

Neutron and x-ray scattering studies of superconductors

lecture 2: unconventional superconductors

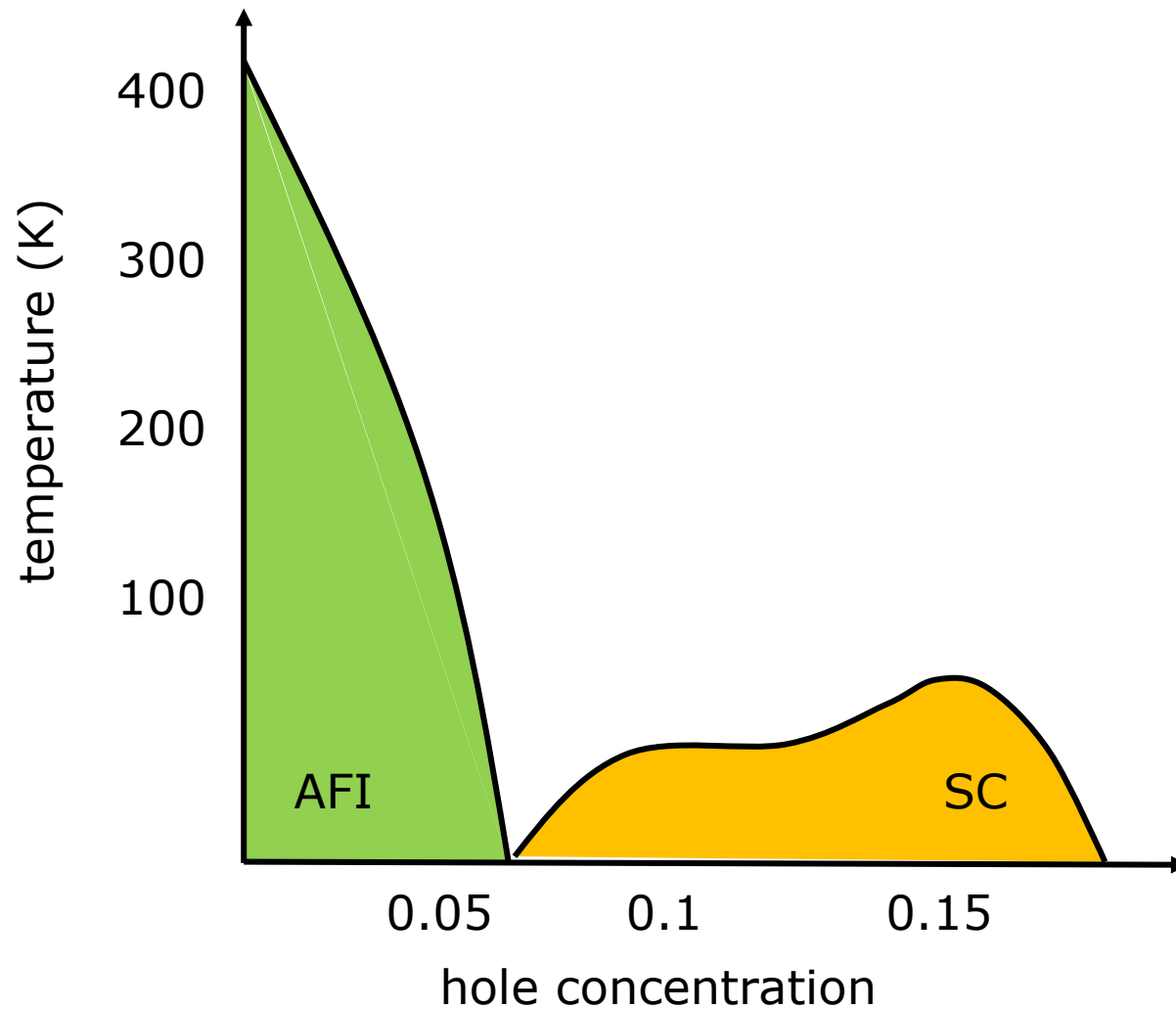
- magnetic neutron scattering continued
- resonant inelastic x-ray scattering from magnons and paramagnons
- resonant elastic x-ray scattering from charge density waves

