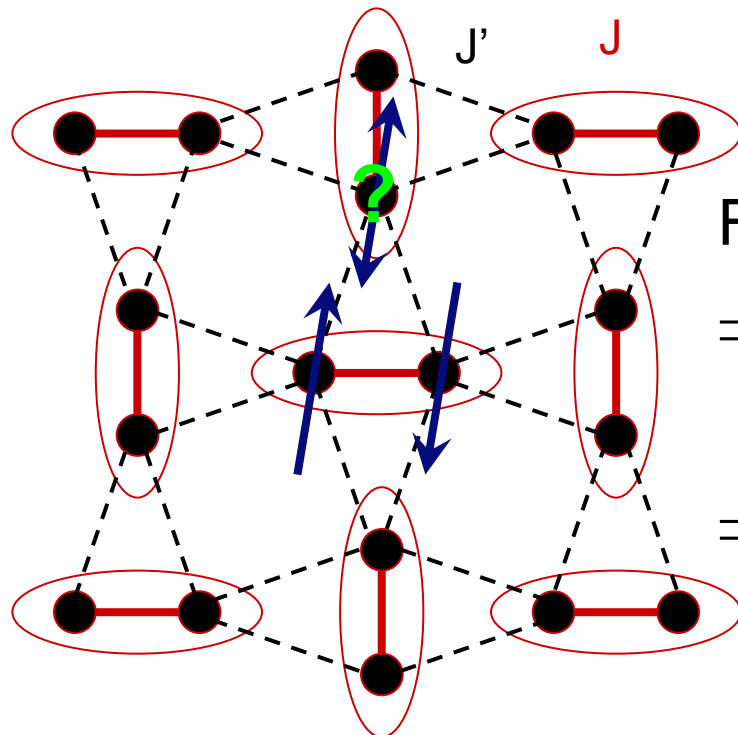
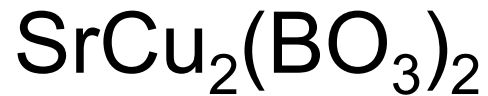




Isolated localised excitations



Localised excitation \Rightarrow Q-indep. Energy spectrum

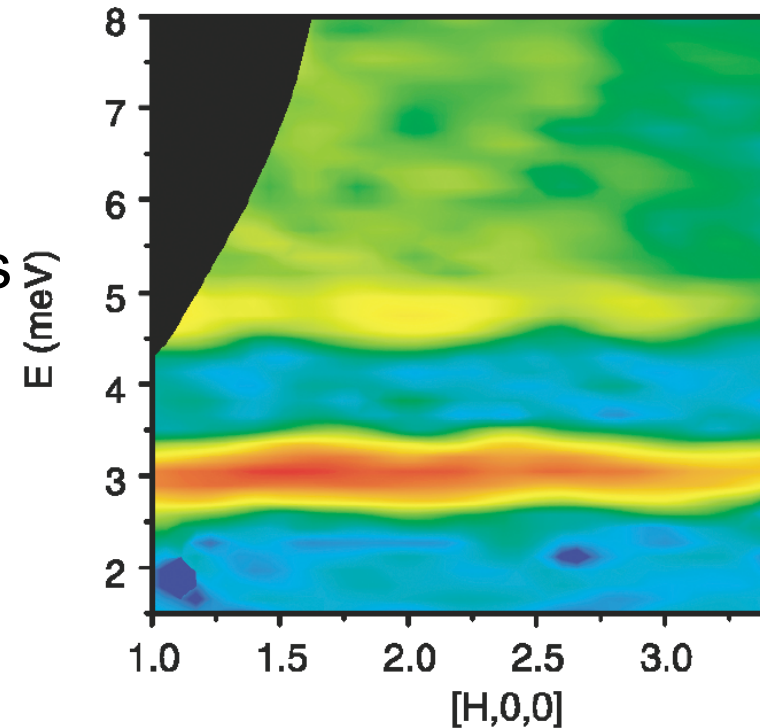


Frustration

\Rightarrow excited triplets
don't move

\Rightarrow No dispersion

Shastry-Sutherland model

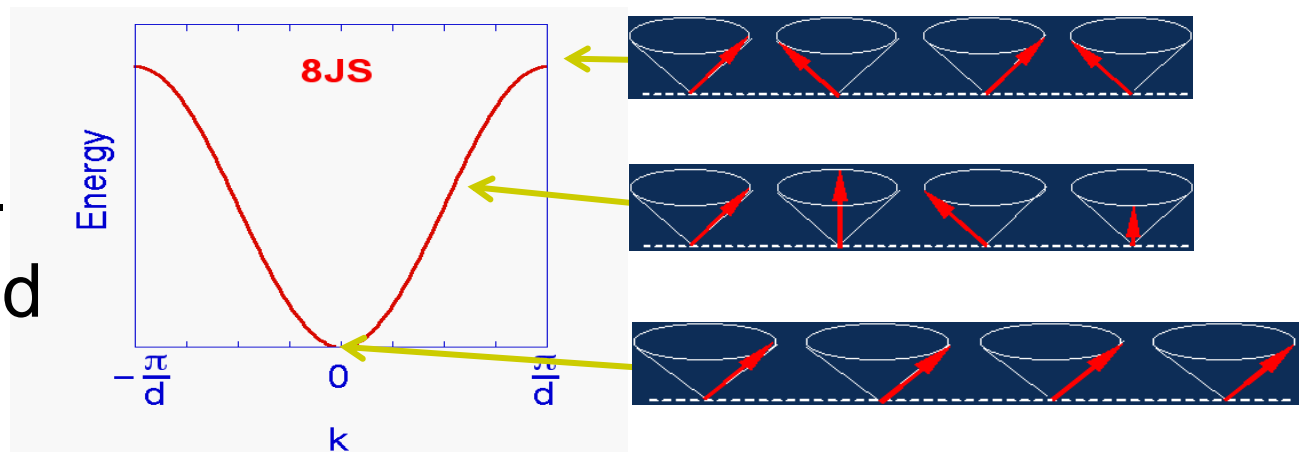


Spin waves (ferromagnet)

- Ordered ground state $H|g\rangle = E_g|g\rangle$
- Single spin flip not eigenstate: $S_r^\pm|g\rangle$
- Periodic linear combination: $|k\rangle = \sum_r e^{ikr} S_r^\pm|g\rangle$ (wave)
- Is (approximately) eigenstate: $H|k\rangle = E_k|k\rangle$
- Time evolution: $|k(t)\rangle = e^{iHt}|k\rangle = e^{iE_k t}|k\rangle$ (sliding wave)

