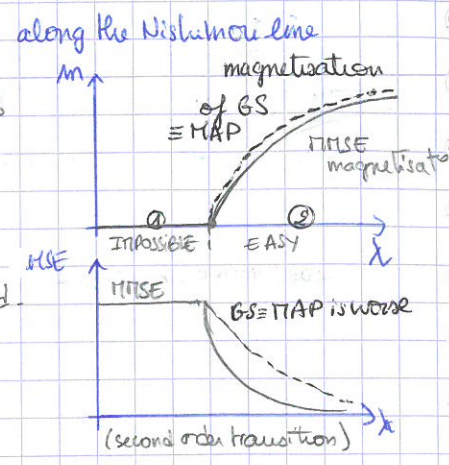
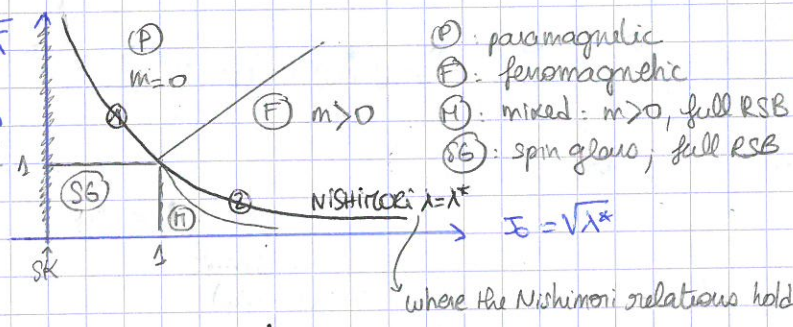
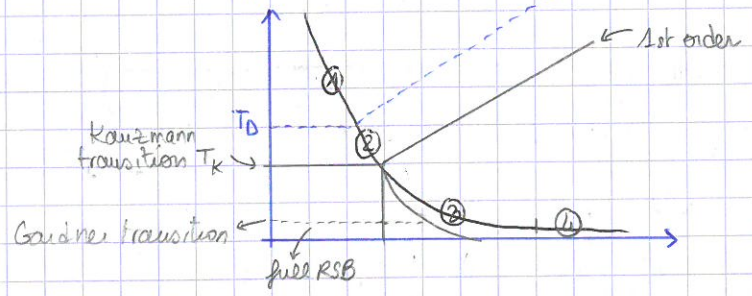


diagram of model parameters PHASE DIAGRAM

$$T = \frac{1}{\beta} = \frac{1}{\lambda}$$

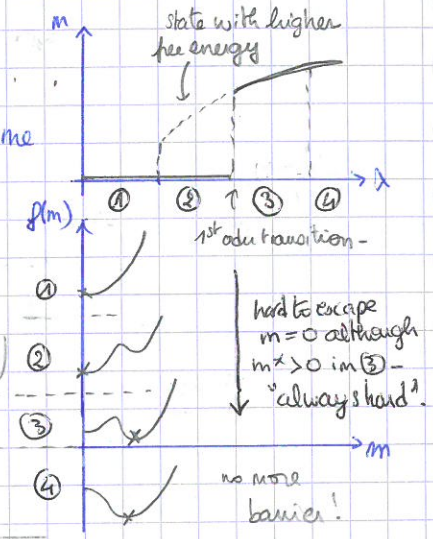


Now in the case of a tensor of order 3: \rightarrow spindled, dynamic transition



now 3 phase along Nishimori line

COMPUTATIONAL COMPLEXITY:
IMPOSSIBLE (indistinguishable)
POSSIBLE but HARD
EASY.