lecture 3: cuprate and nickelate superlattices

- orbital occupation
- magnetic order
- charge density waves

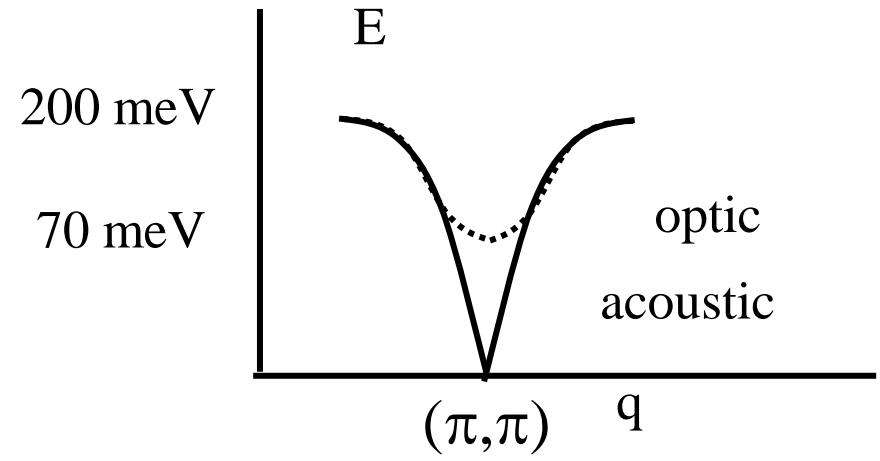
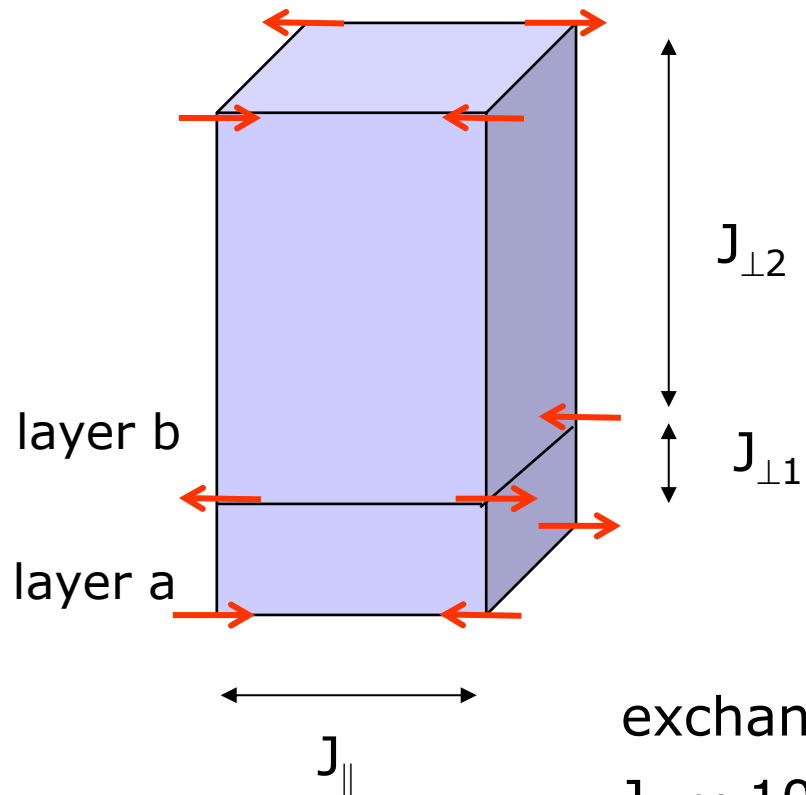


