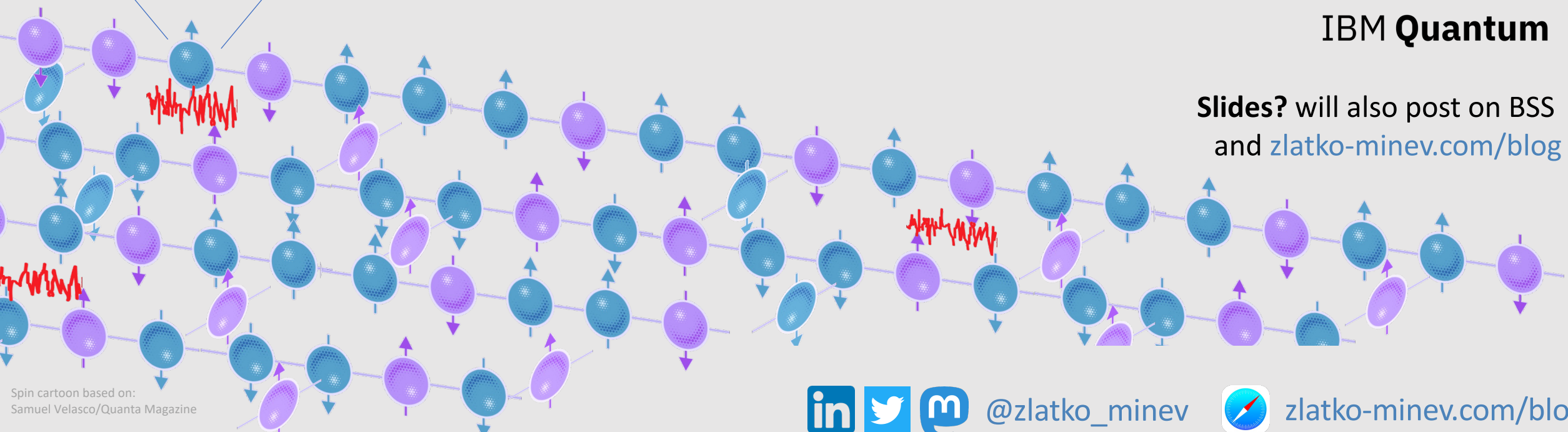
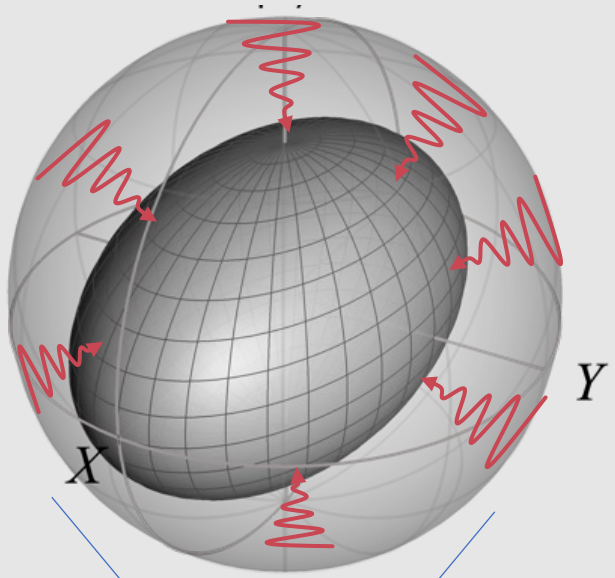


Quantum Error Mitigation for Non-Equilibrium Quantum Dynamics on Quantum Computers

Lecture 2

Zlatko K. Mineev
IBM Quantum

Slides? will also post on BSS
and zlatko-mineev.com/blog



Spin cartoon based on:
Samuel Velasco/Quanta Magazine

Where can you find things?

Lecture Slides

Boulder School for Condensed Matter and Materials Physics

boulderschool.yale.edu/2023/boulder-school-2023-lecture-notes



Will also post on zlatko-minev.com/education

Tutorials and additional lecture notes

Twirling, Measurements and Walsh-Hadamard

Cheat sheets, Videos, ...

zlatko-minev.com/blog

See also lectures on qiskit.org/learn

Tutorials and additional lecture notes

Latest seminar qiskit.org/events/seminar-series

The image shows three article thumbnails from the website zlatko-minev.com. The first thumbnail is titled "7. Digital quantum circuits (pictorial)" and "7A. Basic elements", featuring a diagram of quantum and classical wires and gates. The second thumbnail is titled "Primer on Pauli Twirling" and shows a quantum circuit with P_a , Λ , and P_a^\dagger gates, along with a 2D grid and a 3D plot. The third thumbnail is titled "Learn and cancel quantum noise" and shows a quantum circuit diagram. Each thumbnail includes the author's name, Zlatko K. Minev, and a date.



What is one thing you
learned in Lecture 1?



Review of Lecture 1



Quantum Error Mitigation for Non-Equilibrium Quantum Dynamics

Lecture 1

Big picture

Why quantum computers?
Status and outlook

Why error mitigation?

Noise in quantum computers
Overview of error mitigation

Mitigation fundamentals

Probabilistic error cancelation (PEC)
Introduction
One qubit example



Quantum computers

My experience circa 2010

Maybe 1 or 2 qubits working some small fraction of the time in select labs

Photo with dilution fridge called Sunshine from Michel Devoret's lab at Yale during my Ph.D.



Hopes for a working qubit in here

Mineev, IBM Quantum (9)

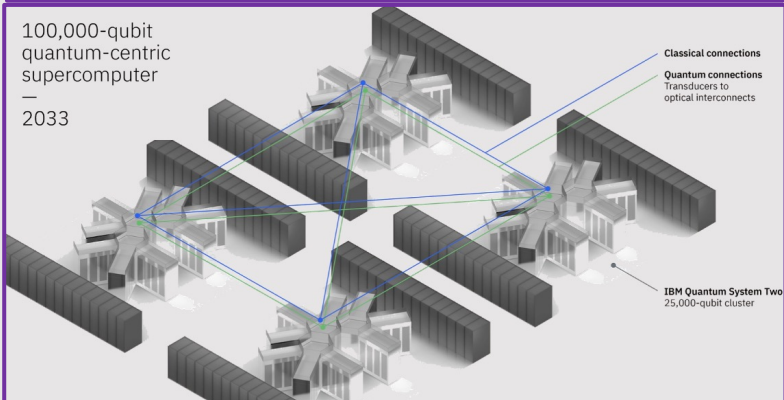
This year 2023

A 127-qubit quantum computer installed in the lobby cafeteria of a research building dutifully executing jobs almost all the time.

at Cleveland Clinic



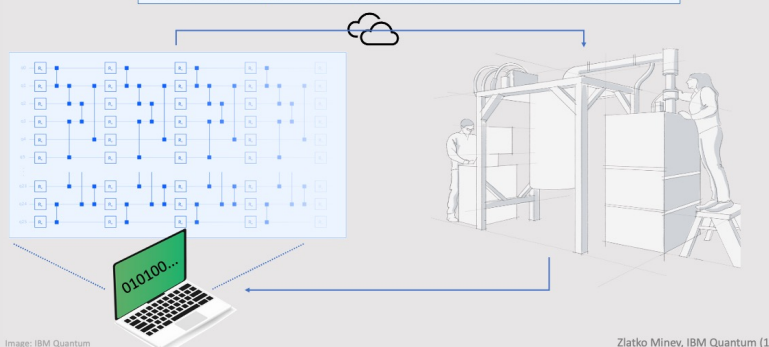
100,000-qubit quantum-centric supercomputer
—
2033



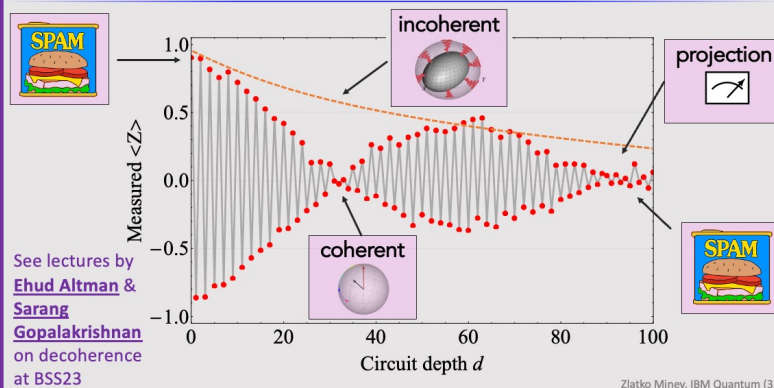
Biggest Problem: Noise

Quantum simulation on a quantum computer

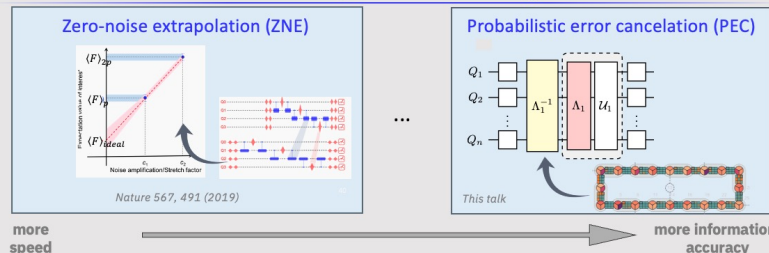
Execute on a real quantum computer device and obtain results as classical data



Elements of noise



Error mitigation landscape



Mitigation

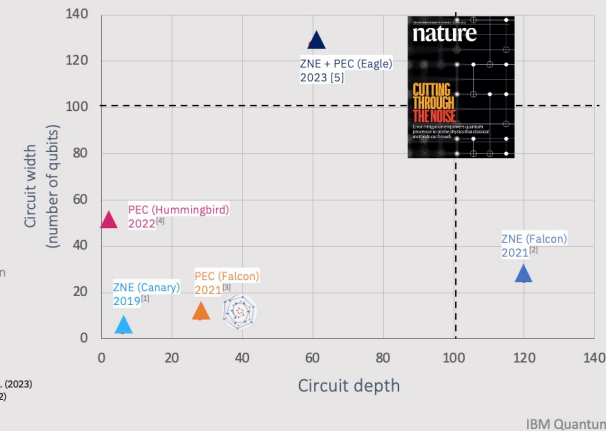
Overview of some key experimental progress in error mitigation:

Error mitigation

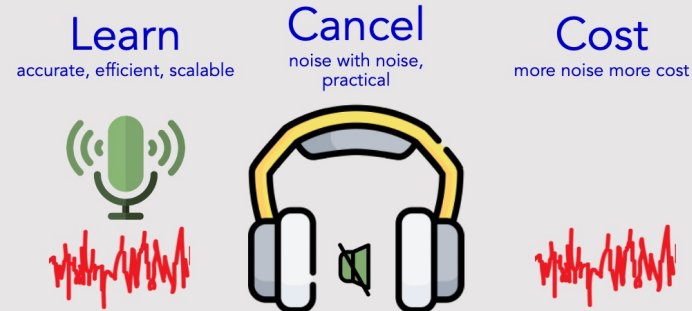
No matter what you do you have to chop it to this graph

PEC: Probabilistic error cancellation

ZNE: Zero-noise extrapolation



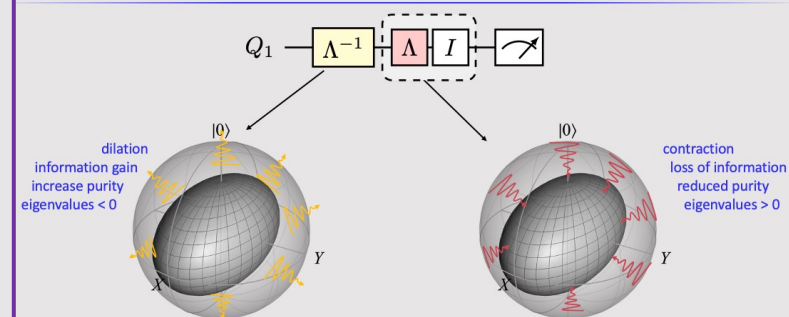
High-level message



from: Flatiron

Zlatko Mineev, IBM Quantum (48)

Inverse of noise map is not physical



Quantum Error Mitigation for Non-Equilibrium Quantum Dynamics

Lecture 2

Mitigation fundamentals

Probabilistic error cancelation (PEC)

Summarize one qubit example

Analogy to random walks

Error bars & confidence

Generalize (optional)

Show unbiased estimator



Learning quantum noise

Challenge

Overcoming: sparse model

Putting it together

Experiments – Ising model

Consequences for the big picture

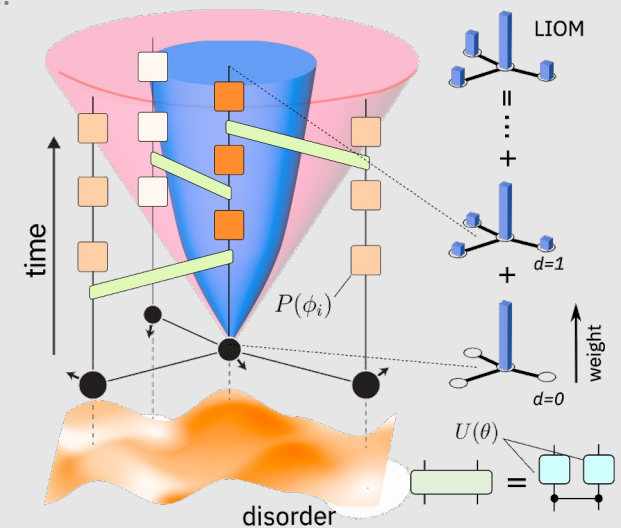
Hardware progress

Lecture 3

Wrap up of Lecture 2, Twirling, ...

State-of-art experiments at the 120Q+, depth 50+: uncovering local integrals of motion

...



Deep dive:
Probabilistic error cancellation (PEC)
To learn and cancel quantum noise



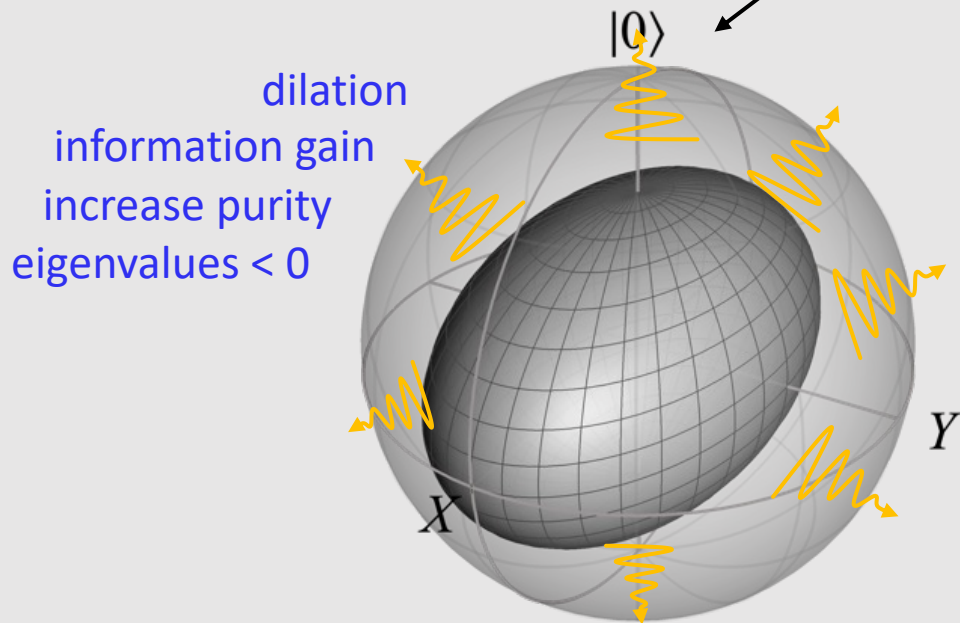
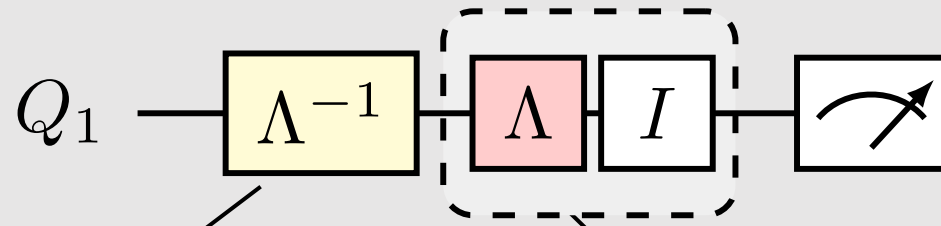
Got Slides?



Paper: [Nature Physics \(2023\)](#)

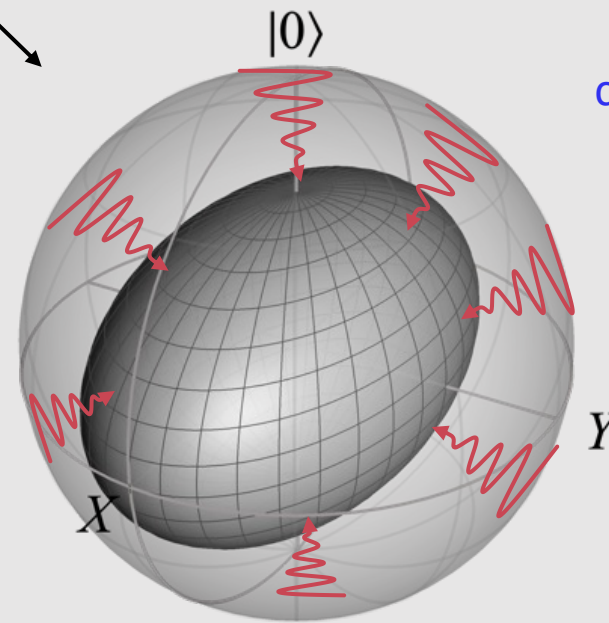
Ewout van den Berg, Zlatko K. Mineev, Abhinav Kandala, Kristan Temme

Inverse of noise map is not physical



dilation
information gain
increase purity
eigenvalues < 0

$$\Lambda^{-1}(\rho) = ?$$



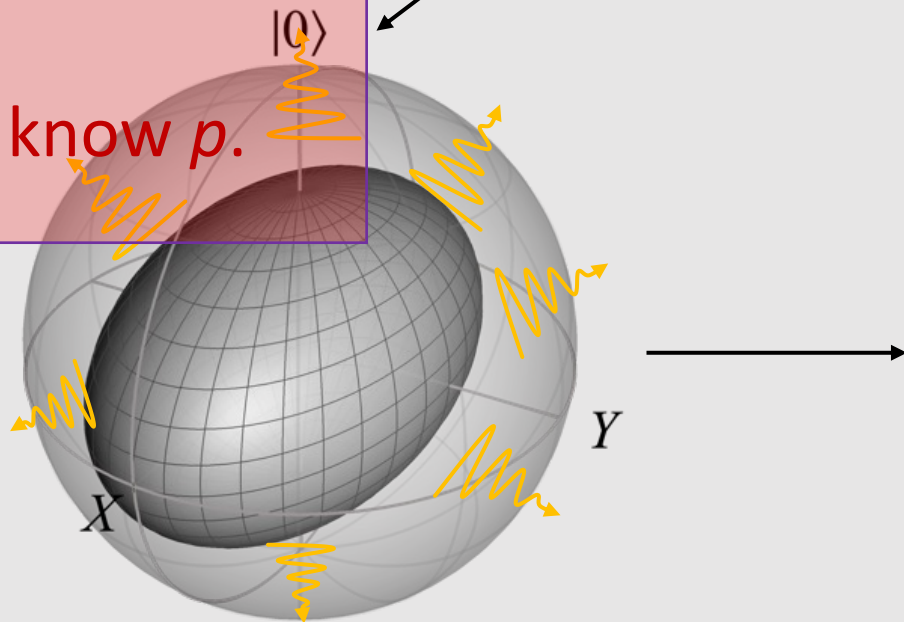
contraction
loss of information
reduced purity
eigenvalues > 0

$$\Lambda(\rho)$$

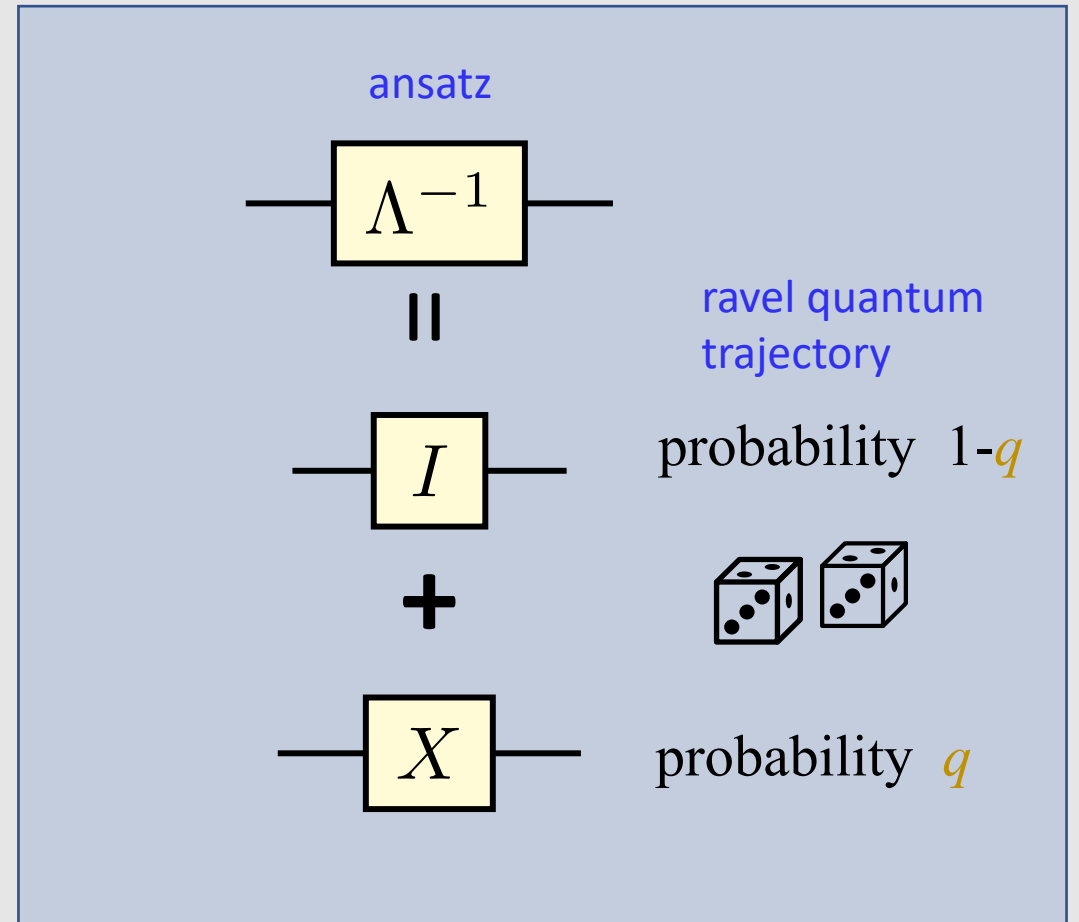
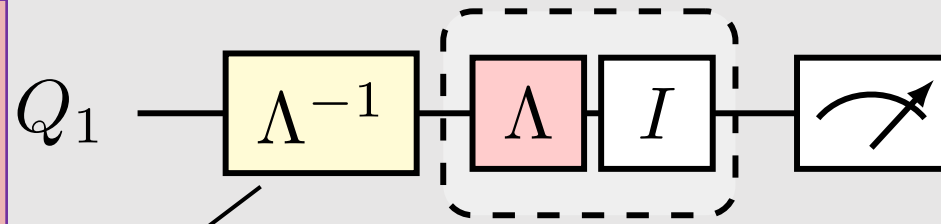
Inverse of noise map is not physical

Critically hinges on knowing the noise exactly!

Need to know p .



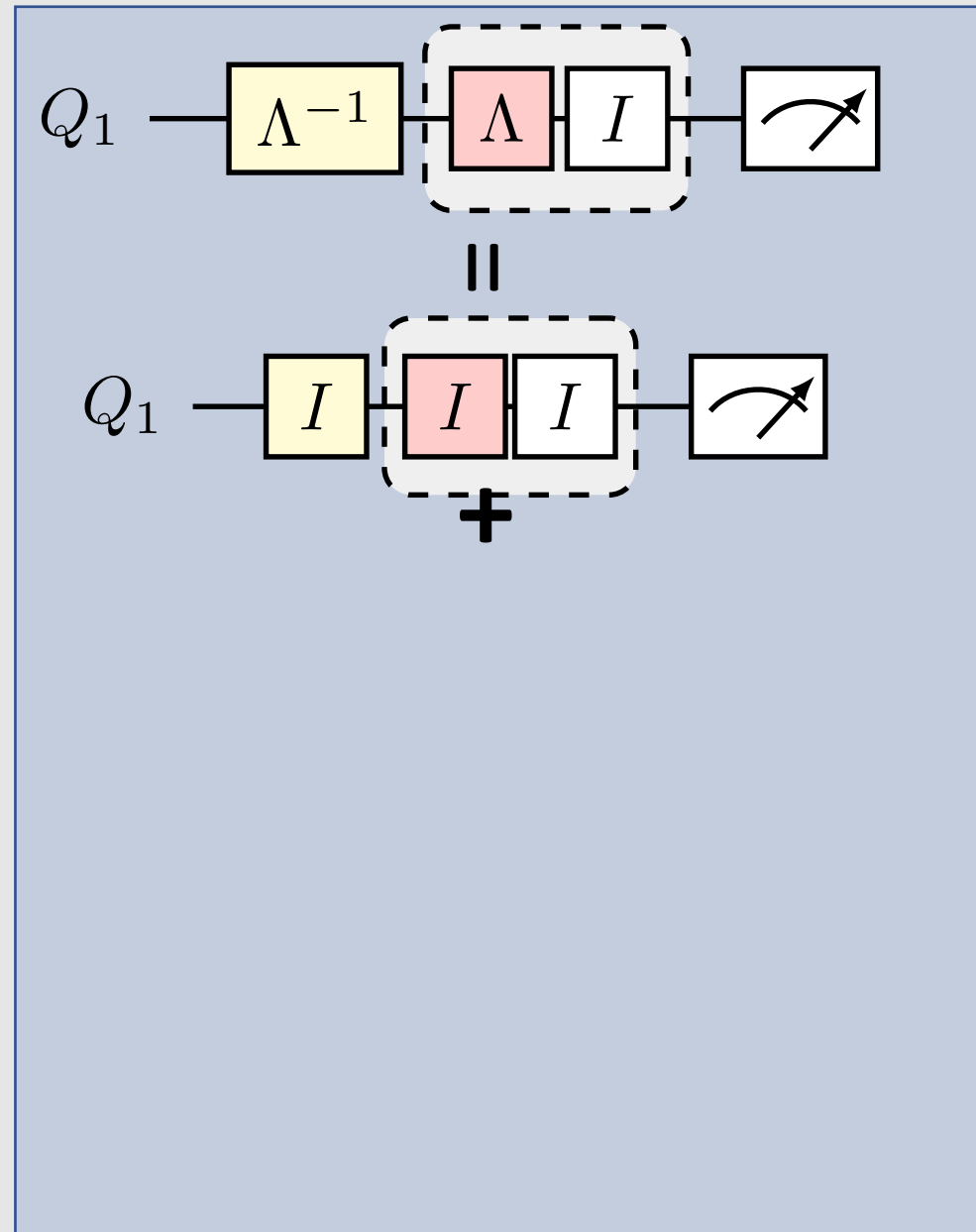
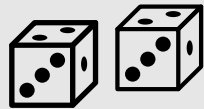
$$\Lambda^{-1}(\rho) = ?$$



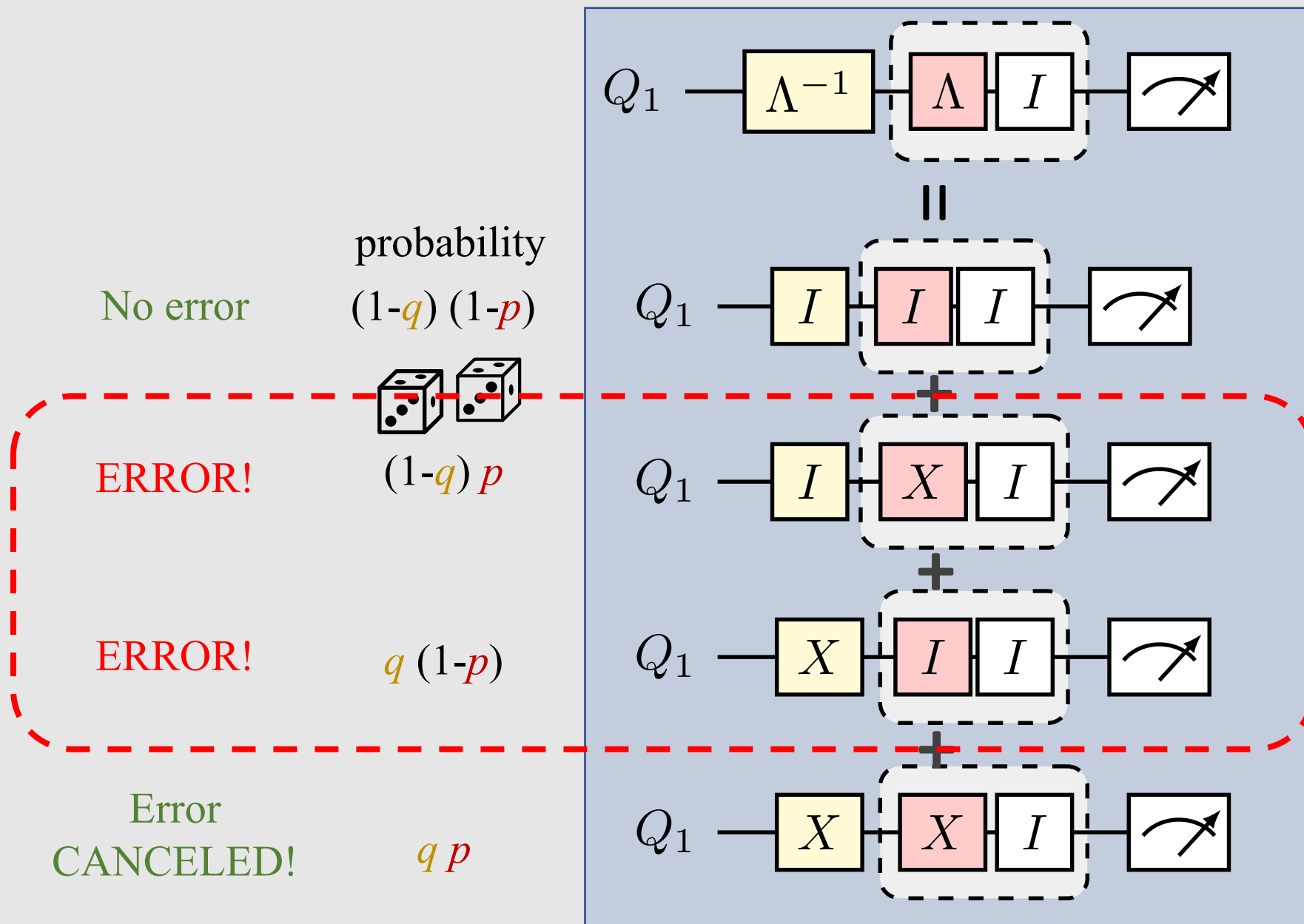
Raveling quantum trajectories to undo noise

No error

probability
 $(1-q)(1-p)$



Raveling quantum trajectories to undo noise

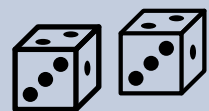


Solution to noise free!

$$q = \frac{-p}{1 - 2p}$$

Interfere
Destructively

Raveling quantum trajectories to undo noise

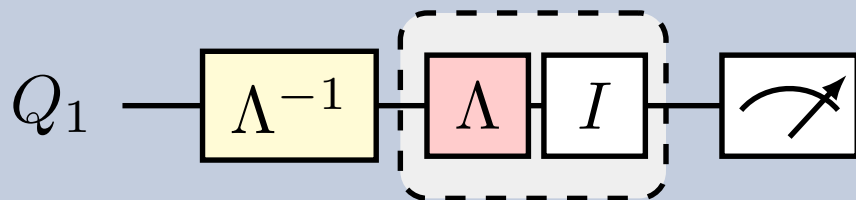


probability

$$P_I = |1 - q|/\gamma$$

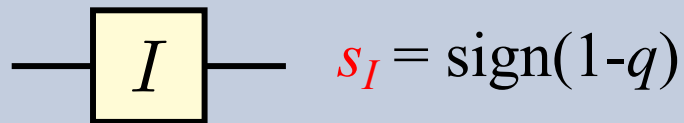
$$P_X = |q|/\gamma$$

$$P_I + P_X = 1$$

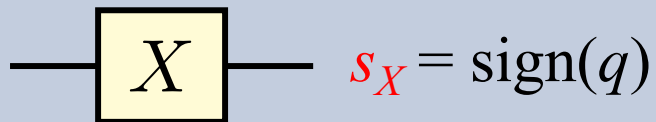


||

sign



+



Solution to noise free!

