

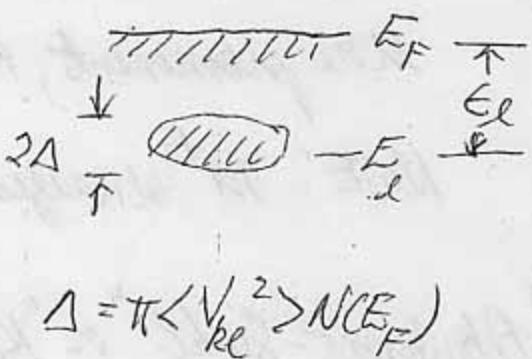
Kondo effectDilute alloy $M_{1-x}T_x$

M - nonmagnetic metallic host

T - impurity ion with partially-filled d- or f-electron shell and spin \downarrow

$$\mathcal{H}_{ex} = -2J \sum \vec{s}_i \cdot \vec{s}_j$$

$$J_{nn} = \frac{\langle V_{ke}^2 \rangle}{\epsilon} < 0 \text{ (AFM)}$$



$$T_K \sim T_F \exp(-1/N(E_F)/|J|)$$

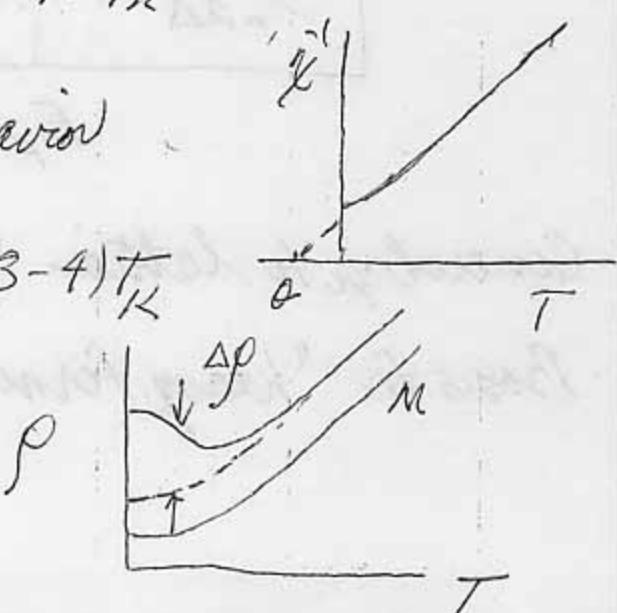
Gradual formation of many body singlet ground state
as T decreases through T_K

Physical properties scale with T_K

$T \gg T_K$ local moment behavior

$$\chi(T) = \frac{N\mu_{eff}^2}{3k_B(T+0)} \cdot 101 \sim (\beta - 4)T_K$$

$$\Delta\rho(T) \sim -\ln(T/T_K)$$



$T \ll T_K$ Nonmagnetic
Local Fermi liquid; effective Fermi temperature

$$\Delta\chi(T), \Delta\chi(T)/T = \gamma \xrightarrow{T \rightarrow 0} \text{Const.} \propto N(E_F)$$

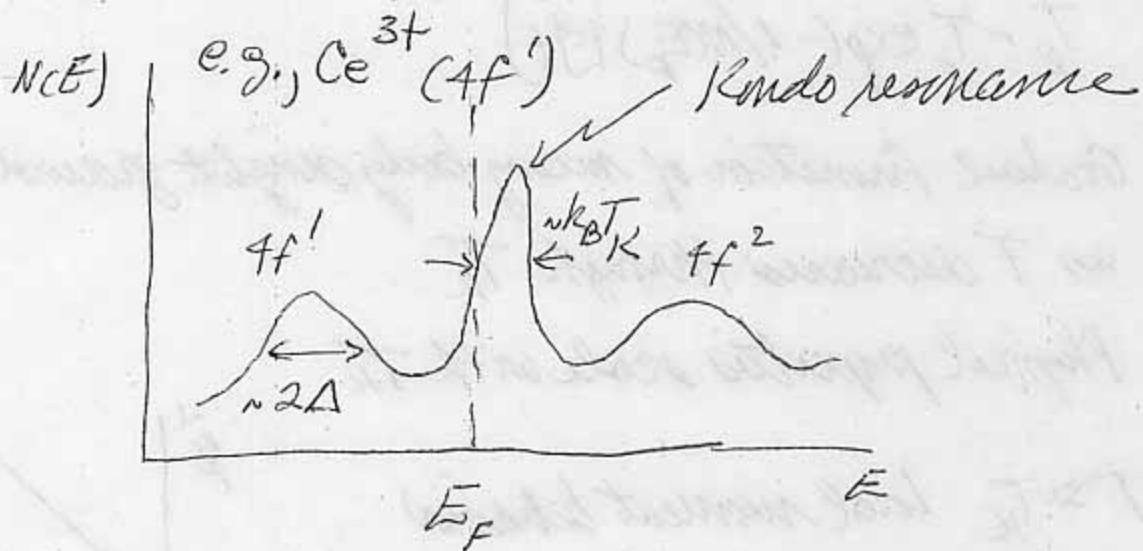
$$T_F^* = T_K$$

$$\rho(T) = \rho_0 \left[1 - \left(\frac{T}{T_K} \right)^2 \right] \xrightarrow{\sim} \propto M^* \propto \frac{1}{T_K}$$

More prominent, the smaller T_K

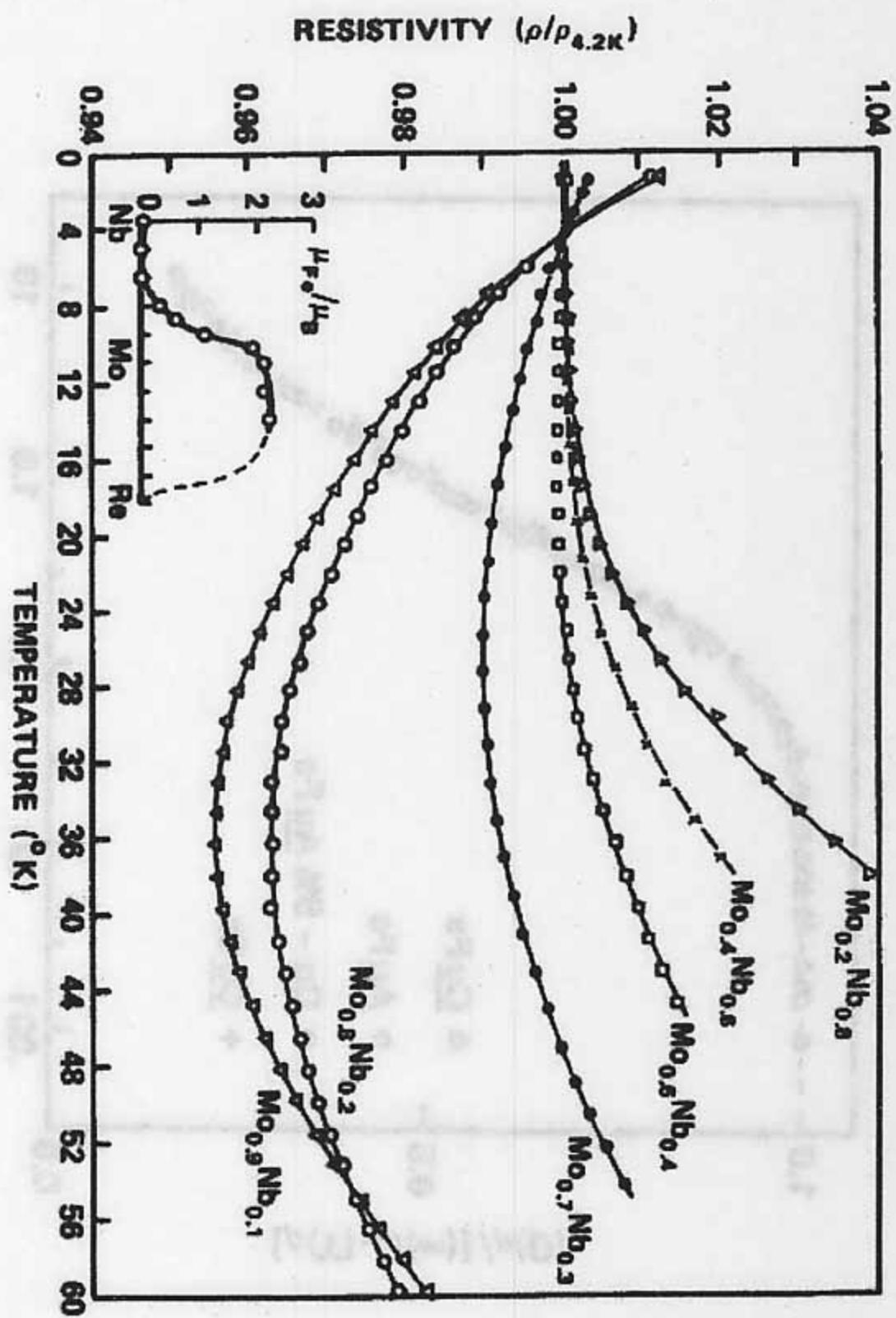
NOTE: M^* diverges as $T_K \rightarrow 0$ ("heavy fermion")