Phase diagram of $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$



YBa₂Cu₃O₆ magnons

$$H = \sum_{ij} (J_{\parallel} \mathbf{S}_i^{(a,b)} \cdot \mathbf{S}_j^{(a,b)}) + \sum_i (J_{\perp 1} \mathbf{S}_i^{(a)} \cdot \mathbf{S}_i^{(b)} + J_{\perp 2} \mathbf{S}_i^{(b)} \cdot \mathbf{S}_i^{(a)})$$



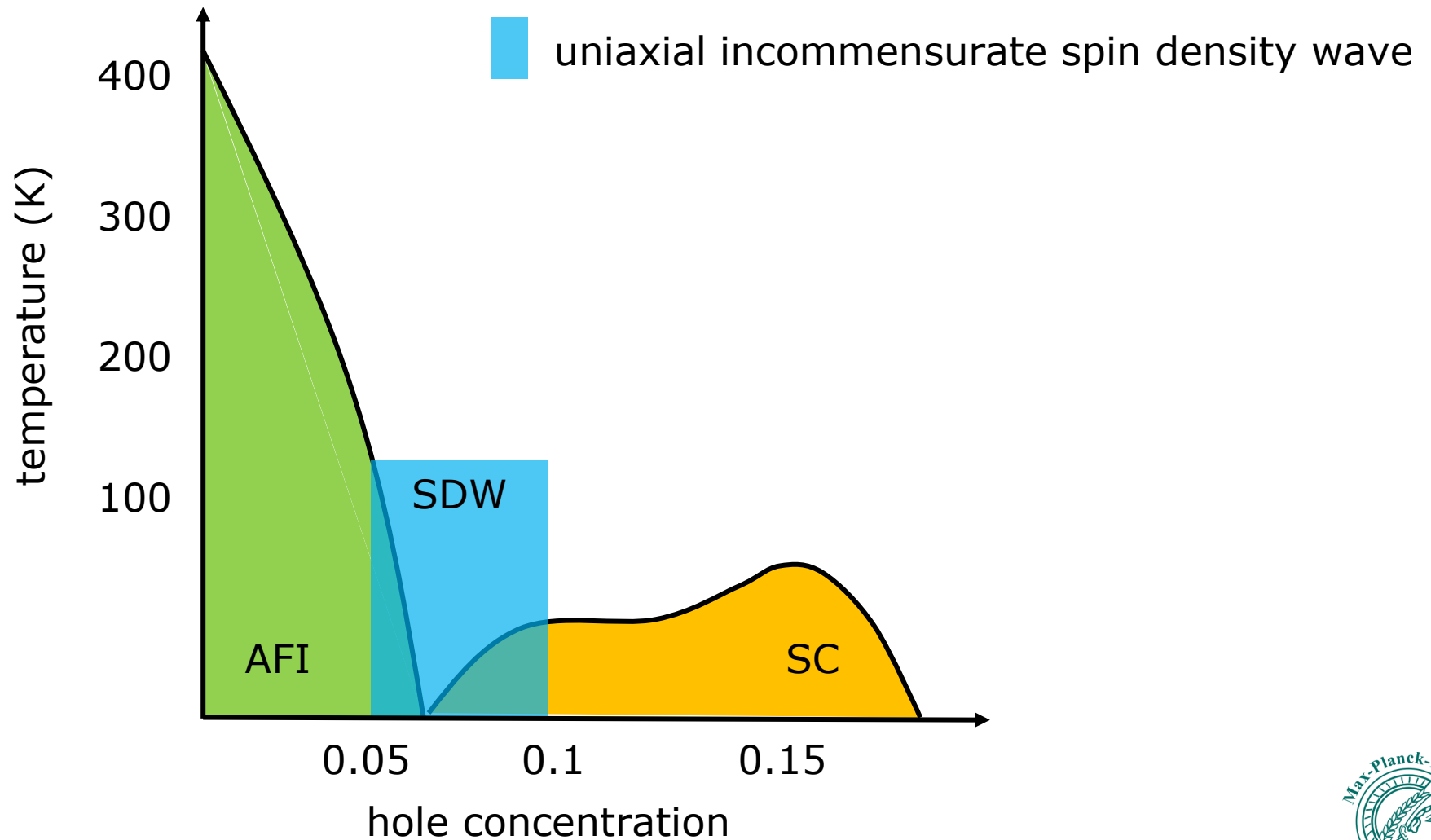
exchange parameters from magnon dispersions