\rightarrow Thus two cases are representative of the different scenario: \rightarrow 1st order with impossible \rightarrow easy
 \rightarrow 2nd order with impossible \rightarrow hard \rightarrow easy

We'll see how to make the necessary computations on the Nishimori line -

REPLICA COMPUTATION ON THE EXERCISE SHEET -

$$\lim_{N \rightarrow \infty} \frac{F}{N} = \min_m \left[\mathcal{L}_{\text{den}}(\lambda_m, \lambda_m) + \frac{\lambda_m^2}{4} \right]$$

\downarrow
meq

\rightarrow replicas lengthy (crazy as well...)
 \rightarrow rigorous proof

Lecture III Free energies and algorithms

III.1 Replica free energy Recall the hamiltonian of the spike matrix model:

$$\mathcal{H} = - \sum_{i,j} \frac{\lambda}{2N} x_i^2 x_j^2 - \frac{\lambda}{N} x_i x_j x_i^* x_j^* + \sqrt{\frac{\lambda}{N}} x_i x_j w_{ij}$$

exercise shows:
$$f = \lim_{N \rightarrow \infty} \frac{\mathbb{E}_{w,x^*} (-\log Z)}{N} = \min_m \left[\mathcal{L}_{\text{den}}(\lambda_m, \lambda_m) + \frac{\lambda_m^2}{4} \right]$$
 (1) for computation later -

$f_{RS}(m)$

meq = $\arg \min_m f_{RS}(m)$

In some cases, the crazy replica method can be actually rigorously proven. The main techniques to do so is the Guerra interpolation.

In our problem we have no replica symmetry breaking \rightarrow pretty simple setting -

\rightarrow idea: interpolate between our problem and the simple model of 1 spin in a field -

simple \rightarrow
$$\mathcal{H}_{\text{sp}} = - \sum_{i=1}^N \frac{\lambda x_i^2}{2} + x_i x_i^* \lambda + x_i \xi \sqrt{\lambda}$$
 \rightarrow independent spins $y_i = \sqrt{\lambda} x_i + \xi_i$

Francesco Guerra trick: introduce a time dependent size λt for \mathcal{H}_{SR}
 $\lambda m(1-t)$ for \mathcal{H}_{DEN}

$$\mathcal{H} = \mathcal{H}_{SR}(\lambda t) + \mathcal{H}_{DEN}(\lambda m(1-t)) \quad t=0 \rightarrow \text{easy denoising problem}$$

$$t=1 \rightarrow \text{pb we would like to solve}$$

basic theorem of calculus = aka integration

↳ defining the time dependent free energy: $F(t) = F(0) + \int_0^t dt \frac{\partial F(t)}{\partial t}$

$$\Rightarrow F_{SR} = N \mathcal{L}_{DEN}(\lambda m, \lambda m) + \int_0^1 dt \mathbb{E}_{w, x^*, \xi} \left\langle \frac{\partial \mathcal{H}}{\partial t} \right\rangle dt$$

$$\sum_{i < j} \left(\frac{\lambda \langle x_i^2 x_j^2 \rangle}{N} - \frac{\lambda}{N} \langle x_i x_j x_i^* x_j^* \rangle - \sqrt{\frac{\lambda}{N}} w_{ij} \langle x_i x_j \rangle \frac{1}{2 \mathbb{E}} \right) + \sum_i \left(\frac{-\lambda \langle x_i^2 \rangle}{2} - \lambda m \langle x_i x_i^* \rangle + \frac{n_i \beta_i \sqrt{\lambda m}}{\sqrt{1-t}} \right)$$

gaussian noise → Stein lemma $\mathbb{E}_w(w f(w)) = \mathbb{E}_w(f'(w))$

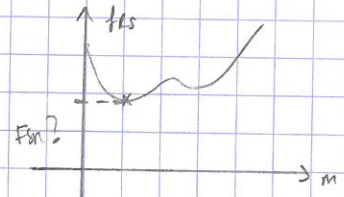
$$\mathbb{E} \left\langle \frac{\partial \mathcal{H}}{\partial t} \right\rangle = \mathbb{E} \left[-\frac{\lambda}{N} \sum_{i < j} \langle x_i x_j x_i^* x_j^* \rangle + \frac{\lambda}{2N} \sum_{i < j} \langle x_i x_j \rangle^2 + \sum_i \langle x_i x_i^* \rangle_t m \lambda - \sum_i \langle x_i^2 \rangle_t \frac{m \lambda}{2} \right]$$