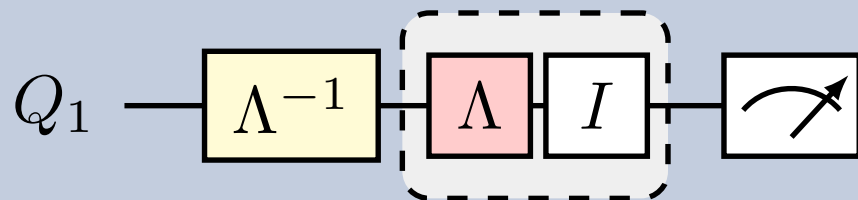
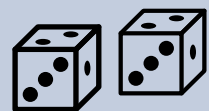
$$q = \frac{-p}{1 - 2p}$$

norm

$$\gamma = |1 - q| + |q|$$

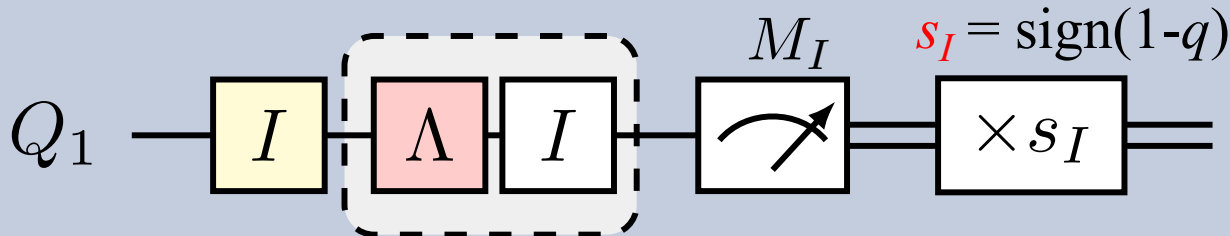
Quasi-probability

Raveling quantum trajectories to undo noise

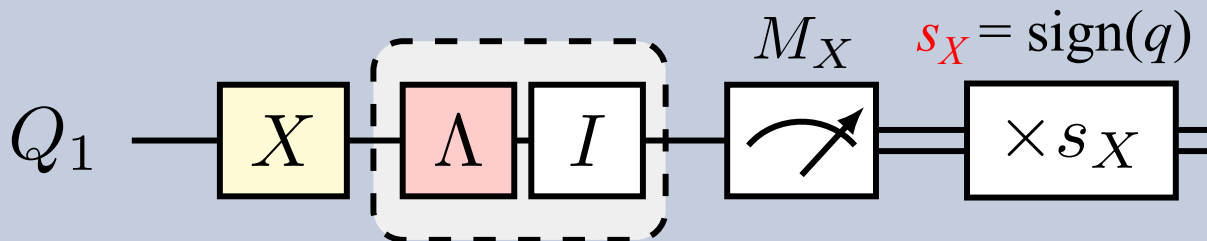


||

$P_I = |1 - q|/\gamma$



$P_X = |q|/\gamma$



mitigated expectation

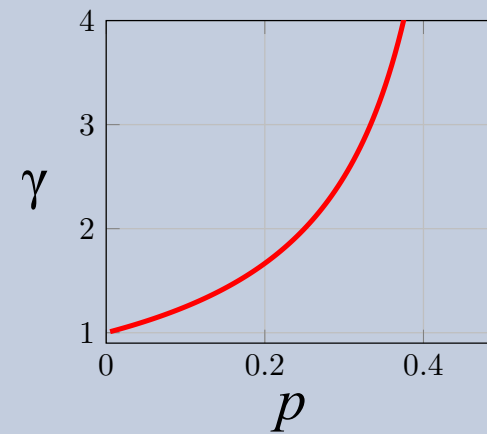
$$\langle M \rangle = \gamma (s_I P_I M_I + s_X P_X M_X)$$

Gain: Bias-free estimate!

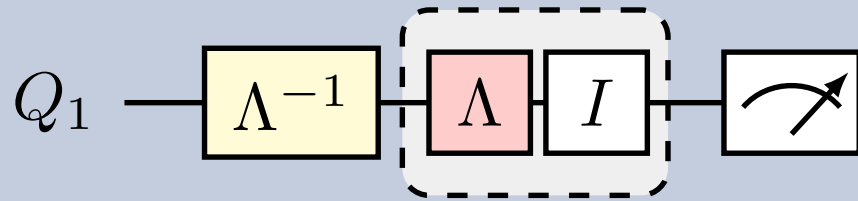
Cost: Variance

Sampling overhead

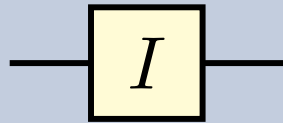
$$\gamma = |1 - q| + |q|$$



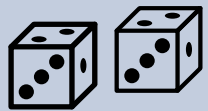
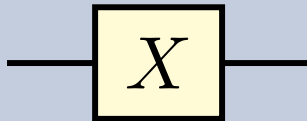
Canceling noise with noise



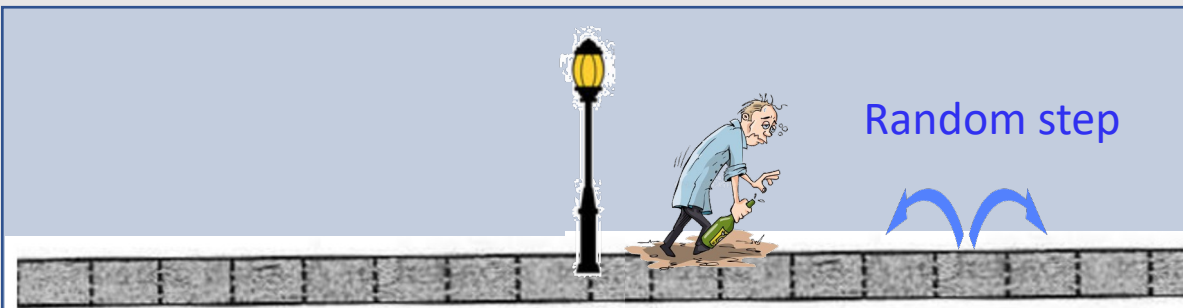
||



+

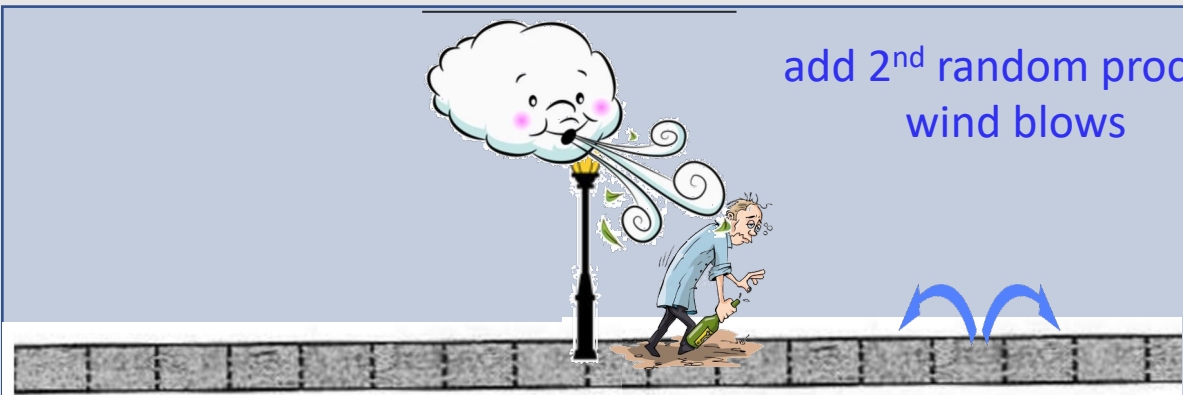
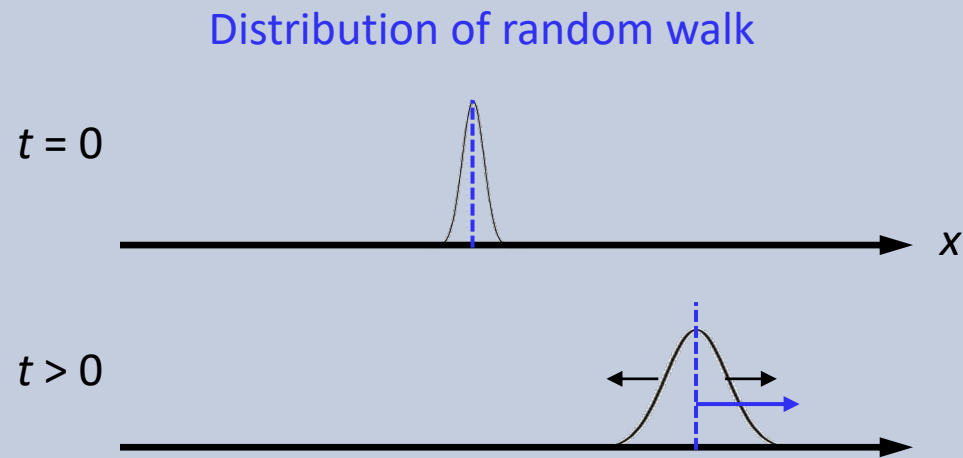


Canceling noise with noise: Drunkard's classical random walk analogy



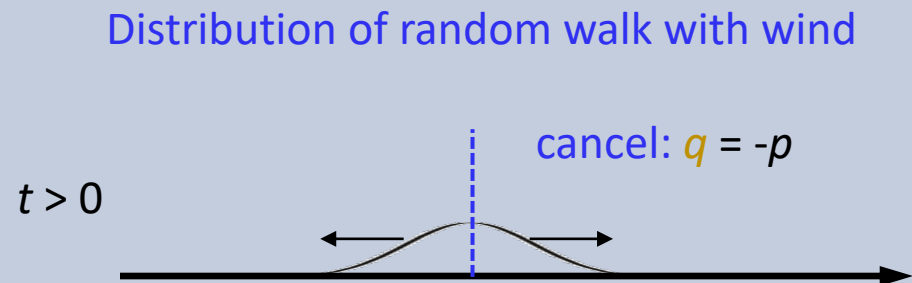
$$P(1 \text{ step left}) = \frac{1}{2} - p$$

$$P(1 \text{ step right}) = \frac{1}{2} + p$$



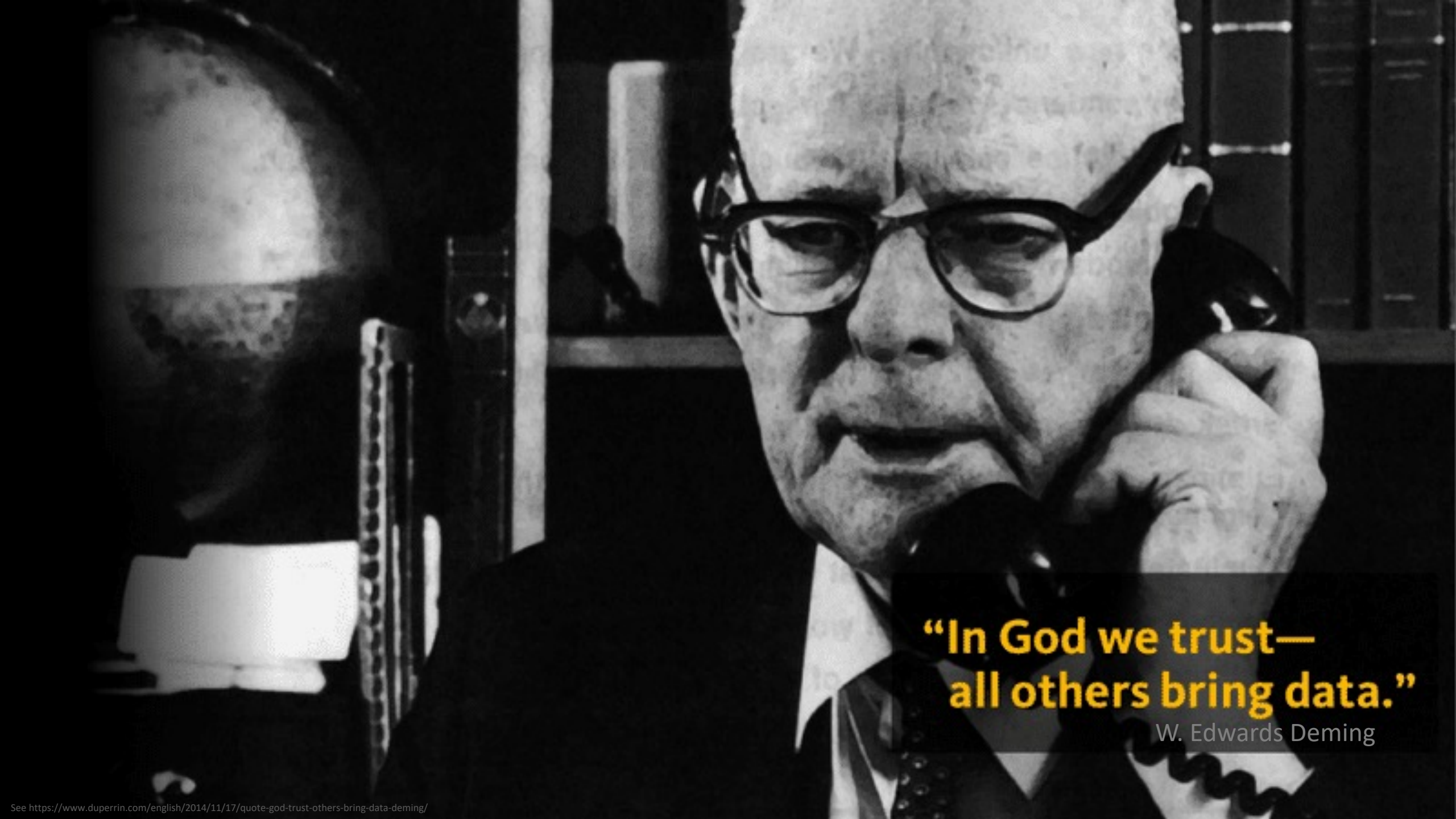
$$P(1 \text{ step left}) = \frac{1}{2} - q$$

$$P(1 \text{ step right}) = \frac{1}{2} + q$$



Gain: Bias-free estimate!

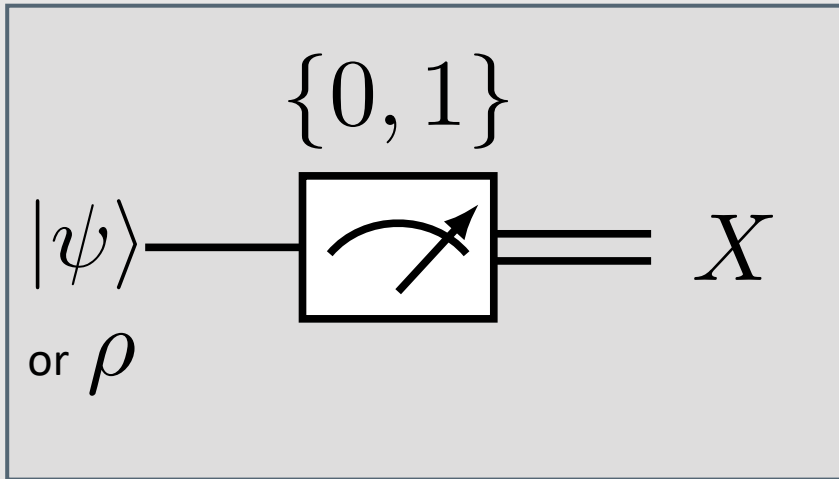
Cost: Variance

A black and white photograph of W. Edwards Deming, an older man with glasses, wearing a suit and tie, talking on a rotary telephone. He is looking slightly to the side with a thoughtful expression. The background shows a bookshelf with several books.

**“In God we trust—
all others bring data.”**

W. Edwards Deming

Summary: Qubit measured in the computational basis



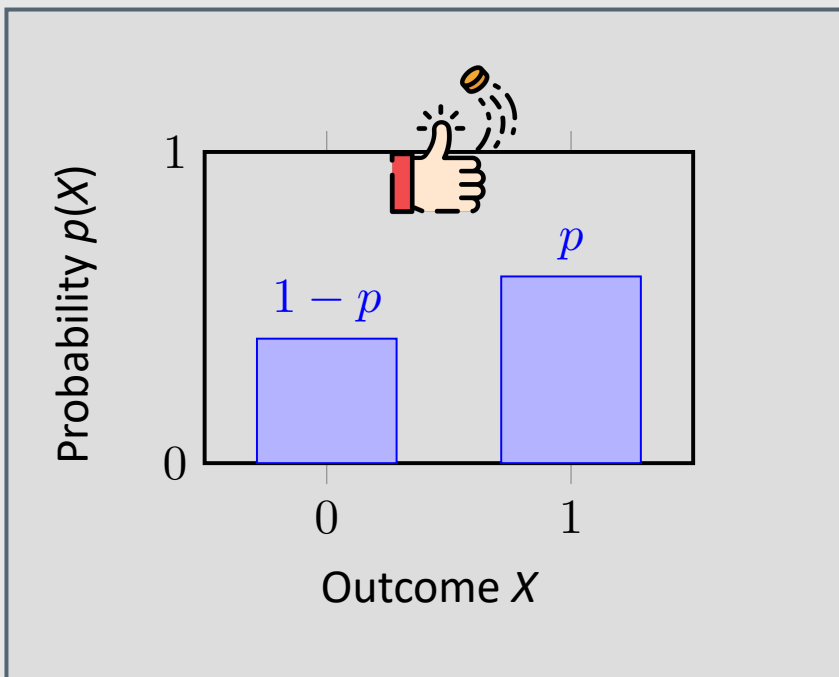
Set of possible outcomes

$$\Sigma = \{0, 1\}$$

$$X \in \Sigma$$

Measurement operators

$$\begin{array}{l} x \quad \hat{\mu}(x) \\ 0 : \quad \hat{\mu}(0) = |0\rangle \langle 0| \\ 1 : \quad \hat{\mu}(1) = |1\rangle \langle 1| \end{array}$$



Probability to measure outcome

$$\begin{cases} X = 0 : & p(X = 0) = \text{Tr}(\hat{\mu}(0)^\dagger \rho) = \text{Tr}(|0\rangle \langle 0| \rho) = \frac{1}{2} (1 + \langle Z \rangle) \\ X = 1 : & p(X = 1) = \text{Tr}(\hat{\mu}(1)^\dagger \rho) = \text{Tr}(|1\rangle \langle 1| \rho) = \frac{1}{2} (1 - \langle Z \rangle) \end{cases}$$

Bernoulli distribution. Single shot outcome follows a Bernoulli distribution:

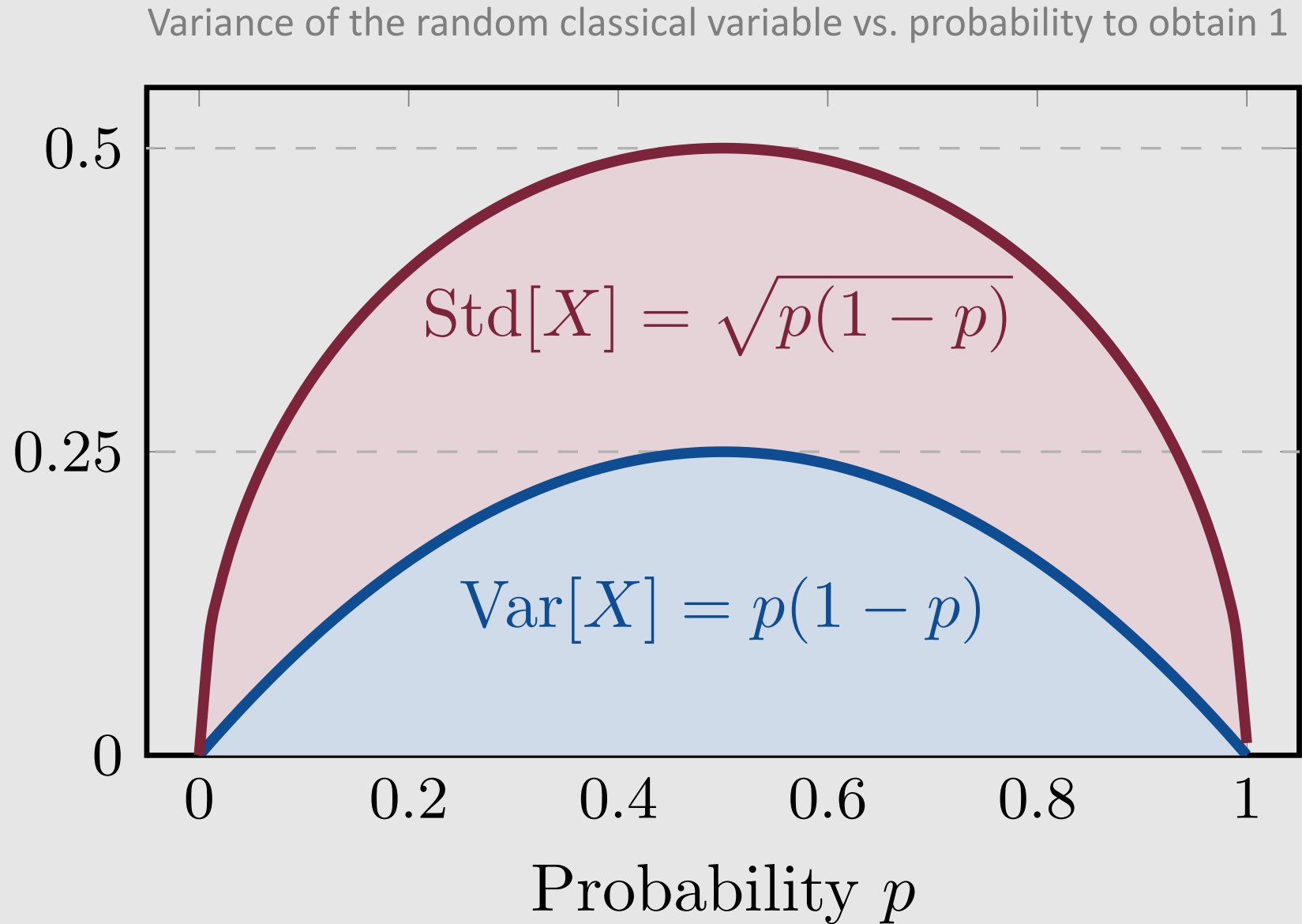
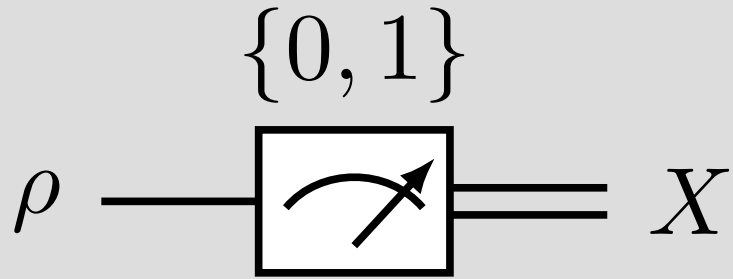
$$X \sim \text{Bernoulli}(p)$$