- $J_{\parallel} \sim 100$ meV
- $J_{\perp 1} \sim 10$ meV
- $J_{\perp 2} \sim 0.01$ meV

Tranquada et al., PRB 1989
Reznik et al., PRB 1996



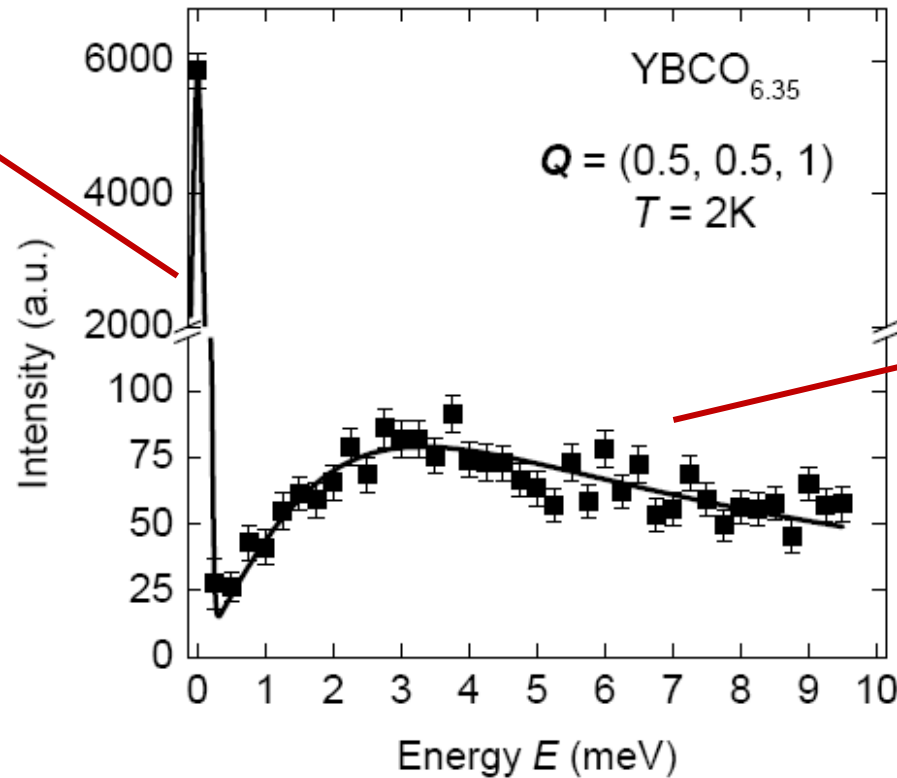
Competing order in $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$



Lightly doped YBCO

YBCO_{6.35} T_c = 10 K p ~ 0.06

quasielastic peak



spin excitations

*Haug et al.
NJP 2010*

analogous to Bragg peak, spin waves of antiferromagnetic Mott insulator
but centered on incommensurate wave vector

