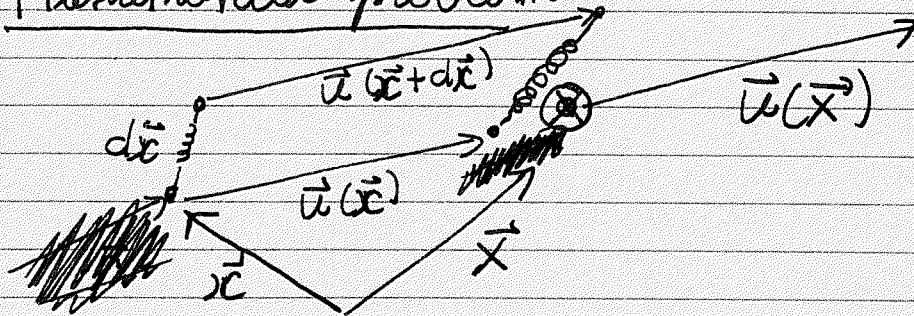


# Finite Strain Elasticity Theory

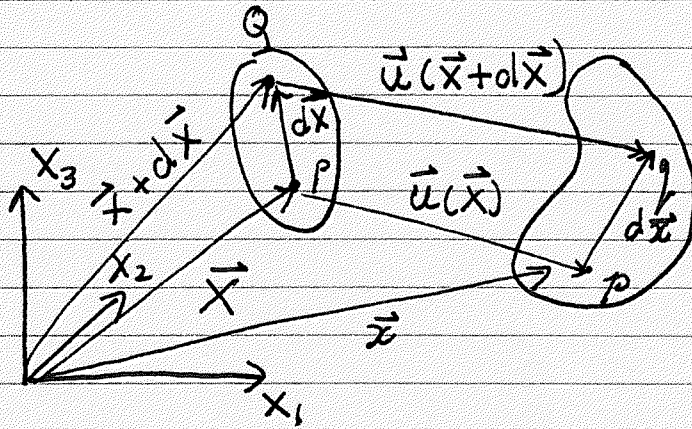
- \* If the strain tensor  $\epsilon_{ij}$  is of the order of one or larger, then classical elasticity theory is mathematically incorrect. ~~It is~~ This is to be distinguished from physical incorrectness due to the fact that higher order invariants must be included.

## \* Mathematical problem



The stress  $\vec{\sigma}$  in the region marked by  $x$  is not related to gradients of  $\vec{u}(\vec{X})$  but instead to gradients of  $\vec{u}(\vec{x})$  originating in mass points originally located far from  $\vec{X}$ .

- \* Define  $\vec{u}(\vec{X})$  as the displacement field in the reference frame of the undeformed body.



P : material point in undeformed body,  
located at  $\vec{X}$

p : <sup>same</sup> material point in deformed body,  
located  $\vec{X} + \vec{u}(\vec{X}) = \vec{x}$

Q : material point in undeformed body,  
located at  $\vec{X} + d\vec{X}$

q : same material point in deformed body,  
located at  $\vec{X} + d\vec{X} + \vec{u}(\vec{X} + d\vec{X}) = \vec{x} + d\vec{x}$

Use  $\vec{x} = \vec{X} + \vec{u}(\vec{X})$ . Define  $d\vec{u} = \vec{u}(\vec{X} + d\vec{X}) - \vec{u}(\vec{X})$

$$d\vec{x} = d\vec{X} + \underbrace{d\vec{u}}_{\text{relative displacement}}$$



Expand

$$\vec{u}(\vec{X} + d\vec{X}) \approx \vec{u}(\vec{X}) + \underbrace{\nabla_{\vec{X}} \vec{u}}_{\frac{\partial u_i}{\partial X_j} \sim \text{strain tensor}} \cdot d\vec{X}$$

$$\begin{aligned} d\vec{x} &= d\vec{X} + d\vec{u} \\ &= d\vec{X} + \nabla_{\vec{X}} \vec{u} \cdot d\vec{X} \\ &= (\mathbf{I} + \nabla_{\vec{X}} \vec{u}) \cdot d\vec{X} \end{aligned}$$