$$p := \text{Tr}(|1\rangle \langle 1| \rho) \in [0, 1]$$

$$\mathbb{E}[X] = p$$

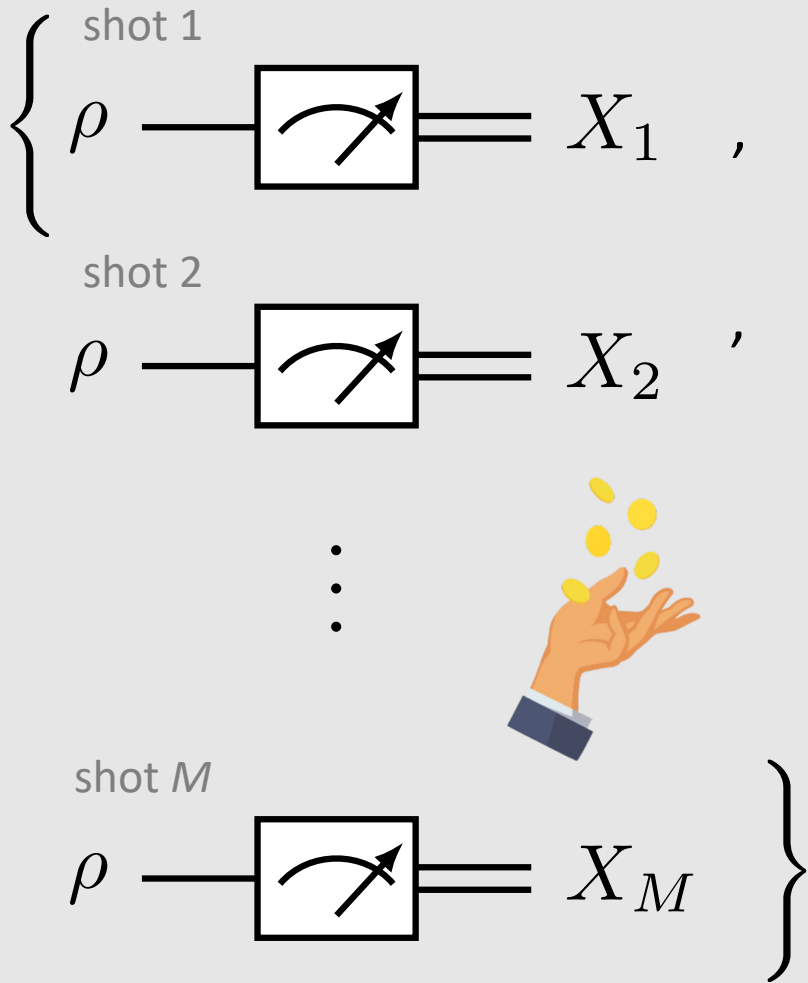
$$\text{Var}[X] = p(1 - p)$$

Quantum projection noise for a single shot





Ideal single qubit measurement with M shots



M shots with IID distribution

M outcomes: independent and identically distributed (iid) random variables

$$X_1, X_2, \dots, X_M \in \Sigma$$

$$X_1, X_2, \dots, X_M \sim \Pr[X = x] = \langle \hat{\mu}(x) \rangle$$

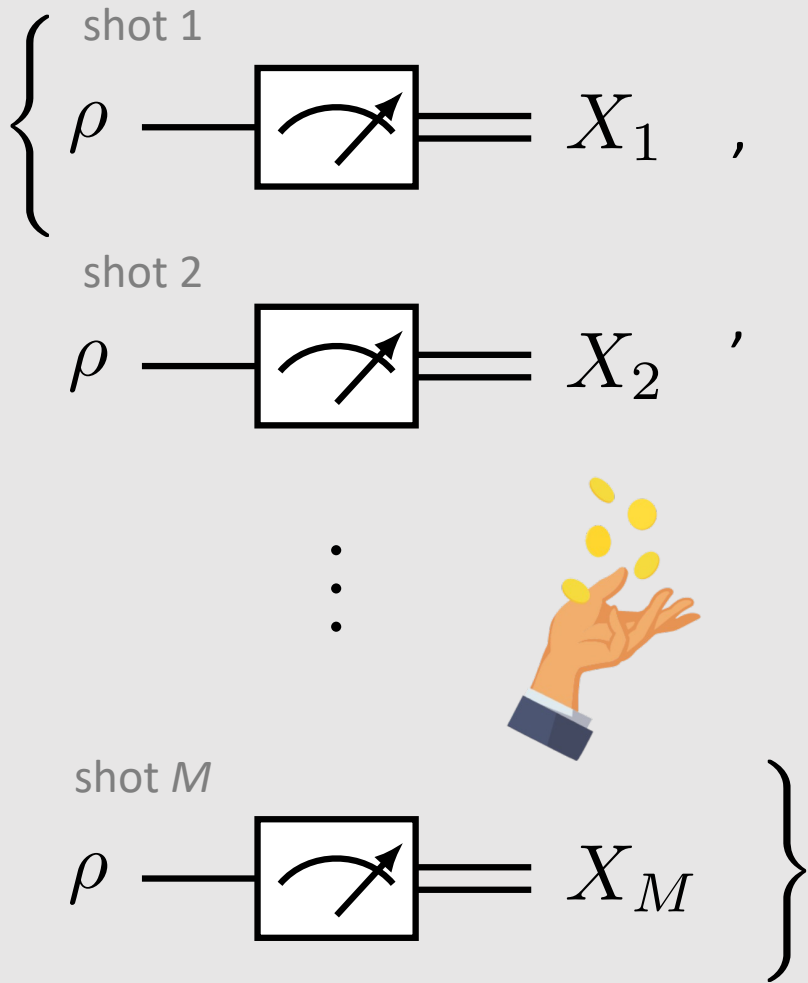
Empirical mean random variable (sample mean statistic)

$$S = \frac{1}{M} \sum_{m=1}^M X_m$$


Find the expectation value and variance of the empirical mean



Empirical mean: an unbiased estimator



$$S = \frac{1}{M} \sum_{m=1}^M X_m$$

Find the expectation value and variance of the empirical mean 

$$\mathbb{E}[S] = \mathbb{E} \left[\frac{1}{M} \sum_{m=1}^M X_m \right]$$

$$\mathbb{E}[X_m] = \mathbb{E}[X] = p \quad \forall m \in \{1, \dots, M\}$$

$$\mathbb{E}[aX_m + bX_n] = a\mathbb{E}[X_m] + b\mathbb{E}[X_n]$$

$$\forall m, n, \quad a, b \in \mathbb{C}$$

$$= \frac{1}{M} \sum_{m=1}^M \mathbb{E}[X]$$

linear functional

$$= \mathbb{E}[X]$$

expectation value of a single shot

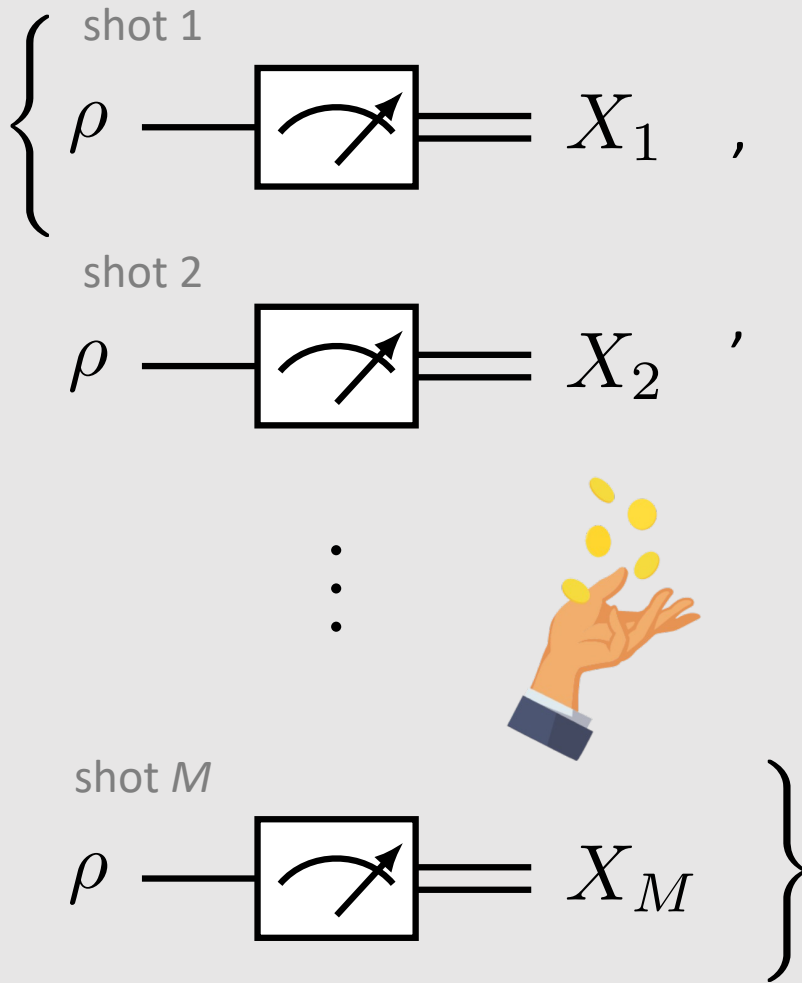
$$= \langle \hat{M} \rangle$$

relation to quantum operator we derived earlier (unbiased estimator)

$$= p$$

relation to probability to measure 1

How noisy is our estimate of the empirical mean?




$$S = \frac{1}{M} \sum_{m=1}^M X_m$$

$$\begin{aligned} \mathbb{V}[S] &= \mathbb{V}\left[\frac{1}{M} \sum_{m=1}^M X_m\right] \\ &= \frac{1}{M^2} \sum_{m=1}^M \mathbb{V}[X] \end{aligned}$$

$$\begin{aligned} &= \frac{1}{M} \mathbb{V}[X] \\ &= \frac{p(1-p)}{M} \end{aligned}$$

$$\begin{aligned} \sigma_S &= \sqrt{\mathbb{V}[S]} \\ &= \sqrt{\frac{\mathbb{V}[X]}{M}} \\ &= \sqrt{\frac{p(1-p)}{M}} \end{aligned}$$

Find the expectation value and variance of the empirical mean 

Use key identity for variance

$$\mathbb{V}[aX_m + bX_n] = a^2\mathbb{V}[X_m] + b^2\mathbb{V}[X_n]$$

(you can derive this from the definition)

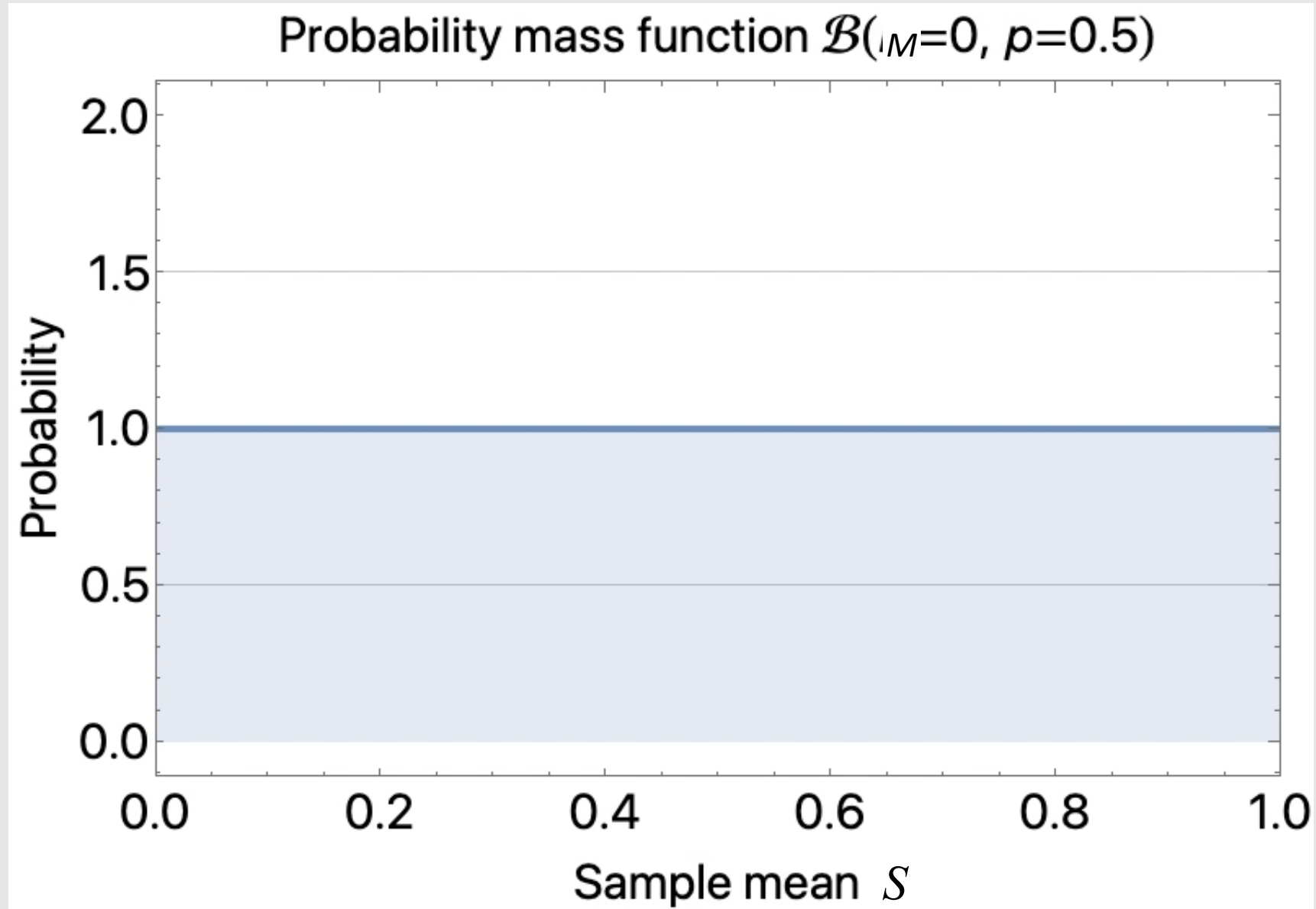
$$\mathbb{V}[X_m] = \mathbb{V}[X] = p(1-p) \quad \forall m \in \{1, \dots, M\}$$

The variance is reduced by the number of sampler we take!

Thus we can suppress the quantum projection noise with enough shots.

The standard deviation drop as one over square root of the number of shots

Animation of convergence of shots expectation value and mean

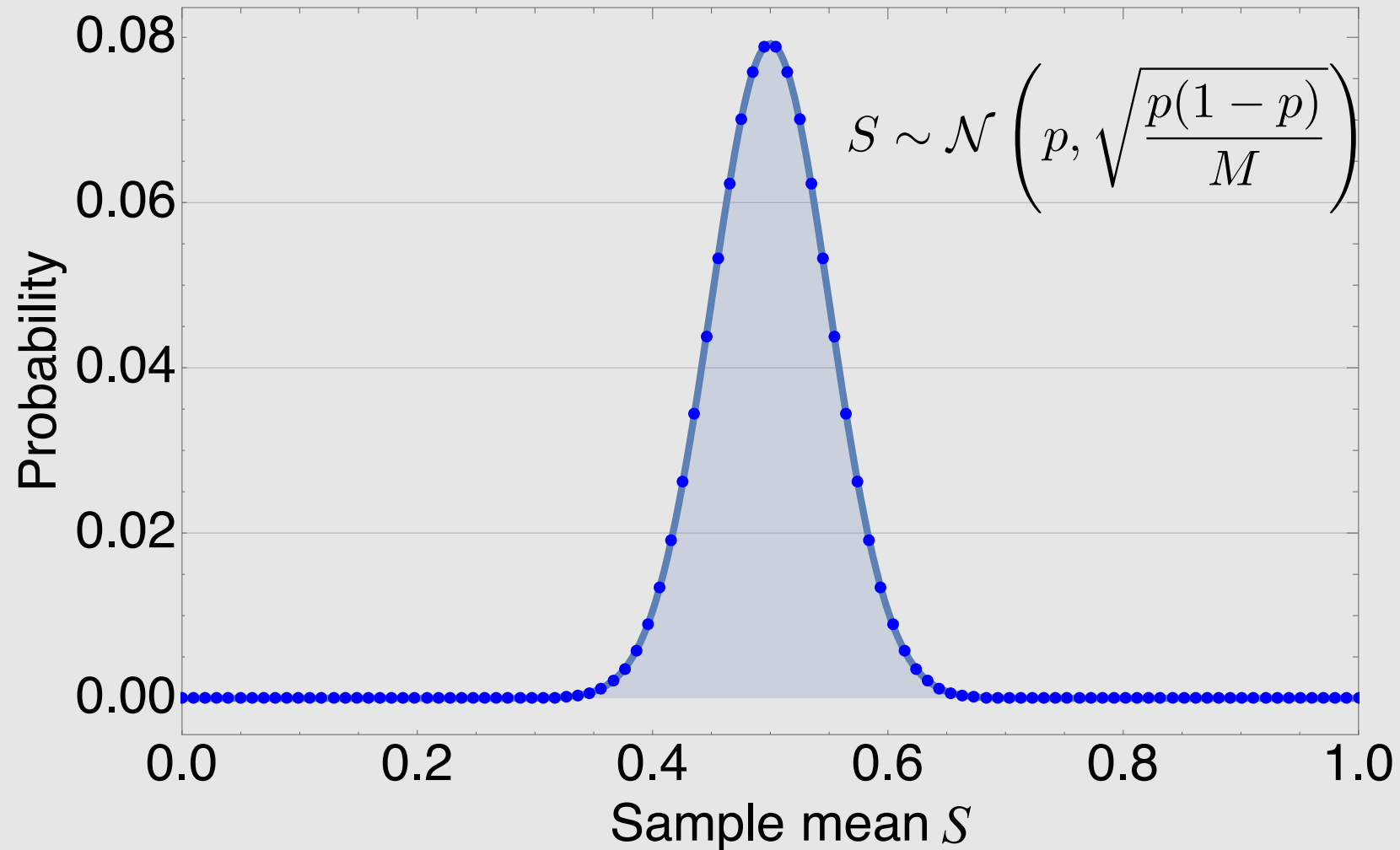


Binomial distribution

$$\binom{M}{k} p^k (1-p)^{M-k}$$

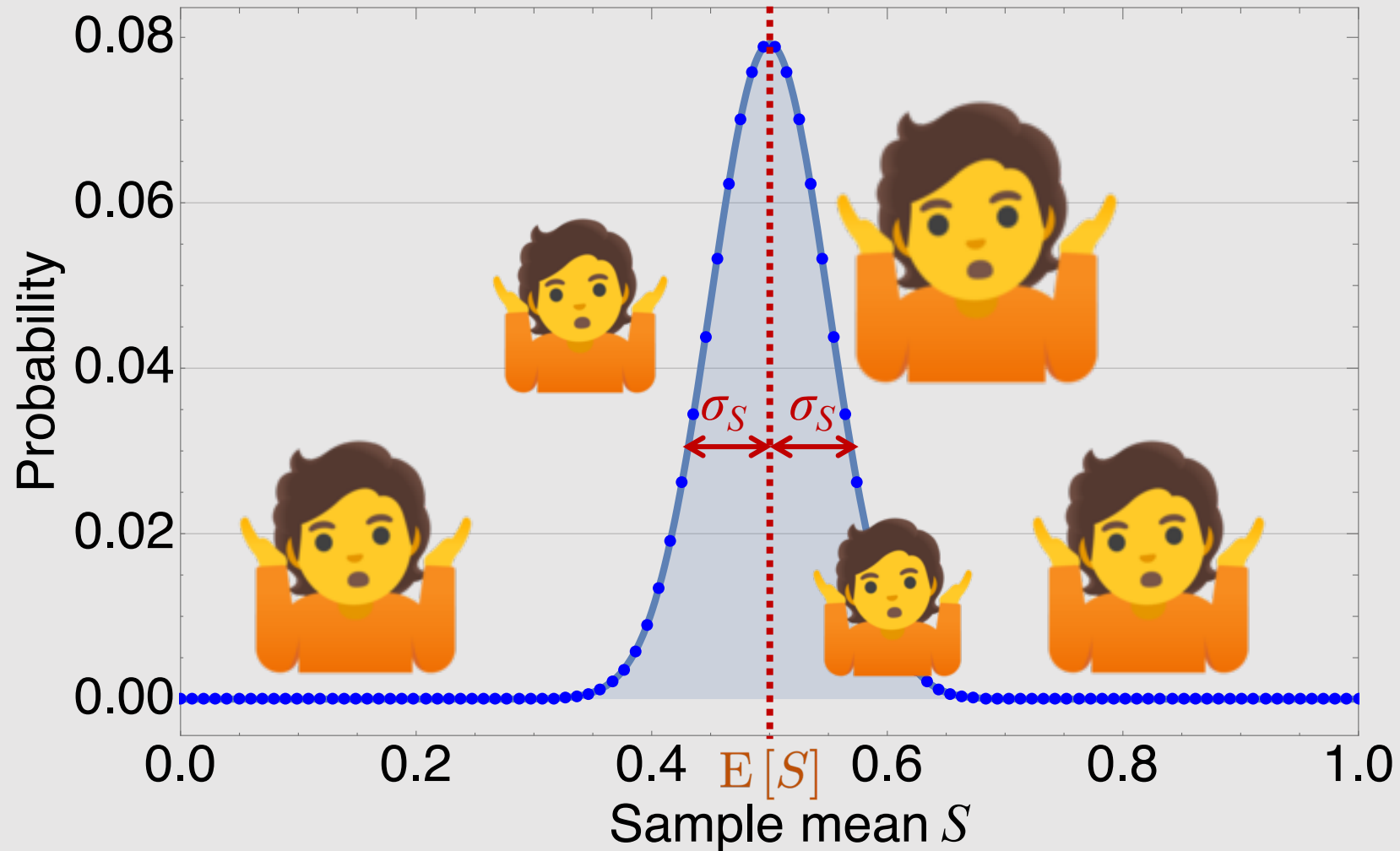
Sampled output distribution

Probability mass function $\mathcal{B}(M = 101, p = 0.5)$



Properties of the output distribution

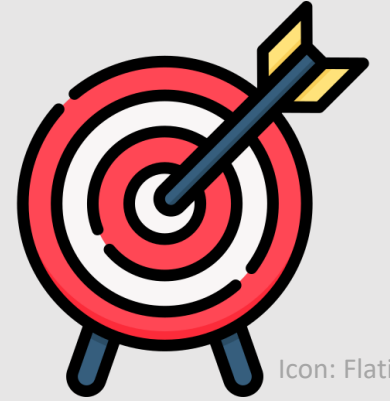
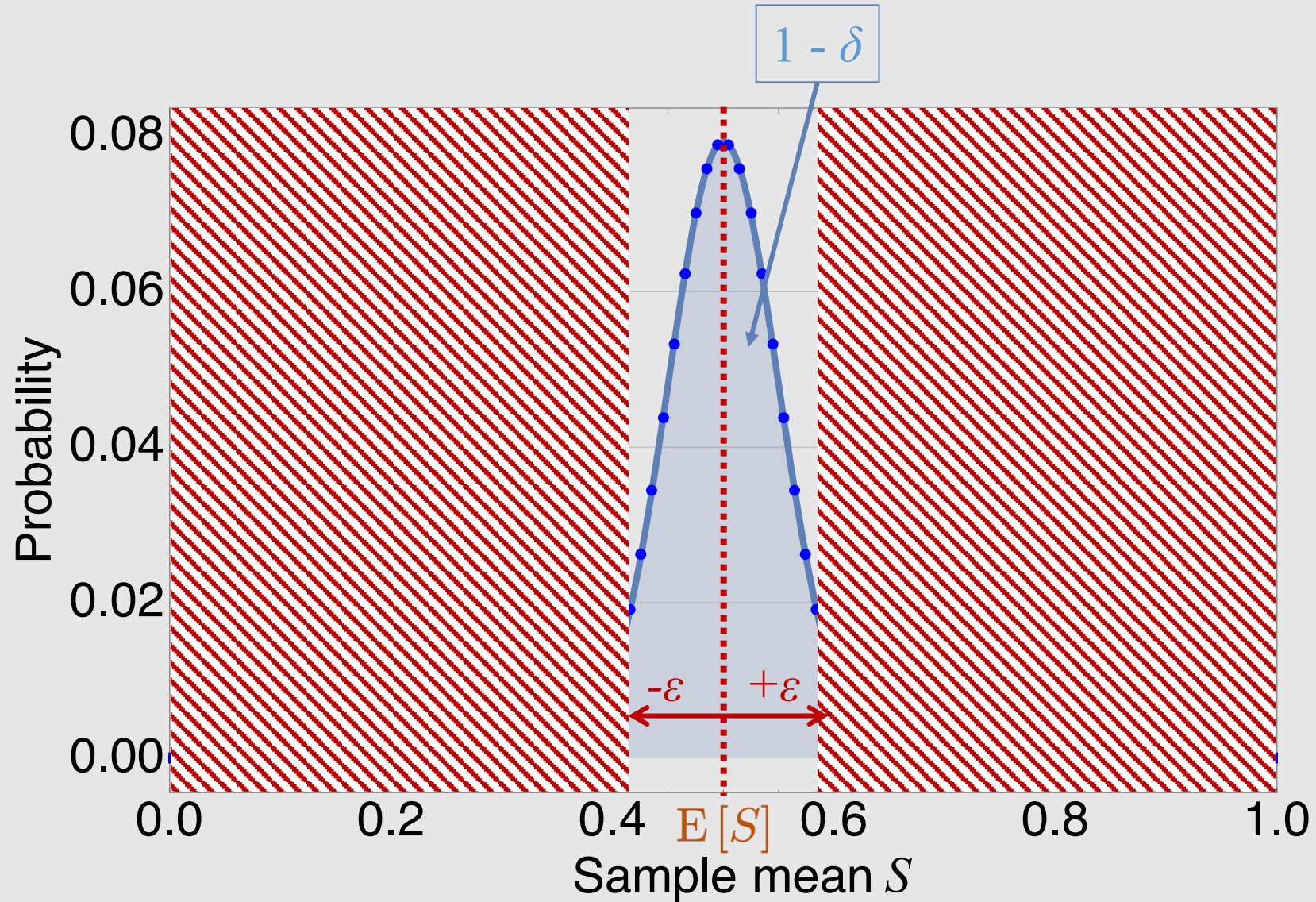
Probability mass function $\mathcal{B}(M = 101, p = 0.5)$



$$E[S] = p$$

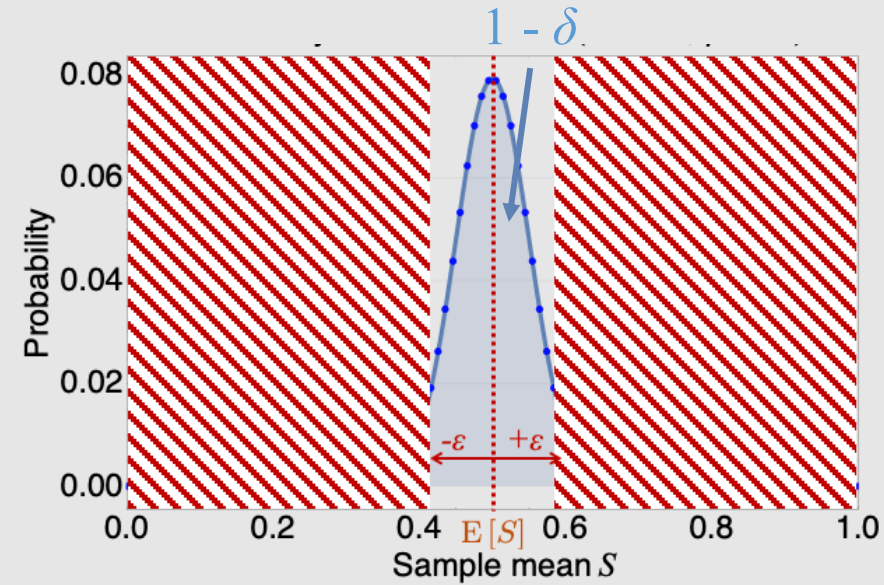
$$\begin{aligned}\sigma_S &= \sqrt{\text{Var}[S]} \\ &= \sqrt{\frac{p(1-p)}{M}}\end{aligned}$$

Concentration Measure for Sampling Expectation Values

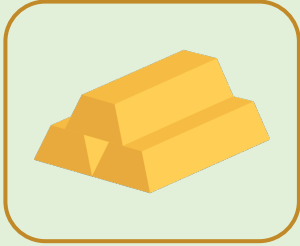


Error Bound on Quantum Expectation Values

Concentration Measure for Sampling Expectation Values



Chernoff-Hoeffding two-sided tail bound, given $|O(x)| \leq 1$ for ϵ specified (additive) precision (worst case additive error) for S with success probability at least $1 - \delta$.



$$\Pr \left[\left| S - \langle \hat{O} \rangle \right| > \epsilon \right] \leq \delta := 2 \exp \left(-\frac{1}{2} M \epsilon^2 \right)$$

↑ δ specific failure probability for meeting precision ϵ empirically.