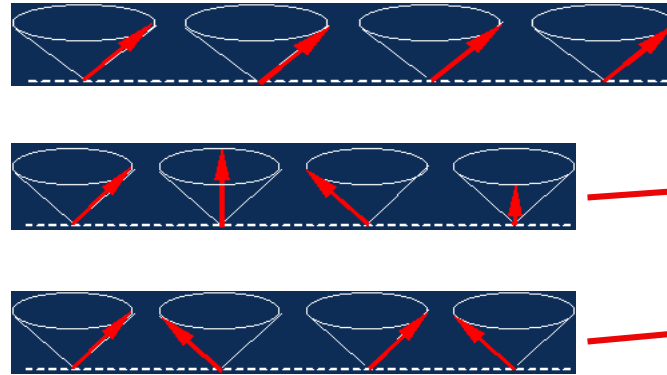
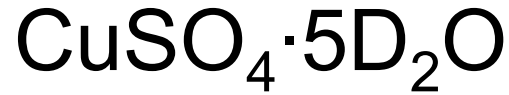


- Dispersion:
relation between
time- and space-
modulation period

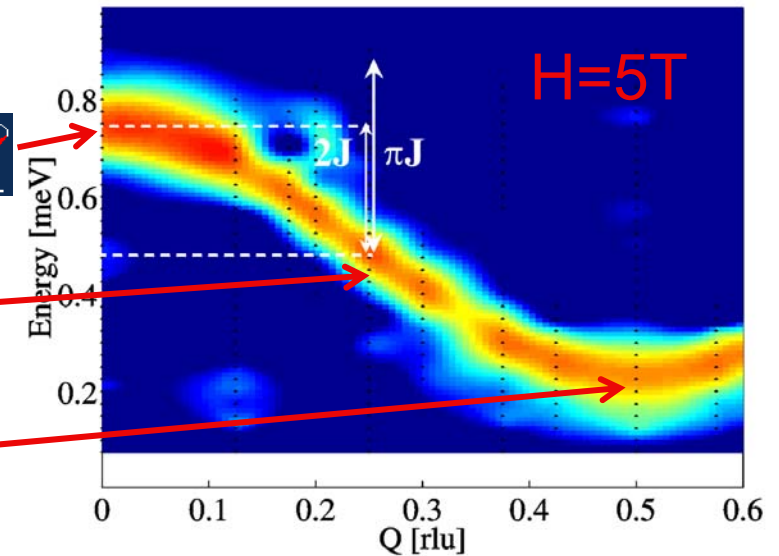




Spin wave example



Cold TAS – point-by-point ...



Antiferromagnet! – forced ferro by magnetic field

1-dimensional, Spin $\frac{1}{2}$, antiferromagnet

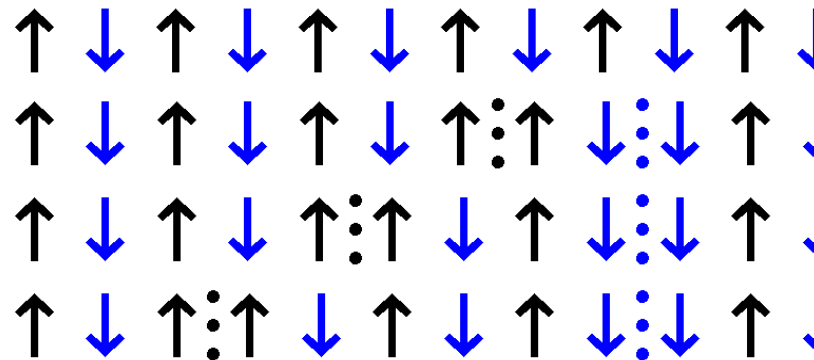
\Rightarrow quantum fluctuations and behaviour !



Spinon excitations

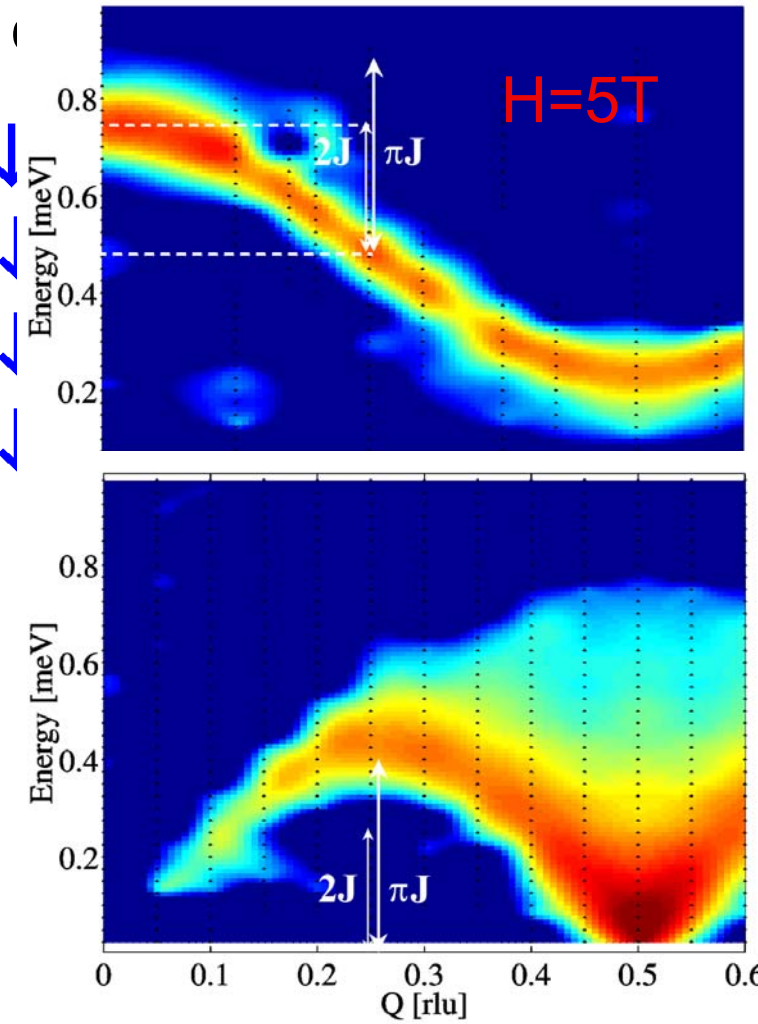


- In 1D a flipped spin ‘unbind’ to two domain walls
- “Spinons”: spin $S = \frac{1}{2}$ domain walls with respect to local AF ‘order’
- Need 2 spinons to form $S=1$ excitation we