"Anderson-Suhl" or "Kondo" resonance

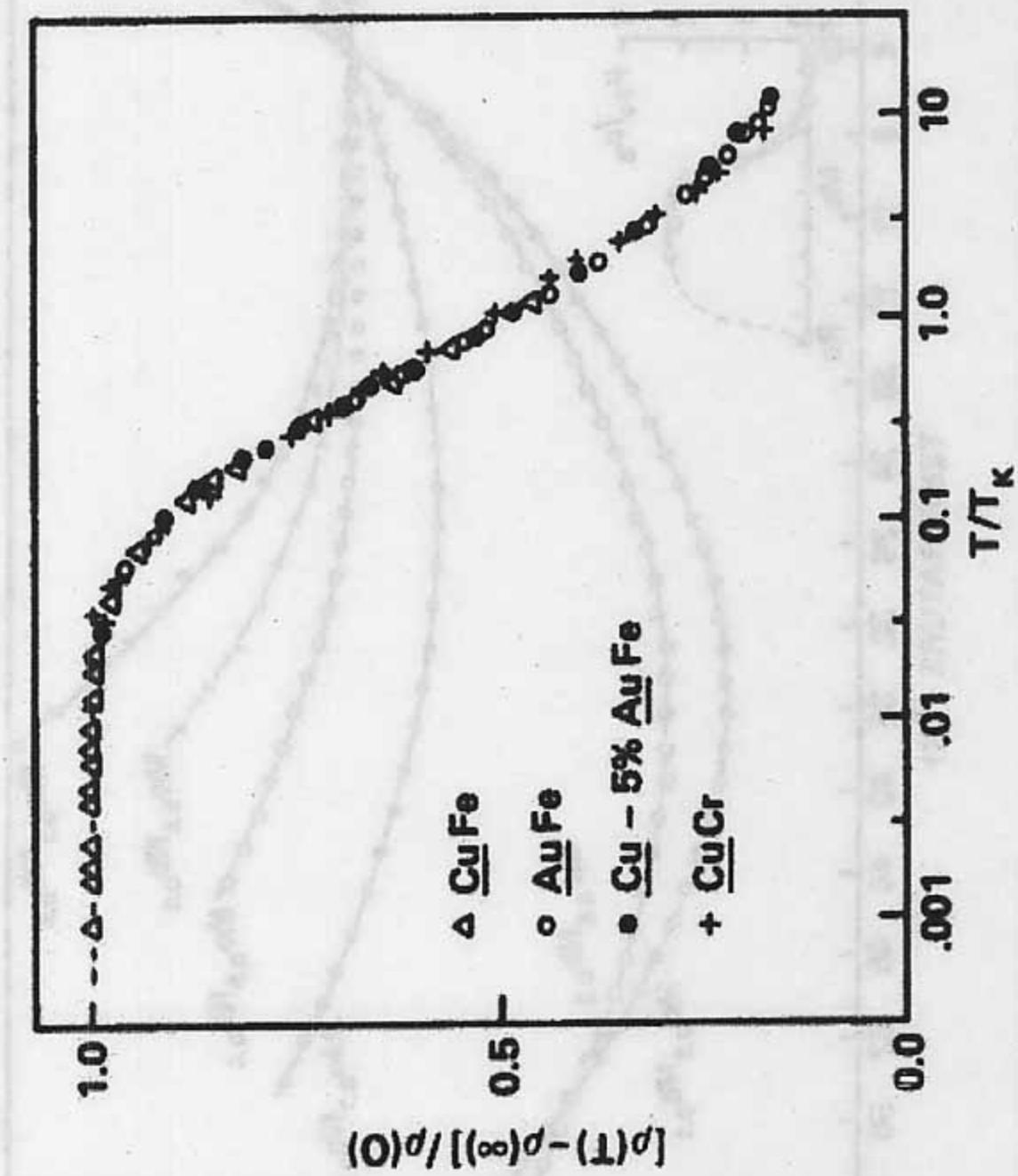


Generalize to lattice "Kondo lattice"
Base for "heavy fermion" metal

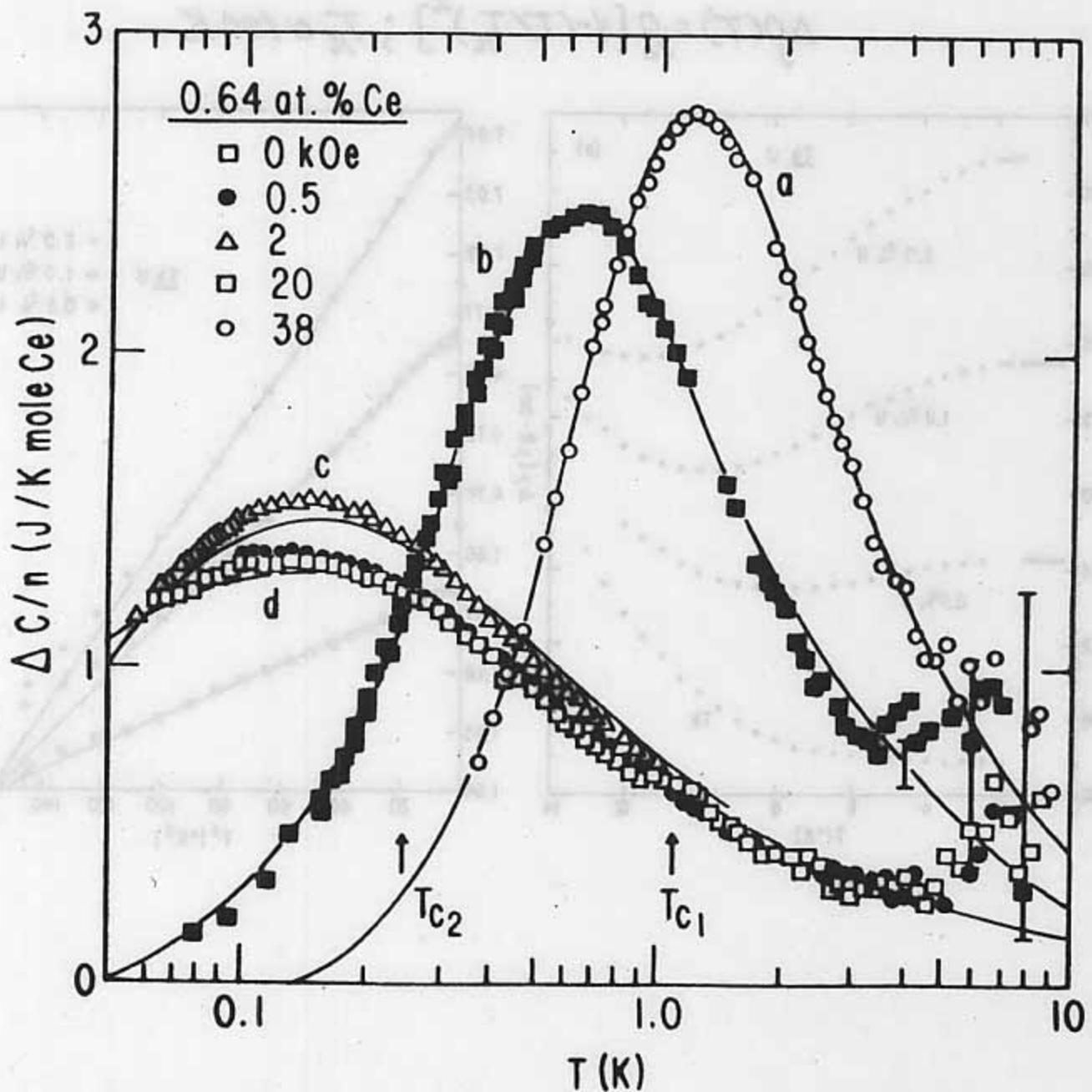
Characteristic Kondo resistance minimum associated with Fe dissolved in Mo-Nb alloys. After Sarachik *et al.* '64.
Inset: μ_{eff} vs composition. After Clogston *et al.* '62.



Excess resistivity associated with scattering from the impurity spin. The resistivity saturates at low temperatures because all of the d-wave is being scattered. After White and Geballe '79.



$(La, Ce)Al_2$



S.D. Bader, N.E. Phillips, M.B. Maple & C.A. Luengo (1975)

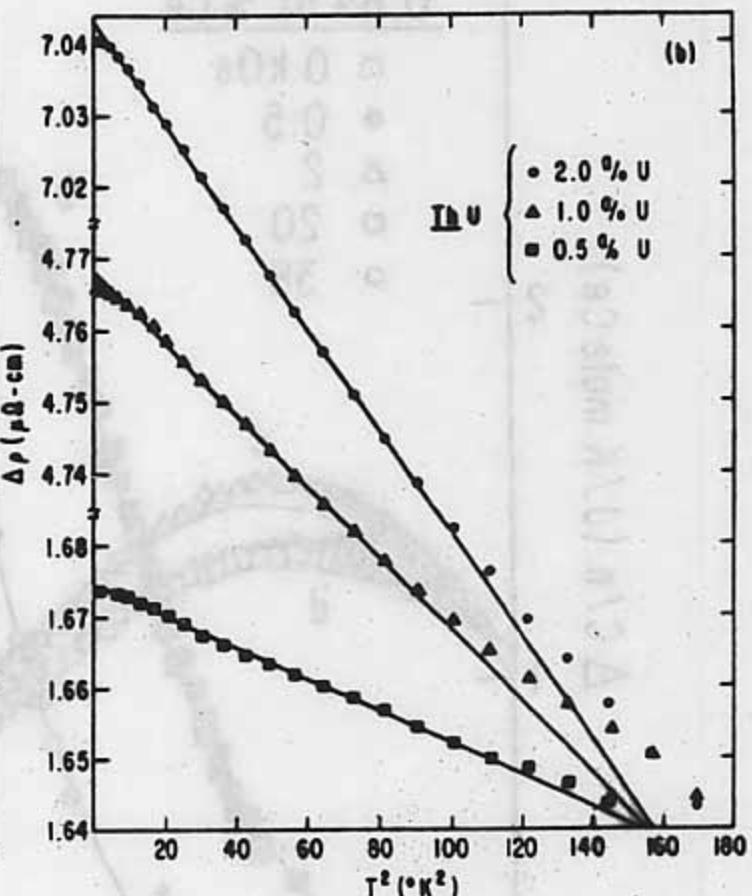
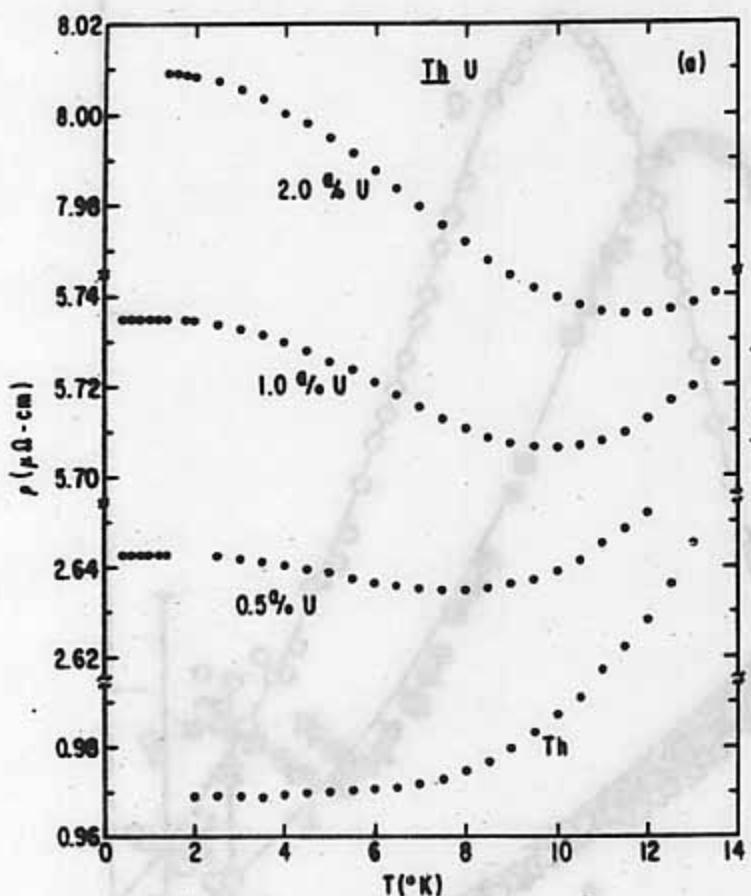
Curve d - Bloomfield-Hamann theory for $S = \frac{1}{2}$, $T_k = 0.42 K$
(also consistent with Bethe Ansatz solution)

Rajan et al. - 1983

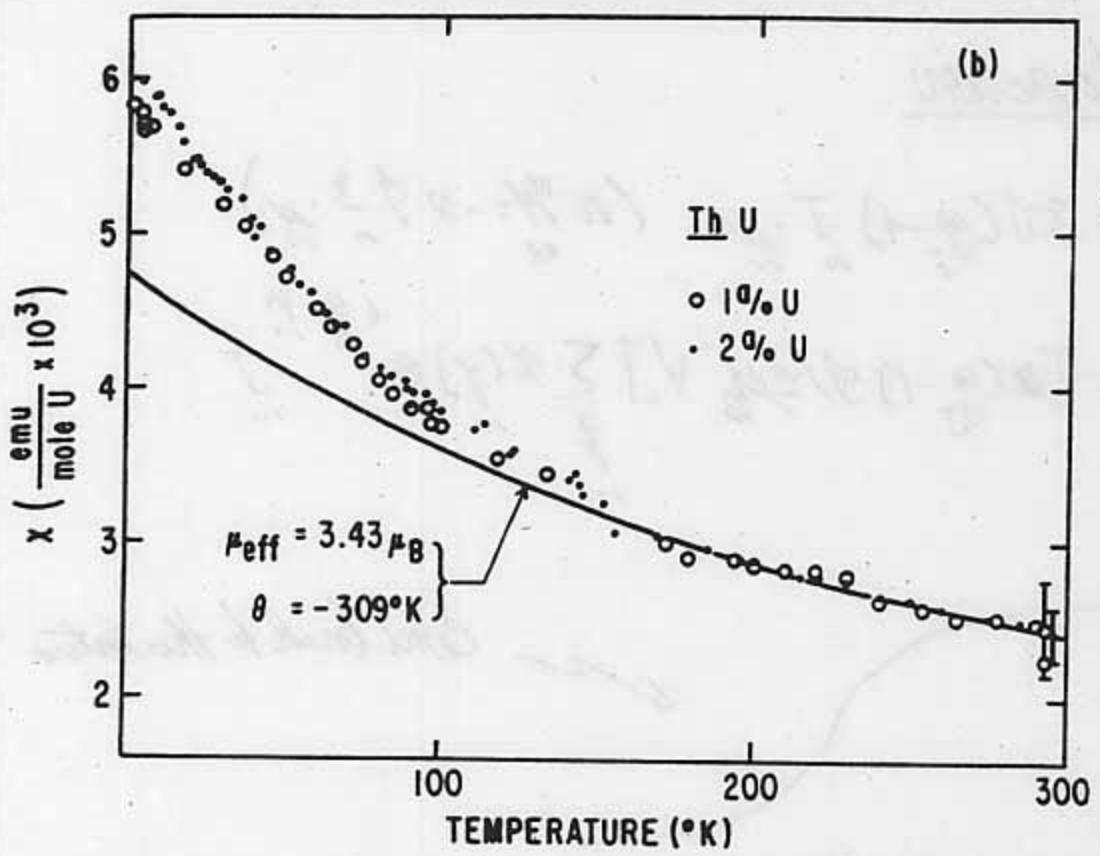
$Th_{1-x}U_x$ Conventional Kondo effect (Fermi liquid-low T)