Empirical mean & sample properties

$$S = \frac{1}{M} \sum_{m=1}^M O(X_m)$$

Tf. required number of shots is at least [or with high probability (greater than 2/3)]

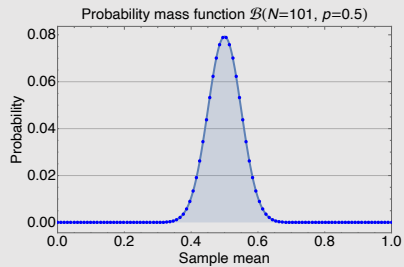
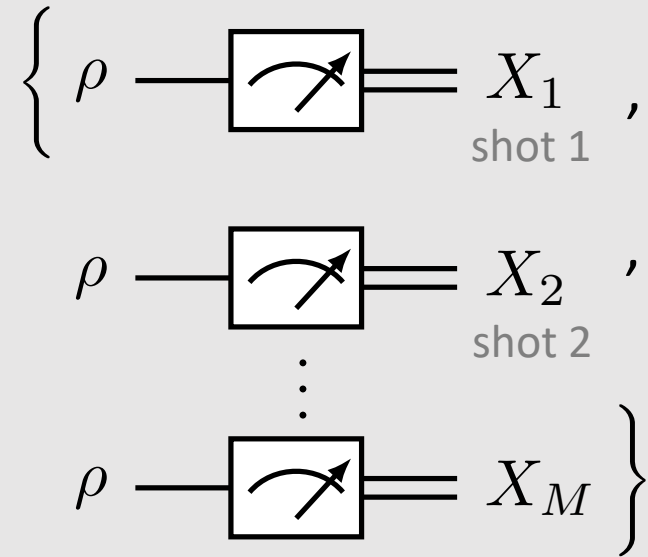
$$M \geq 2\epsilon^{-2} \log(2\delta^{-1}) \quad \left[M \gtrsim 4\epsilon^{-2} \right] .$$

* Can find even tighter bound here owing to smaller [0,1] range

Note that this scales same way (mod δ) as the variance bound with $\epsilon = \sigma$:

$$M \geq \frac{1}{4} \epsilon^{-2}$$

Ideal single qubit measurement with M shots



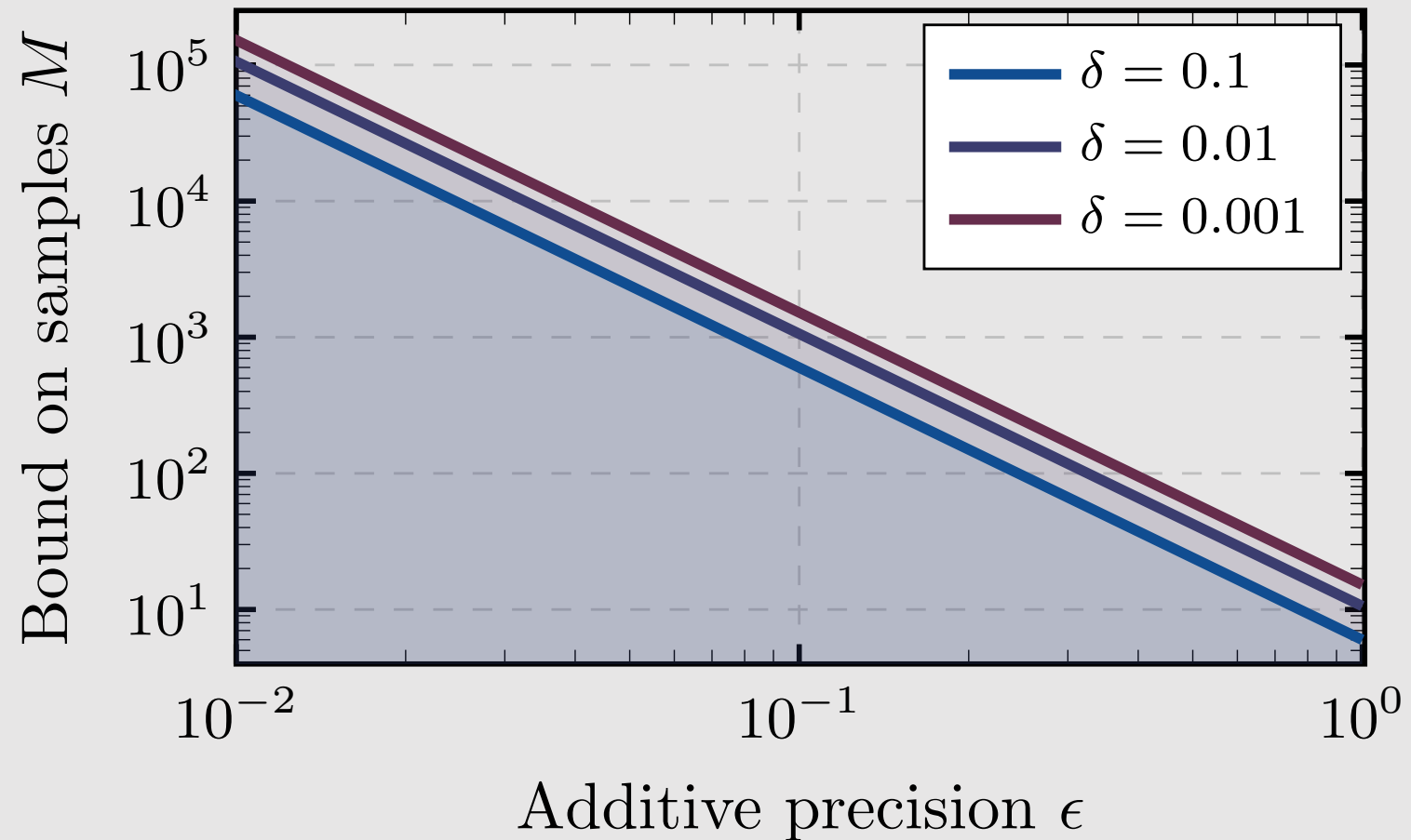
Observe that the **probability** δ is much cheaper than the **precision** ϵ .

Observe, n is not part of the equation.

“Knowing you are *not* wrong is cheaper than knowing you are right.” - Derek

Tail bound on sample complexity

$$M \geq 2\epsilon^{-2} \log(2\delta^{-1}) \quad [M \gtrsim 4\epsilon^{-2}]$$



Concentration inequalities and tail bounds



*Making a list,
checking it twice,
going to see
which inequality
is nice!*

*Markov? Hoeffding?
Jensen? Chebyshev?
Chernoff?*

<https://www.zlatko-minev.com/blog/inequalities>

1. Probability (Technical note 11.9 v0.6)

1A. Concentration inequalities and tail bounds

Unless otherwise specified, all variables are real \mathbb{R} . Inequalities come as one-sided $\Pr(\dots \leq \dots)$ and two-sided $\Pr(|\dots| \leq \dots)$. Notation: X is a random variable, $\mu := \mathbb{E}[X]$, $\sigma^2 := \text{Var}[X]$, $S_n := X_1 + \dots + X_n$.

Inequality	Conditions	Common form	Notes / Alternate form
Single random variable			
Markov¹	Non-negative $X \geq 0$	$\Pr[X \geq a] \leq \frac{\mathbb{E}[X]}{a}$	$\forall a > 0$ $\Pr[X \geq k\mathbb{E}[X]] \leq \frac{1}{k}$ $k > 1$ [3, Sec. 5.1][6, Thm 1.13]
extension	+ non-negative, strictly increasing func Φ $X \geq 0$ $\Phi(X) \geq 0$ increasing	$\Pr[X \geq a] = \Pr[\Phi(X) \geq \Phi(a)] \leq \frac{\mathbb{E}[\Phi(X)]}{\Phi(a)}$	$\forall a > 0$ Wiki
Reverse Markov	upper-bounded by U $\max X = U$ (can be positive)	$\Pr[X \leq a] \leq \frac{U - \mathbb{E}[X]}{U - a}$	$\forall a > 0$ [1, Sec. 3.1]
Chebyshev²	Finite mean and variance $\mathbb{E}[X], \text{Var}[X]$ finite	$\Pr[X - \mathbb{E}[X] \geq a] \leq \frac{\sigma^2}{a^2}$	$\Pr[X - \mathbb{E}[X] \geq a \cdot \sigma] \leq \frac{1}{a^2}$ $\forall a > 0, \sigma^2 = \text{Var}[X]$ [3, Sec. 5.1][2, Thm 18.11]
Cantelli	Improved Chebyshev (same; but one-sided)	$\Pr[X - \mathbb{E}[X] \geq a] \leq \frac{\sigma^2}{\sigma^2 + a^2}$	$\forall a > 0, \sigma^2 = \text{Var}[X]$ Wiki
Chernoff³	Generic	$\Pr[X \geq a] = \Pr[e^{tX} \geq e^{ta}]$	$\forall t > 0, a \in \mathbb{R}$ [1, Sec. 3.3]
Jensen		$f: \mathbb{R} \rightarrow \mathbb{R}; f$ convex $f(\mathbb{E}[X]) \leq \mathbb{E}[f(X)]$	[3, Prob. 5.3][6, Thm 1.14]
Hoeffding's lemma		$\mathbb{E}[X] = \mu$ $a \leq X \leq b$ $\mathbb{E}[e^{\lambda X}] \leq e^{\lambda\mu} e^{\frac{\lambda^2(b-a)^2}{8}}$	$\lambda \in \mathbb{R}$ [1, Sec. 3.4]
Sum of random variables			
Chernoff-Hoeffding (one-sided)	n independent random vars X_1, \dots, X_n indep $S_n = X_1 + \dots + X_n$ $X_i \in [a_i, b_i] \forall i$	$\Pr[S_n - \mathbb{E}[S_n] \geq t] \leq \exp\left(\frac{-2t^2n^2}{\sum_{i=1}^n (b_i - a_i)^2}\right)$	[1, Sec. 3.5]
(two-sided) ⁴	(same as above)	$\Pr[S_n - \mathbb{E}[S_n] > t] \leq 2 \exp\left(\frac{-2t^2}{\sum_{i=1}^n (b_i - a_i)^2}\right)$	$\forall t \in (0, \frac{1}{2})$ [5, Thm.1.1]
(two-sided iid)	same plus iid, range, mean μ for each $X_1, \dots, X_n \in [0, 1]$ $\mathbb{E}[X_i] = \mu$ iid	$\Pr\left[\left \frac{S_n}{n} - \mu\right \geq \epsilon\right] \leq 2 \exp(-2n\epsilon^2)$	$\forall \epsilon > 0$ [6, Thm 1.16]
Thm 1.3	n independent random vars X_1, \dots, X_n indep $S_n = X_1 + \dots + X_n$	$\Pr[S_n - \mathbb{E}[S_n] > \epsilon] \leq 2 \exp\left(\frac{-\epsilon^2}{4 \sum_{i=1}^n \text{Var}[X_i]}\right)$	$\epsilon \in (0, 2 \text{Var}[S_n] / (\max_i X_i - \mathbb{E}[X_i]))$ [5, Thm. 1.3]
Azuma			
Weak law of large numbers	n independent iid random vars X_1, \dots, X_n indep $\mathbb{E}[X_i] = \mu$ iid	$\lim_{n \rightarrow \infty} \Pr\left[\left \frac{1}{n}S_n - \mu\right \geq \epsilon\right] = 0$	$\forall \epsilon > 0$ [3, Sec. 5.2][6, Thm 1.15]
Strong law of large numbers	(same)	(same)	$\Pr\left[\lim_{n \rightarrow \infty} \frac{1}{n}S_n = \mu\right] = 1$ [3, Sec. 5.5]
Advanced			
Bennett	n independent zero-mean X_1, \dots, X_n indep $\mathbb{E}[X_i] = 0$ iid	$\Pr[S_n > \epsilon] \leq \exp\left(-n\sigma^2 h\left(\frac{\epsilon}{n\sigma^2}\right)\right)$	$\sigma^2 := \frac{1}{n} \sum_{i=1}^n \text{Var}[X_i], \forall \epsilon > 0,$ $h(a) := (1+a) \log(1+a) - a$ for $a \geq 0$ [1, 4.1]
Bernstein	(same)	(same)	(same) [1, 4.2]
Efron-Stein	scalar func of vars $f: \mathcal{X}^n \rightarrow \mathbb{R}$	X_1, \dots, X_n indep w/ values in set \mathcal{X} $\text{Var}[Z] \leq \sum_{i=1}^n \mathbb{E}\left[\left(Z - \mathbb{E}_i[Z]\right)^2\right]$	$Z := g(X_1, \dots, X_n)$ $\mathbb{E}_i[Z] := \mathbb{E}[Z X_1, \dots, X_{i-1}, X_{i+1}, \dots, X_n]$ [1, 4.3]
McDiarmid's	scalar func of vars $f: \mathcal{X}^n \rightarrow \mathbb{R}$	X_1, \dots, X_n indep w/ values in set \mathcal{X} $\Pr[f(X_1, \dots, X_n) - \mathbb{E}[f(X_1, \dots, X_n)] \geq \epsilon] \leq \exp\left(\frac{-2\epsilon^2}{\sum_{i=1}^n c_i^2}\right)$	condition: c -bounded difference property $\forall \epsilon > 0$ $ f(X_1, \dots, X_i, \dots, X_n) - f(X_1, \dots, X_i', \dots, X_n) \leq c_i$ [1, 4.4]

¹Markov's inequality bounds the first moment of random variable. Use it when a constant probability bound is sufficient [1, Sec. 3.3].

²Chebyshev is derived from Markov. It bounds the second moment. It is the appropriate one when the variance σ is known. If σ is unknown, we can use the bounds of $X \in [a, b]$.

³Chernoff bound is used to bound the tails of the distribution for a sum of independent random variables. By far the most useful tool in randomized algorithms [1, Sec. 3.3].

⁴This probability can be interpreted as the level of significance ϵ (probability of making an error) for a confidence interval around the mean of size 2ϵ . Therefore, we require at least $\log(2\alpha)/2t^2$ samples to acquire $1 - \alpha$ confidence interval $\mathbb{E}[X] \pm t$.

Scaling PEC to n qubits and larger circuits (ADVANCED – OPTINAL)

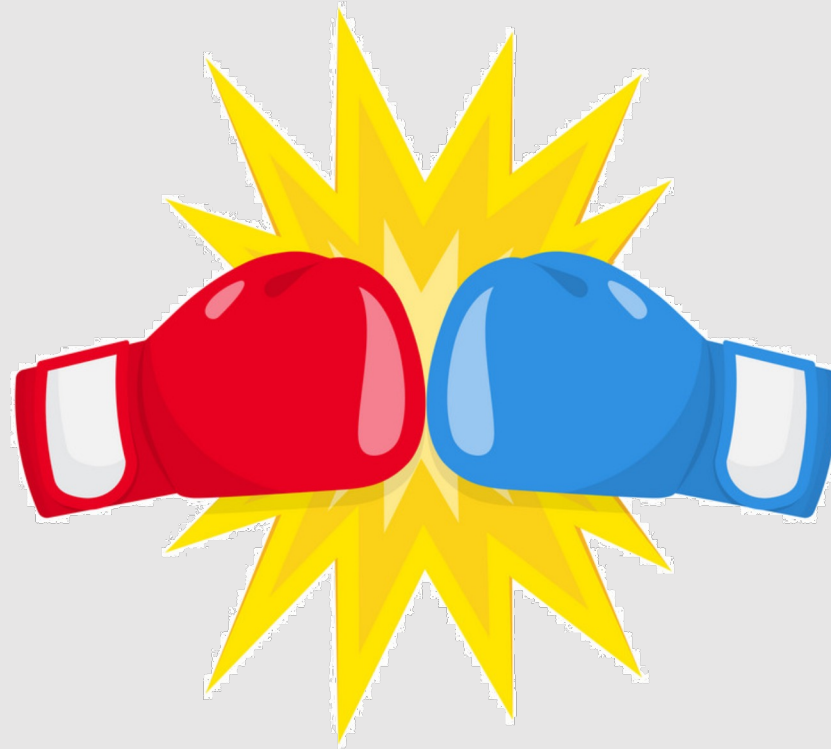


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Language of errors and error mitigation

$$\frac{d}{dx} \frac{1}{x} = -\frac{1}{x^2}$$

@formulas_for_your_comfort

Notation: Super bra ket

$$\rho \in L(\mathcal{H})$$

$$\hat{A}, \hat{B} \in L(\mathcal{H})$$

$$\mathcal{S} \in L(L(\mathcal{H}))$$

$$|\hat{A}\rangle\rangle \leftrightarrow \hat{A}$$

$$\mathcal{S}|\hat{A}\rangle\rangle \leftrightarrow \mathcal{S}(\hat{A})$$

Vectorization isomorphism

$$\begin{aligned} \langle\langle \hat{A} | \hat{B} \rangle\rangle &\leftrightarrow \langle \hat{A}, \hat{B} \rangle \\ &= \text{Tr}(\hat{B} \hat{A}^\dagger) \end{aligned}$$

$$\begin{aligned} \langle\langle \hat{A} | \mathcal{S} | \hat{B} \rangle\rangle &= \langle\langle \hat{A} | \mathcal{S}(\hat{B}) \rangle\rangle \\ &= \langle \hat{A}, \mathcal{S}(\hat{B}) \rangle \\ &= \text{Tr}(\hat{A}^\dagger \mathcal{S}(\hat{B})) \end{aligned}$$

$$\begin{aligned} \langle\langle \hat{A} | \cdot \rangle\rangle &\leftrightarrow \langle \hat{A}, \cdot \rangle \\ &= \text{Tr}(\hat{A}^\dagger \cdot) \end{aligned}$$

$$\begin{aligned} |\hat{A}\rangle\rangle \langle\langle \hat{B} | &\leftrightarrow \hat{A} \langle \hat{B}, \cdot \rangle \\ &\leftrightarrow \hat{A} \text{Tr}(\hat{B}^\dagger \cdot) \end{aligned}$$

Familiar ideas revisited with super notation

$$\langle \hat{P}_a | \hat{P}_b \rangle = \text{Tr} \left(\hat{P}_a^\dagger \hat{P}_b \right) = d \delta_{ab}$$

$$\Gamma = \{I, X, Y, Z\}^{\otimes n}$$

$$|\Gamma| = 4^n = d^2, \quad d = 2^n$$

$$a, b \in \Gamma$$

$$\mathcal{I} = \sum_{a \in \Gamma} \frac{|\hat{P}_a\rangle \langle \hat{P}_a|}{\langle \hat{P}_a | \hat{P}_a \rangle}$$

$$\Lambda = \mathcal{I} \Lambda \mathcal{I} = \sum_{a, b \in \Gamma} \frac{\langle \hat{P}_a | \Lambda | \hat{P}_b \rangle}{\langle \hat{P}_a | \hat{P}_a \rangle \langle \hat{P}_b | \hat{P}_b \rangle} |\hat{P}_a\rangle \langle \hat{P}_b| = \frac{1}{d} \sum_{a, b \in \Gamma} \Lambda_{ab} |\hat{P}_a\rangle \langle \hat{P}_b| ,$$

Stochastic Pauli channel

$$\Lambda : L(\mathcal{H}) \rightarrow L(\mathcal{H}) \quad 0 \leq p_a \leq 1 ,$$

$$\Lambda(\rho) = \sum_{a \in \Gamma} p_a P_a \rho P_a , \quad \sum_{a \in \Gamma} p_a = 1 ,$$

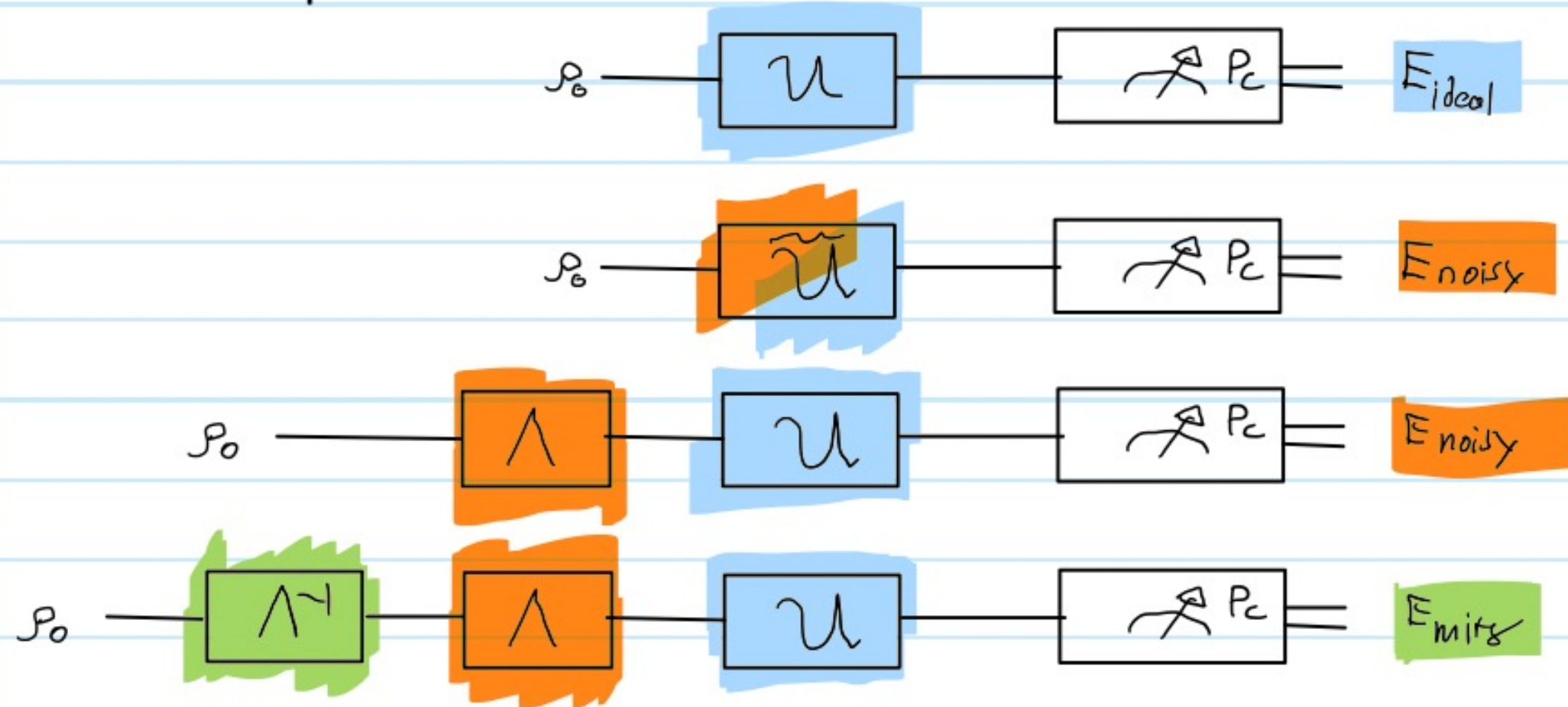
$$\Lambda(\rho) = \frac{1}{\sqrt{|\Gamma|}} \sum_{b \in \Gamma} f_b \text{Tr}(P_b \rho) P_b = \sum_{b \in \Gamma} f_b \frac{|P_b\rangle\rangle \langle\langle P_b|}{\langle\langle P_b|P_b\rangle\rangle} ,$$

$$\Lambda(P_b) = f_b P_b , \quad \forall a \in \Gamma ,$$

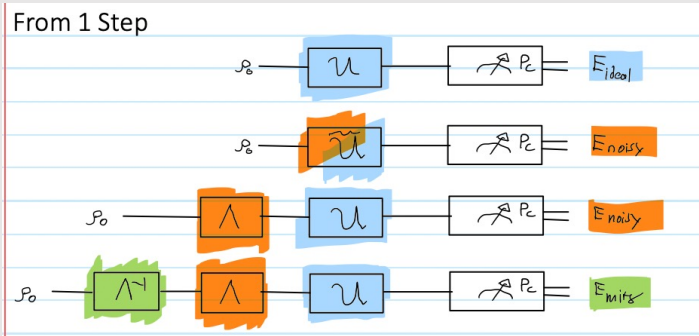
$$-1 \leq f_b \leq 1 \quad f_I = 1 ,$$

Probabilistic error cancellation: Derivation

From 1 Step



Probabilistic error cancellation: Derivation



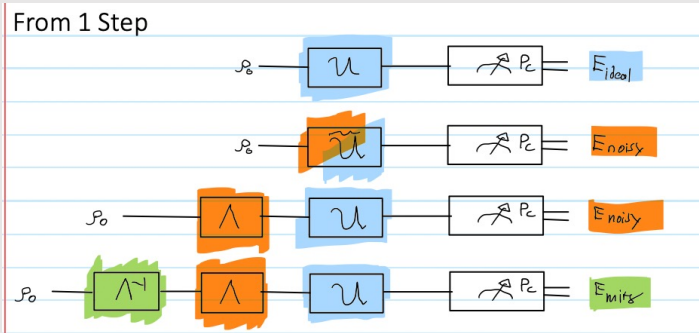
Channel definitions

$$U = U \cdot U^\dagger$$

$$\Lambda = \sum_a f_a |P_a\rangle\langle P_a|$$

$$= \sum_b c_b P_a$$

Probabilistic error cancellation: Derivation



Channel definitions

$$U = U \cdot U^\dagger$$

$$\langle \hat{P}_c \rangle (\cdot) = \langle \langle P_c | \cdot \rangle \rangle \quad \text{Need}$$

$$\begin{aligned} \Lambda &= \sum_a f_a |P_a\rangle \langle P_a| \\ &= \sum_b c_b P_b \end{aligned}$$

$$-1 \leq f_a \leq 1$$

$$c_b \geq 0, \sum_0 c_b = 1$$

$$c_b = \frac{1}{2^n} \sum_a (-1)^{\langle a, b \rangle_{sp}} f_a$$

$$\vec{c}_b = W^{-1} \vec{f}_a \quad \vec{f}_a = W c_b$$

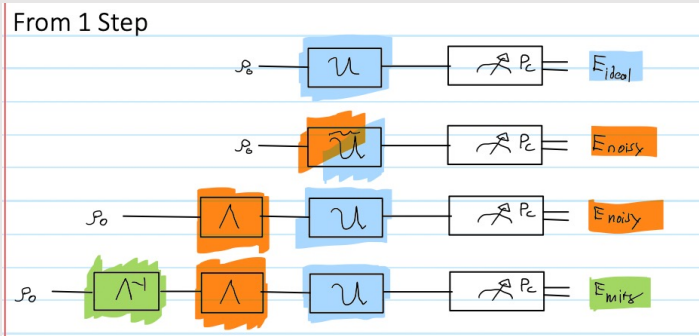
$$\begin{aligned} \Lambda^{-1} &= \sum_a f_a^{-1} |P_a\rangle \langle P_a| \\ &= \sum_b c_b^{inv} P_b \end{aligned}$$

$$\vec{c}_b^{inv} = W \vec{f}_a^{-1}$$

$$c_b^{inv} = \frac{1}{2^n} \sum_a (-1)^{\langle a, b \rangle_{sp}} f_a$$

$$c_b^{inv} \in \mathbb{R}$$

Probabilistic error cancellation: Derivation



Channel definitions

$$U = U \cdot U^\dagger$$

$$\langle \hat{P}_C \rangle(\cdot) = \langle P_C | \cdot \rangle \quad \text{Need}$$

$$\Lambda = \sum_a f_a |P_a\rangle \langle P_a|$$

$$-1 \leq f_a \leq 1$$

$$= \sum_b c_b P_a$$

$$c_b \geq 0, \sum_0 c_b = 1$$

$$c_b = \frac{1}{2^n} \sum_a (-1)^{\langle a, b \rangle_{sp}} f_a$$

$$\vec{c}_b = W^{-1} \vec{f}_a \quad \vec{f}_a = W c_b$$

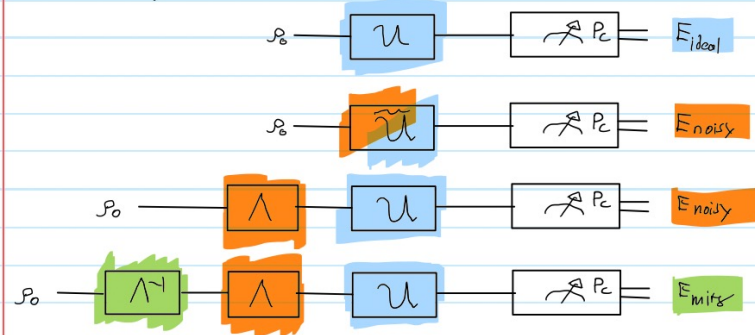
Circuit expectation values

$$E_{ideal} = \langle \hat{P}_c \rangle_{\rho_0}$$

$$= \langle\langle P_c | U | \rho_0 \rangle\rangle$$

ideal exp value with noiseless unitary

From 1 Step



Channel definitions

$$U = U \cdot U^\dagger$$

$$\Lambda = \sum_a f_a |P_a\rangle\langle P_a|$$

$$= \sum_b c_b P_b$$

$$\Lambda^{-1} = \sum_a f_a^{-1} |P_a\rangle\langle P_a|$$

$$= \sum_b c_b^{inv} P_b$$

$$\langle \hat{P}_c \rangle(\cdot) = \langle\langle P_c | \cdot \rangle\rangle$$

$$-1 \leq f_a \leq 1$$

$$c_b \geq 0, \sum_b c_b = 1$$

$$c_b = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

$$\vec{c}_b = W^{-1} \vec{f}_a \quad \vec{f}_a = W c_b$$

$$c_b^{inv} = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

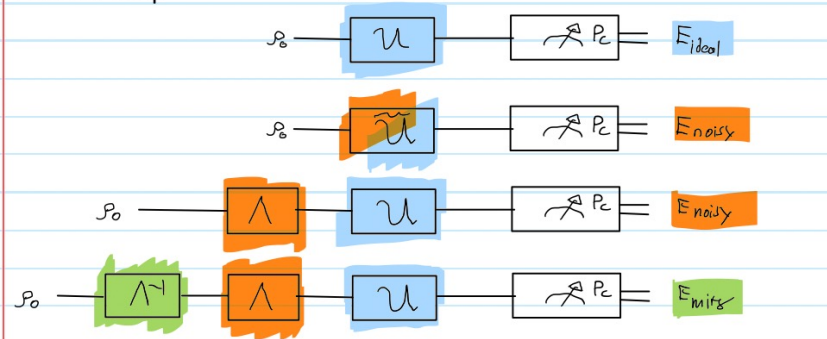
$$c_b^{inv} \in \mathbb{R}$$

Need to introduce

- Super-op. bra-ket notation
- Pauli channel $\rho \rightarrow P \rho P$
- separ-op notation
- unit
- Pauli check representation
- SQ
- \mathbb{Z}_2^n
- $\mathbb{C}^{0,0,2^n}$