~~$\vec{x} = \vec{X} + \vec{u}$~~

 $d\vec{x} = \overset{\leftrightarrow}{F} d\vec{X}$

Deformation Gradient Tensor ~~metric tensor~~

Right Cauchy-Green deformation tensor

$\overset{\leftrightarrow}{C} = \overset{\leftrightarrow}{F}^T \overset{\leftrightarrow}{F}$

 $C_{IJ} = \frac{\partial x_k}{\partial X_I} \frac{\partial x_k}{\partial X_J}$

Metric Tensor

$$d\vec{x}^2 = d\vec{X} \overset{\leftrightarrow}{C} d\vec{X} \quad \vec{x} \rightarrow \vec{X}$$

# Invariants of $\overset{\leftrightarrow}{C}$ ( $d=2$ )

$$\text{Tr } \overset{\leftrightarrow}{C}, \text{Det } \overset{\leftrightarrow}{C}, \frac{1}{2}(\text{Tr } \overset{\leftrightarrow}{C}^2 - (\text{Tr } \overset{\leftrightarrow}{C})^2) \dots$$

\*  $\overset{\leftrightarrow}{E} = \frac{1}{2}(\overset{\leftrightarrow}{C} - \overset{\leftrightarrow}{I})$  strain Tensor \* ("Green-Lagrange" strain Tensor)

$$E_{KL} = \frac{1}{2} \left( \frac{\partial u_K}{\partial x_L} + \frac{\partial u_L}{\partial x_K} + \frac{\partial u_M}{\partial x_K} \frac{\partial u_M}{\partial x_L} \right)$$

Let  $\lambda_1^2$  &  $\lambda_2^2$  be the eigenvalues of  $\overset{\leftrightarrow}{C}$

\* Show that  $(\text{Det } \overset{\leftrightarrow}{C})^{1/2} = 1 = \lambda_1 \lambda_2 - 1$

measures area changes. (In general  $dA' = \det \overset{\leftrightarrow}{F} dA$ )

\* Show that  $\frac{\text{Tr } \overset{\leftrightarrow}{C}}{\text{Det } \overset{\leftrightarrow}{C}^{1/2}} - 2 = \frac{\lambda_1}{\lambda_2} + \frac{\lambda_2}{\lambda_1} - 1$

measures shear.



# Cauchy Stress Tensor

## Elastic Energy

$$F = 2\mu E_{ij}^2 + \lambda E_{ll}^2$$

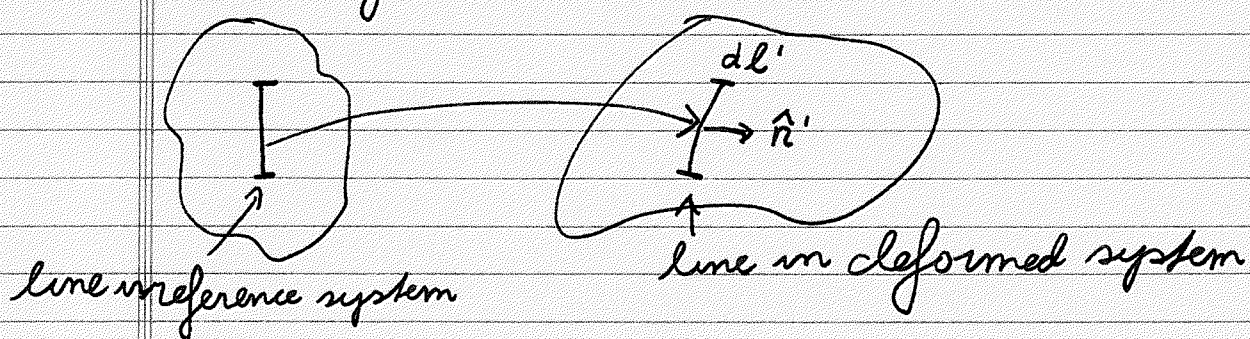
"Saint-Venant"

$E_{ij}$  = Green-Lagrange Strain

Analogy of usual stress tensor

$$S_{ij} = \partial F / \partial E_{ij} = 2\mu E_{ij} + \lambda E_{ll} \delta_{ij}$$

"Cauchy Stress Tensor"



$\vec{S} \cdot \hat{n}'$  is the force/unit length in deformed configuration on deformed line element