$$\Delta\gamma \approx 270 \text{ mJ/mol U-K}^2$$

$$\Delta\rho(T) = \rho_0 [1 - (T/T_K)^2]; T_K \approx 100 \text{ K}$$



M.B. Maple et al. '70

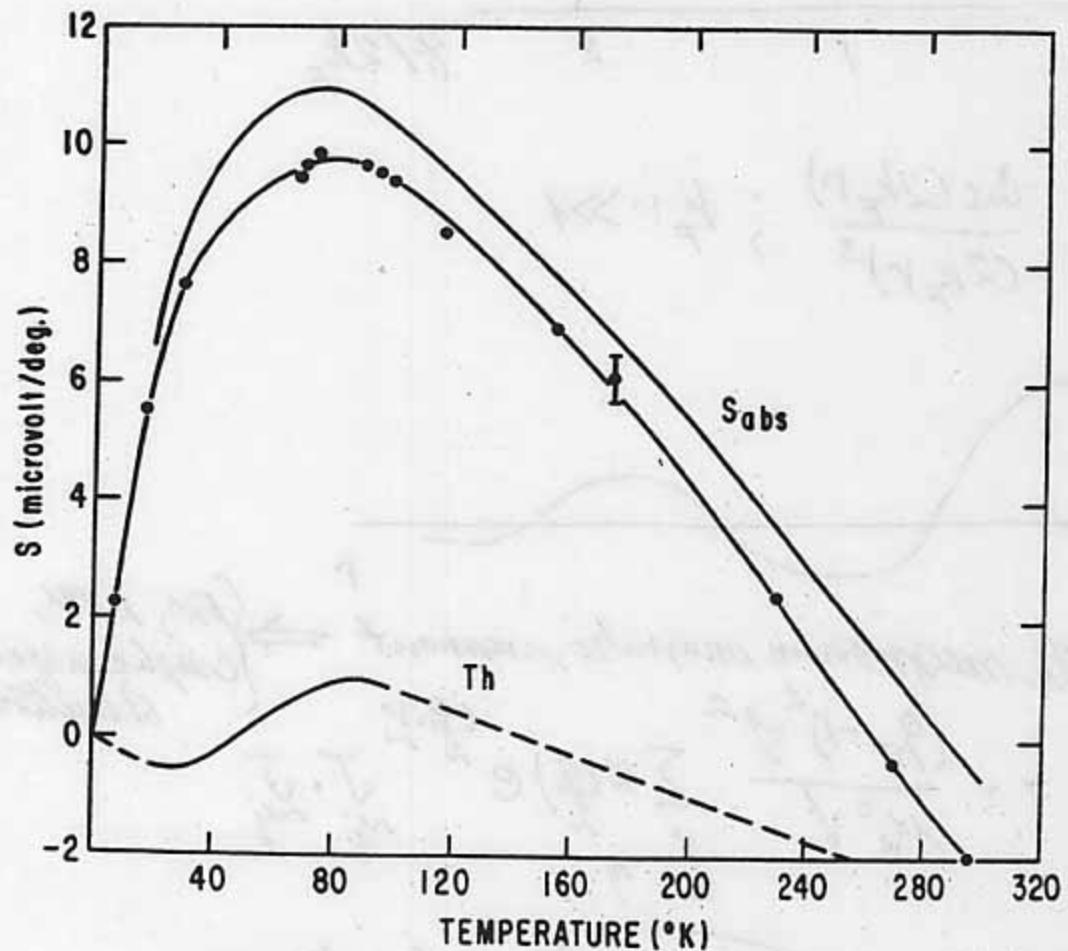


$$\frac{T h_{1-x} U_x}{}$$

$$\chi(T) = C/(T-\theta)$$

$$\theta \approx -3 T_K$$

$$C = N \mu_{\text{eff}}^2 / 3 k_B$$

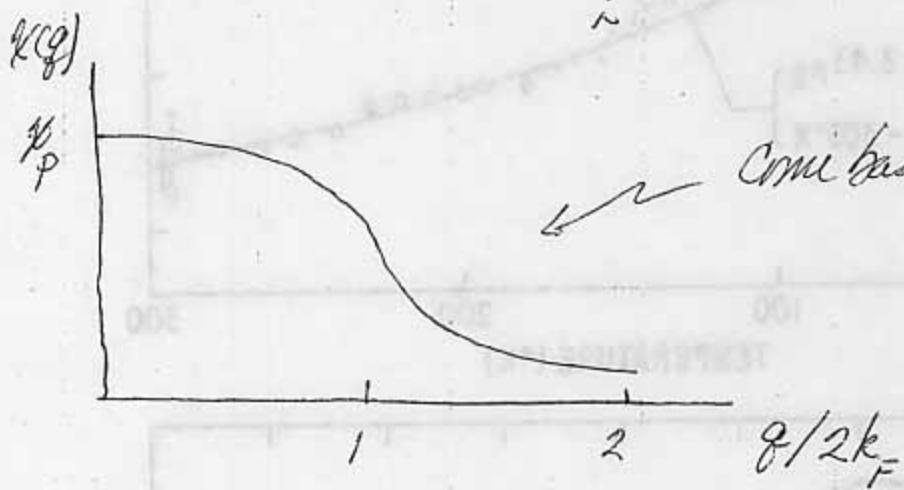


M.B. Maple et al. '70

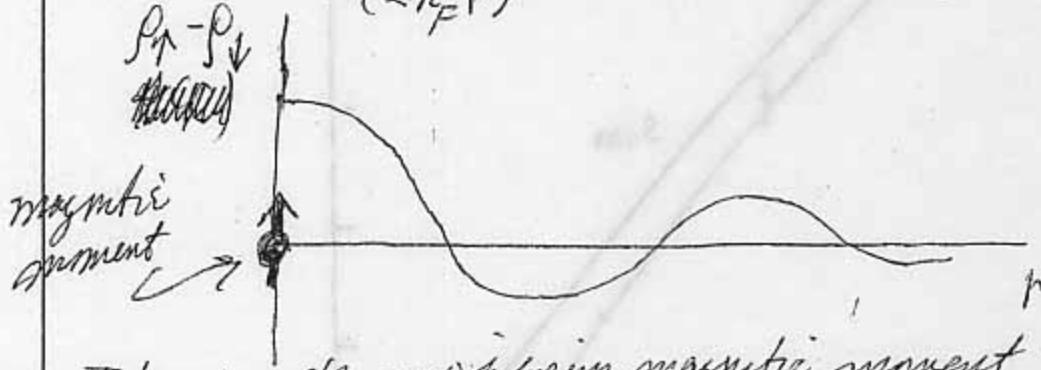
RKKY interaction

$$\mathcal{M}_{\text{ex}} = -2J(g_J - 1) \frac{J}{V} \sum_{\mathbf{q}} \chi(\mathbf{q}) \quad (\text{or } \mathcal{M}_{\text{ex}} = -2g_J S \cdot \mathcal{P})$$

$$\chi(\mathbf{q}) = [J(g_J - 1) \frac{J}{2\mu_B^2 V}] \sum_{\mathbf{q}} \chi(\mathbf{q}) e^{i\mathbf{q} \cdot \mathbf{r}}$$



$$\chi(\mathbf{q}) \sim \frac{\cos(2k_F r)}{(2k_F r)^3}; k_F r \gg 1$$



Interacts with neighboring magnetic moment \Rightarrow {FM, AFM, complex magnetic structures}

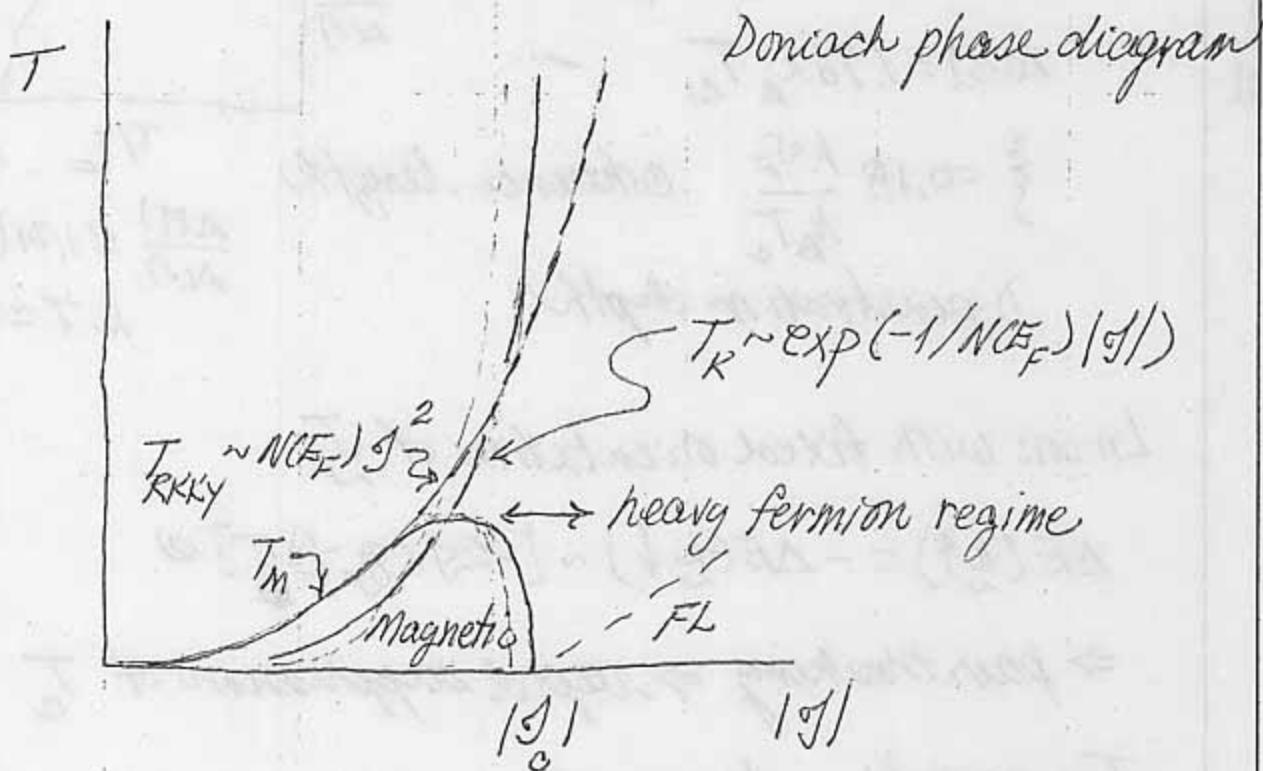
$$\mathcal{M}_{\text{RKKY}} = -\frac{(g_J - 1)^2 J^2}{\mu_B^2 V} \sum_{\mathbf{q}} \chi(\mathbf{q}) e^{i\mathbf{q} \cdot \mathbf{r}} \sum_{i,j} J_i J_j$$

Electron band structure

Competition between magnetic and nonmagnetic ground states

Local singlet (Kondo) : $T_K \sim \exp(-1/N\langle E_F \rangle / |J|)$

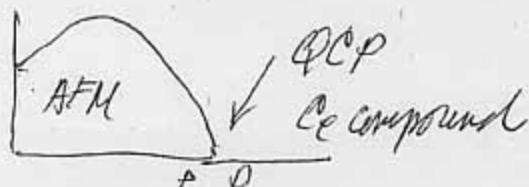
Magnetic order (RKKY) : $T_{RKKY} \sim N\langle E_F \rangle J^2$