Circuit expectation values

From 1 Step



Channel definitions

$$U = u \cdot u^\dagger$$

$$\Lambda = \sum_a f_a |P_a\rangle\langle P_a|$$

$$= \sum_b c_b P_b$$

$$\Lambda^{-1} = \sum_a f_a^{-1} |P_a\rangle\langle P_a|$$

$$= \sum_b c_b^{inv} P_b$$

$$\langle \hat{P}_c \rangle(\cdot) = \langle P_c | \cdot$$

$$-1 \leq f_a \leq 1$$

$$c_b \geq 0, \sum c_b = 1$$

$$c_b = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

$$\vec{c}_b = W^{-1} \vec{f}_a \quad \vec{f}_a = W c_b$$

$$\vec{c}_b^{inv} = W \vec{f}_a^{-1}$$

$$c_b^{inv} = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

$$c_b^{inv} \in \mathbb{R}$$

Need to introduce

- Separable bra-ket notation
- Pauli channel to PTM
- separable notation
- unit
- Pauli check representation
- SG
- \mathbb{Z}_2^n
- $\mathbb{C}^{0,0,2^n}$

$$E_{ideal} := \langle \hat{P}_c \rangle_{\rho}$$

$$= \langle\langle P_c | U | \rho_0 \rangle\rangle$$

ideal exp value with noiseless unitary

$$E_{noisy} := \langle \hat{P}_c \rangle_{\tilde{\rho}}$$

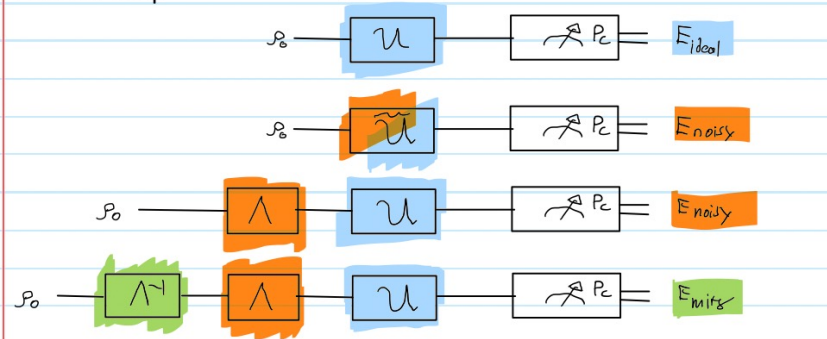
$$= \langle\langle P_c | U \Lambda | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | \tilde{U} | \rho_0 \rangle\rangle$$

noisy-gate expectation value

Circuit expectation values

From 1 Step



Channel definitions

$$U = U \cdot U^\dagger$$

$$\Lambda = \sum_a f_a |P_a\rangle\langle P_a|$$

$$= \sum_b c_b P_b$$

$$\Lambda^{-1} = \sum_a f_a^{-1} |P_a\rangle\langle P_a|$$

$$= \sum_b c_b^{inv} P_b$$

$$\langle \hat{P}_c \rangle(\cdot) = \langle\langle P_c | \cdot \rangle\rangle$$

$$-1 \leq f_a \leq 1$$

$$c_b \geq 0, \sum_b c_b = 1$$

$$c_b = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

$$\vec{c}_b = W \vec{f}_a \quad \vec{f}_a = W c_b$$

$$\vec{c}_b^{inv} = W \vec{f}_a^{-1}$$

$$c_b^{inv} = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

$$c_b^{inv} \in \mathbb{R}$$

Need to introduce

- Super-op. bra-ket notation
- Pauli channel vs PTM
- separ-op notation
- unit
- Pauli check representation
- SQ
- CNOT

$$E_{ideal} := \langle \hat{P}_c \rangle_{\rho_0}$$

$$= \langle\langle P_c | U | \rho_0 \rangle\rangle$$

ideal exp value with noiseless unitary

$$E_{noisy} := \langle \hat{P}_c \rangle_{\tilde{\rho}_0}$$

$$= \langle\langle P_c | U | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | \tilde{U} | \rho_0 \rangle\rangle$$

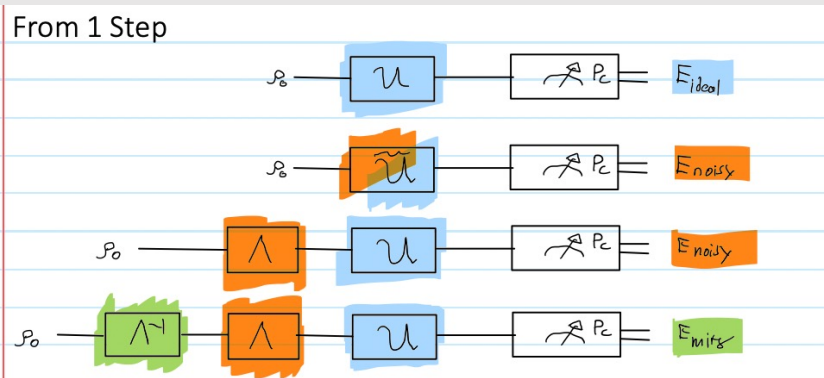
noisy-gate expectation value

$$E_{mitg} := \langle\langle P_c | U \Lambda \Lambda^{-1} | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | U \Lambda \left(\sum_a f_a^{-1} |P_a\rangle\langle P_a| \right) | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | U \Lambda \left(\sum_b c_b^{inv} P_b \right) | \rho_0 \rangle\rangle$$

Circuit expectation values



Channel definitions

$U = u \cdot u^\dagger$

$\Lambda = \sum_a f_a |P_a\rangle\langle P_a|$
 $= \sum_b c_b P_b$

$\Lambda^{-1} = \sum_a f_a^{-1} |P_a\rangle\langle P_a|$
 $= \sum_b c_b^{inv} P_b$

$\langle \hat{P}_c \rangle(\cdot) = \langle\langle P_c | \cdot \rangle\rangle$

$-1 \leq f_a \leq 1$

$c_b \geq 0, \sum_b c_b = 1$

$c_b^{inv} = \frac{1}{c_b}$

$\vec{c}_b = W \vec{f}_a, \vec{f}_a = W c_b$

$\vec{c}_b^{inv} = W \vec{f}_a^{-1}$

$c_b^{inv} = \frac{1}{c_b}$

$c_b^{inv} \in \mathbb{R}$

Need to introduce:

- Super-op. bra-ket notation
- Pauli channel Λ & PTM
- separ-op notation
- unit
- Pauli trace representation
- SQ
- c_b^{inv}
- c_b^{inv}

$$E_{ideal} := \langle \hat{P}_c \rangle_u$$

$$= \langle\langle P_c | U | P_0 \rangle\rangle$$

ideal exp value with noiseless unitary

$$E_{noisy} := \langle \hat{P}_c \rangle_{\tilde{u}}$$

$$= \langle\langle P_c | U \Lambda | P_0 \rangle\rangle$$

$$= \langle\langle P_c | \tilde{u} | P_0 \rangle\rangle$$

noisy-gate expectation value

$$E_{mitg} := \langle\langle P_c | U \Lambda \Lambda^{-1} | P_0 \rangle\rangle$$

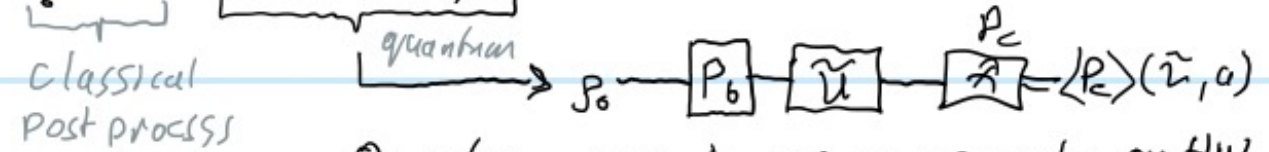
$$= \langle\langle P_c | U \Lambda \left(\sum_a f_a^{-1} |P_a\rangle\langle P_a| \right) | P_0 \rangle\rangle$$

$$= \langle\langle P_c | U \Lambda \left(\sum_b c_b^{inv} P_b \right) | P_0 \rangle\rangle$$

$$= \sum_b c_b^{inv} \langle\langle P_c | U \Lambda P_b | P_0 \rangle\rangle$$

$$= \sum_b c_b^{inv} \langle P_c \rangle(\tilde{u}, b)$$

sum of trajectories with weight $c_b^{inv} \in \mathbb{R}$



$$E_{mitg} := \langle\langle P_c | U \Lambda U^\dagger | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | U \Lambda \left(\sum_a f_a^{-1} |P_a\rangle \langle P_a| \right) | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | U \Lambda \left(\sum_b c_b^{inv} P_b \right) | \rho_0 \rangle\rangle$$

$$= \sum_b c_b^{inv} \langle\langle P_c | U \Lambda P_b | \rho_0 \rangle\rangle$$

sum of trajectories
with weight $c_b^{inv} \in \mathbb{R}$

$$= \underbrace{\sum_b c_b^{inv}}_{\text{classical post process}} \underbrace{\langle P_c \rangle(\tilde{u}, b)}_{\text{quantum}}$$

classical
post process



Quantum circuit we can execute on FHW
and find exp. value from.

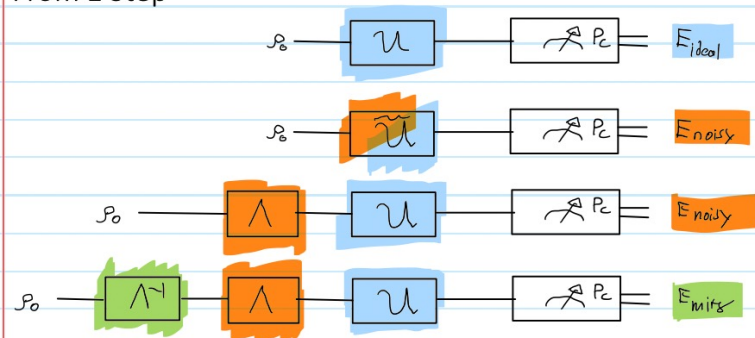
\therefore To find noise-free val all we have to do is to compute exp-val
of all 4^n b-modified circuits! This would give us ideal
exp value.

However, $| \{b\} | = 4^n$ grows exponentially, hence, infeasible.

but what if we could sample from it to approximate
full sum. But... cant sample directly from c_b^{inv} which

does not form a valid prob-distribution. let's solve:

From 1 Step



Channel definitions

$$U = u \cdot u^\dagger$$

$$\langle P_c \rangle(\cdot) = \langle\langle P_c | \cdot \rangle\rangle$$

- Need to introduce
- Super-op. bra-ket notation
 - Pauli channel $\rho \rightarrow P \rho P^\dagger$
 - separ-op notation
 - unit
 - Pauli check representation
 - SQ
 - $C \circ O_{sq}$

$$\Lambda = \sum_a f_a |P_a\rangle \langle P_a|$$

$$-1 \leq f_a \leq 1$$

$$= \sum_b c_b P_b$$

$$c_b \geq 0, \sum_b c_b = 1$$

$$c_b = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a$$

$$\vec{c}_b = W^{-1} \vec{f}_a \quad \vec{f}_a = W \vec{c}_b$$

$$\Lambda^{-1} = \sum_a f_a^{-1} |P_a\rangle \langle P_a|$$

$$\vec{c}_b^{inv} = W \vec{f}_a^{-1}$$

$$= \sum_b c_b^{inv} P_b$$

$$c_b^{inv} = \frac{1}{2^n} \sum_a (-1)^{a \cdot b} f_a^{-1}$$

$$c_b^{inv} \in \mathbb{R}$$

Quasi probability distribution

c_b^{inv} can be outside $[0,1]$

$\sum_b c_b^{inv} =: \gamma \geq 1$ generally for Λ not unitary

eg 6A th chapter

$$\Lambda = (1-p)[I + pX \cdot X]$$

$$\Lambda^{-1} = (1 + \frac{p}{1-2p})I + \frac{-p}{1-2p} X \cdot X$$

$$c_{\uparrow}^{inv} = 1 + \frac{p}{1-2p}$$

$b = (0,0)$
chose vector

$$c_{\downarrow}^{inv} = \frac{-p}{1-2p}$$

$b = (1,0)$

Turn into probability

$$c_b^{inv} = \underbrace{\text{sgn}(c_b^{inv})}_{\text{sign} \in \{-1, +1\}} \underbrace{\frac{|c_b^{inv}|}{\gamma}}_{\substack{\text{prob} \in [0,1] \\ \text{scale}}}$$

$$\bar{c}_b^{inv} := \frac{|c_b^{inv}|}{\|c_b^{inv}\|_1} \rightarrow \|c_b^{inv}\|_1 = \sum_b |c_b^{inv}| \Leftrightarrow L_1 \text{ norm}$$

$$= \frac{|c_b^{inv}|}{\gamma}$$

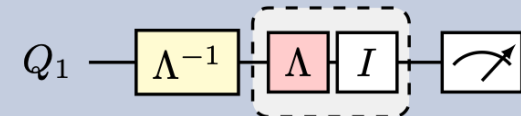


probability

$$P_I = |1 - q|/\gamma$$

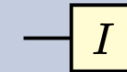
$$P_X = |q|/\gamma$$

$$P_I + P_X = 1$$



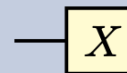
||

sign



$$s_I = \text{sign}(1-q)$$

+



$$s_X = \text{sign}(q)$$

E mitigated

$$E_{\text{mitg}} = \sum_b c_b^{\text{inv}} \langle \hat{P}_c \rangle(\tilde{u}, b)$$

$$= \sum_b \text{sgn}(c_b^{\text{inv}}) \frac{|c_b^{\text{inv}}|}{\gamma} \gamma \langle \hat{P}_c \rangle(\tilde{u}, b)$$

$$= \gamma \sum_b \underbrace{\text{sgn}(c_b^{\text{inv}})}_{\text{scale}} \underbrace{\bar{c}_b^{\text{inv}}}_{\text{classical post-processing}} \underbrace{\langle \hat{P}_c \rangle(\tilde{u}, b)}_{\text{valid QC circuit can run & find value on HW}}$$

scale classical post-processing. valid QC circuit can run & find value on HW

In this form, the decomposition of the error error mitigated expectation value is simply a sum over expectation values of Pauli-gate modified circuits, whose value can be obtained from direct quantum computer execution, weighted by a probability, c bar b inverse, and the sign. The elements that perform the weighing and rescaling can all be done in classical post-processing.

c_b^{inv} can be outside $\{0,1\}$

$\sum_b c_b^{\text{inv}} =: \gamma \geq 1$ generally, for Λ not unitary

eg b4 th chapter $\Lambda = (1-p)I + pX \cdot X$
 $\Lambda^\dagger = (1+p)I + \frac{-p}{1-2p} X \cdot X$

$c_b^{\text{inv}} = 1 + \frac{p}{1-2p}$ for $b=(0,0)$ choice vector
 $c_b^{\text{inv}} = \frac{-p}{1-2p}$ for $b=(1,0)$

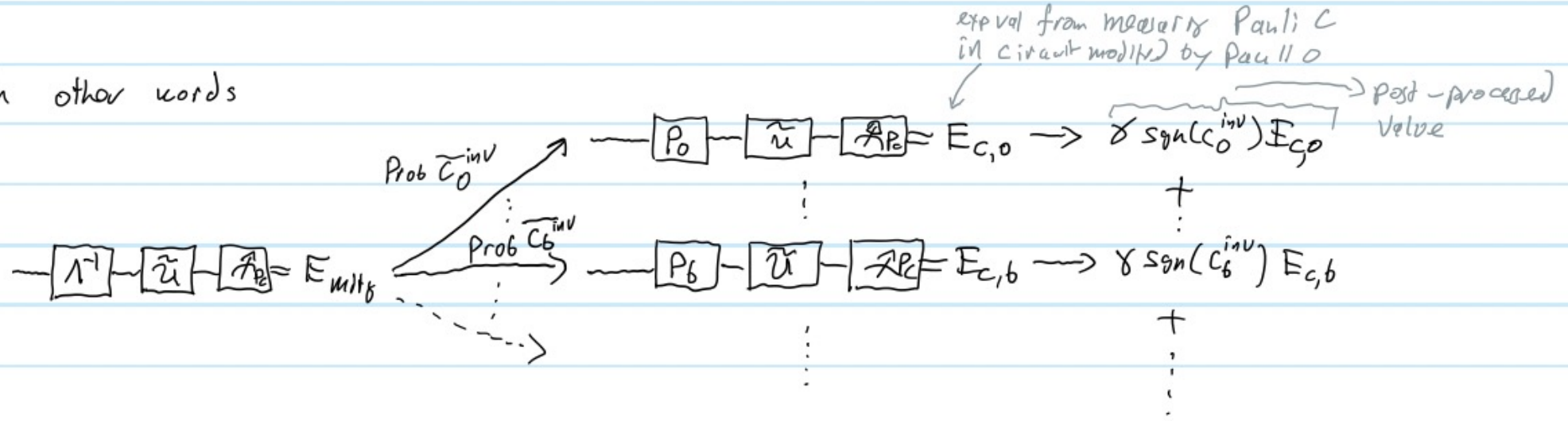
Turn into probabilities

$$c_b^{\text{inv}} = \underbrace{\text{sgn}(c_b^{\text{inv}})}_{\text{sign in } \{-1, +1\}} \underbrace{\frac{|c_b^{\text{inv}}|}{\gamma}}_{\substack{\text{prob} \in \{0,1\} \\ \text{scale}}} \gamma$$

$$\bar{c}_b^{\text{inv}} := \frac{|c_b^{\text{inv}}|}{\|c_b^{\text{inv}}\|_1} \xrightarrow{\text{abs}} \|c_b\|_1 = \sum_b |c_b^{\text{inv}}| \Leftrightarrow L_1 \text{ norm}$$

$$= \frac{|c_b^{\text{inv}}|}{\gamma}$$

In other words



Quasi-Probability Distribution
 c_b^{inv} can be outside $\{0,1\}$
 $\sum_b c_b^{inv} = \gamma \geq 1$ generally for Λ not unitary
 eg. bit-flip channel $\Lambda = (1-p)I + pxX$
 $\Lambda^{-1} = (1+p)I + \frac{p}{1+p}X$
 $c_b^{inv} = \frac{1+p}{1-p}$ $c_x^{inv} = -\frac{p}{1-p}$
 Turn into probability
 $C_b^{inv} = \frac{\text{sgn}(c_b^{inv}) |c_b^{inv}|}{\sum_b \text{sgn}(c_b^{inv}) |c_b^{inv}|} \gamma$
 $\tilde{c}_b^{inv} = \frac{|c_b^{inv}|}{\sum_b |c_b^{inv}|} \gamma$ $\|c_b^{inv}\| = \sum_b |c_b^{inv}| \rightarrow L_1$ norm
 $= \frac{|c_b^{inv}|}{\sum_b |c_b^{inv}|} \gamma$

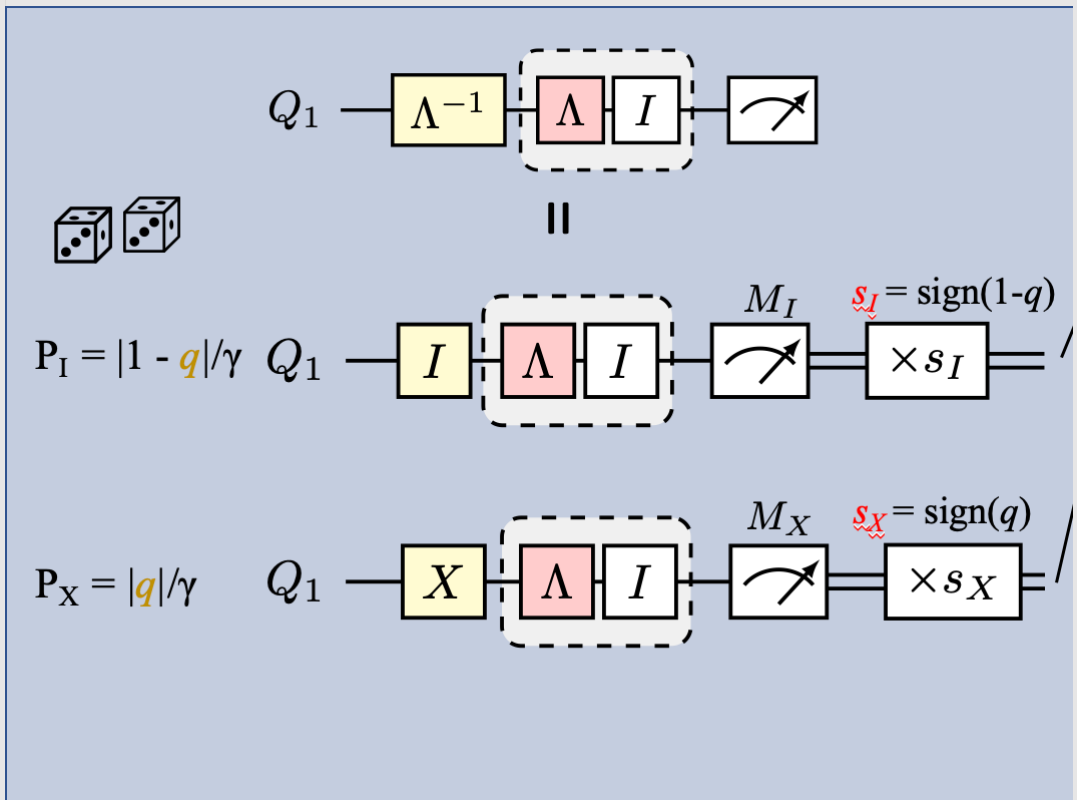
$$E_{mit,b} = \sum_b c_b^{inv} \langle \hat{P}_b \rangle (\tilde{U}, b)$$

$$= \sum_b \text{sgn}(c_b^{inv}) \frac{|c_b^{inv}|}{\gamma} \langle \hat{P}_b \rangle (\tilde{U}, b)$$

$$= \gamma \sum_b \text{sgn}(c_b^{inv}) \tilde{c}_b^{inv} \langle \hat{P}_b \rangle (\tilde{U}, b)$$

scale classical post-processing valid QC circuit
 can't run forward with Pauli

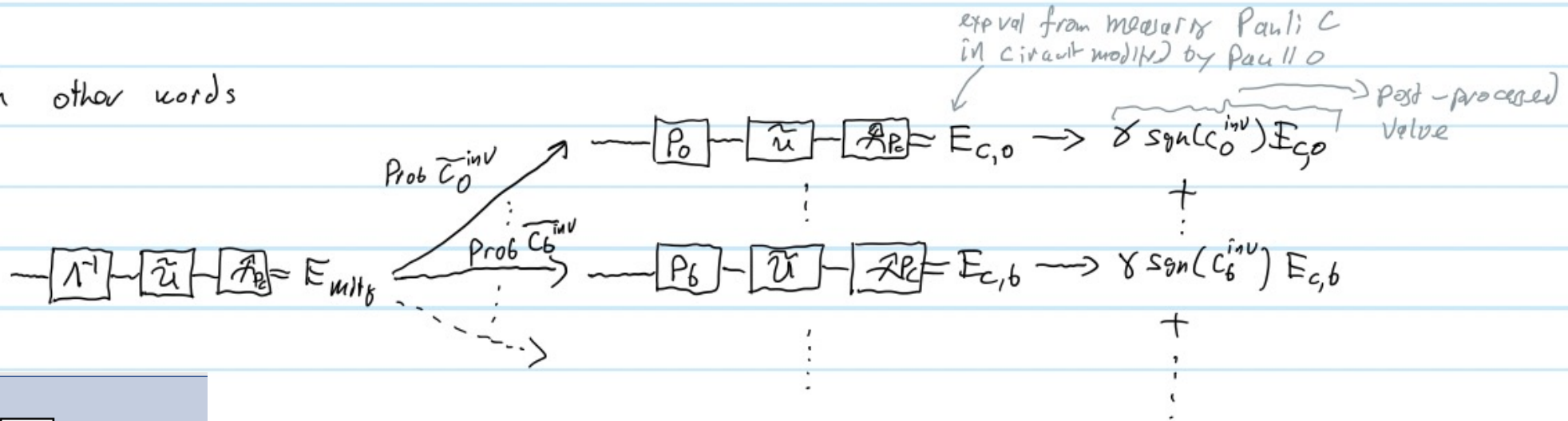
In this form, the decomposition of the error error mitigated expectation value is simply a sum over expectation values of Pauli-gate modified circuits, whose value can be obtained from direct quantum computer execution, weighted by a probability, c bar is inverse, and the sign. The elements that perform the weighing and rescaling can all be done in classical post-processing.



Estimator

post processing.

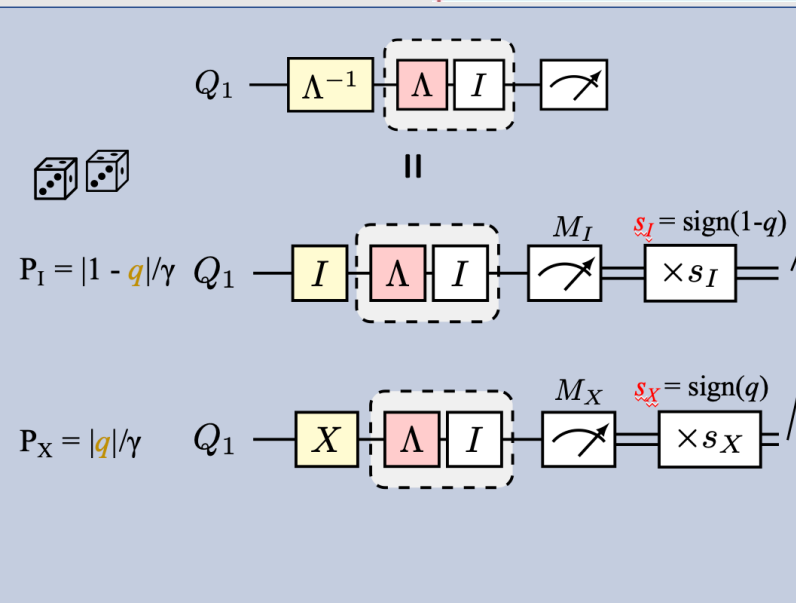
In other words



$$E_{c,mitg} = \sum_b \gamma \text{sgn}(C_b^{inv}) E_{c,b}$$

mitigated value for Pauli C obtained from the quasi-prob distribution

From above, we know this is an unbiased estimator, but what about the error and sampling

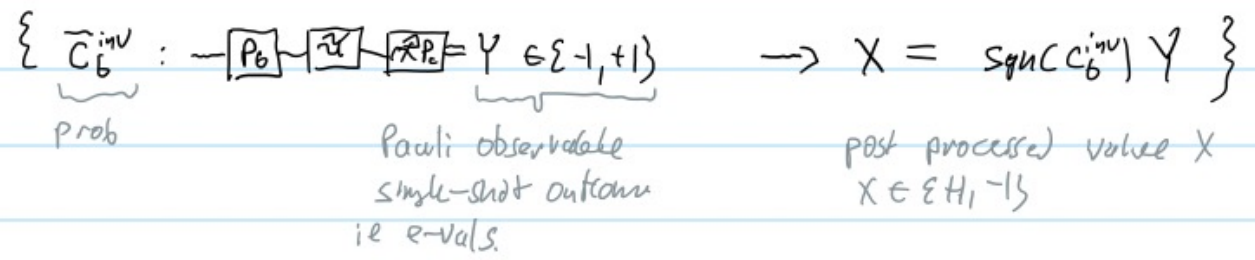


mitigated value for Pauli C obtained from the quasi-prob distribution

From above, we know this is an unbiased estimator,
but what about the error and sampling

Estimator, Sampling, and Error Bounds

Sample circuits of the form



Let's say we sample M instances, randomly sample a single value for b and obtain one-shot value on the QC, ie one random instance of $Y = 1$ or $Y = -1$, which we then post-process.

The results are thus the classical random variables

$$\{X_1, X_2, \dots, X_M\} \quad \text{or} \quad \{X_m : m = 1, \dots, M\}$$

Quasi-Probability Distribution

c_b^{inv} can be outside $[0, 1]$

$\sum_b c_b^{inv} = 1$ generally for N not unitary

eg bit flip circuit $M = (I + \sigma_x)^n + \frac{1 - i^n}{2} \sigma_x$

$c_b^{inv} = \frac{1 + P}{1 - 2P}$ $C_b^{inv} = \frac{P}{1 - 2P}$

Turn into probability

$C_b^{inv} = \underbrace{\text{sgn}(c_b^{inv})}_{\substack{\text{sgn} \in \{-1, 1\} \\ \text{prob} \in [0, 1]}} \underbrace{|c_b^{inv}|}_{\text{scale}} \gamma$

$\bar{c}_b^{inv} = \frac{|c_b^{inv}|}{\sum_b |c_b^{inv}|} \rightarrow \|c_b^{inv}\| = \sum_b |c_b^{inv}| \rightarrow L_1 \text{ norm}$

$= \frac{|c_b^{inv}|}{\gamma}$

$$E_{mitg} = \sum_b c_b^{inv} \langle \hat{P}_b \rangle_{\psi, b}$$

$$= \sum_b \text{sgn}(c_b^{inv}) \frac{|c_b^{inv}|}{\gamma} \langle \hat{P}_b \rangle_{\psi, b}$$

$$= \gamma \sum_b \underbrace{\text{sgn}(c_b^{inv})}_{\text{scale}} \underbrace{\frac{|c_b^{inv}|}{\gamma}}_{\text{classical post-processing}} \langle \hat{P}_b \rangle_{\psi, b}$$

scale = classical post-processing
 valid QC circuit can run forward and backward

In this form, the decomposition of the error error mitigated expectation value is simply a sum over expectation values of Pauli-gate modified circuits, whose value can be obtained from direct quantum computer execution, weighted by a probability, a bar b inverse, and the sign. The elements that perform the weighing and rescaling can all be done in classical post-processing.

Probabilistic error cancellation: Derivation

where each $X_m \in \{1, -1\}$ and is distributed to model a Bernoulli distribution with some probability which can be any valid value and can vary from shot-to-shot m .