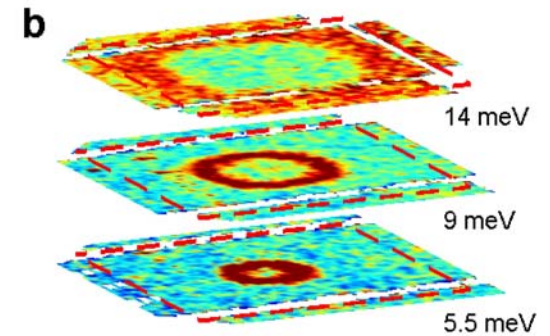
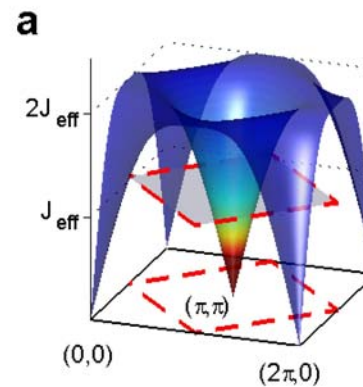
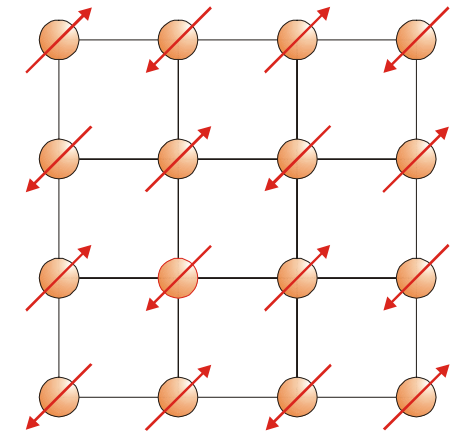
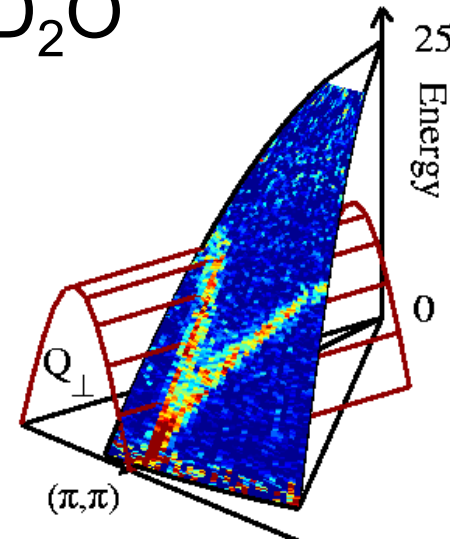
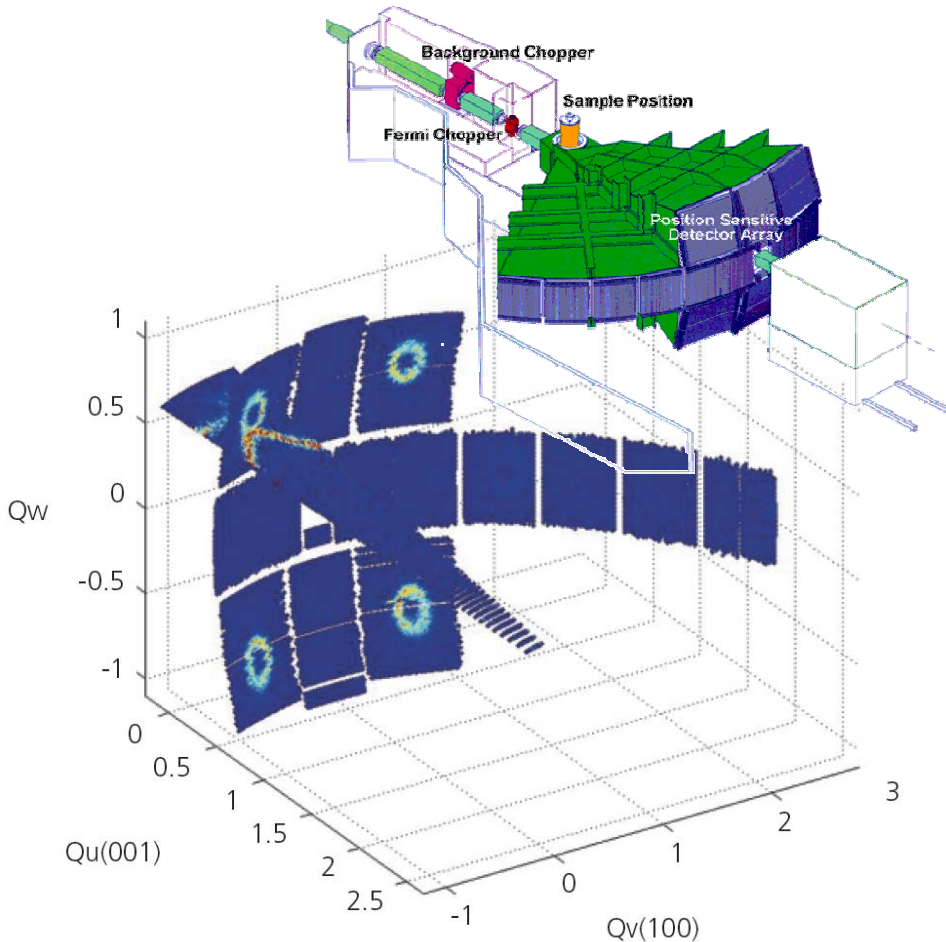


Energy: $E(\mathbf{q}) = E(k_1) + E(k_2)$
 Momentum: $\mathbf{q} = \mathbf{k}_1 + \mathbf{k}_2$
 Spin: $S = \frac{1}{2} \pm \frac{1}{2}$

Many possibilities for each \mathbf{q}
 Continuum of scattering \Rightarrow



2D Heisenberg antiferromagnet – long range ordered Spin waves in $\text{Cu}(\text{DCOO})_2 \cdot 4\text{D}_2\text{O}$



LQM 2-spin wave scattering



Ordered moment $\langle S \rangle = 60\%$,
 Spin wave amplitude $Z_\chi = 51\%$,
 Where is the remaining weight?

Polarised neutrons:

Spin-flip / non-spin-flip channels

- Separate **transverse** and **longitudinal**

$$|G\rangle = |\text{Neel}\rangle + \sum_k a_k |1\text{-spin waves}\rangle + \dots$$

spin-wave

2-spin-waves:

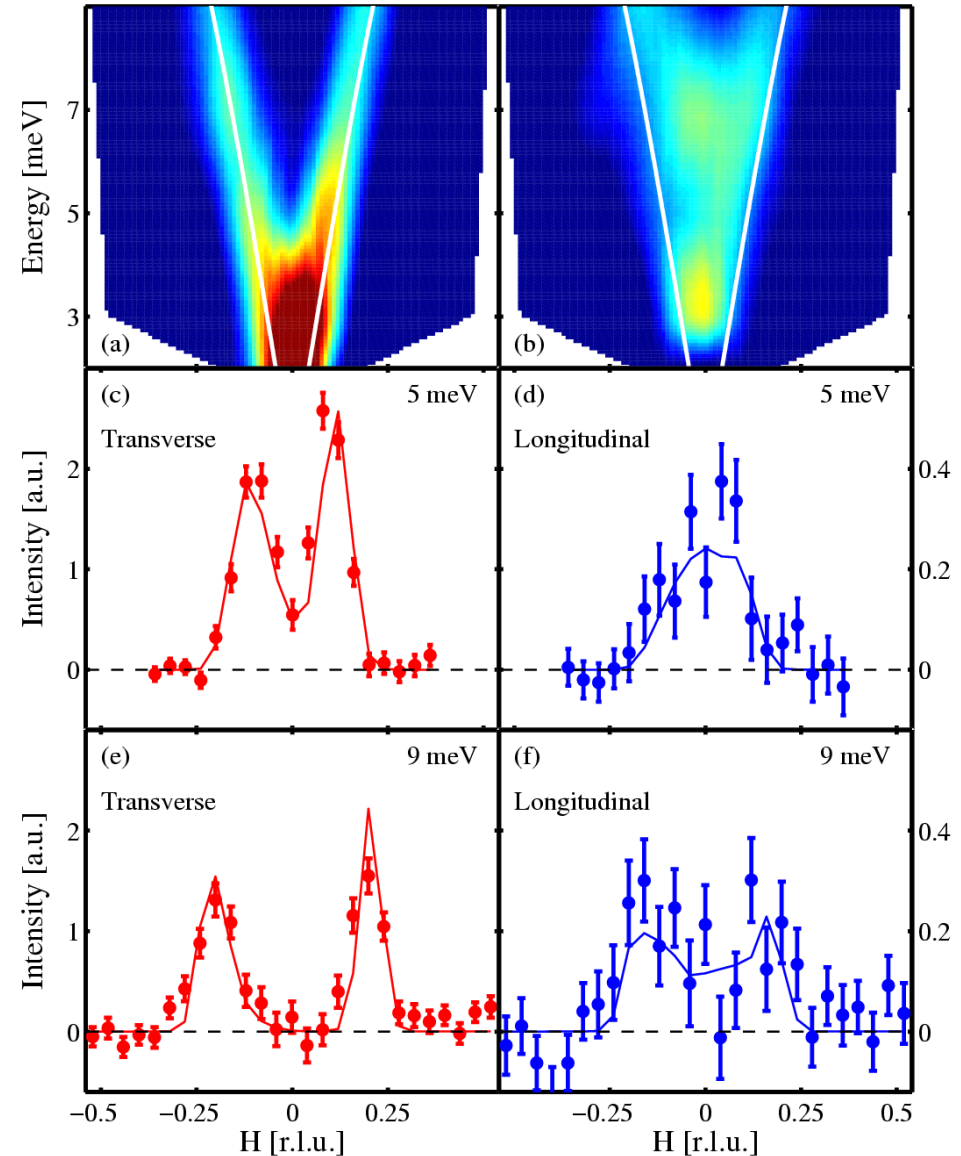
$$k = k_1 + k_2$$

$$E = E_1 + E_2$$

transverse

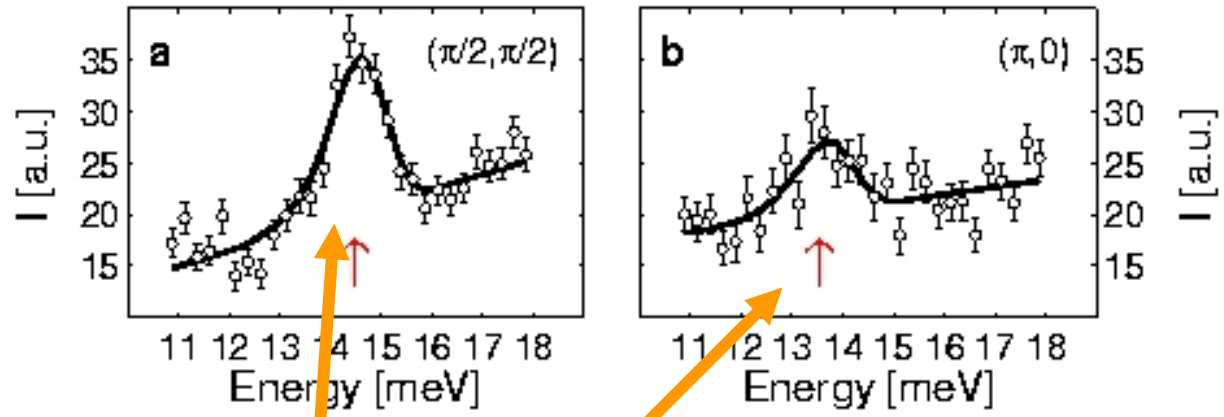
longitudinal

(spin-waves ok for long wavelengths)



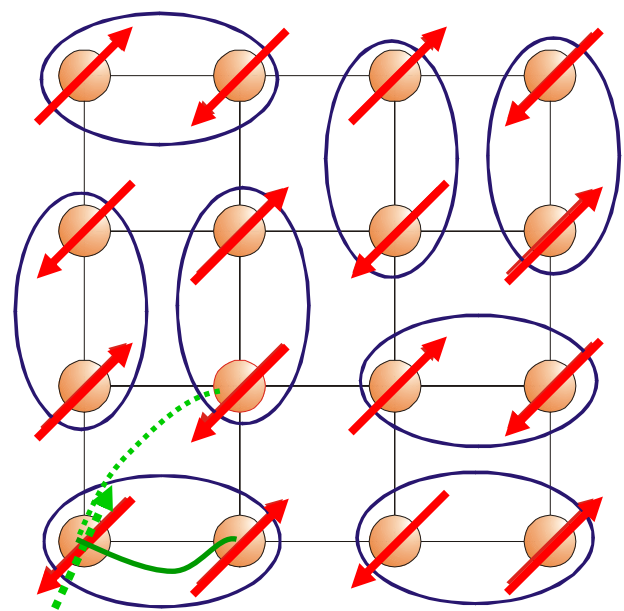
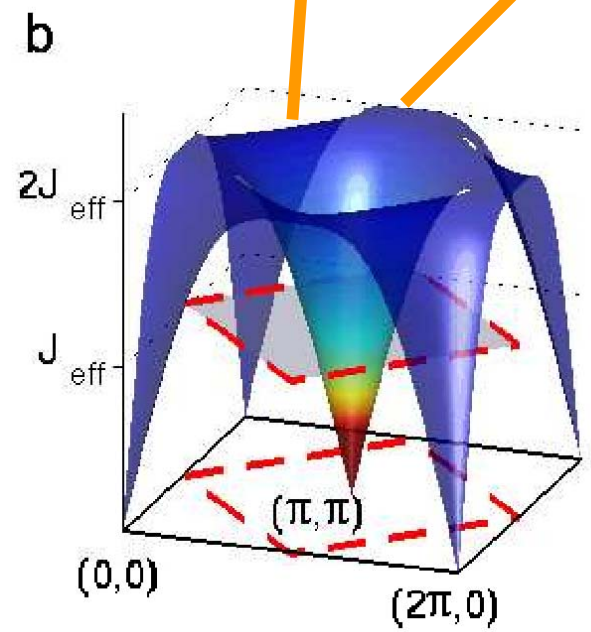


Surprise: zone boundary anomaly!