$0 \leq |J| < |J_c|$: Local moment - magnetic order
"small" Fermi volume

$|J| > |J_c|$: Local singlet - nonmagnetic
"large" Fermi volume

NOTE: $P \Rightarrow$ measure $|J|$ fr. T_N



Paramagnetic impurities in conventional superconductors

SC: (\downarrow , \uparrow) Cooper pairs

$$T_c \sim \theta_D \exp(-1/N\langle E_F \rangle V)$$

$$F_n(0) - F_s(0) = \frac{1}{2} N\langle E_F \rangle \Delta^2(0)$$

$$\Delta(0) = 1.76 k_B T_c$$

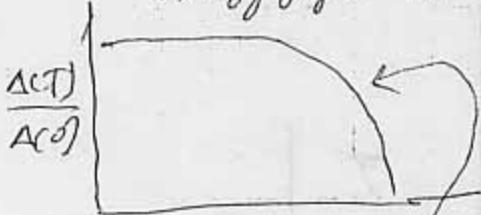
$$\xi = 0.18 \frac{\hbar v_F}{k_B T_c}$$

coherence length

λ penetration depth

BCS

Energy gap (0 P)



T/T_c

$$\frac{\Delta(T)}{\Delta(0)} \approx 1.74 \left(1 - \frac{T}{T_c}\right)^{1/2}$$

$\ln T \approx T_c$

In in with fixed orientation of \vec{J}

$$\Delta E(\downarrow, \uparrow) = -\Delta E(\downarrow, \downarrow) \sim [-2J(g_f - 1)J]^{1/2}$$

\Rightarrow pair breaking \Rightarrow rapid suppression of T_c

Two cases: $J > 0$ FM

$J < 0$ AFM (Kondo effect)

$\Delta > 0 \text{ FM}$

$$T_c/T_{c_0} = \ln(\alpha/\alpha_{cr}) \quad \text{Abrikosov & Gor'kov (AG) '60}$$

α pair breaking parameter

$$\alpha = \frac{\gamma^{-1}}{n\hbar N(E_F)} = n\hbar N(E_F) \gamma^2 \underbrace{(q_J - 1)^2}_{\rightarrow \text{lifetime}} \overline{J(J+1)} \quad \begin{matrix} \gamma \rightarrow \text{lifetime} \\ \text{gap opening} \end{matrix}$$

$$\alpha_{cr} = k_B T_c / 4\pi\delta^4 \quad (\ln \delta \text{ Euler's constant})$$

$$\alpha \rightarrow 0$$

$$T_c/T_{c_0} = 1 - 0.691(\alpha/\alpha_{cr}) = 1 - 0.691(n/n_{cr})$$

Linear region $n/n_{cr} \ll 1$

$$\left. \frac{dT_c}{dn} \right|_{n=0} = -(\pi^2/2) k_B^{-1} N(E_F) \gamma^2 \underbrace{(q_J - 1)^2}_{\text{de Gennes' factor}} \overline{J(J+1)}$$

Other predictions

$$\Delta C/\Delta C_0 = \sqrt{n} (T_c/T_{c_0})$$

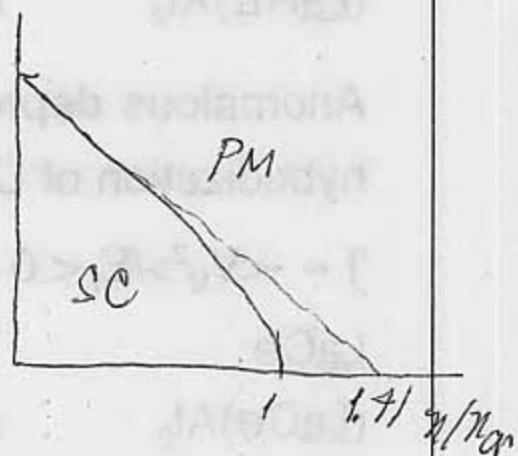
Depression relative to BCS
law of corresponding states

$$(\Delta C/\Delta C_0 = T_c/T_{c_0})$$

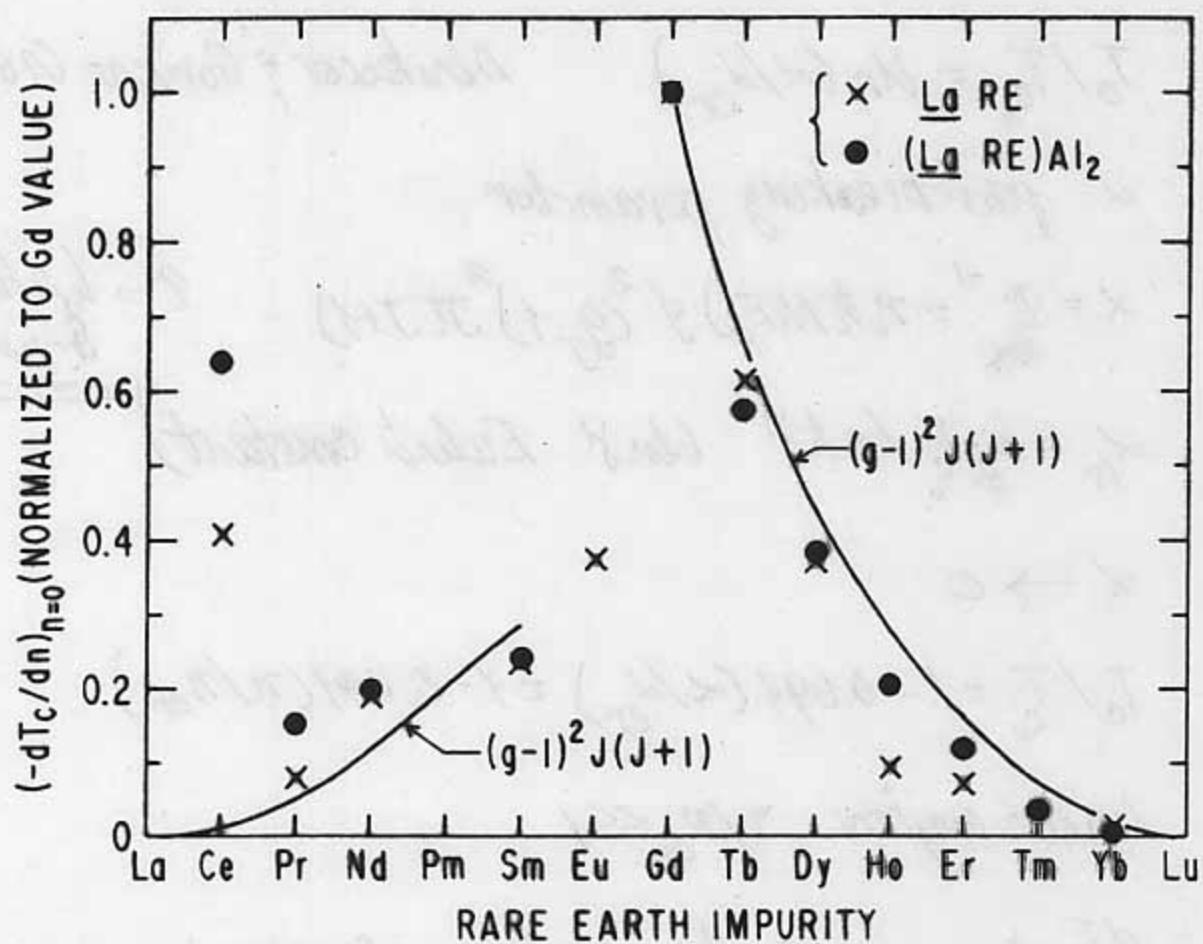
Gapless SC

$$\Delta \leftrightarrow \alpha$$

$\alpha \rightarrow 0$ faster than $\Delta \rightarrow 0$ with α



Paramagnetic impurities in superconductors



La: $T_\infty = 6$ K, LaAl_2 : $T_\infty = 3.3$ K, $J \sim 0.1$ eV

LaRE *B. T. Matthias, H. Suhl, E. Corenzwit '58*

(LaRE)Al₂ *M. B. Maple '70*

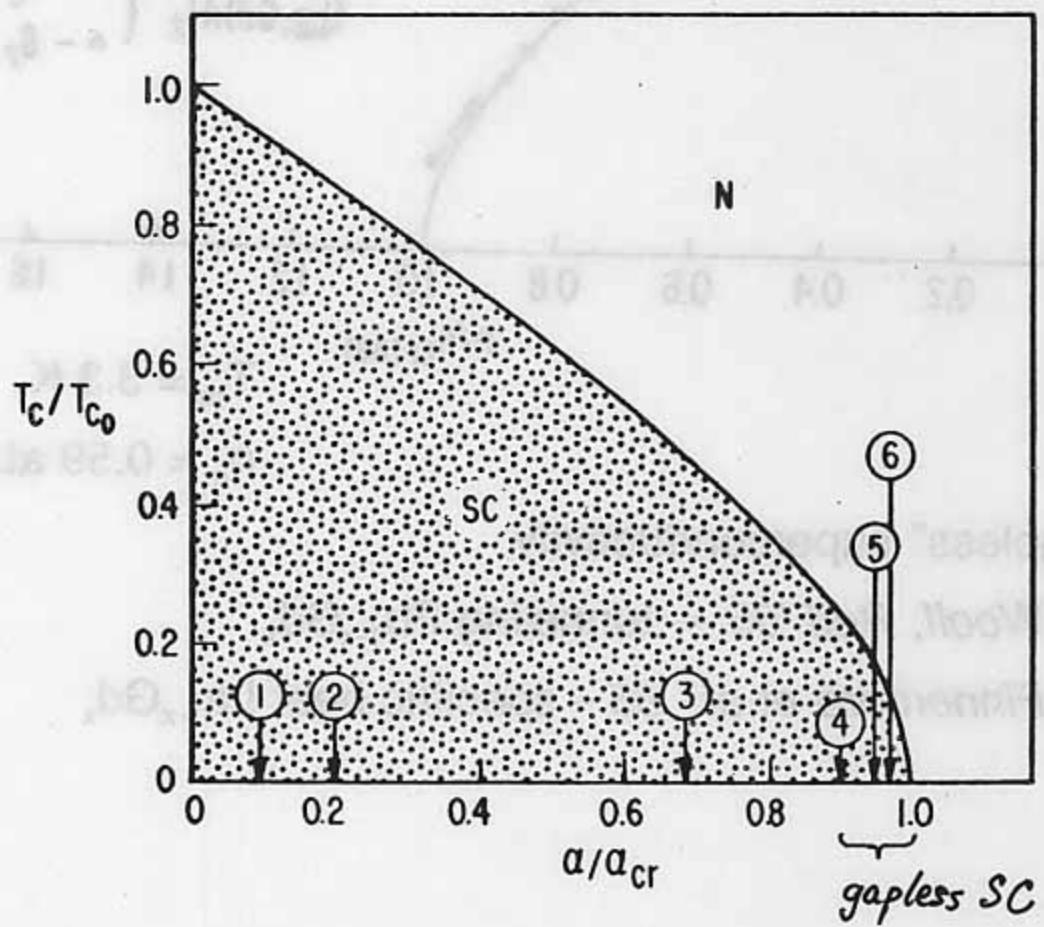
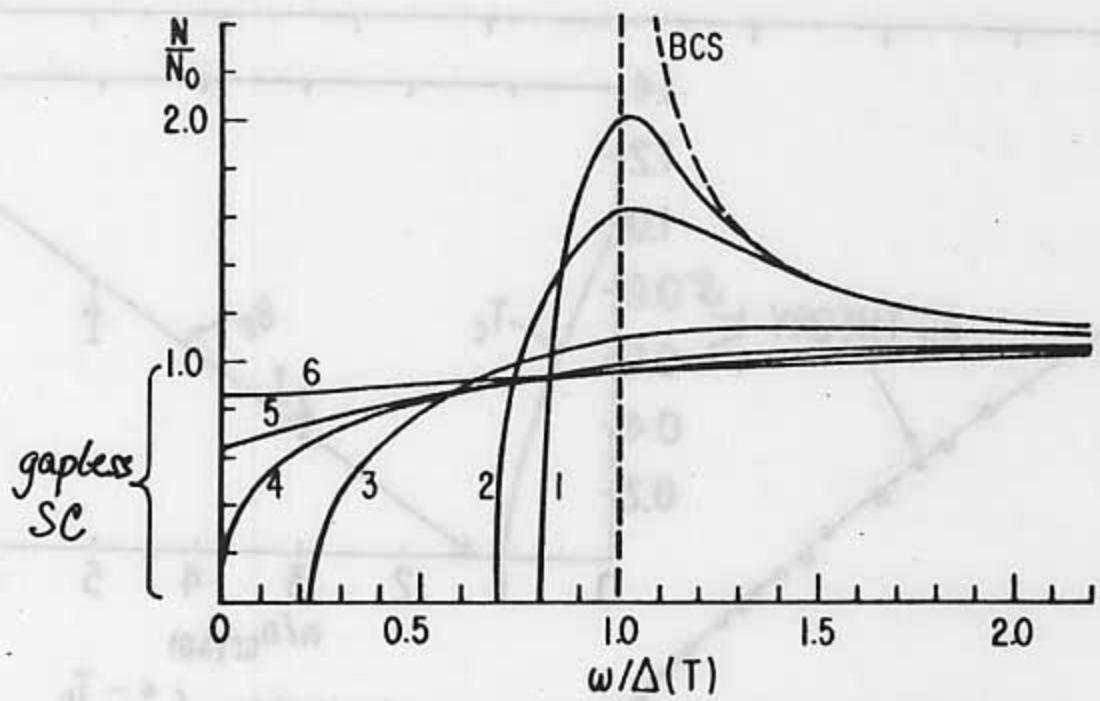
Anomalous depression of T_c for Ce \Rightarrow