Our mitigation estimator is then for M shots:

$$E_M := \gamma \frac{1}{M} \sum_{m=1}^M X_m = \frac{1}{M} \sum_{m=1}^M \gamma \text{sgn}(c_{b_m}^{\text{inv}}) \underbrace{Y_{m,b_m}}_{\substack{\text{rand outcome of } m\text{-th shot for } b_m \text{ circ.} \\ \text{noisy circuit}}} (\vec{u}, b_m)$$

Pauli chosen for m -th shot

There are now 2 random processes:

b_m : which pauli b we pick for shot m

Y_{b_m} : which outcome ± 1 we get for b_m circuit of shot m

Unbiased estimator of the ideal, noise-free circuit expectation

$$\mathbb{E}[E_M] = \frac{1}{M} \sum_{m=1}^M \mathbb{E}[\gamma X_m] \quad \text{iid rand vars}$$

Unbiased estimator

$$\mathbb{E}[F_M] = \frac{1}{M} \sum_{m=1}^M \mathbb{E}[\delta X_m] \quad \text{iid rand vars}$$

$$= \mathbb{E}[\delta X_m] \quad \text{no } X_m \text{ is different}$$

$$= \mathbb{E}[\delta \text{sgn}(c_{b_m}^{\text{inv}}) Y_{b_m}(\tilde{u}, b_m)] \quad \text{where rand var is } b_m \text{ now, not just } m, \text{ so}$$

$$\approx \mathbb{E}_{b_m}[\delta \text{sgn}(c_b^{\text{inv}}) Y_b(\tilde{u}, b)] \quad \text{Prob}[b] = \bar{c}_b^{\text{inv}}$$

drop notation

$$= \sum_b \mathbb{E}_Y[\delta \text{sgn}(c_b^{\text{inv}}) Y_b] \text{Prob}[b]$$

$$= \sum_b \underbrace{\delta \text{sgn}(c_b^{\text{inv}})}_{\text{post-process of outcome}} \underbrace{\bar{c}_b^{\text{inv}}}_{\text{sample prob}} \underbrace{\mathbb{E}[Y_b]}_{\text{rand outcome}} \longrightarrow$$

$$\approx \sum_b \delta \text{sgn}(c_b^{\text{inv}}) \frac{\bar{c}_b^{\text{inv}}}{\delta} \langle\langle P_C | \mathcal{U} \wedge \mathcal{P}_b | \rho_0 \rangle\rangle$$

$$\text{note: } \langle\langle \hat{P}_C \rangle\rangle(\tilde{u}, b) = \langle\langle P_C | \mathcal{U} \wedge \mathcal{P}_b | \rho_0 \rangle\rangle = \mathbb{E}[Y_b]$$

↑ Rand variable ± 1 for output of the b -th pauli circuit.

\therefore for some classical func of b $f(b)$ which does not depend on the value Y_b but only on the label b :

$$\mathbb{E}[f(b) Y_b] = f(b) \langle\langle \hat{P}_C \rangle\rangle(\tilde{u}, b)$$

Probabilistic error cancellation: Derivation

$$= \langle\langle P_c | U \Lambda \left(\sum_b c_b^{inv} P_b \right) | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | U \Lambda \Lambda^{-1} | \rho_0 \rangle\rangle$$

$$= \langle\langle P_c | U | \rho_0 \rangle\rangle$$

$$= \langle \hat{P}_c \rangle (U_{ideal}) \quad \text{without noise!}$$

unbiased estimator of the true noise-free, ideal value of the circuit



the value γ_b but only on the label b :

$$\mathbb{E}[f(c_b) \gamma_b] = f(c_b) \langle \hat{P}_c \rangle (U, b)$$

(Optional step) Variance

Variance of E_M

$$\underbrace{V_{\{b_m, x_m\}}}_{\text{with respect to}} [E_M] = \frac{\gamma^2}{M^2} \sum_{m=1}^M V[X_m]$$

$$= \frac{\gamma^2}{M} V[X(U, b)]$$

X_m iid
can drop subscript m
and emphasize b and value X

note the same variance is just rescaled by γ^2 due to γ

Generalizing: Raveling trajectories with quasiprobabilities

Channel we *want*
to implement

CPTP operation we
can implement

$$\mathcal{C}(\cdot) = \sum_i a_i \mathcal{F}_i(\cdot)$$

Real coefficients, turn
into *quasi-probability*

Putting the following techniques all on the same footing

Technique

Prob. error cancelation (PEC)

Circuit cutting (knitting) of gates

Circuit cutting of wires

Classical sim. algorithms (QP)

Channel \mathcal{C}

noise inverse

non-local gate

large unitary

unitary

Quantum Error Mitigation for Non-Equilibrium Quantum Dynamics

Lecture 1

Big picture

Why quantum computers?

Status and outlook

Why error mitigation?

Noise in quantum computers
Overview of error mitigation



Mitigation fundamentals

Probabilistic error cancelation (PEC)

Introduction

One qubit example

General derivation



Next lectures

Learning noise

State-of-art PEC experiments

Key techniques: Twirling

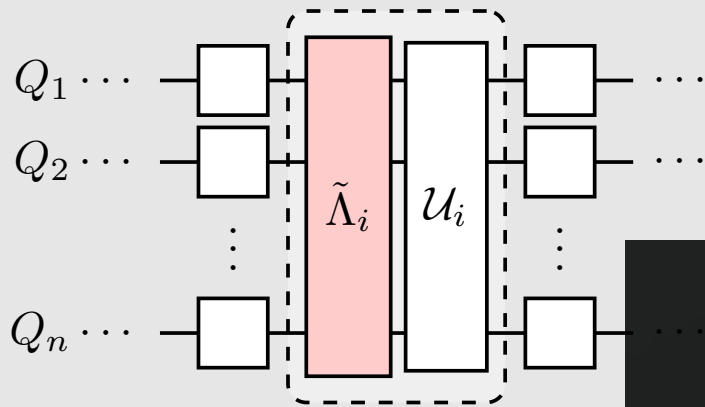
T-REX mitigation

...

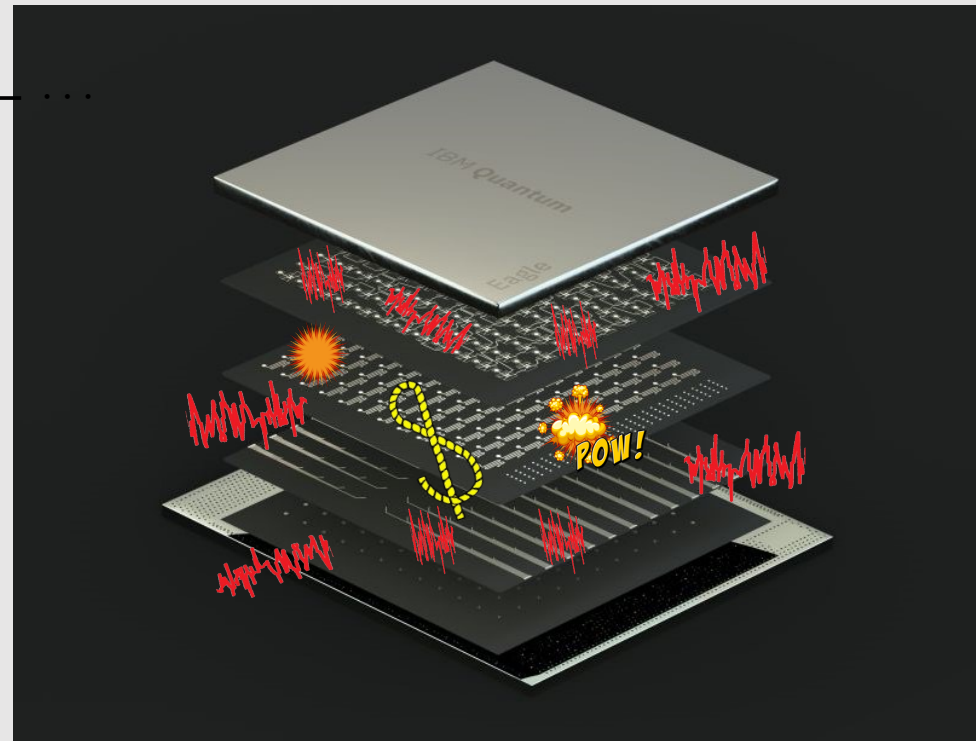
State-of-art experiments at the 100Q+, depth 50+:
uncovering local integrals of motion

...

Is it possible to learn the noise with accuracy, efficiency, and scalability?



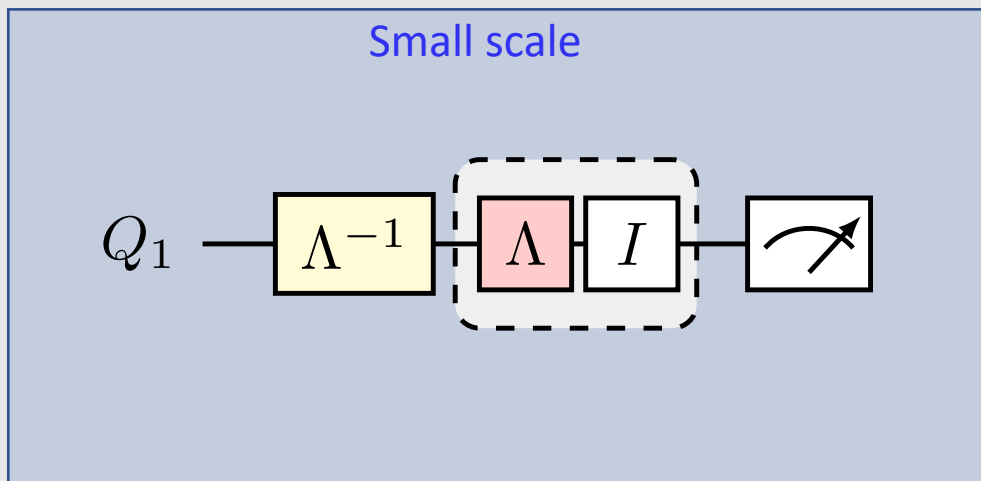
Energy relaxation T_1
Dephasing T_2
Coherent errors ZZ
Classical crosstalk
Quantum crosstalk
State preparation error
Measurement correlated errors
...



Control errors
Photon shot noise
1/f charge noise
1/f flux noise
Nonequilibrium quasiparticles
Leakage
Cosmic rays
...

PEC: Nice, **but** why hasn't worked so far for experiments?

Practical challenges



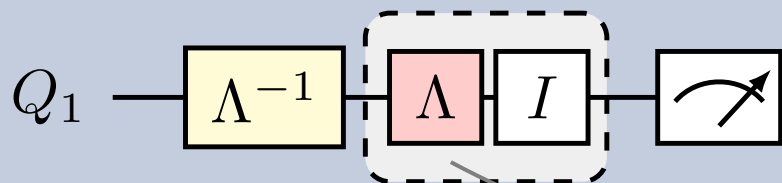
Critically hinges on knowing the full noise near perfectly

Despite the method's theoretical appeal (1-10), practical challenges have limited its demonstration to the one and two-qubit level (2, 3)

1. S. Endo, S. C. Benjamin, Y. Li, *Physical Review X* 8, 031027 (2018).
2. C. Song, et al., *Science Advances* 5, arXiv:2109.04457(2019).
3. S. Zhang, et al., *Nature Communications* 11, 587 (2020).
4. C. Piveteau, D. Sutter, S. Woerner, arXiv:2101.09290 (2021).
5. S. Endo, et al., *J. Phys. Soc. of Japan* 90, 032001 (2021).
6. C. Piveteau, et al., arXiv:2103.04915 (2021).
7. R. Takagi, *Phys. Rev. Research* 3, 033178 (2021).
8. R. Takagi, S. Endo, S. Minagawa, M. Gu, arXiv:2109.04457 (2021).
9. Y. Guo, S. Yang, arXiv preprint arXiv:2201.00752 (2022).
10. ...

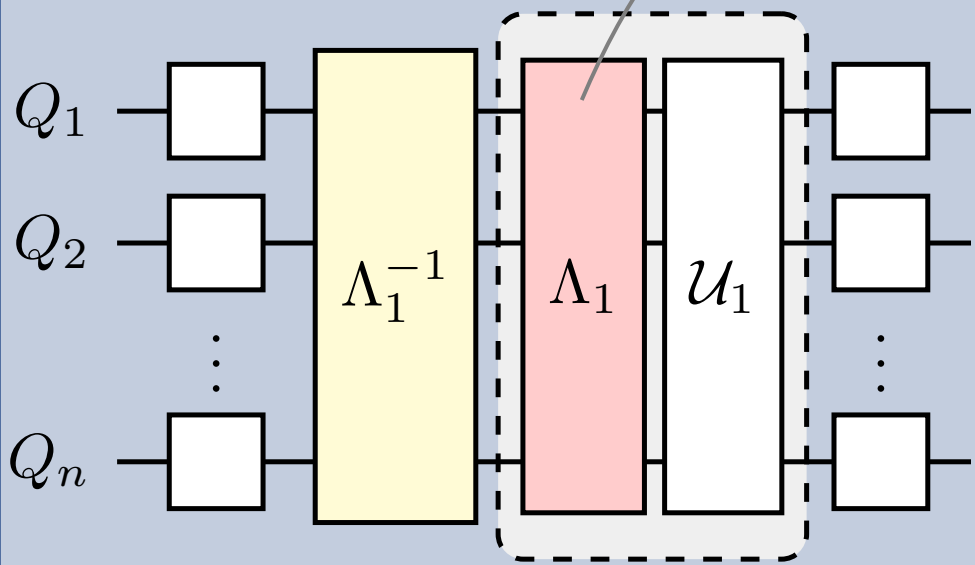
PEC: Nice, **but** why hasn't worked so far? Challenges

Small scale



2 qubits 255 parameters
10 qubits 10^{12} parameters
50 qubits **10^{60} parameters**
noise param values $10^{-2} - 10^{-5}$
additive error sampling cost ($>10^2 - 10^{10}$)

Large scale



Challenges

learning complexity

- efficient
- scalable
- accurate
- compact, tractable representation

noise in full device

- cross-talk
- correlated errors
- parallel gates

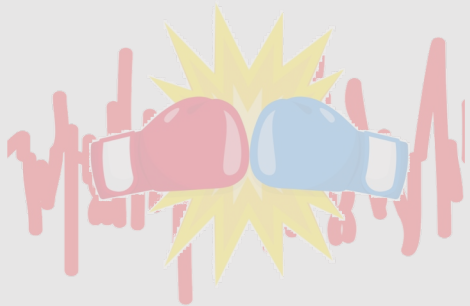
Outline



Idea

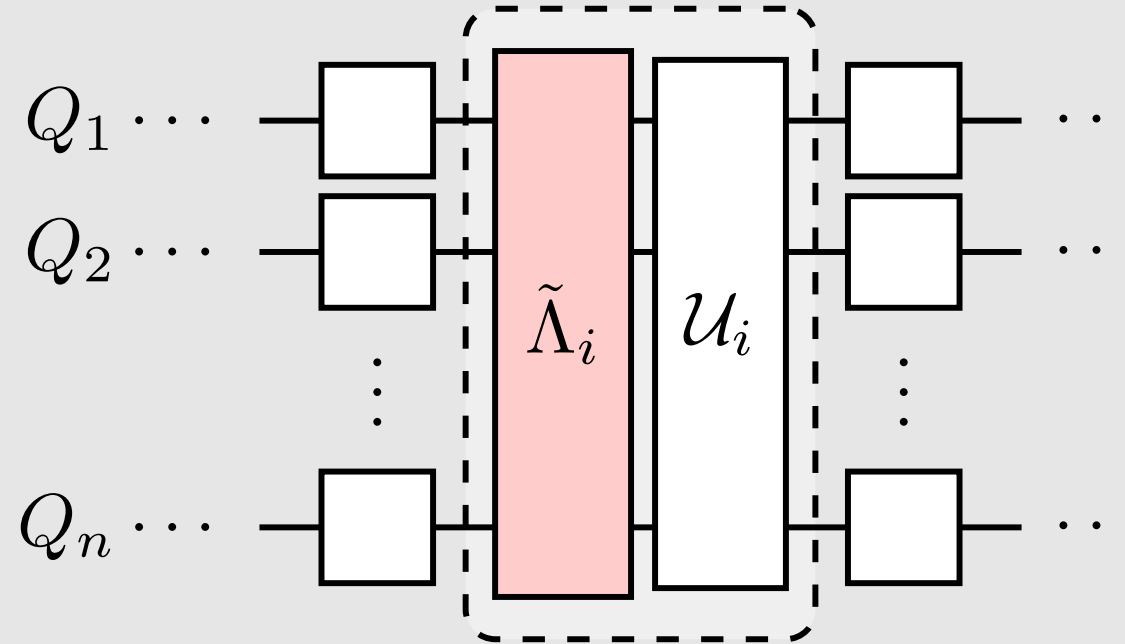


Learn



Cancel
(realization)

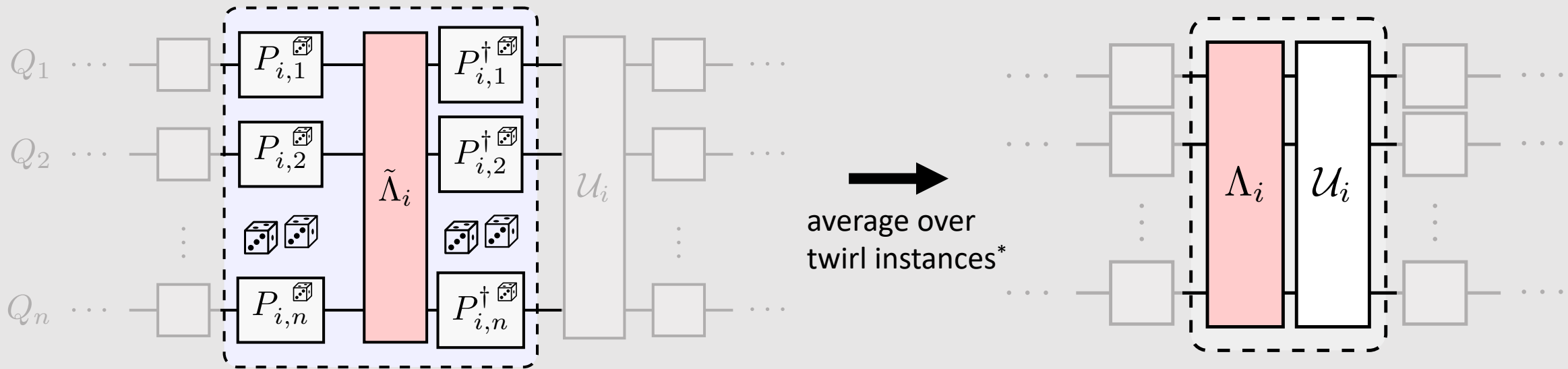
Step 1: Simplify the noise



noise that includes cross-talk errors, etc.
characterized by some $4^n \times 4^n$ matrix

Step 1: Simplify the noise: twirl

twirl reduces to noise $4^n \times 4^n$ matrix to diagonal one with 4^n entries in Pauli basis



Twirling references

1. C. H. Bennett, G. Brassard, S. Popescu, B. Schumacher, J. A. Smolin, **Tutorial** W. K. Wootters, et al., Phys. Rev. Lett. 76, 722 (1996).
2. E. Knill, arXiv:0404104 (2004).
3. O. Kern, G. Alber, D. L. Shepelyansky, EPJ D 32, 153 (2005).
4. M. R. Geller, Z. Zhou, Physical Review A 88, 012314 (2013).
5. J. J. Wallman, J. Emerson, Physical Review A 94, 052325 (2016)
6. Hashim *et al.*, Phys. Rev. X 11, 041039 (2021)
7. Tutorial: zlatko-minev.com/blog/twirling (2022)
8. ...



Stochastic Pauli channel

$$\Lambda_i(\rho) = \sum_{a=0}^{4^n-1} c_{ia} P_a \rho P_a^\dagger$$

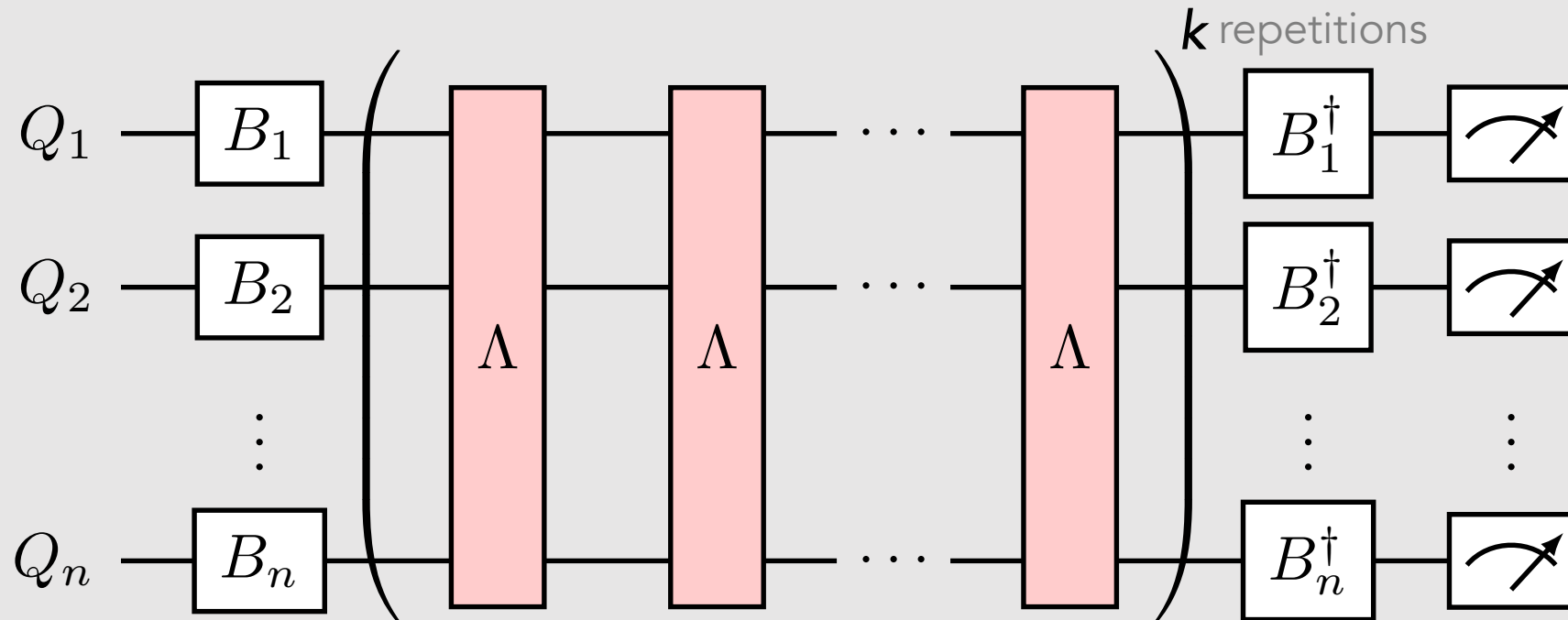
$$\Lambda(P_a) = f_a P_a \quad \text{eigenvecs are Paulis}$$

* some sub-Clifford twirl group (use Paulis)

Zlatko Mineev, IBM Quantum (69)

Step 2 wish: amplify noise

for the i -th layer



prepare circuit in pre-determined Pauli basis

Since diagonal channel will amplify eigenvalues learn with multiplicative precision

measure circuit in same pre-determined Pauli basis

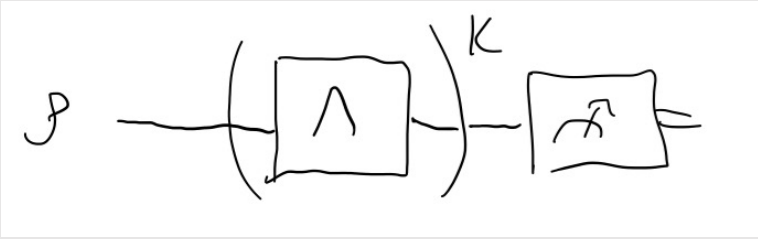
Akin to:

RB, Cycle RB, K-body noise reconstruction, ...

S.T. Flammia and J.J. Wallman ACM Trans QC 1, 3 (2020), ...

For something of a review of protocols, see Helsen, *et al.*, *A general framework for randomized benchmarking* (arXiv:2010.07974)

Notes



$$\Lambda_i(\rho) = \sum_{a=0}^{4^n-1} c_{ia} P_a \rho P_a^\dagger$$

$$\Lambda(P_a) = f_a P_a$$

For a qubit

$$\rho = \frac{1}{2} \left(\hat{I} + \rho_X \hat{X} + \rho_Y \hat{Y} + \rho_Z \hat{Z} \right)$$

In general:

$$\rho = \frac{1}{2^n} \sum_b \mathcal{P}_b \hat{P}_b$$

$$\Lambda(\rho) = \rho'$$

$$\rho' = \frac{1}{2^n} \sum_b \mathcal{P}'_b \hat{P}_b$$

Pauli decomposition of a density matrix is a powerful tool - offering a versatile representation. It expresses a density matrix as a linear combination of Paulis, often referred to as the Pauli basis.

Posterior state: Action of the channel on the input state

Pauli decomposition of posterior state

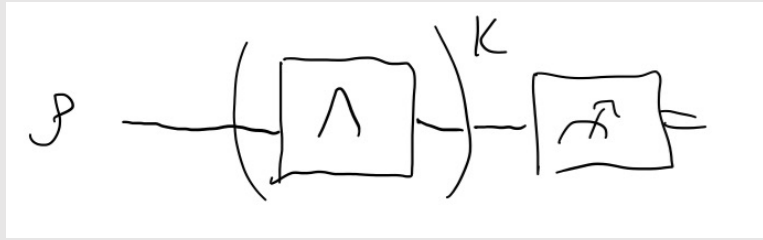
Since the channel is linear

$$\Lambda(a \hat{A} + b \hat{B}) = a \Lambda(\hat{A}) + b \Lambda(\hat{B})$$

$$\Lambda(\rho) = \Lambda\left(\frac{1}{2^n} \sum_b \mathcal{P}_b \hat{P}_b\right)$$

$$= \frac{1}{2^n} \sum_b \mathcal{P}_b \underbrace{\Lambda(\hat{P}_b)}$$

Just need to know action on P_b



$$\Lambda_i(\rho) = \sum_{a=0}^{4^n-1} c_{ia} P_a \rho P_a^\dagger$$

$$\Lambda(P_a) = f_a P_a$$

or written another way:

$$\begin{aligned} f_c &= \sum_b f_b \frac{1}{2^n} \text{Tr}(\hat{P}_c^\dagger \Lambda(\hat{P}_b)) \\ &= \sum_b f_b \frac{\langle\langle P_c | \Lambda | P_b \rangle\rangle}{\langle\langle P_c | P_c \rangle\rangle} \end{aligned}$$

$$\Lambda(\rho) = \Lambda\left(\frac{1}{2^n} \sum_b f_b \hat{P}_b\right)$$

$$= \frac{1}{2^n} \sum_b f_b \Lambda(\hat{P}_b)$$

just need to know action on P_b

$$= \rho'$$

$$= \frac{1}{2^n} \sum_c f'_c \hat{P}_c$$

since $\text{Tr}(\hat{P}_a^\dagger \hat{P}_b) = \delta_{ab} 2^n$
orthogonal

Solve for f'_c by equate & orthogonal case:

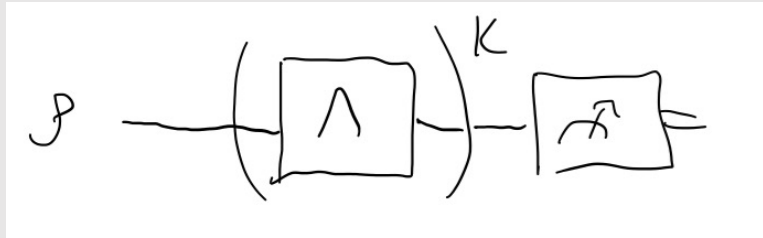
$$\text{Tr}\left(\hat{P}_c^\dagger \frac{1}{2^n} \sum_b f_b \Lambda(\hat{P}_b)\right) = \text{Tr}\left(\hat{P}_c^\dagger \frac{1}{2^n} \sum_{c'} f'_{c'} \hat{P}_{c'}\right)$$

$$\sum_b f_b \text{Tr}(\hat{P}_c^\dagger \Lambda(\hat{P}_b)) = \sum_{c'} f'_{c'} \text{Tr}(\hat{P}_c^\dagger \hat{P}_{c'})$$

$$= \sum_{c'} f'_{c'} \delta_{cc'} 2^n$$

$$\frac{1}{2^n} \sum_b f_b \text{Tr}[\hat{P}_c^\dagger \Lambda(\hat{P}_b)] = f'_c \quad \checkmark$$

Notes



$$\Lambda_i(\rho) = \sum_{a=0}^{4^n-1} c_{ia} P_a \rho P_a^\dagger$$

$$\Lambda(P_a) = f_a P_a$$

The measurement outcome then is (say measured in the \hat{P}_c basis):

$$\langle \hat{P}_c \rangle = \text{Tr}(\hat{P}_c \rho')$$

For a single step

$$\approx \text{Tr}(\hat{P}_c \sum_{c'} \frac{1}{2^n} \rho_{c'} \hat{P}_{c'})$$

$$= \sum_{c'} \rho_{c'} \frac{1}{2^n} \overbrace{\text{Tr}(\hat{P}_c \hat{P}_{c'})} \rightarrow 2^n \delta_{cc'}$$

$$= \rho_c$$

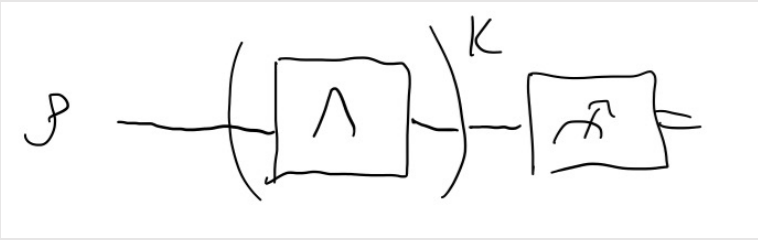
now sub in

$$= \sum_b \rho_b \frac{\langle\langle P_c | \Lambda | P_b \rangle\rangle}{2^n}$$

$$\rightarrow \text{for a Paul channel } \langle\langle P_c | \Lambda | P_b \rangle\rangle = f_c \delta_{bc} 2^n$$

$$= \underbrace{f_c}_{\text{fidelity } c \text{ of } \Lambda} \rho_c \leftarrow \text{initial state coefficient}$$