↳ the explicit dependence on t has vanished, but recall that $\langle \cdot \rangle_t$

$$\begin{aligned} \frac{F_{SR}}{N} &= \mathcal{L}_{DEN}(\lambda m, \lambda m) + \int dt \frac{\lambda}{4} \mathbb{E} \left\langle \left(\frac{x x^*}{N} - m \right)^2 \right\rangle_t - \frac{\lambda}{2} \mathbb{E} \left[\left\langle \left(\frac{x x^{*T}}{N} - m \right)^2 \right\rangle_t \right] - \frac{\lambda m^2}{4} + \frac{\lambda m^2}{2} \\ &\stackrel{(1)}{=} f_{RS}(m) + \int dt \frac{\lambda}{4} \mathbb{E} \left[\left\langle \left(\frac{x x^*}{N} - m \right)^2 \right\rangle_t \right] - \frac{\lambda}{2} \mathbb{E} \left[\left\langle \left(\frac{x x^{*T}}{N} - m \right)^2 \right\rangle_t \right] \\ &= f_{RS}(m) - \int dt \frac{\lambda}{4} \mathbb{E} \left[\left\langle \left(\frac{x x^*}{N} - m \right)^2 \right\rangle \right] \end{aligned}$$

Nishimori

⇒ which gives us a bound: $F_{SR}/N \leq f_{RS}(m) \quad \forall m \geq 0$



To find the converse bound -

we consider $\mathcal{H}_{DEN}(q, m, \lambda, t) = -\sum_i -\frac{\lambda}{2} q(1-t) x_i^2 + x_i x_i^* \lambda m(1-t) + x_i \xi_i \sqrt{\lambda q(1-t)}$

$$\mathcal{H}_{DEN}^t + \mathcal{H}_{SR}^t \leftarrow \frac{F'(t=1)}{N} = \mathcal{L}_{DEN}(\lambda q, \lambda m) + \int dt \frac{\lambda}{4} \mathbb{E} \left[\left\langle \left(\frac{x^{(1)} x^{(2)}}{N} - q \right)^2 \right\rangle \right] - \frac{\lambda}{2} \mathbb{E} \left[\left\langle \left(\frac{x x^{*T}}{N} - m \right)^2 \right\rangle \right] - \frac{\lambda q^2}{4} + \frac{\lambda m^2}{2}$$

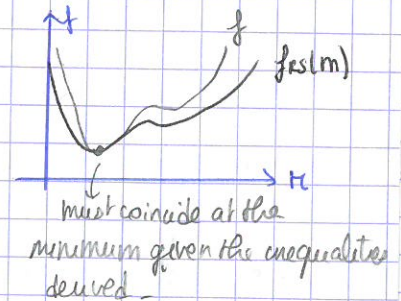
↳ consider the restricted free at fixed values of overlap $\mathcal{H} = \frac{x x^*}{N}$, $F(\mathcal{H}) = -\mathbb{E} \log Z(\mathcal{H})$

with $Z(\mathcal{H}) = \sum_{x/x^*} e^{-\mathcal{H}}$ $\ll Z_{\text{all config}}$

$$\frac{F}{N}[\mathcal{H}=m] = \mathcal{L}_{DEN}(\lambda q, \lambda m) + \int dt \frac{\lambda}{4} \mathbb{E} \left[\left\langle \left(\frac{x x^*}{N} - q \right)^2 \right\rangle \right] - \frac{\lambda q^2}{4} + \frac{\lambda m^2}{2}$$

$$\Rightarrow f(\mathcal{H}=m) \geq f_{RS}(m, q) \quad \forall q \Rightarrow \underset{m=q}{f(\mathcal{H}=m)} \geq f_{RS}(m)$$

now, what's the true free energy?
 $f(m, q) / m, q = \text{argmin}_m f(m)$!



must coincide at the minimum given the inequalities derived -

$$\Rightarrow \underline{f = \min_m f_{RS}(m)}$$

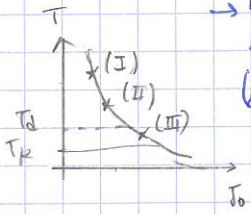
$f(\mathcal{H}) \equiv$ fixed overlap with truth free energy
 $\equiv f(q) \equiv$ fixed overlap with equilibrium configuration
 \equiv FRANZ-PARISI POTENTIAL

Rk: We can use this same method for the p spin: $f_{RS}(m) = \mathcal{L}_{DEN}(\lambda m, \lambda m) + \frac{\lambda m^p}{2p} (p-1)$
 as long as p is even.