hybridization of Ce 4f and conduction electron states \Rightarrow

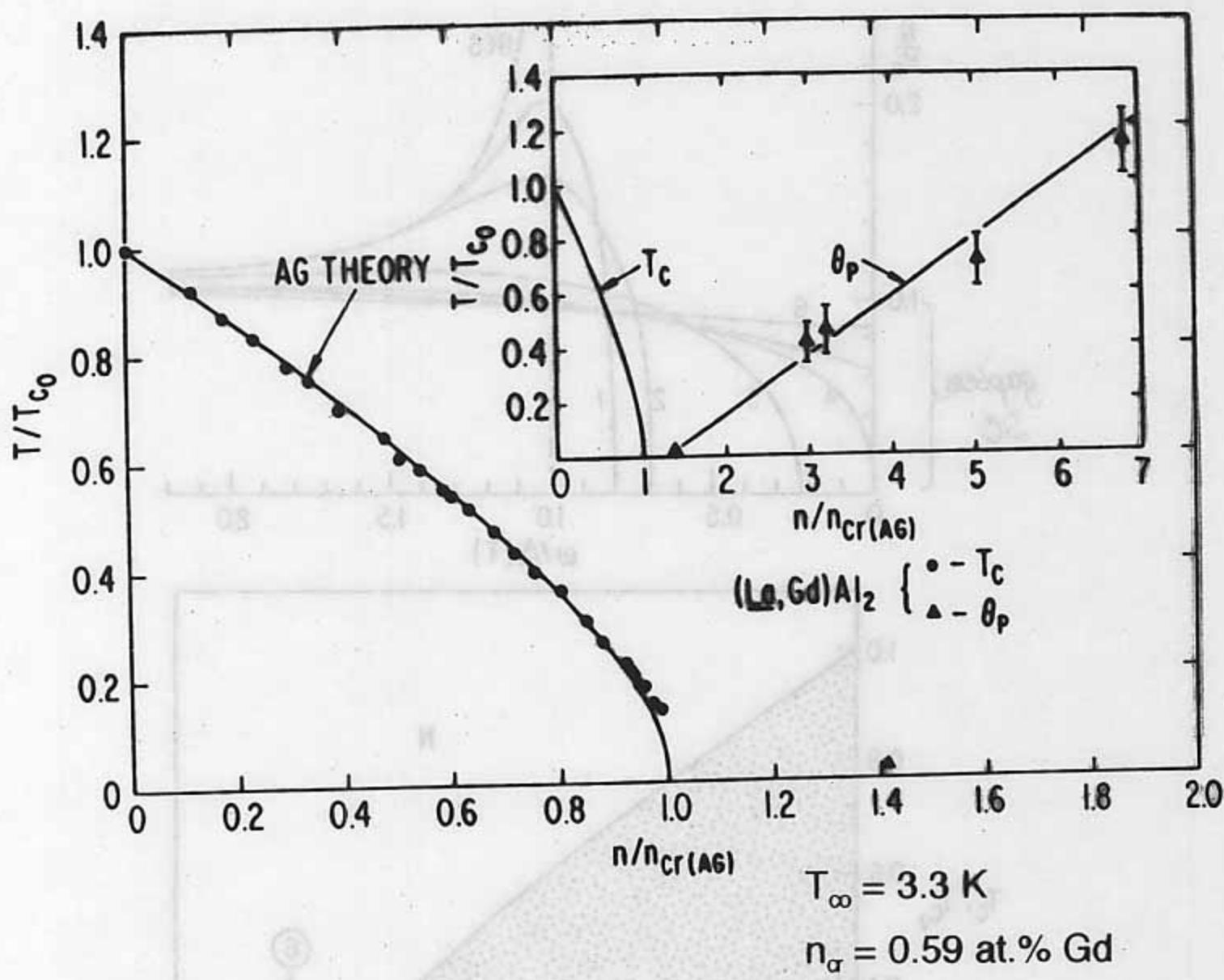
$J \sim -\langle V_{kf}^2 \rangle / E_f < 0 \Rightarrow$ Kondo effect

LaCe *T. Sugawara, H. Eguchi '66*

(LaCe)Al₂ *M. B. Maple, Z. Fisk '68*



Paramagnetic impurities in superconductors

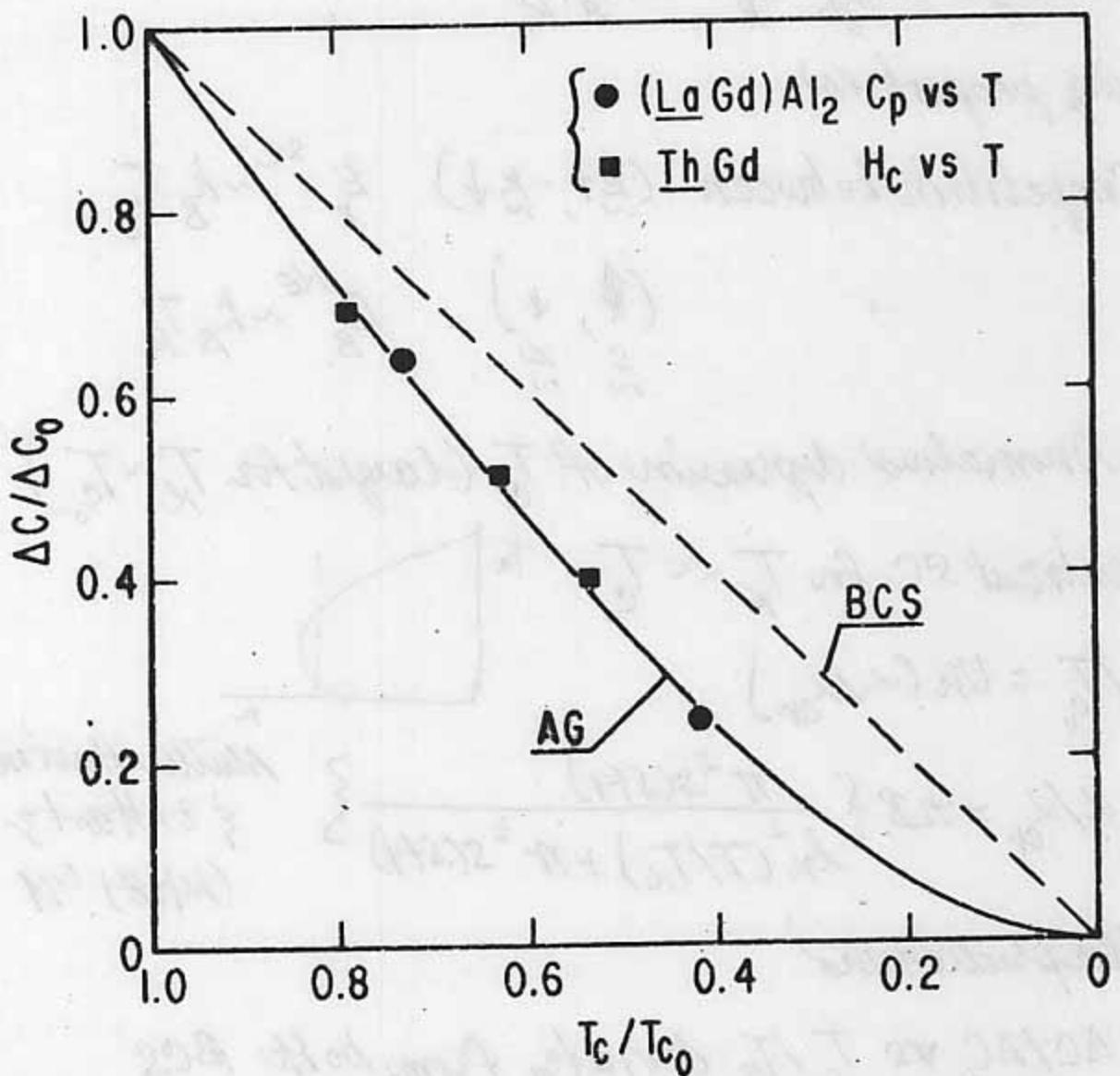
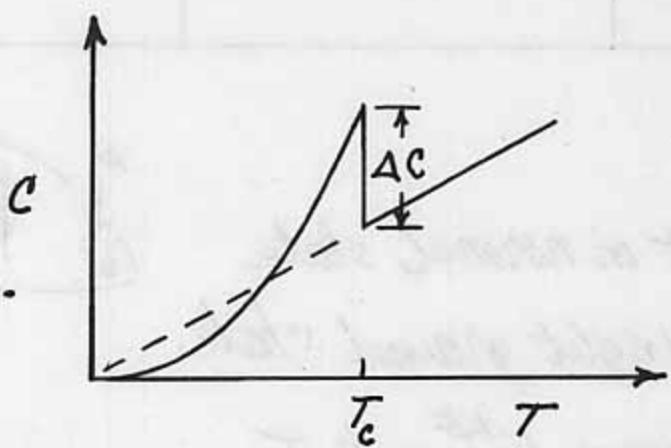


"Gapless" superconductivity

Woolf, Reif '65 — tunneling $Pb_{1-x}Gd_x$

Finnemore et al. '65 — specific heat $La_{1-x}Gd_x$

M. B. Maple '68

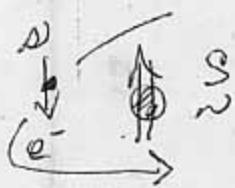


W. R. Decker & D. K. Finnemore (1968)

C. A. Luengo & M. B. Maple (1973)

$T < 0$ AFM

Romolo effect in normal state



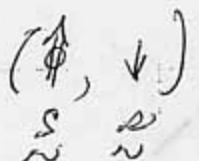
Many body singlet ground state

Binding energy $E_B^{RE} \sim k_B T_K$

SCing properties

Competition between $(k\uparrow, -k\downarrow)$

$$E_B^{SC} \sim k_B T_C$$

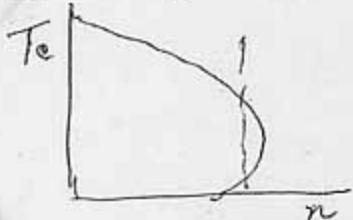


$$E_B^{RE} \sim k_B T_K$$

\Rightarrow Anomalous depression of T_c (largest for $T_K \sim T_{C_0}$)