Notes



$$\Lambda_i(\rho) = \sum_{a=0}^{4^n-1} c_{ia} P_a \rho P_a^\dagger$$

$$\Lambda(P_a) = f_a P_a$$

What is ρ_c for an eigenstate of $\hat{P}_c \rightarrow 1$.

eg. for $c=z$: $\hat{P}_c = \hat{Z}$ and $\rho = |0\rangle = \frac{1}{2}(I + Z)$

$$\Rightarrow \rho_{c=z} = \text{Tr}(Z \rho) = +1$$

For repeating the channel K times



$$\rho' = \Lambda^K \rho$$

$$\Rightarrow \rho'_c = \sum_b f_b \frac{\langle\langle P_c | \Lambda^K | P_b \rangle\rangle}{2^n}$$

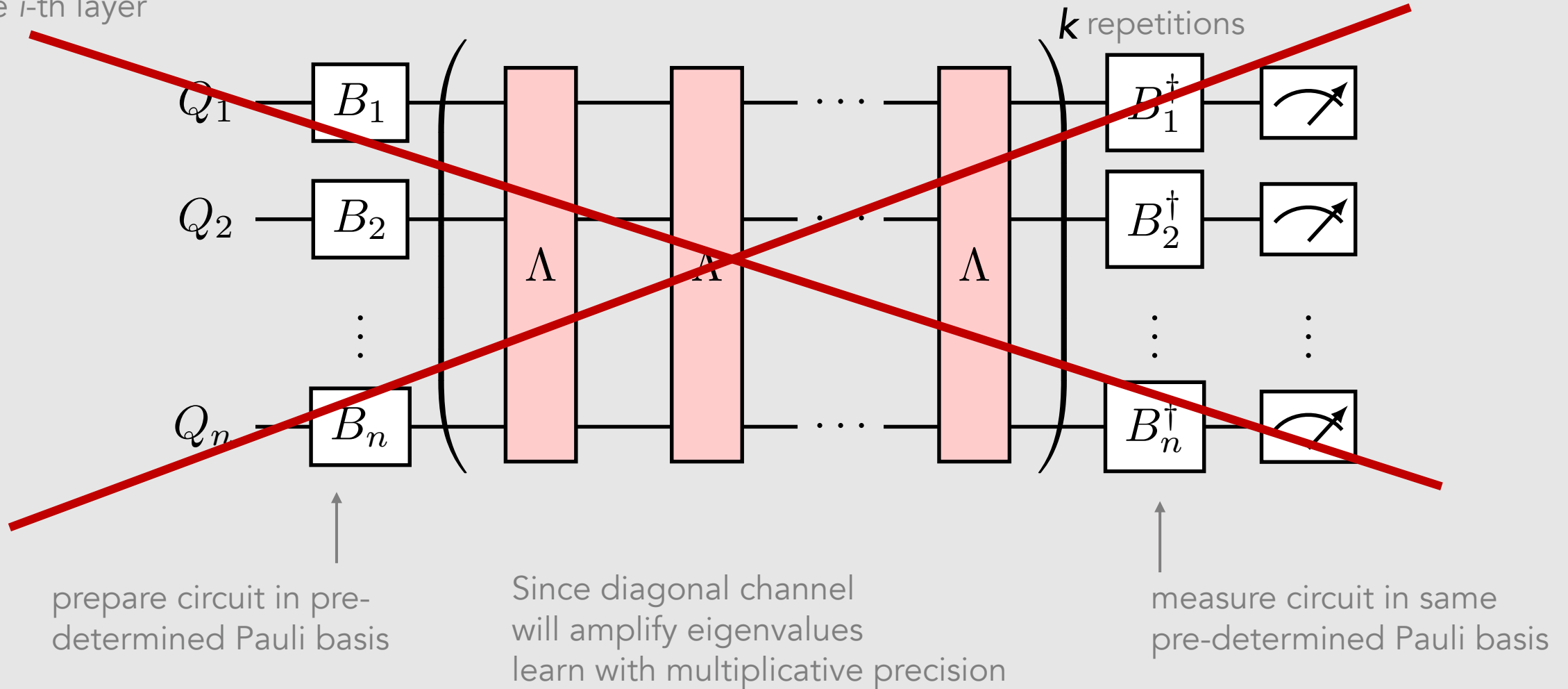
$$= \sum_b f_b^K \rho_b$$

$$\Rightarrow \langle \hat{P}_c \rangle = f_c^K \rho_c$$

Can include all diagonal SPAM

Step 2 wish: amplify noise

for the i -th layer



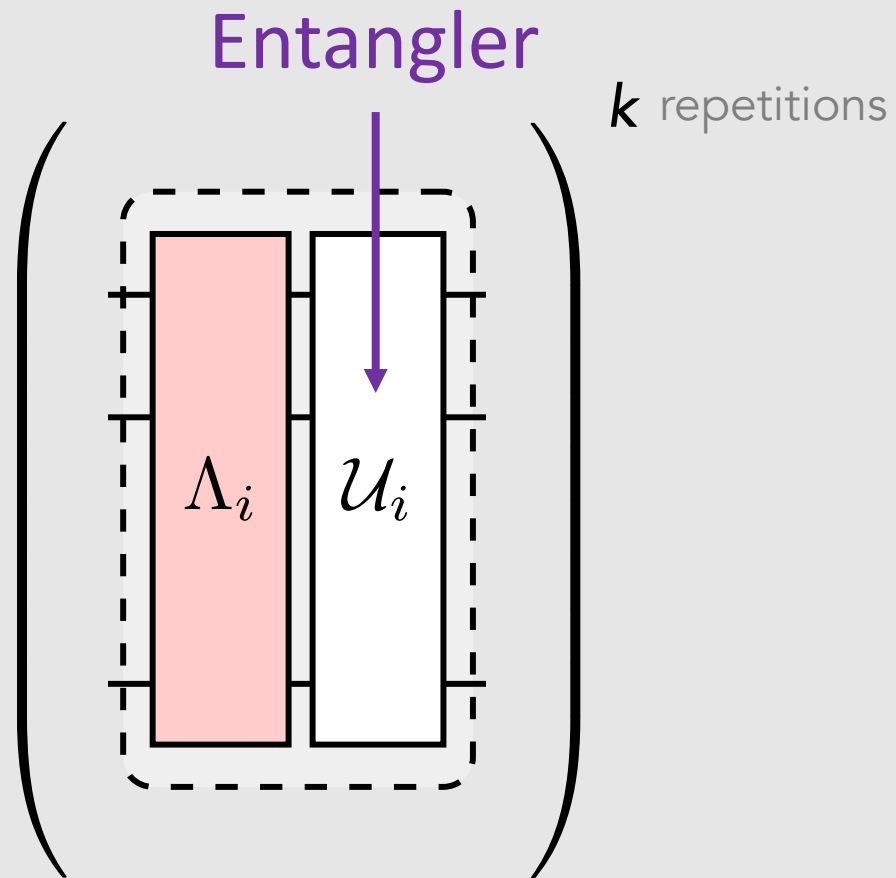
Akin to:

RB, Cycle RB, K-body noise reconstruction, ...

S.T. Flammia and J.J. Wallman ACM Trans QC 1, 3 (2020), ...

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Step 2: Ideally, amplify the noise and learn



Ideally wish

$$\Lambda_i^k(P_a) = f_{ia}^k P_a$$

Akin to:

RB, Cycle RB, K-body noise reconstruction, ...

S.T. Flammia and J.J. Wallman ACM Trans QC 1, 3 (2020), ...

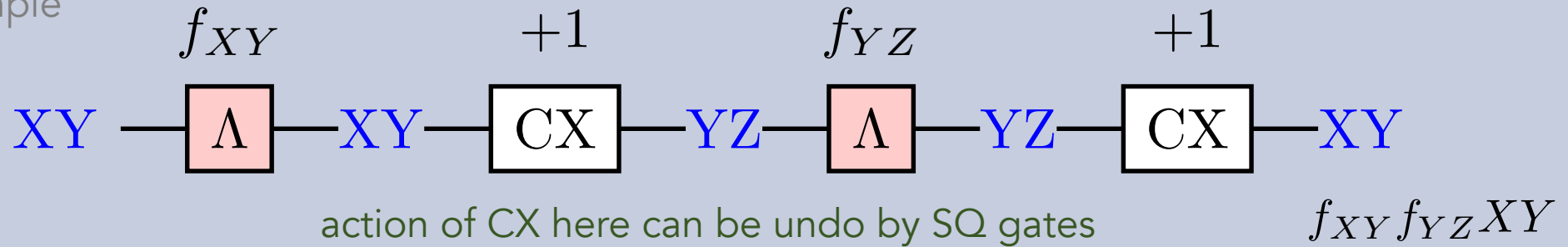
Erhard *et al.*, arXiv:1902.08543; Ferracin *et al.*, arXiv:2201.10672, ...

For something of a review of protocols, see Helsen, *et al.*, arXiv:2010.07974

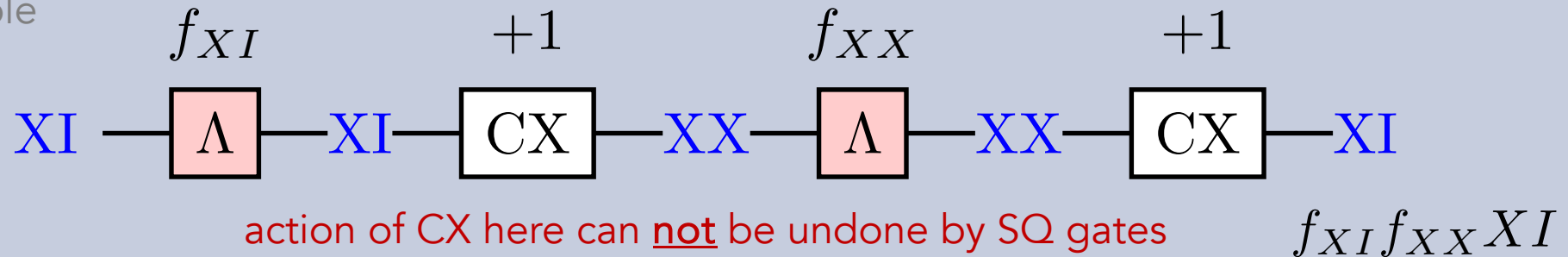
Let's see how the amplification works with gates: no-go theorem

$$\Lambda(P_a) = f_a P_a$$

2Q example



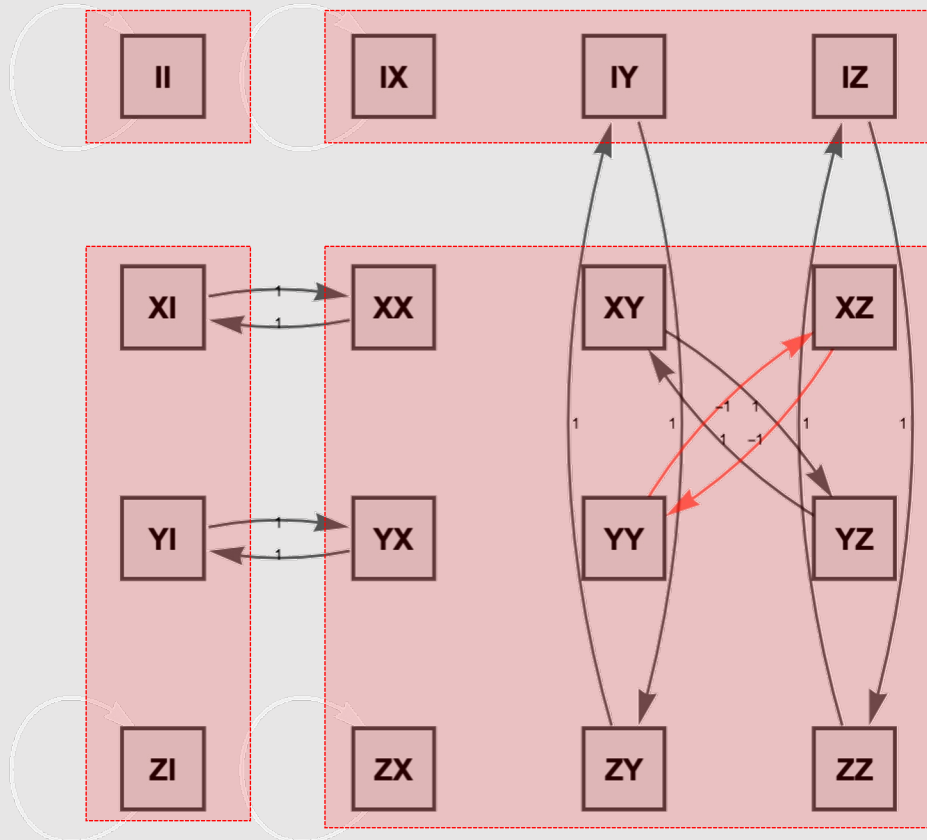
2Q example



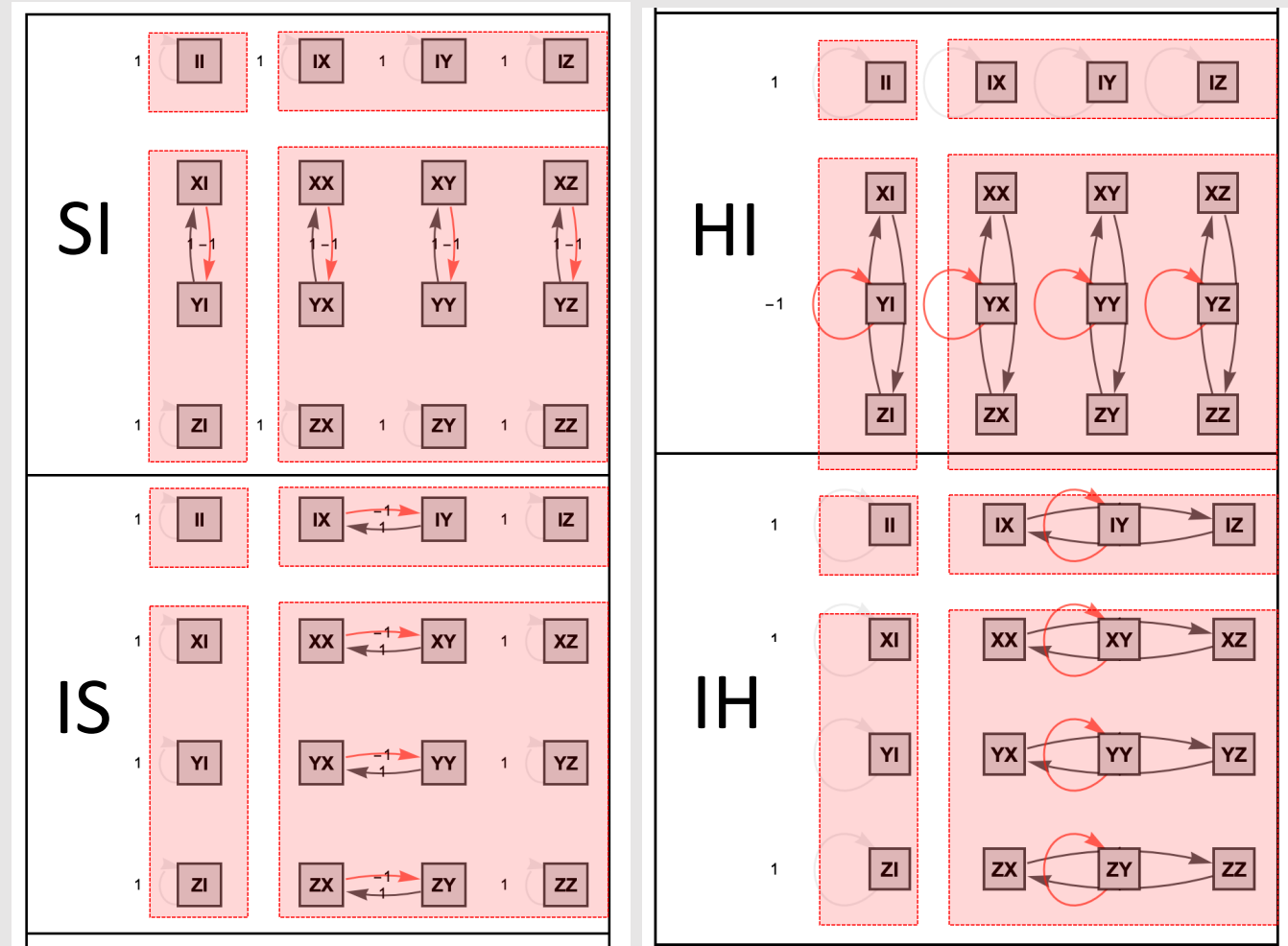
fundamental degeneracy – can not undo some non-local – need entangling operation

How to gates move state Paulis around?

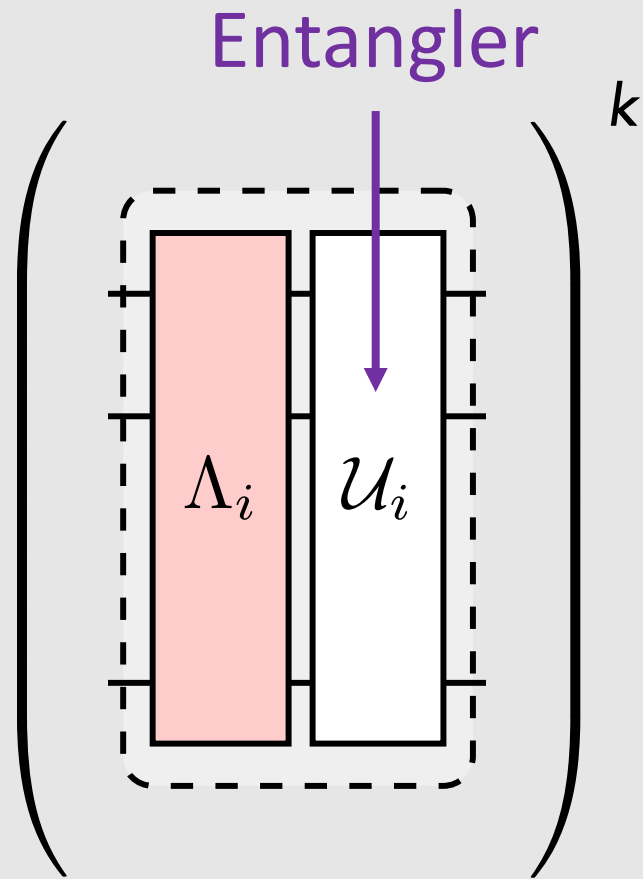
cNOT



Example single qubit gates



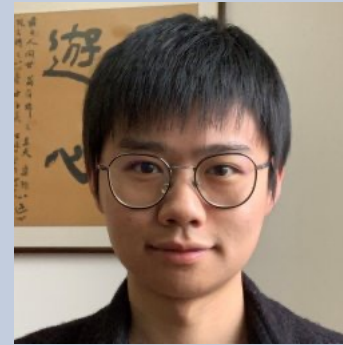
Not so simple to learn noise of entangling gates



Ideally wish

$$\cancel{\Lambda_i^k(P_a) = f_{ia}^k P_a}$$

Fundamental **no-go theorem on learning**



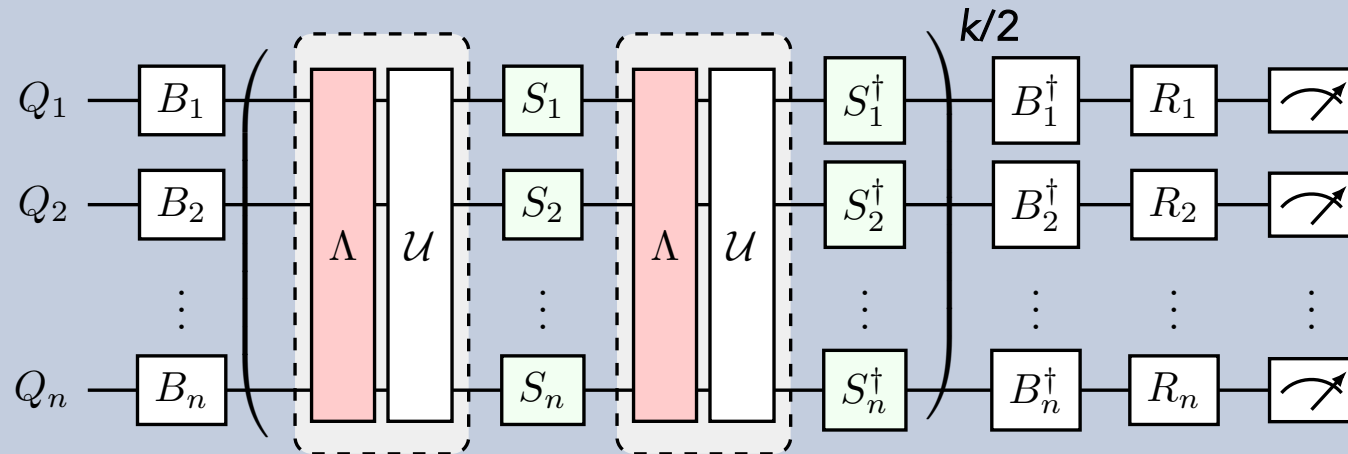
For general and in-depth

Senrui Chen, Y. Liu, M. Otten, A. Seif, B. Fefferman, L. Jiang
arXiv:2206.06362 (2022)

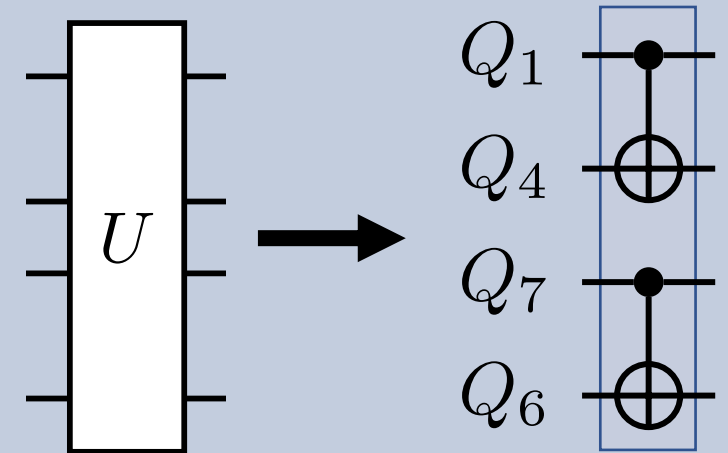
or supplement of our paper for qubit version and work by
S. Flammia, S Benjamin, and teams.

Solution: Custom protocol + weak assumption

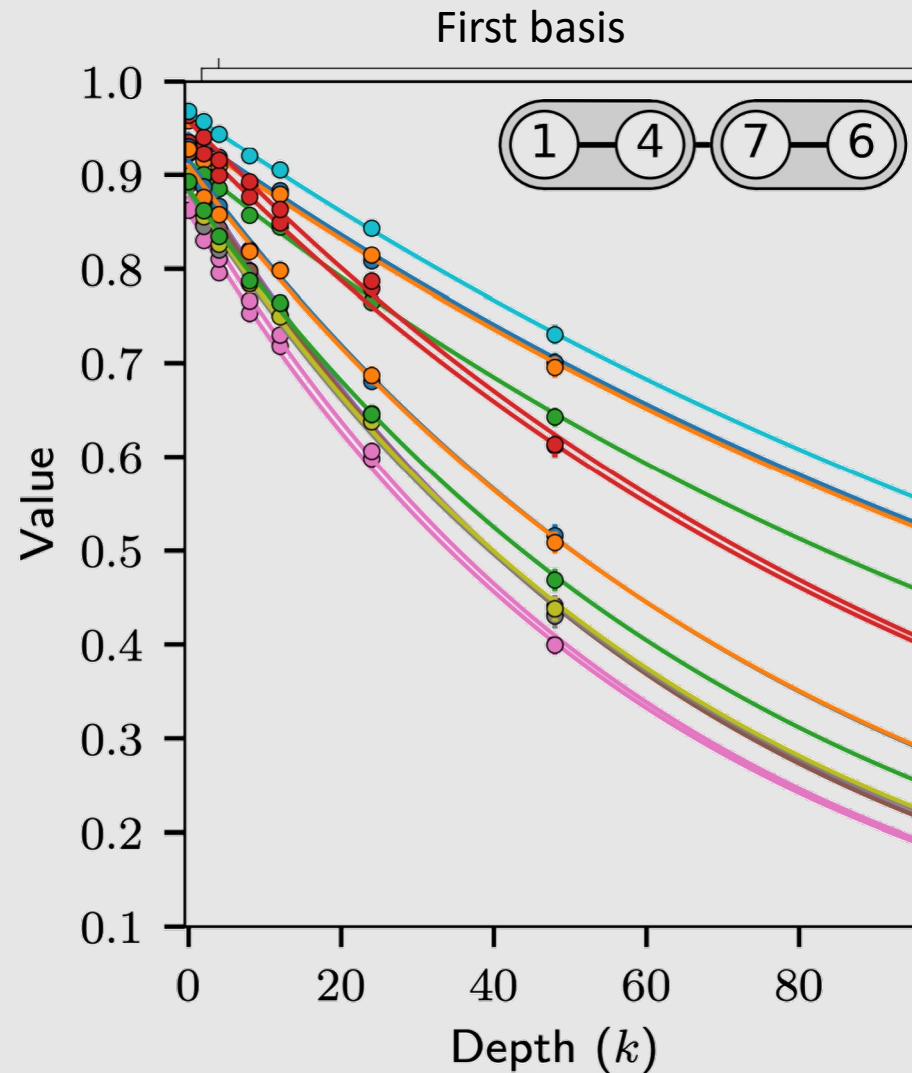
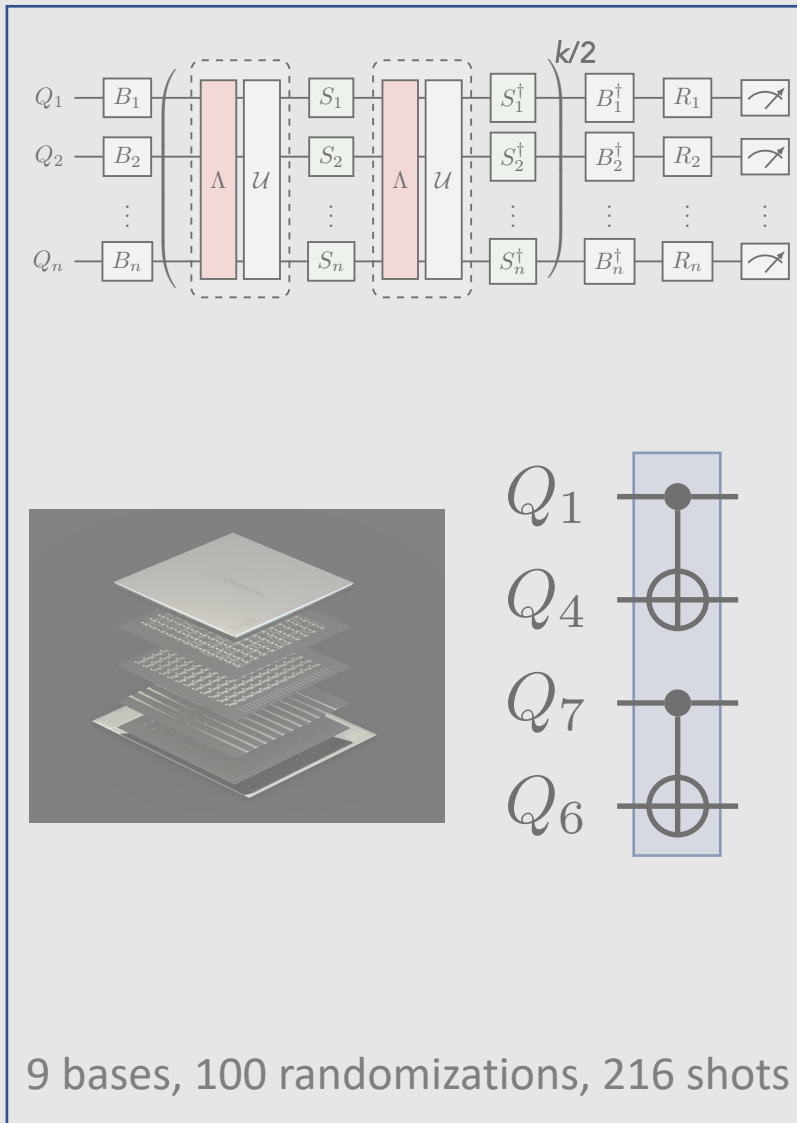
Learning circuits