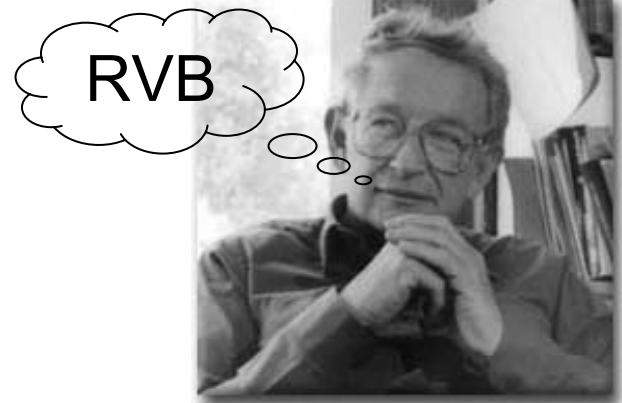


Should be same:
 7% energy effect
 50% intensity effect

What is nature of
 $\checkmark \pi 0n$ [$\pi 0n$] excitation



Hidden singlet correlations
 doping: afm order dies
 singlets survive \Rightarrow SC?





S(Q) instantaneous correlations



S(Q) is Fourier transform of 'snap-shot'

$$S(\mathbf{Q}) = \sum_{\mathbf{r}' \neq \mathbf{r}} e^{-i\mathbf{Q} \cdot (\mathbf{r}' - \mathbf{r})} \langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(t) \rangle$$

Real space exponential decay:

$$\langle S_{\mathbf{r}'}(t) S_{\mathbf{r}}(t) \rangle \propto e^{-|\mathbf{r} - \mathbf{r}'|/\xi}$$