Reentrant SC for $T_K \ll T_{C_0}$

$$T_c/T_{C_0} = V_n(\lambda/\lambda_{cr})$$



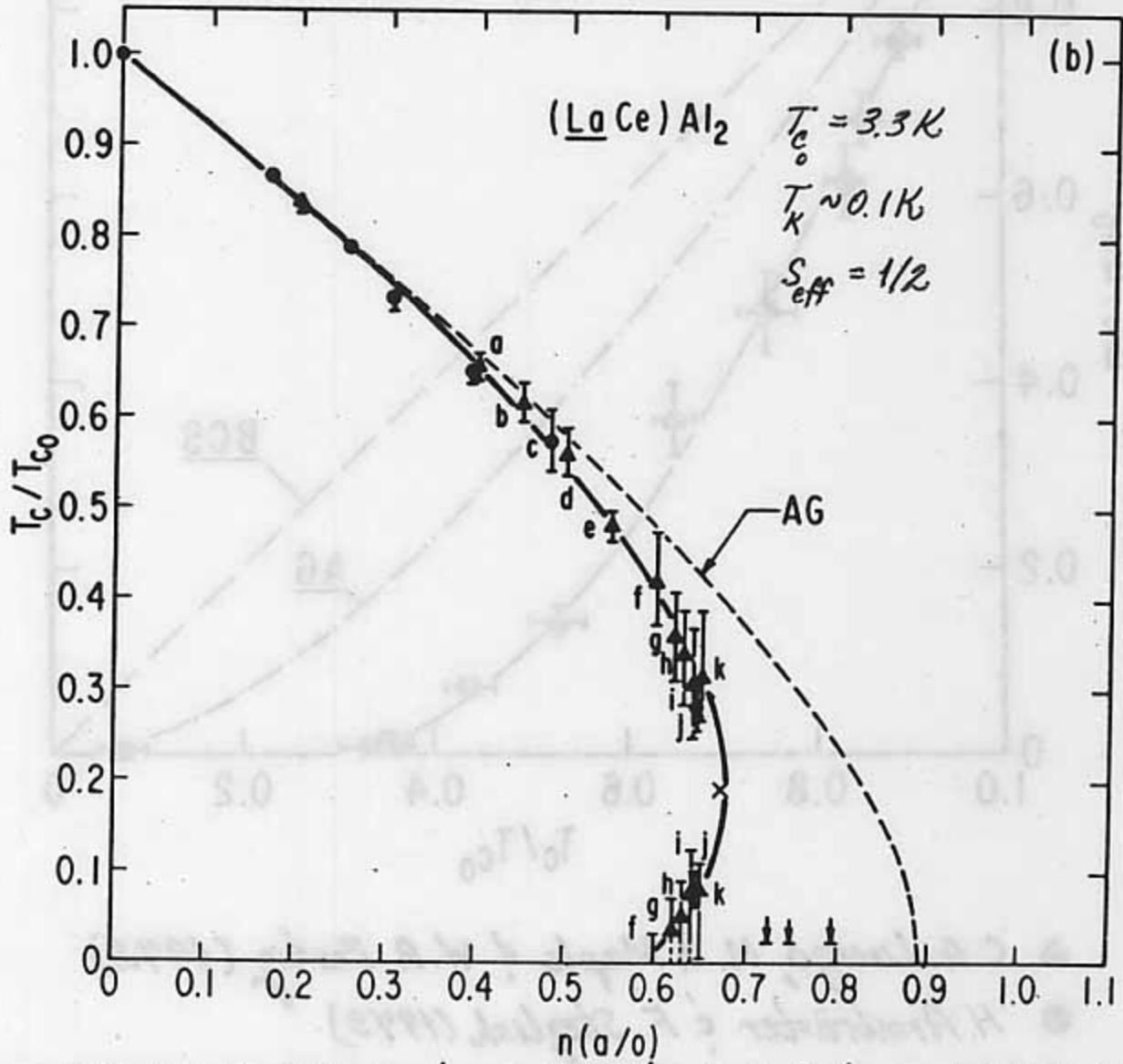
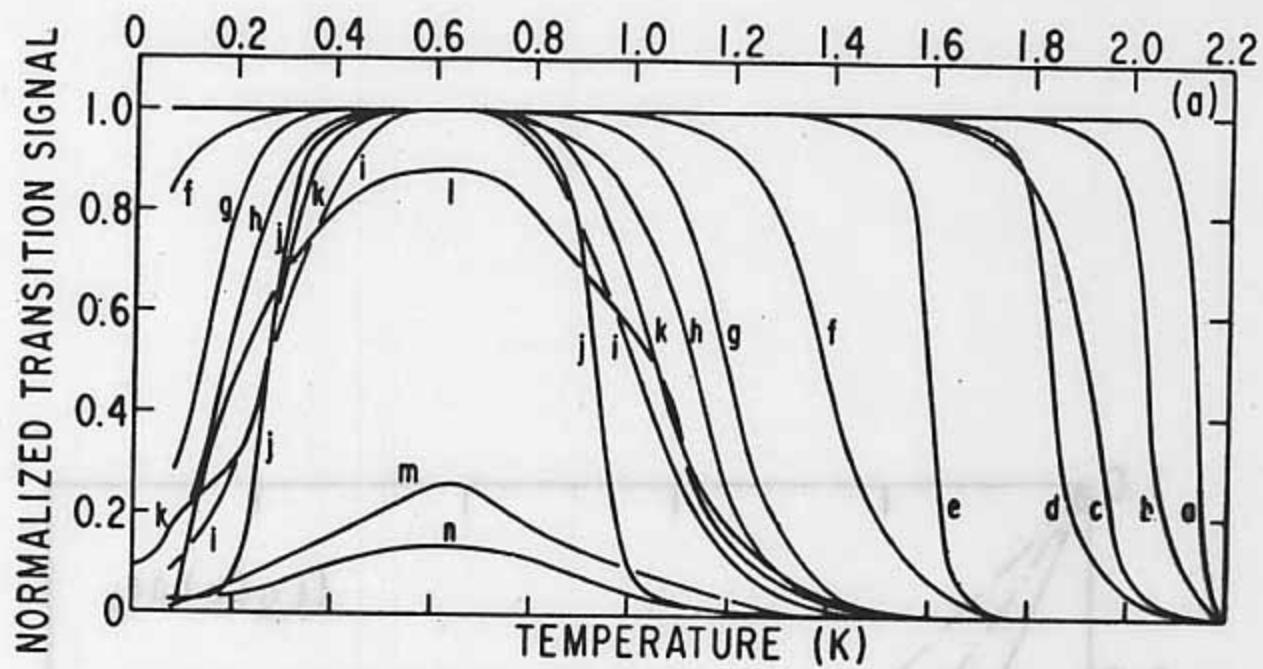
$$\lambda/\lambda_{cr} = nB \left\{ \frac{\pi^2 S(S+1)}{\ln(T/T_K) + \pi^2 S(S+1)} \right\}$$

Müller Hartmann
& Zi Hartz
(MHZ) '71

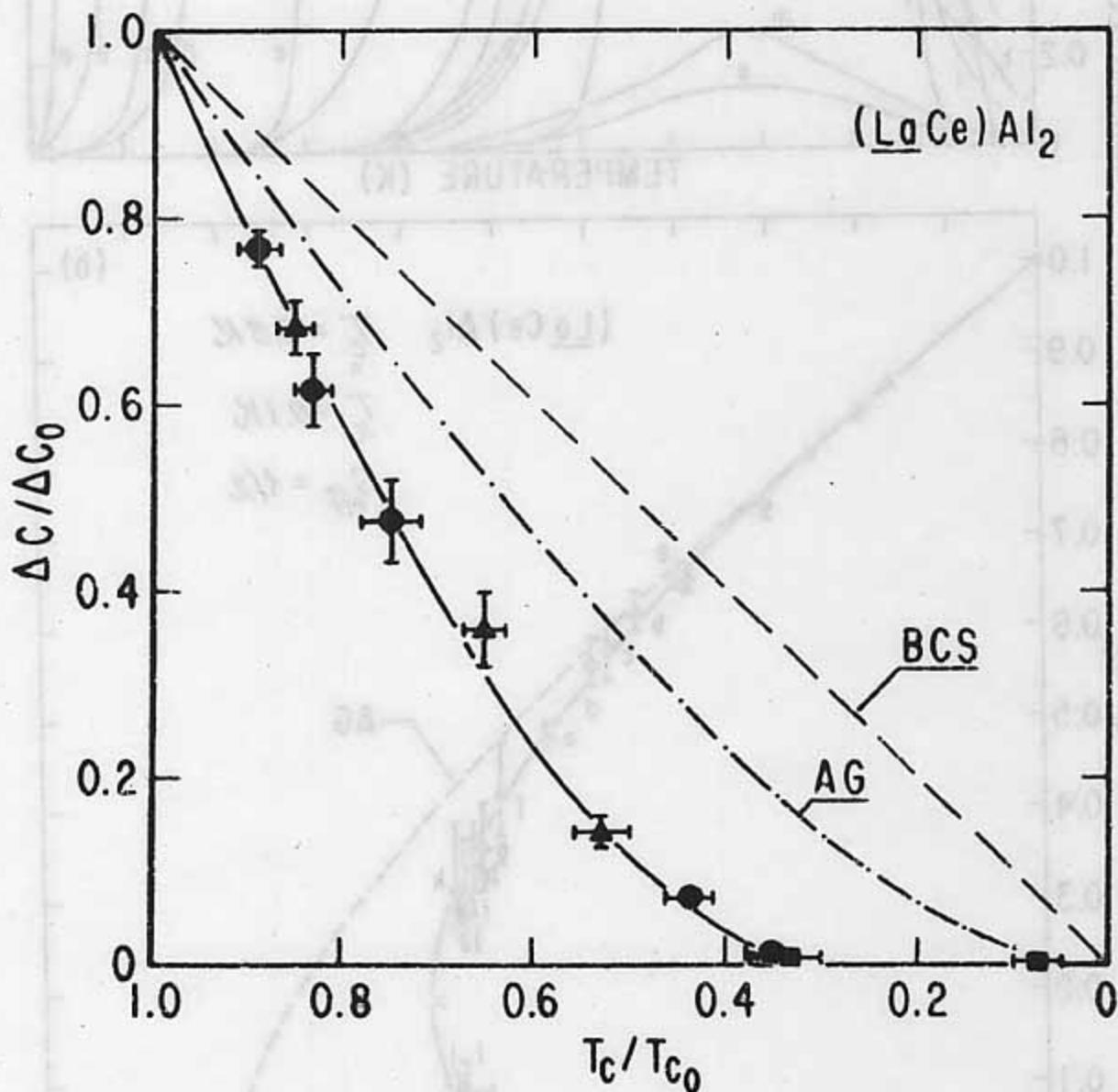
Other predictions

$\Delta C/\Delta C_0$ vs T_c/T_{C_0} deviates from both BCS law of corresponding states and SC theory

Bound state in gap



M. B. Maple, W. A. Fertig, A. C. Mota, L. E. DeLong, D. Wohlleben
 & R. Fitzgerald (1972)



- ▲ C. A. Luengo, M. B. Maple & W. A. Fertig (1972)
- H. Armbrüster & F. Steglich (1973)
- S. D. Bader, N. E. Phillips, M. B. Maple & C. A. Luengo (1973)

$$T_K \gg T_C$$

$$T_c = T_c^* \exp[-An/(1-Dn)] \quad A.B. Kaiser '7$$

Based on nonmagnetic resonant state at E_F

- ① Reduction of NCE_F) - one body dilution effect
- ② Reduction of V by Coulomb repulsion between electrons that scatter into resonant state - two body effect