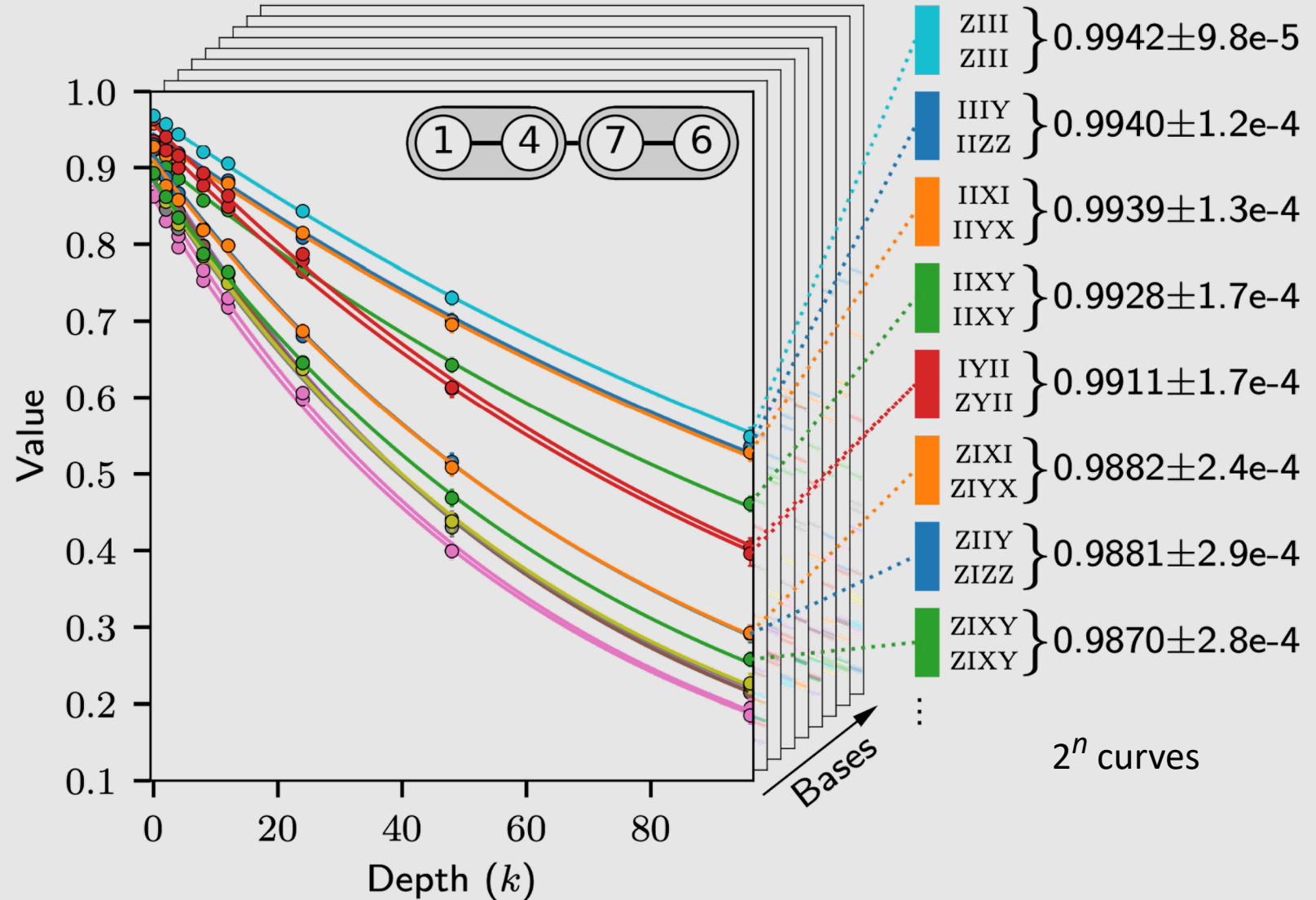
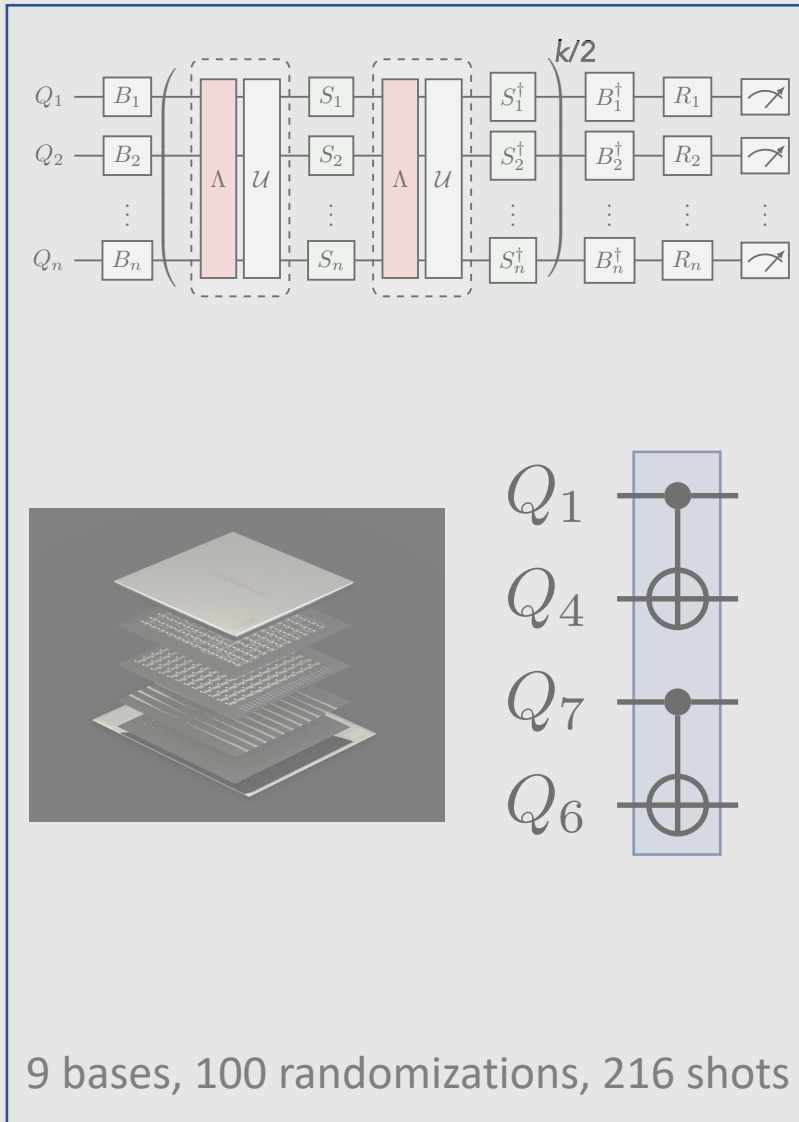
Example



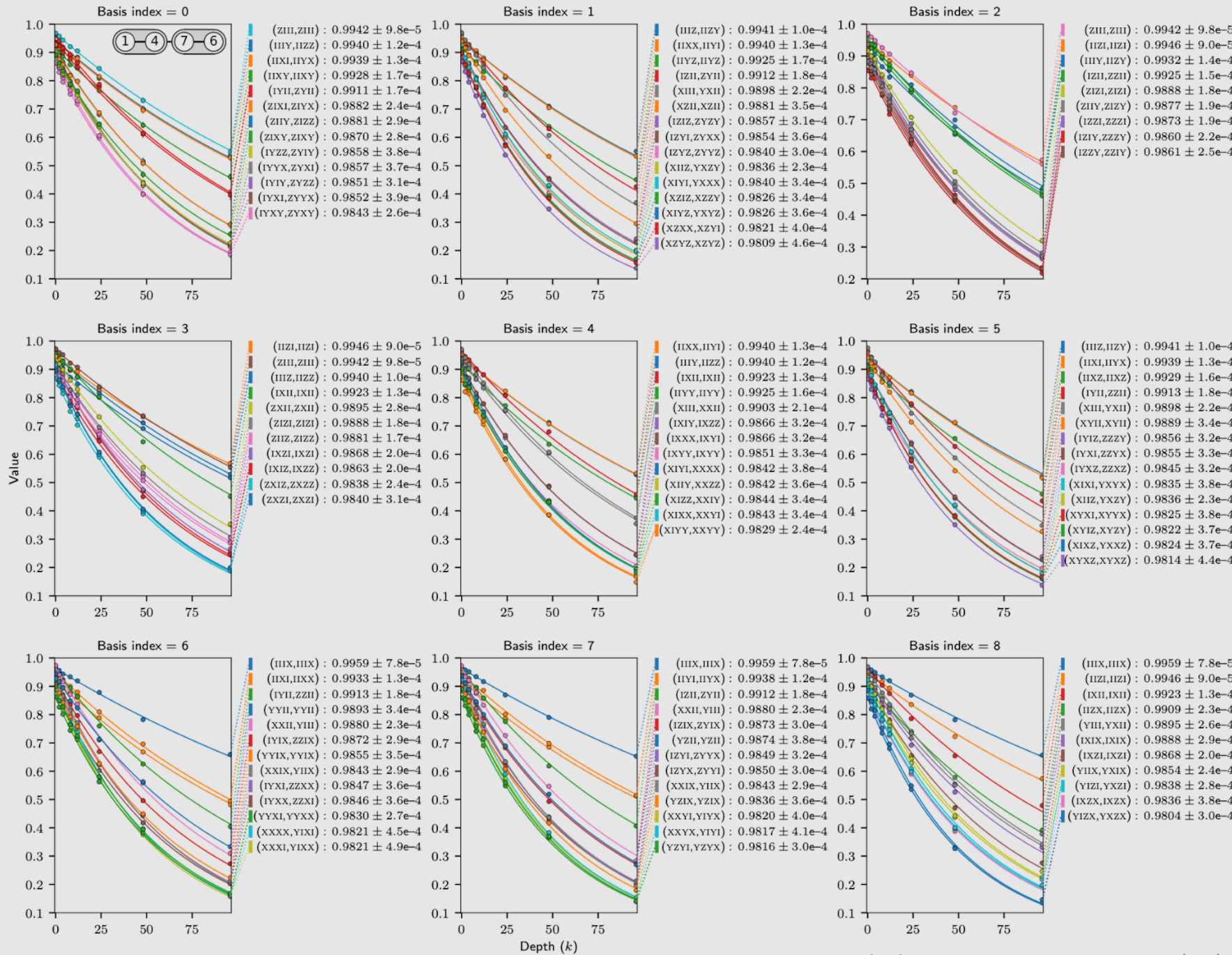
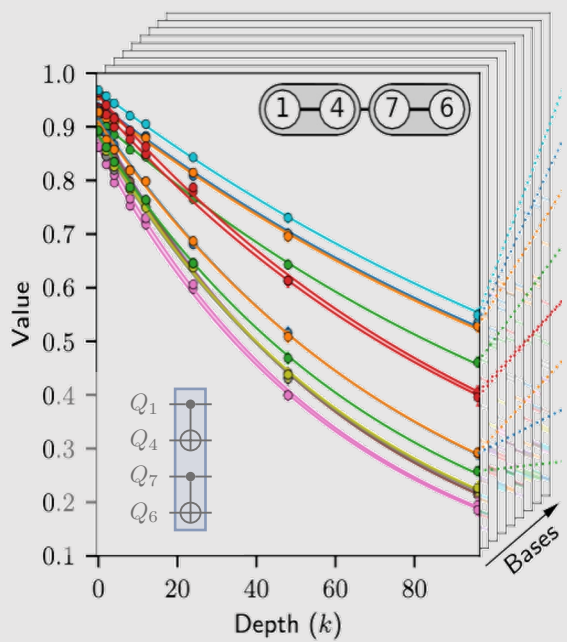
Learning the noise: raw data



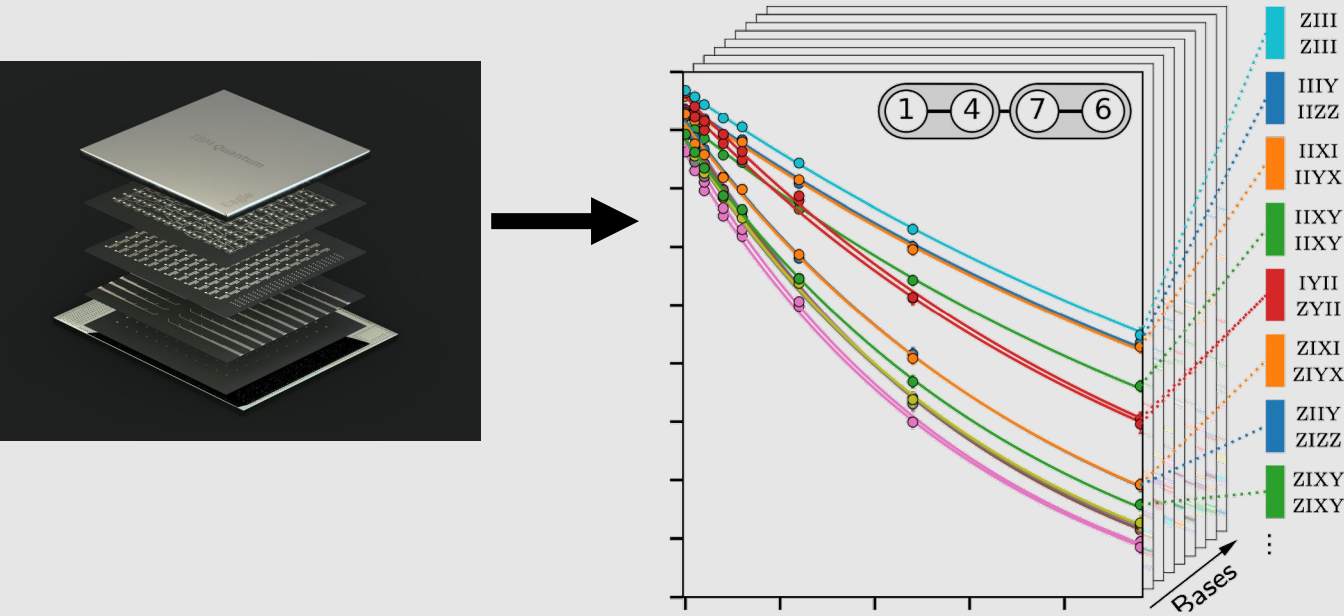
Learning the noise: raw data



Raw data



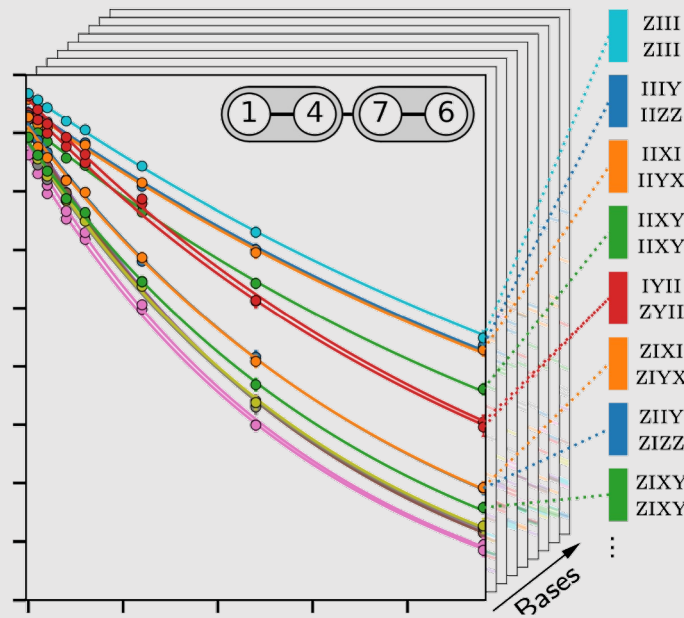
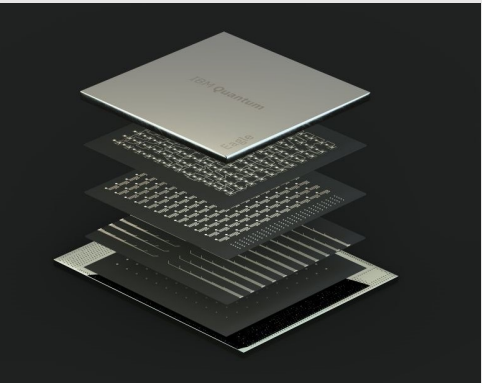
Reconstructing quantum channel from measurement data



$$\Lambda(\rho) = \sum_{a=0}^{4^n - 1} c_a P_a \rho P_a^\dagger$$

Still 4^n

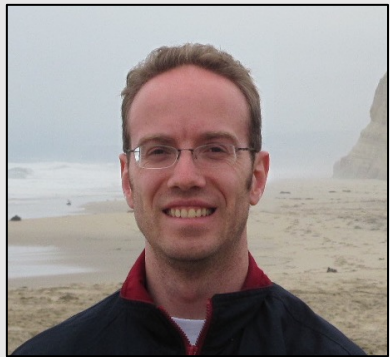
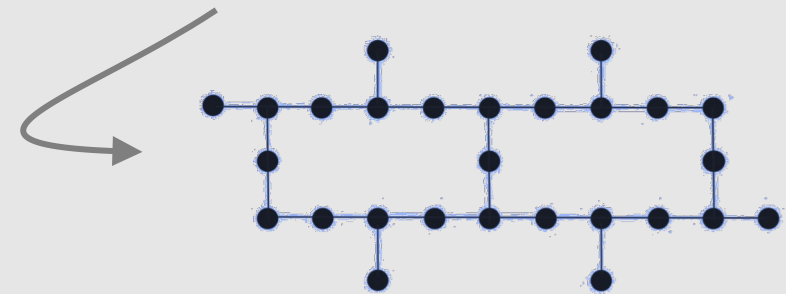
Sparse Pauli-Lindblad model



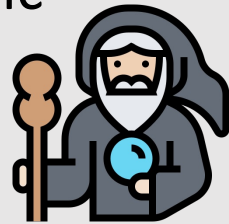
$$\Lambda(\rho) = \sum_{a=0}^{4^n - 1} c_a P_a \rho P_a^\dagger$$

$$\Lambda(\rho) = \exp[\mathcal{L}](\rho)$$

$$\mathcal{L}(\rho) = \sum_{k \in \mathcal{K}} \lambda_k (P_k \rho P_k - \rho)$$



Magic



icon: Eucalyp

Highlight: Ewout van den Berg

Notes

Dissipator for a given Pauli

$$\begin{aligned} \mathcal{D}[\rho]_P &= P \rho P^\dagger - \frac{1}{2} (P \rho P^\dagger + P^\dagger \rho P) \\ &= P \rho P^\dagger - \rho \\ &= (P \cdot P - I) \rho \end{aligned}$$

$$\begin{aligned} \hat{\mathcal{L}} &= \sum_a \gamma_a \left(\hat{P}_a - I \right) \\ &= \sum_a \gamma_a \mathcal{D}[\hat{P}_a] \end{aligned}$$

$$e^{\hat{\mathcal{L}}} = \prod_a \exp(\gamma_a \mathcal{D}[\hat{P}_a])$$

$$= \prod_a \Lambda[\hat{P}_a]$$

$$\text{but } [\mathcal{D}[\hat{P}_a], \mathcal{D}[\hat{P}_b]] = 0$$

Each sub channel

$$\Lambda[\rho_a] =$$

$$\begin{aligned} & \exp(\gamma \mathcal{D}[\rho_a]) \\ &= \exp(\gamma (\hat{P}_a + \hat{I})) \\ &= \exp(\gamma \hat{P}_a) \exp(\gamma \hat{I}) \exp\left(\frac{\gamma^2}{2} [\hat{P}_a, \hat{I}]\right) \dots \\ &= \exp(\gamma \hat{P}_a) \exp(\gamma \hat{I}) \end{aligned}$$

$$\hookrightarrow \sum_n \frac{\gamma^n \hat{I}^n}{n!} = \sum_n \frac{\gamma^n}{n!} \hat{I}$$

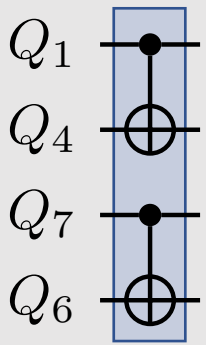
$$= \exp(\gamma) \hat{I} = \exp(\gamma) \hat{I} \cdot \hat{I}$$

Each sub channel

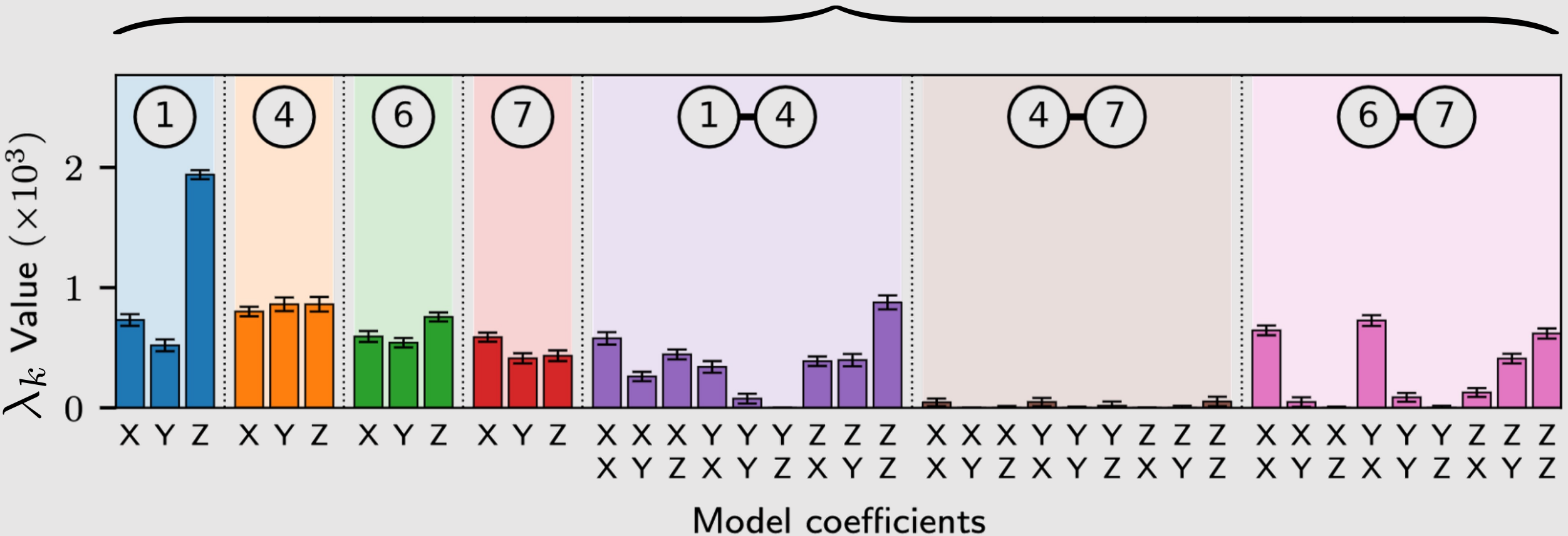
$$\begin{aligned}
 &= \exp(\gamma \hat{P}_a) \\
 &\quad \hookrightarrow \sum_{n=0}^{\infty} \frac{\gamma^n}{n!} (\hat{P}_a)^n \qquad \text{note } \hat{P}_a^{n-\text{even}} = \hat{P}_a^{\hat{1}} \\
 &\qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \hat{P}_a^{n-\text{odd}} = \hat{P}_a^{\hat{1}} \hat{P}_a \\
 &= \left(\sum_{\substack{n=0 \\ \text{even}}}^{\infty} \frac{\gamma^n}{n!} \right) \hat{P}_a^{\hat{1}} + \left(\sum_{\substack{n=1 \\ \text{odd}}}^{\infty} \frac{\gamma^n}{n!} \right) \hat{P}_a^{\hat{1}} \hat{P}_a \\
 &= \cosh(\gamma) \hat{P}_a^{\hat{1}} + \sinh(\gamma) \hat{P}_a^{\hat{1}} \hat{P}_a
 \end{aligned}$$

$$\Lambda[\hat{P}_a] = \cosh(\gamma_a) \hat{P}_a^{\hat{1}} + \sinh(\gamma_a) \hat{P}_a^{\hat{1}} \hat{P}_a$$

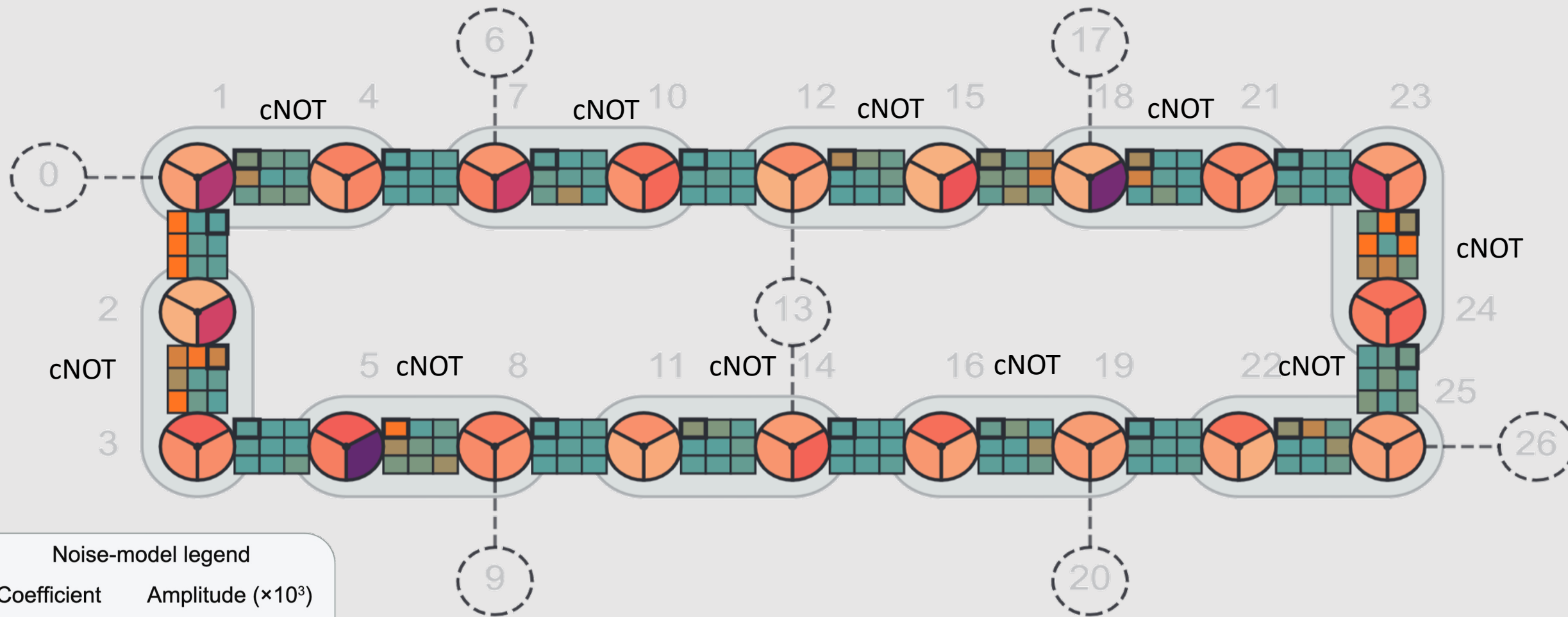
Sparse Lindblad tomogram



$$\mathcal{L}(\rho) = \sum_{k \in \mathcal{K}} \lambda_k \left(P_k \rho P_k^\dagger - \rho \right)$$




Noise tomogram for 20Q Ising-ring Trotter layer



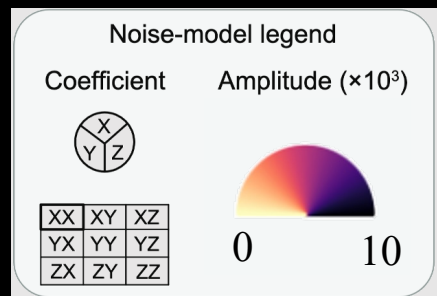
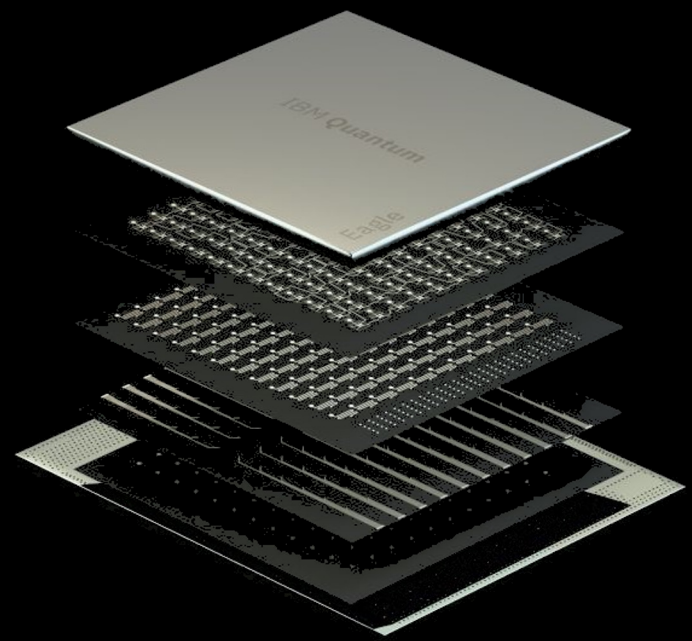
Noise-model legend

Coefficient Amplitude ($\times 10^3$)



XX	XY	XZ
YX	YY	YZ
ZX	ZY	ZZ

Noise tomogram for 127Q Trotter layer



Same number of learning circuits as for 4Q