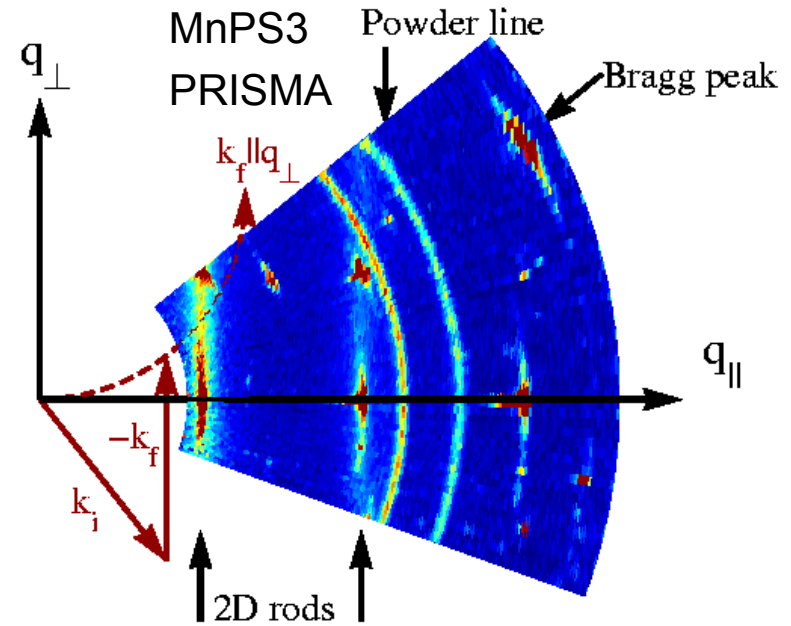
Reciprocal space: Lorentzian

$$S(\mathbf{Q}) \propto \frac{1}{1 + Q^2 \xi^2}$$

Correlation length ξ

In 2D, S(Q) gives rods along Q_{\perp}

Integrate by having $k_f \parallel Q_{\perp}$



Wildes et al. \Uparrow



S(Q) instantaneous correlations



$$S(Q) \propto \frac{1}{1 + Q^2 \xi^2}$$

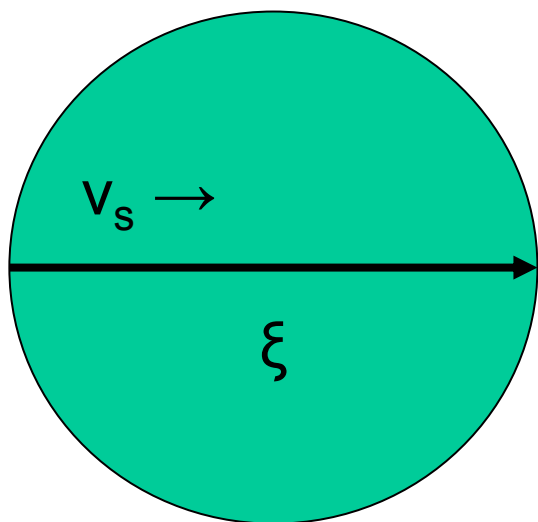
Width \Rightarrow Correlation length ξ

Softening

$$v_s \star Z_c$$

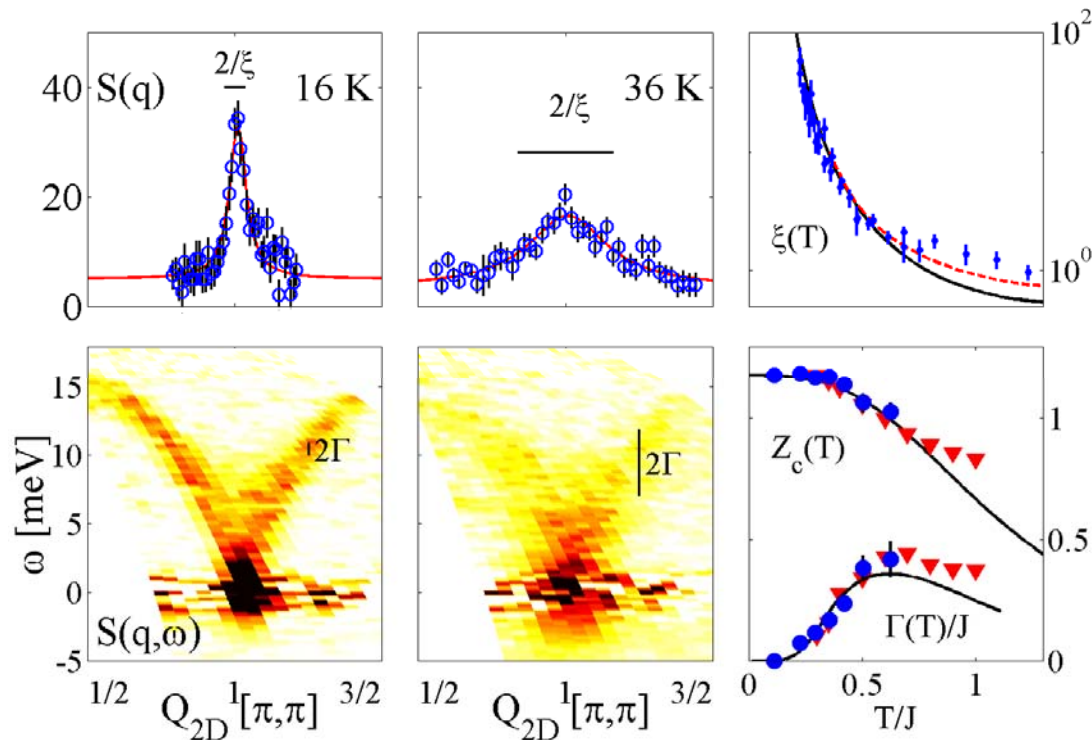
Damping

$$\Gamma = v_s / \xi$$



Life-time

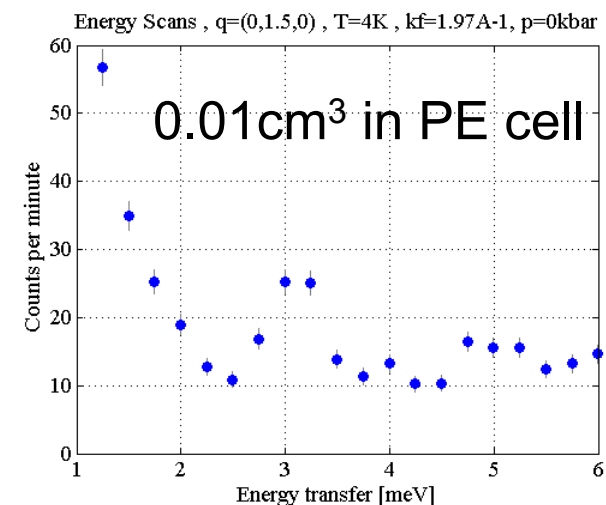
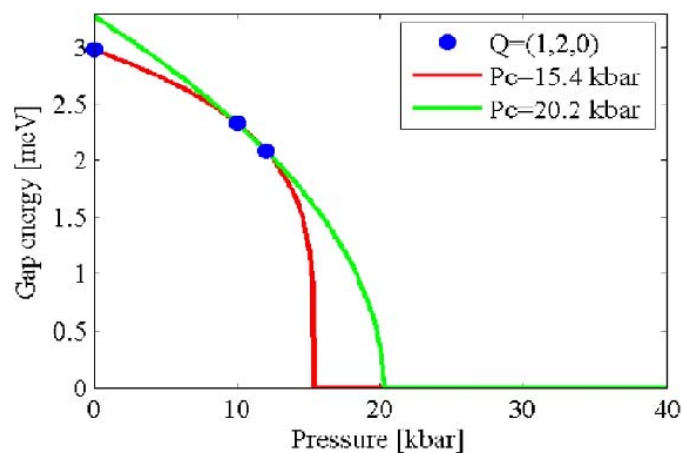
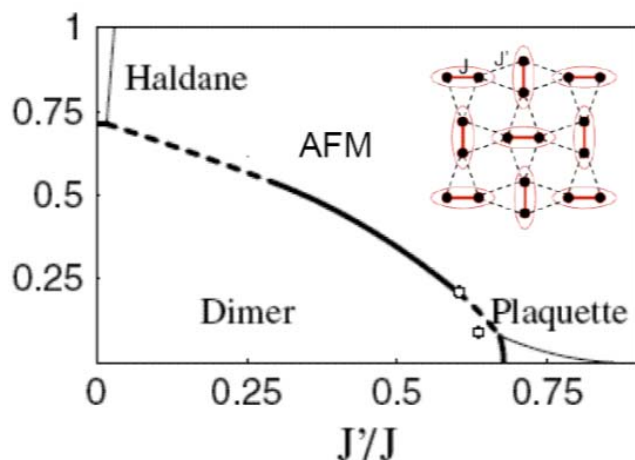
$$\tau = \xi / v_s = 1 / \Gamma$$



LQM Parameters



- E from micro eV to eV
- T from few mK to 1000K
- H to 15T
- P to 17kbar / 50kbar
example: $\text{SrCu}_2(\text{BO}_3)_2$





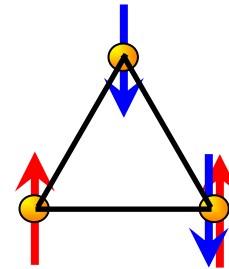
Recap.: What do we measure



- Localised excitations
- Dispersive excitations
- Dynamic correlations
- Polarisation analysis
- Dependencies:
Temperature, Magnetic field, Pressure field



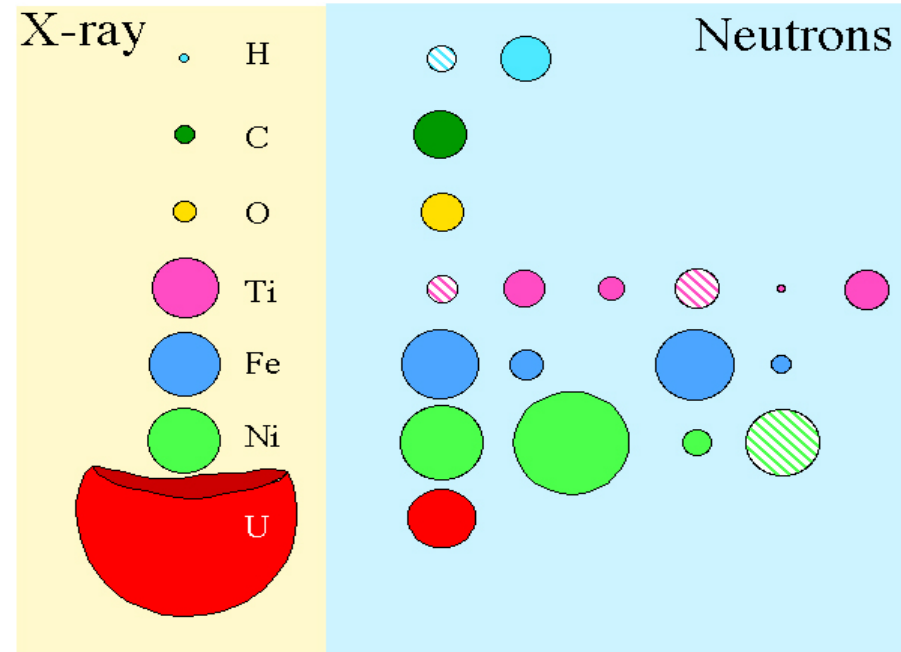
- It's a battle: samples, statistics, resolution