BCS law of corresponding states

$$\Delta C/\Delta C_0 = T_c/T_c^*$$

Magnetic nonmagnetic transition

$$|f|/n = \frac{\langle V_{kf}^2 \rangle}{E_F}$$

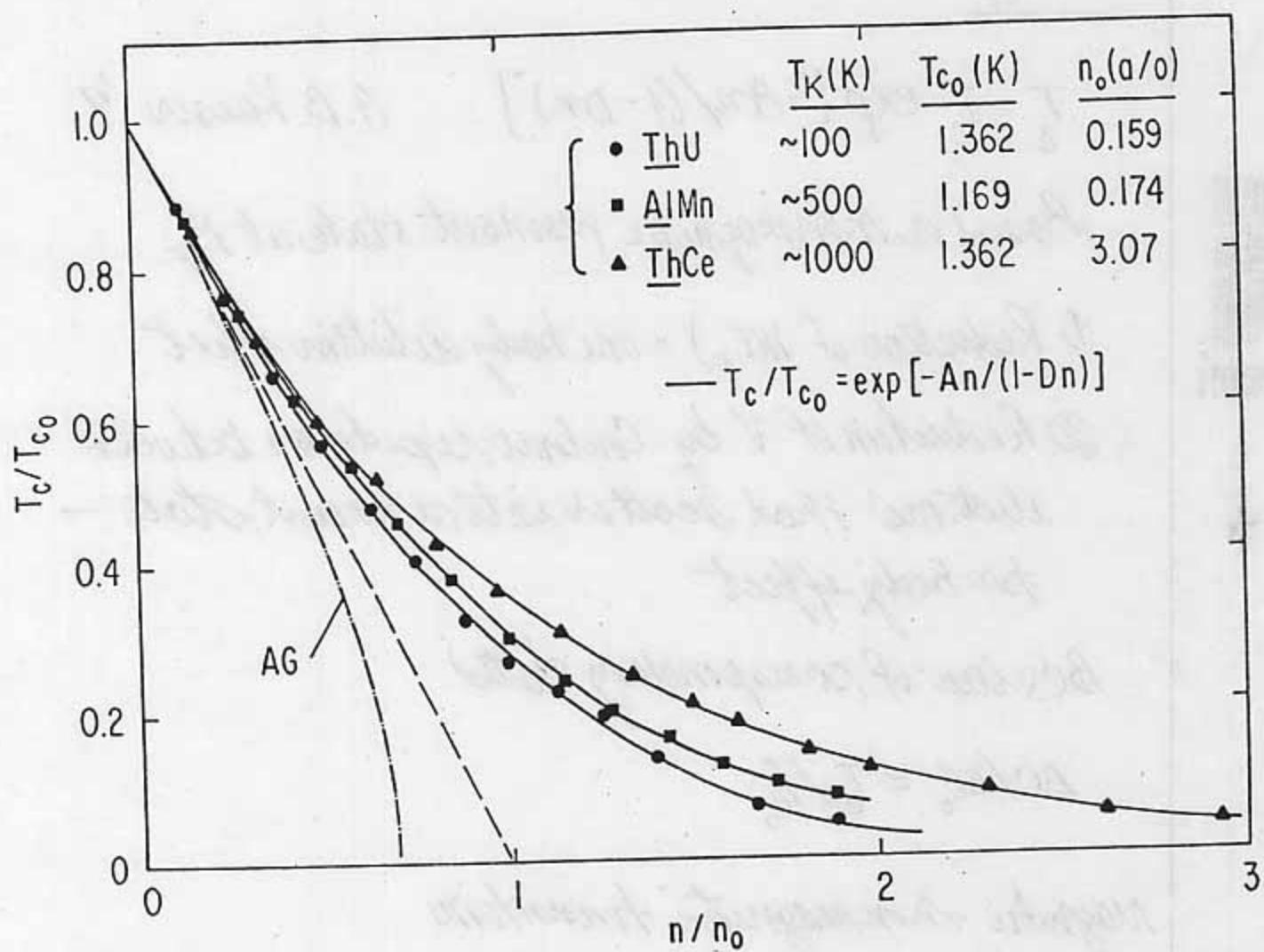
$P \Rightarrow$ increase of $\langle V_{kf}^2 \rangle$ and/or decrease of E_F

\Rightarrow increase of T_K

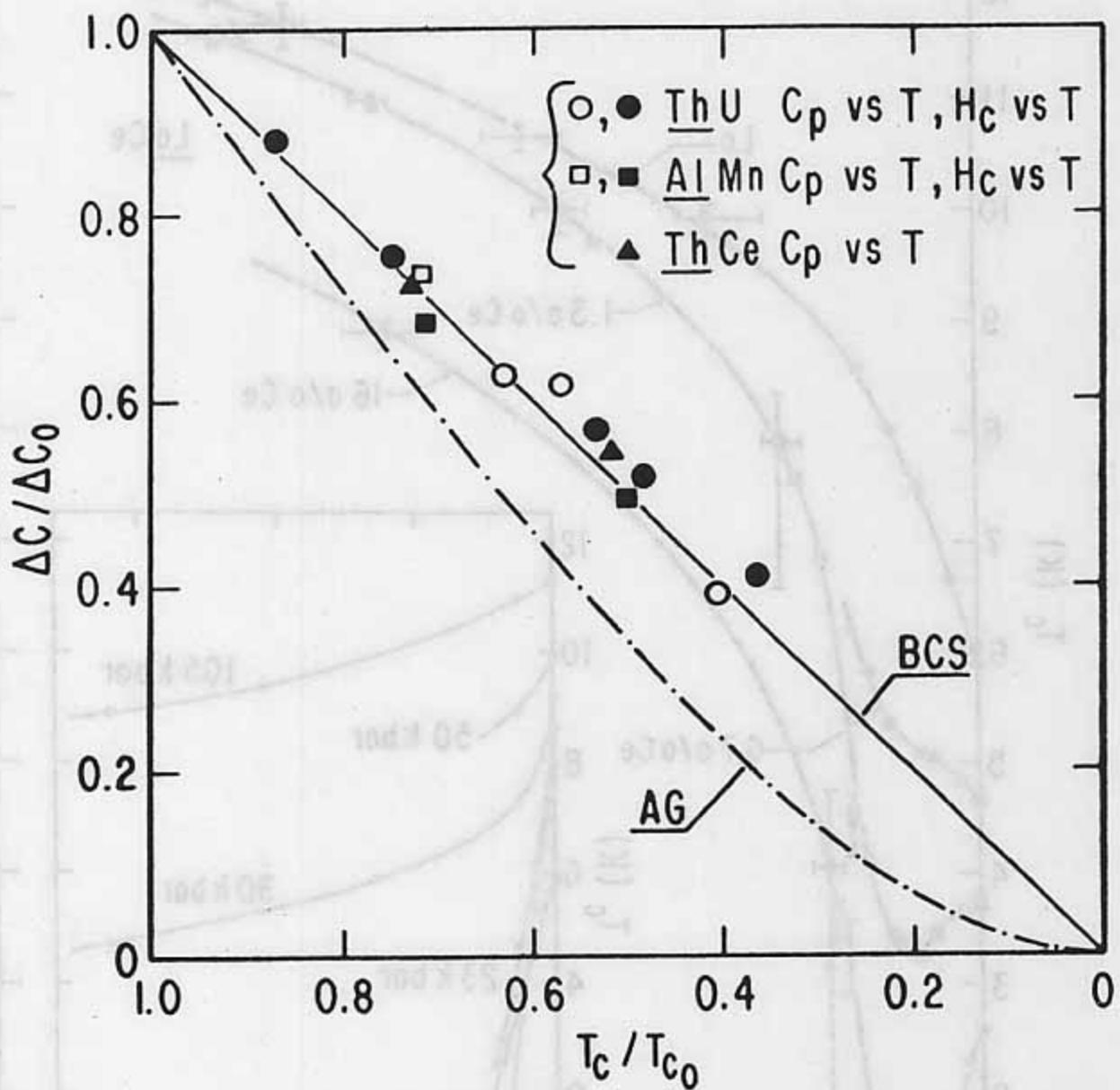
$(dT_c/dn)_{n=0}$ passes through maximum at $T_K \approx T_c$

Paramagnetic impurities in superconductors

Kondo effect: $T_K \gg T_c$

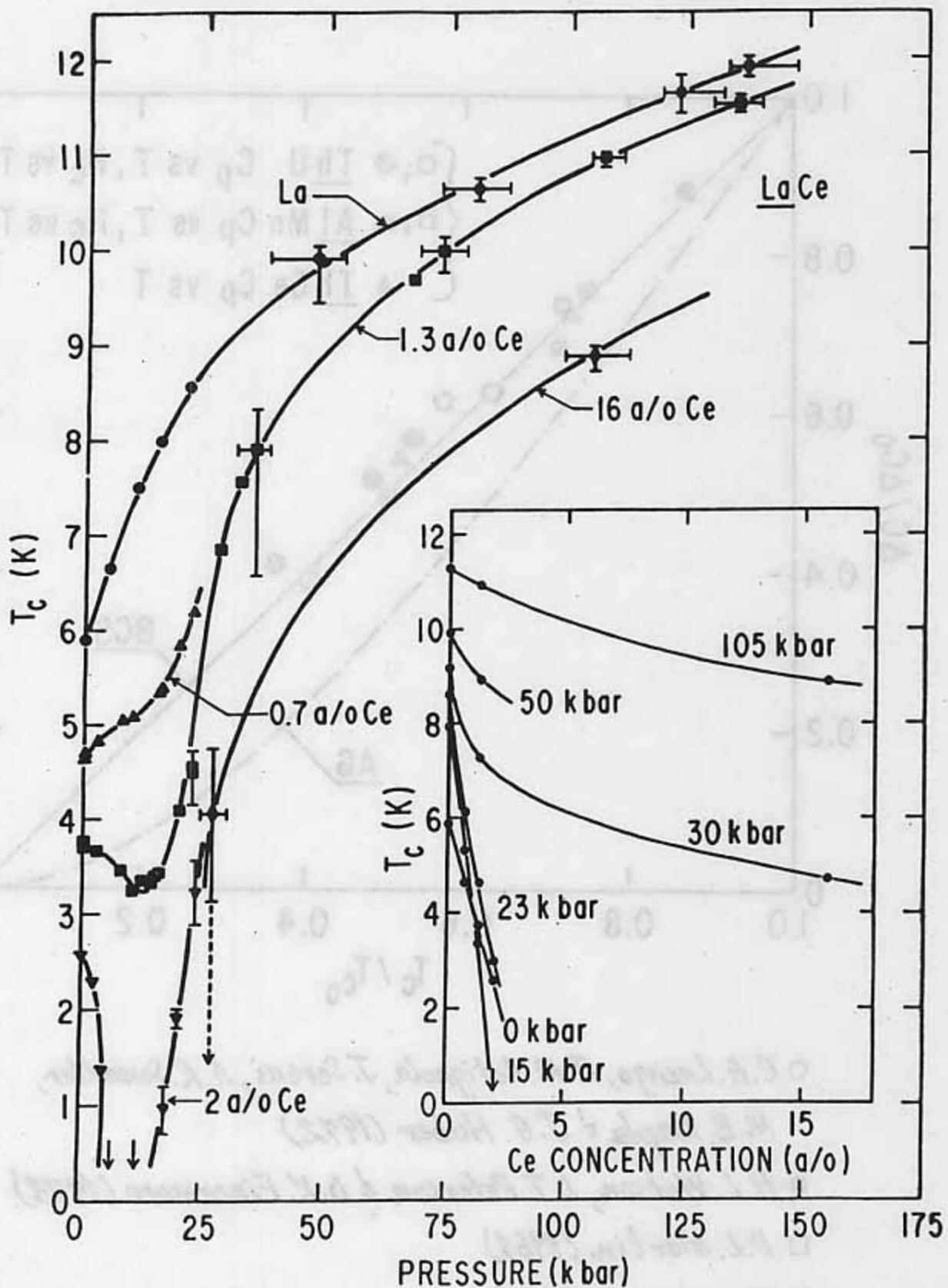


M.B. Maple, J.G. Huber, B.R. Coles, A.C. Lawson '70
 J.G. Huber & M.B. Maple '70



- C.A. Luengo, J.M. Cotignola, J. Screni, A.R. Sweeney,
M.B. Maple & J.G. Huber (1972)
- H.L. Watson, D.T. Peterson & D.K. Finnemore (1973)
- D.L. Martin (1961)
- F.W. Smith (1972)
- ▲ C.W. Dempsey (1970)

