

Quasiparticle in d-wave Superconductor [Ong]

- By doing charge-heat excitation, one can differentiate between quasiparticles, vortex, phonons, etc.

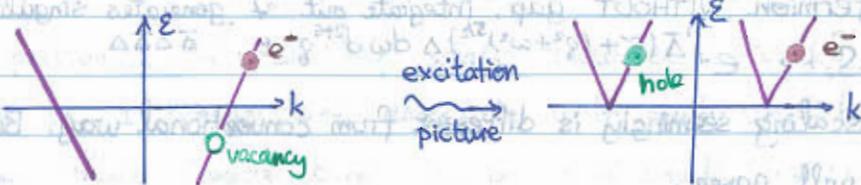


Vortex

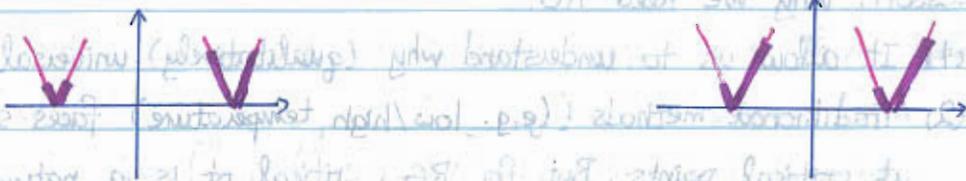
Dirac quasiparticle

- Theoretically, $\vec{\sigma}$ in $\vec{J} = \vec{\sigma} \cdot \vec{E}$ is easier to calculate, but experimentally it is \vec{p} is $\vec{E} = \vec{p} \cdot \vec{J}$ that's measured.

- Including heat current, $\vec{J} = \vec{\sigma} \cdot \vec{E} + \alpha (-\nabla_x T)$
 $\vec{\sigma}$ and α can be computed by Boltzmann Egn.



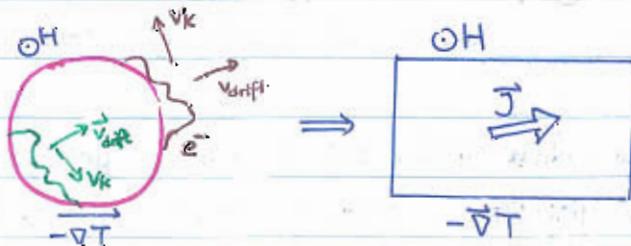
- Electric field : vs. Heat Gradient



Large charge current

Small mass current

- Thermal Hall



- Nernst effect

Setting $\vec{J} = \sigma \cdot \vec{E} + \vec{\alpha} \cdot (\nabla T)$ with $\vec{J} = 0$

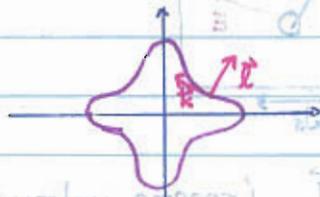
$\Rightarrow \vec{E} = -\vec{\rho} \cdot \vec{\alpha} \cdot (\nabla T)$

Effect is small since α_{xy} and σ_{xy} cancels

- The Hall effect coeff. can be written as:

$\sigma_{xy} = e^3 B \cdot \frac{1}{2} \oint d\vec{l} \times \vec{l} \Rightarrow$ area swept out in k -space.

\vec{l} mean free path



- Note that Hall effect is 2nd-order: need both \vec{E} and \vec{B} presence.

- In ordinary metal, Hall coeff. (ρ) should be approx. const. at Kelvin scale. But in cuprates it varies drastically. When convert to Hall angle it shows a T^2 (rather than T) dependence.

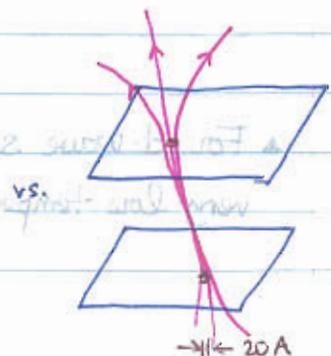
$\sim \sigma_{xy} / \sigma_{xx}$

- In Cuprates, Fermi arc ends abruptly, and its length is roughly proportion to T .

\Rightarrow Take a model where only states on Fermi arc have long mean free path.

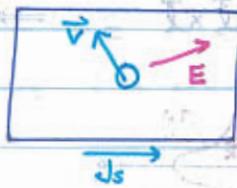
- The model fit data well EXCEPT at small temperature, where observed R_H becomes small.

- Since Cuprate is 2D, vortex forms "pancakes". Vortex size much smaller than in 3D.



Burdick - Stephen Theory

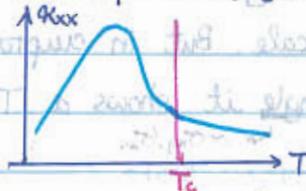
- ▲ Supercurrent exert forces on vortex core
- ▲ As vortex moves, an internal \vec{E} -field is produced by Faraday effect
- ▲ The internal \vec{E} -field attract pairs into core, where they break into normal particles, disipate, and recombine.
- ▲ A macroscopic, observable \vec{E} -field is then generated by Josephson effect.



- ▲ In clean limit $\vec{v} \propto -\vec{J}_s$ (reason unknown)

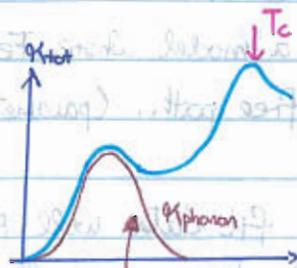
Problem — separates quasiparticles from vortex.

- ▲ Consider thermal power:



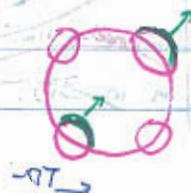
Is peak caused by phonon or quasiparticle? (vortex has small contribution, supercurrent carry no entropy)

- ▲ Case of old SC (Pb):



For low T , phonon- e^- scattering vanishes, giving 2nd peak.

- ▲ For d-wave superconductor, quasiparticle exists even down to very low temperature, because of nodes.



- Quasiparticle and phonon can be distinguished since scattering of q.p. by vortex is asymmetric, while phonon-vortex scattering is symmetric.
- Cause by Doppler shift



- Conclusion: peak is caused mostly by quasiparticles.

$$\frac{(1+2)Z^2 \mu^2 \rho^2}{2E} = A$$



$$\frac{(1+2)Z^2 \mu^2 \rho^2}{2E} = A \Rightarrow E = \frac{(1+2)Z^2 \mu^2 \rho^2}{2A}$$

- For high temperatures, the superconducting transition temperature T_c is low.
- For single cuprate, $T_c^{\text{max}} = 140\text{K}$.
- Empirically, frustration occurs if (Bohrer) .
- $f = \frac{\hbar v_F}{\hbar v} \gg 1$

Conceptually, frustration \leftrightarrow competing interactions

• Control magnetic dipole-dipole interactions

• Many states are as good as each other

(i.e. many degenerate ground states)

• All states are equally likely to be occupied

to be equally likely