- But often *'the final kill'* for problems in magnetism

Thank you

- Similar in technique, but
 - X-rays much more intense
 - Different scattering lengths:
 - Neutrons simpler for magnetism
 - Neutrons better for energy resolution



- Mostly, only one will work for a given problem

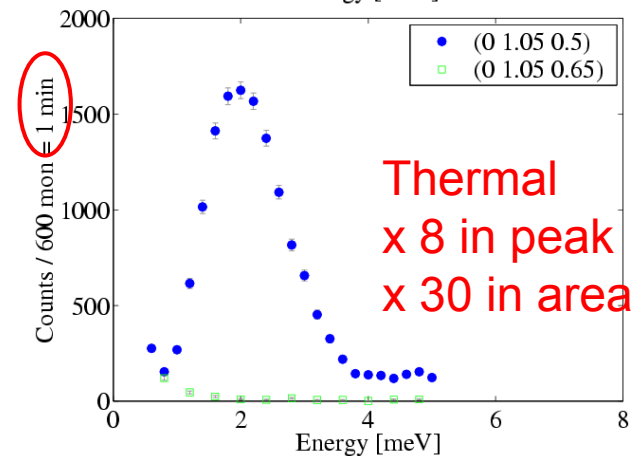
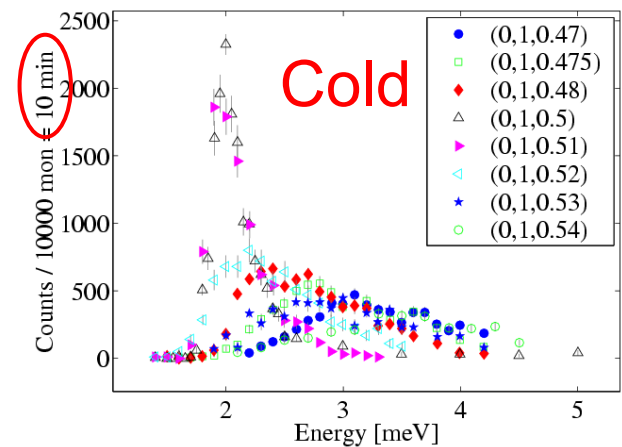
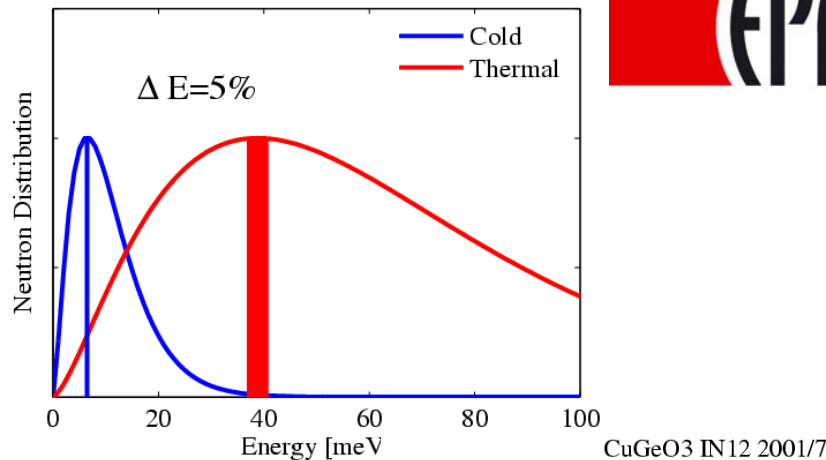
LQM TAS resolution



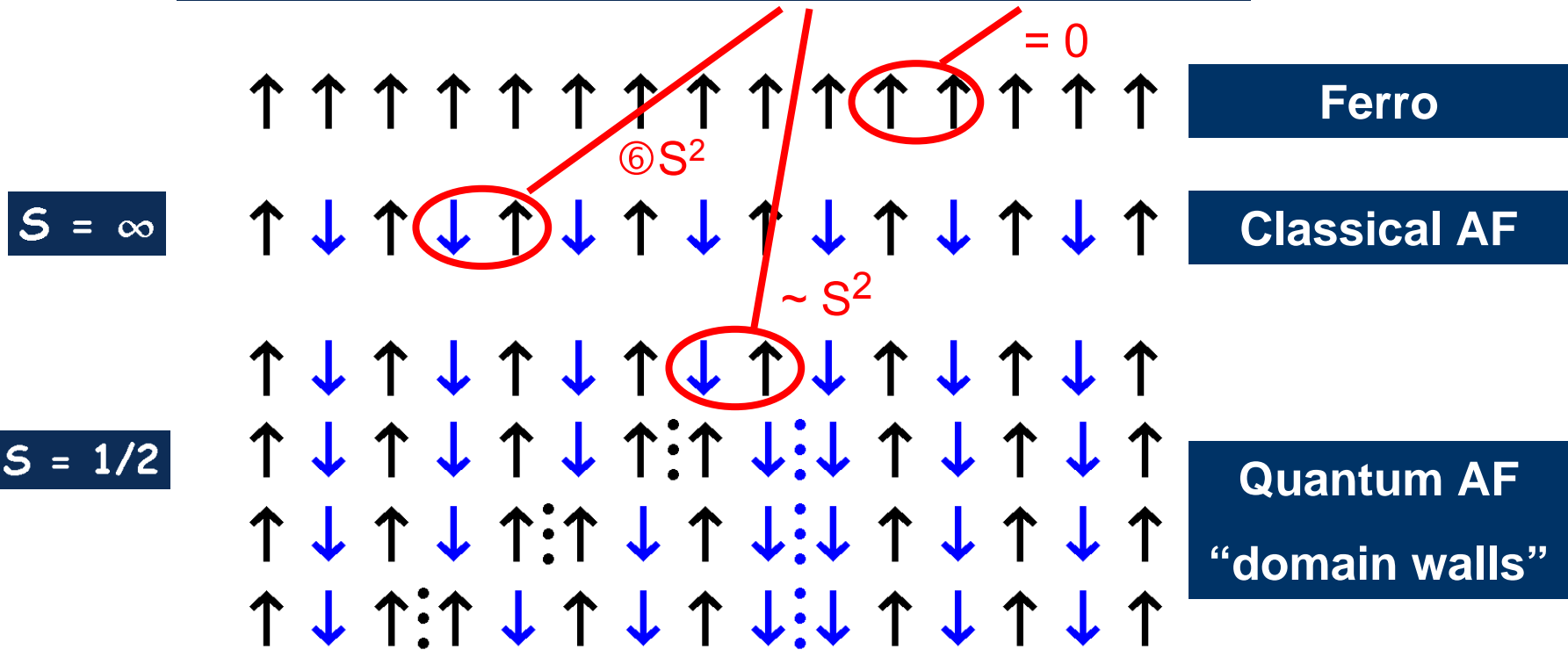
- Source spectrum:
 - Cold: $E_i=2-15$ meV
 - Thermal $E_i=15-100$ meV
 - Epithermal $E_i \Rightarrow 1-2$ eV
- Intensity / resolution:

Some typical numbers (PG002):

E_i, E_f	resolution (fwhm)
2 meV	30 μ eV
5 meV	0.15 meV
14.7 meV	1.2 meV
- Typical counting $\sim 1-10$ min/pt



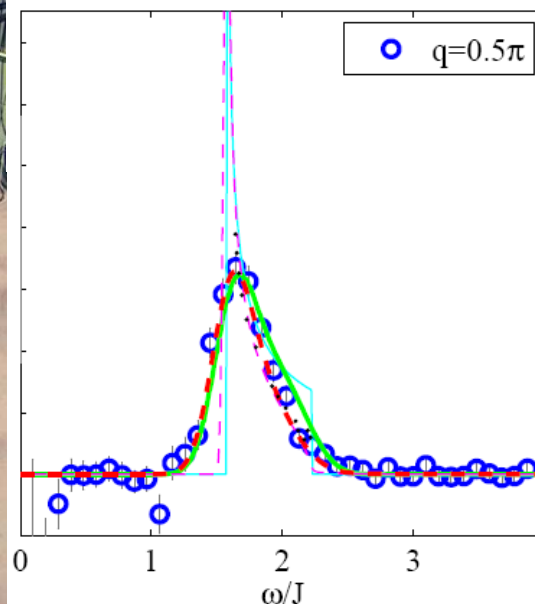
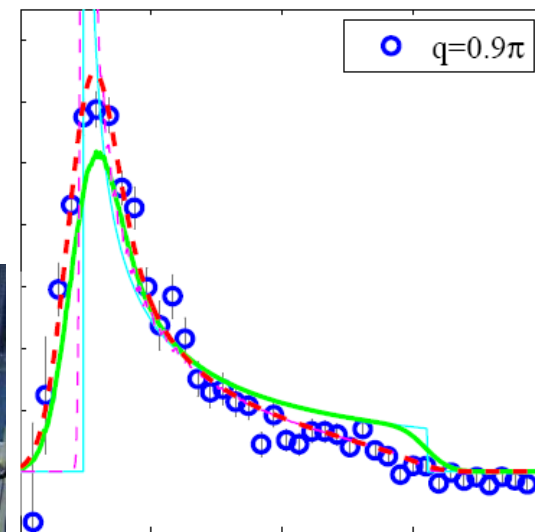
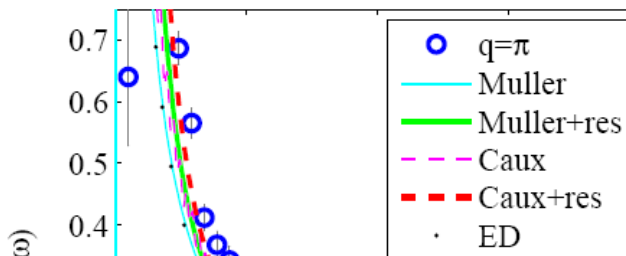
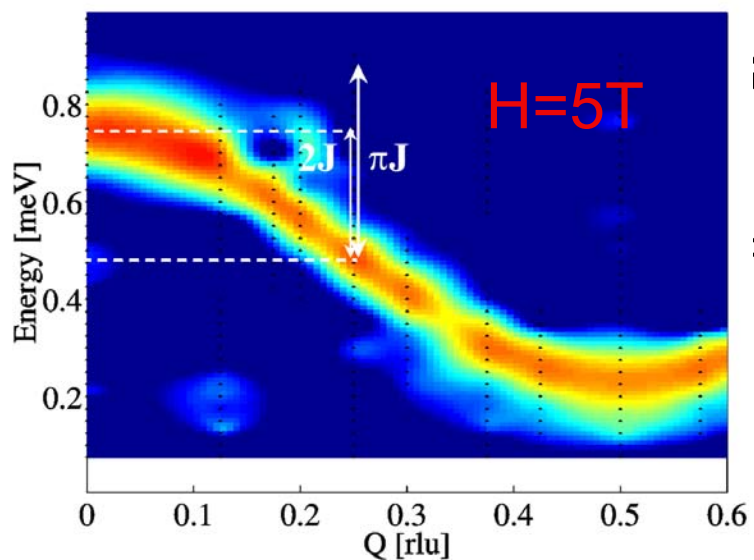
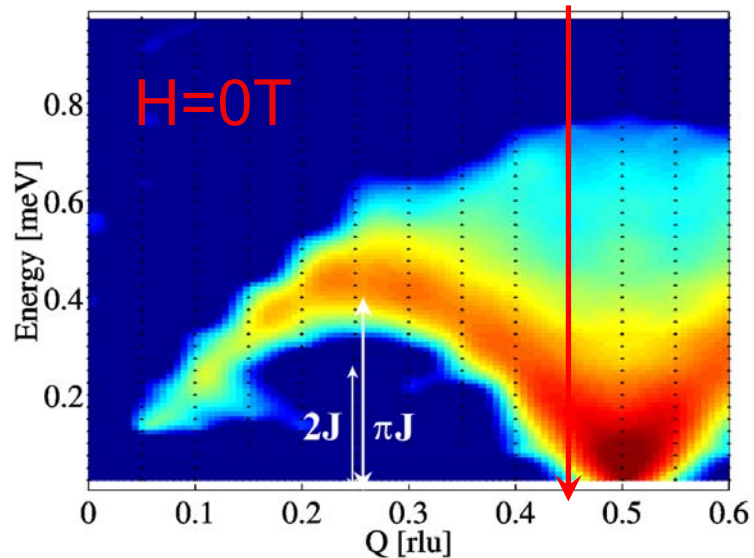
$$\mathcal{H} = J \sum S_n^z S_{n+1}^z + \frac{1}{2} (S_n^+ S_{n+1}^- + S_n^- S_{n+1}^+)$$



Ground state (Bethe 1931) disordered by quantum fluctuations



Continuum in 1D



Cold TAS: IN14



2-spin wave scattering



Ordered moment $\langle S \rangle = 60\%$,
 Spin wave amplitude $Z_\chi = 51\%$,
 Extra 25% dip at $(\pi, 0)$

Where is the remaining weight?

Ground state:

$$|0\rangle = |\text{Neel}\rangle + \sum_k a_k |\text{spin wave with momentum } k\rangle + \dots$$

spin-wave
transverse

2 spin-waves: $k_1 + k_2 = k$ and $E_1 + E_2 = E$
longitudinal