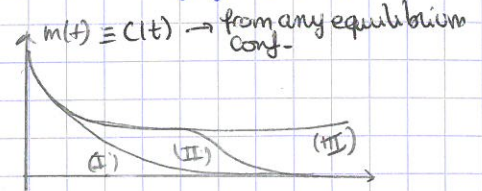
looking back at the phase diagram \equiv spin - lower factorization

For $T > T_K$, on the Nishimori line, paramagnetic phase \rightarrow indistinguishable from random problem, same as if $J_0 = 0 \equiv$ unplanted

\rightarrow but if you planted the problem, you know one equilibrium configuration - which is of great help in many circumstances.



Monte Carlo simulation from large magnetizations on the Nishimori line

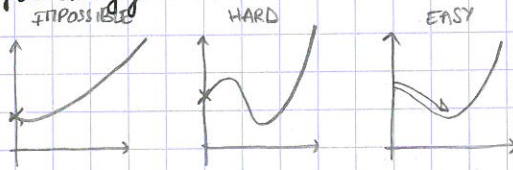


\equiv simulation from any equilibrium conf, which would be very costly to obtain otherwise at low temperatures, especially $T < T_d \rightarrow$ GREAT TRICK!

14/07/07

III.2 Belief propagation, TAP free energy

We have this situation



to find the solution we can run monte carlo but alternative, faster algorithms \rightarrow Belief Propagation

Belief propagation was rediscovered many times, but originally Belief-Perls in the 30s.

The idea is the following, (see exercise sheet):

hamiltonian: $H = - \sum_{i,j} J_{ij} S_i S_j$

1. write recursive BP "tree with high connectivity":
 $m_{i \rightarrow j} = f(\{m_{a \rightarrow i}\}) = \tanh\left(\sum_{a \in \mathcal{N}(i)} \tanh(\beta J_{ia}) m_{a \rightarrow i}\right)$

2. Expand in J (small J)

large connectivity assumption, marginals \leftarrow

3. Close the equations on the marginals \equiv magnetizations
 $m_i = \tanh\left(\beta \sum_j J_{ij} m_j - \beta^2 \sum_j J_{ij}^2 (1 - m_j^2) m_i\right)$
 \hookrightarrow time indexing is crucial for convergence!!

This algorithm can be analyzed: it is basically doing a gradient descent along the replica symmetric potential \rightarrow TAP magnetization get to (local) minima of the RS free energy.

\hookrightarrow still stuck for the hard phase
 \hookrightarrow finds optimal solution in every phase

Rk: Belief in stat mech community that message passing algorithms are somewhat optimal: there are not polynomial algorithm that can succeed in the hard phase.

Lecture IV: GENERALISED LINEAR MODELS

The situation is still signal recovery from noisy observations:

signal $x^* \in \mathbb{R}^N$ distributed $P_x(x^*) \rightarrow$ observation in \mathbb{R}^M $Y = \Phi X^*$
 known measurement matrix $\mathbb{R}^{M \times N}$ KNOWN!

- 1) $\alpha = M/N$
- 2) $\rho_0 =$ non zeros in x^*/N
- 3) $\Phi_{ij} \sim W(0, \frac{1}{N})$ iid.

component-wise: $y_{\mu} = f[\Phi_{\mu} \cdot x^*]$

different situations for different f : GENERALISED LINEAR MODELS!

- 1) $f(x) = x \rightarrow$ linear model \rightarrow stat phys community Tanaka 2000 something
- 2) stochastic version: $f(x) = x + \xi \rightarrow W(0, \Delta)$
- 3) non-linearities: $f(x) = \text{sign}(x + \xi)$ - perceptron classification problem