Pressure-induced demagnetization of Ce impurities in La



Analogue of δ - α transition in Ce

M.B.Maple, J.Wittig, K.S.Kim '69

► Magnetic - nonmagnetic transition of Ce impurities in
(La, Th) alloys -

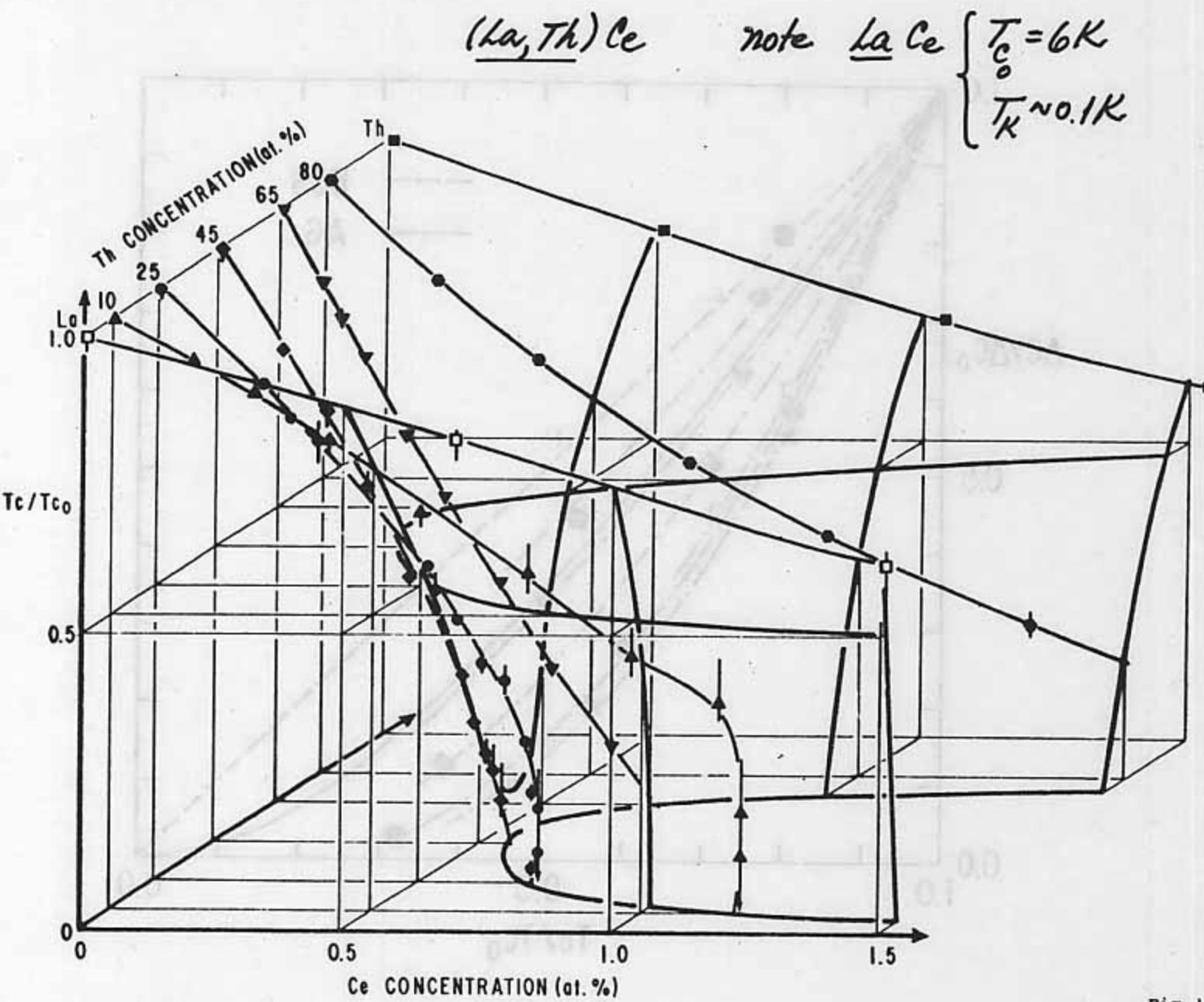
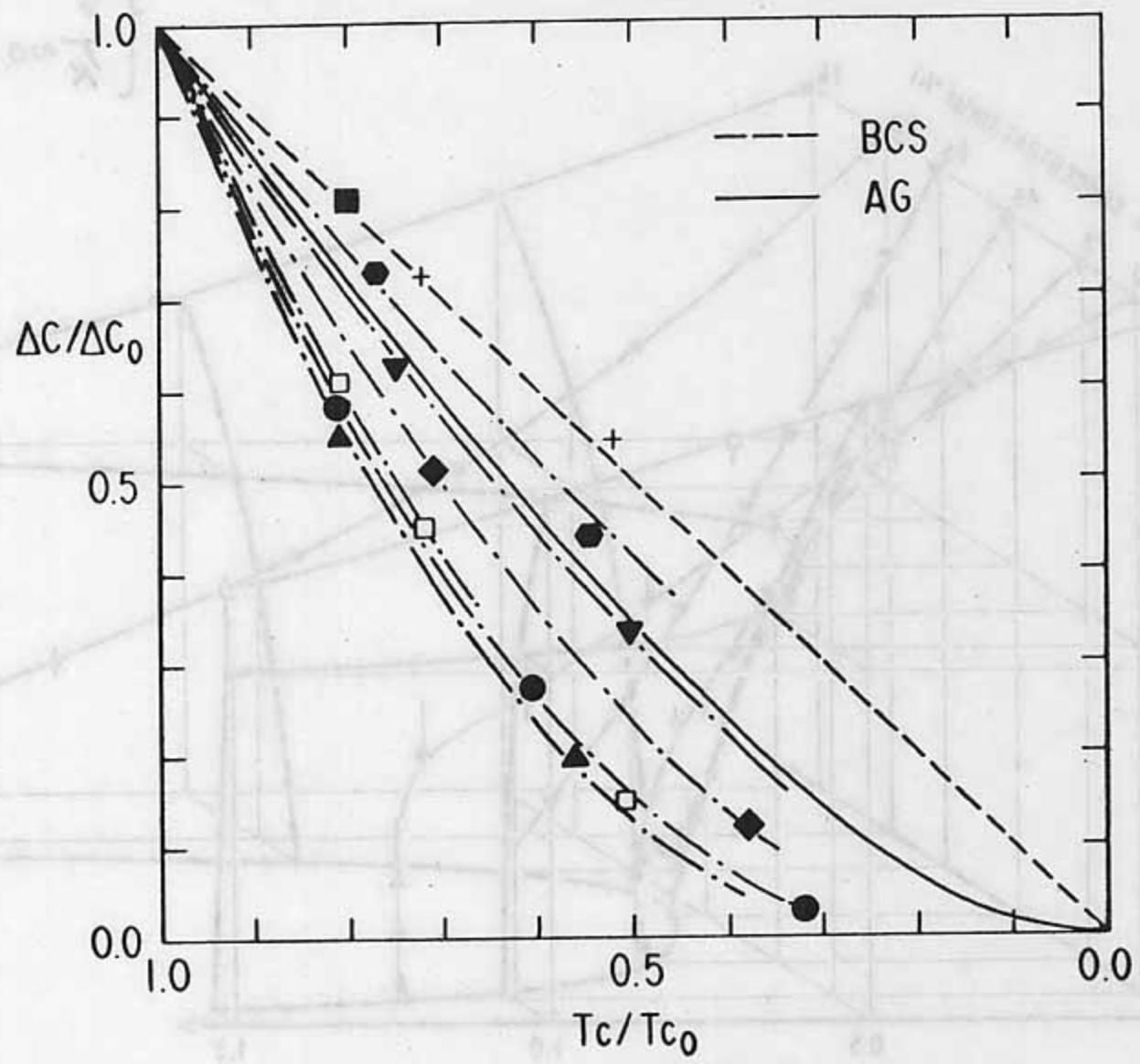


Fig. 11

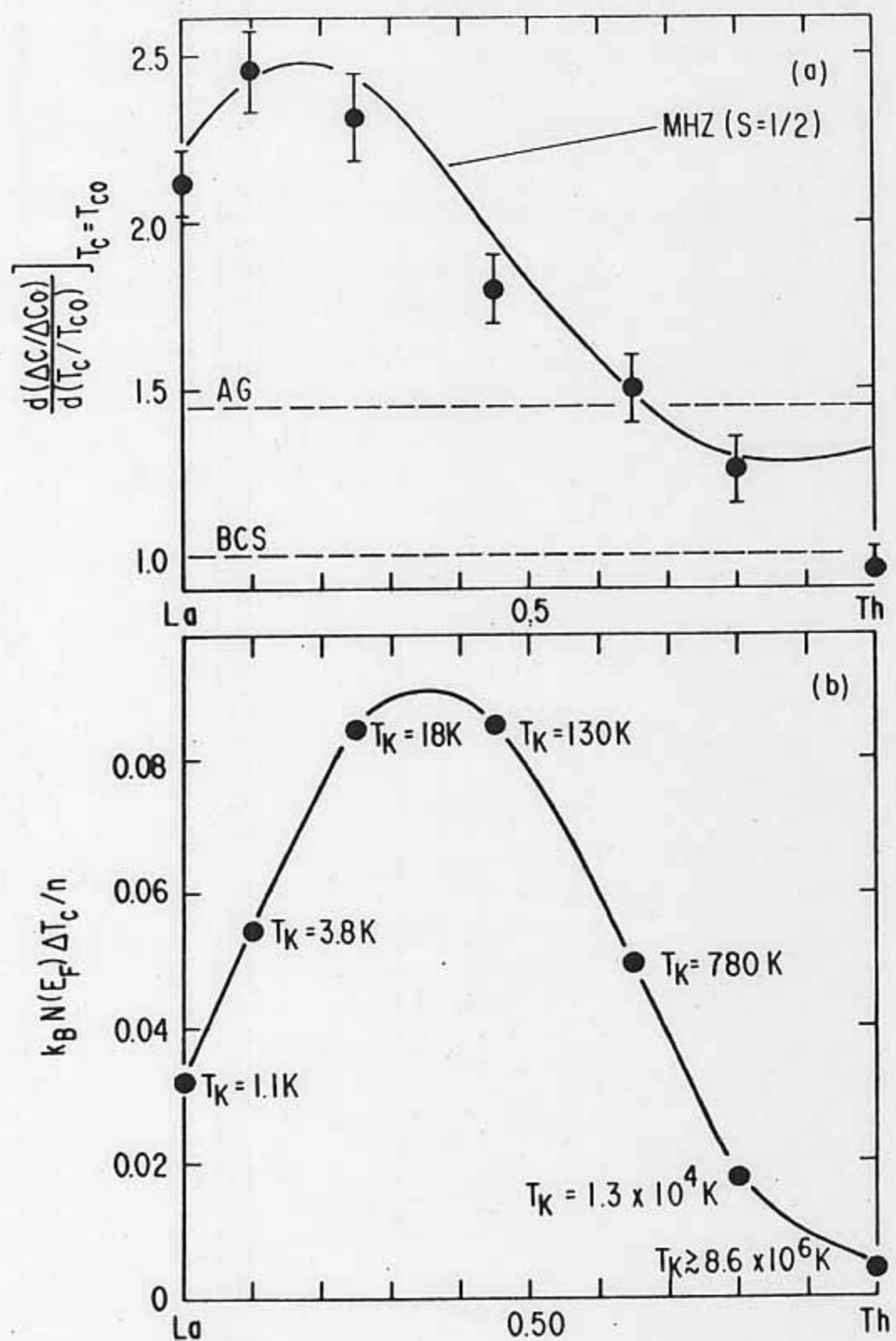
J.G. Huber, W.A. Fertig & M.B. Maple (1974)

$$T_K \sim T_F \exp(-1/N(E_F)|\beta|) ; \quad \beta \sim -\frac{\langle V_{kf}^2 \rangle}{E_F}$$

$P \Rightarrow$ increase. $|\beta| \Rightarrow$ increase. $\langle V_{kf}^2 \rangle$ and/or decrease. E_F .



C.A. Luengo, J.G. Huber, M.B. Maple & M. Roth (1974)



C.A. Luengo, J.G. Huber, M.B. Maple, & M. Roth (1974)