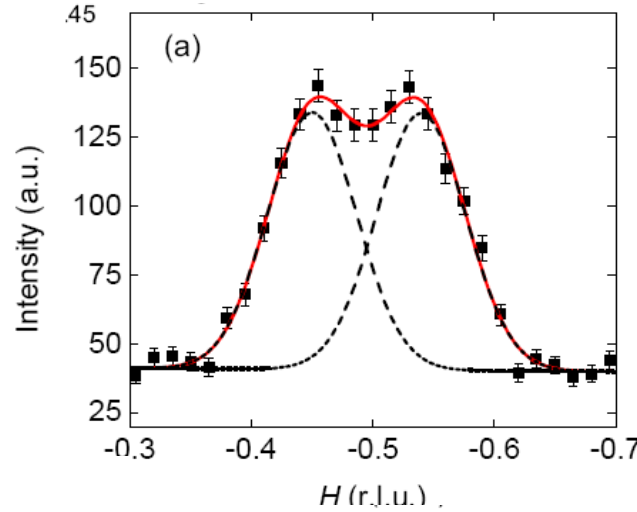
Low-energy spin fluctuations

neutron scattering

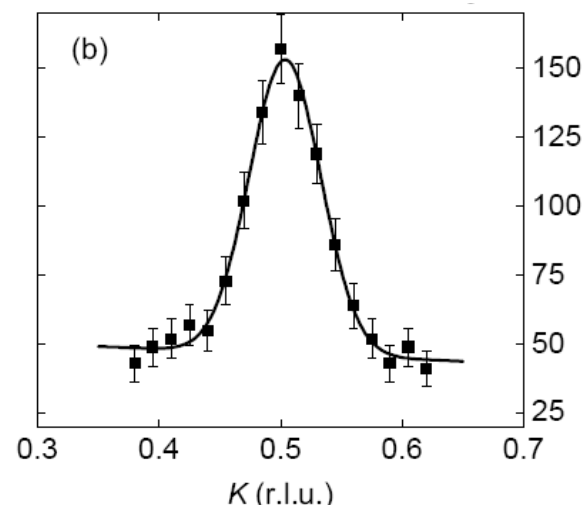
YBCO_{6.45}

T_c = 35 K

along a*

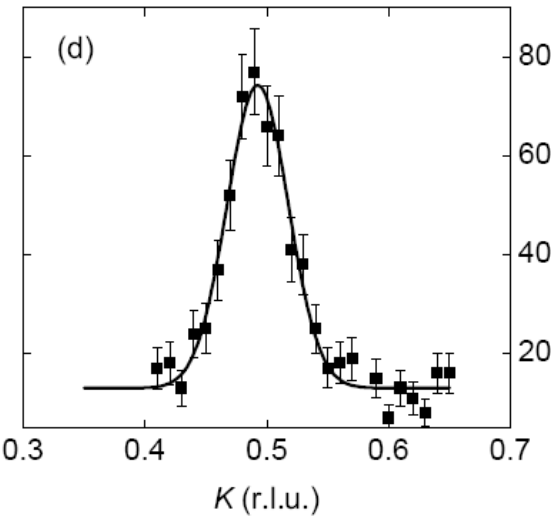
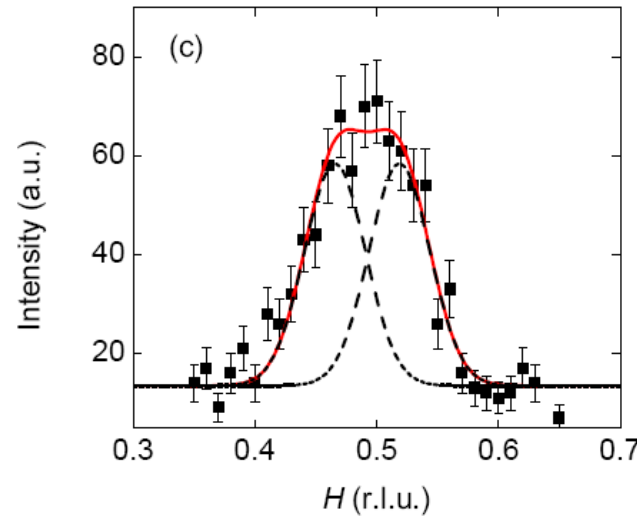


along b*



YBCO_{6.35}

T_c = 10 K



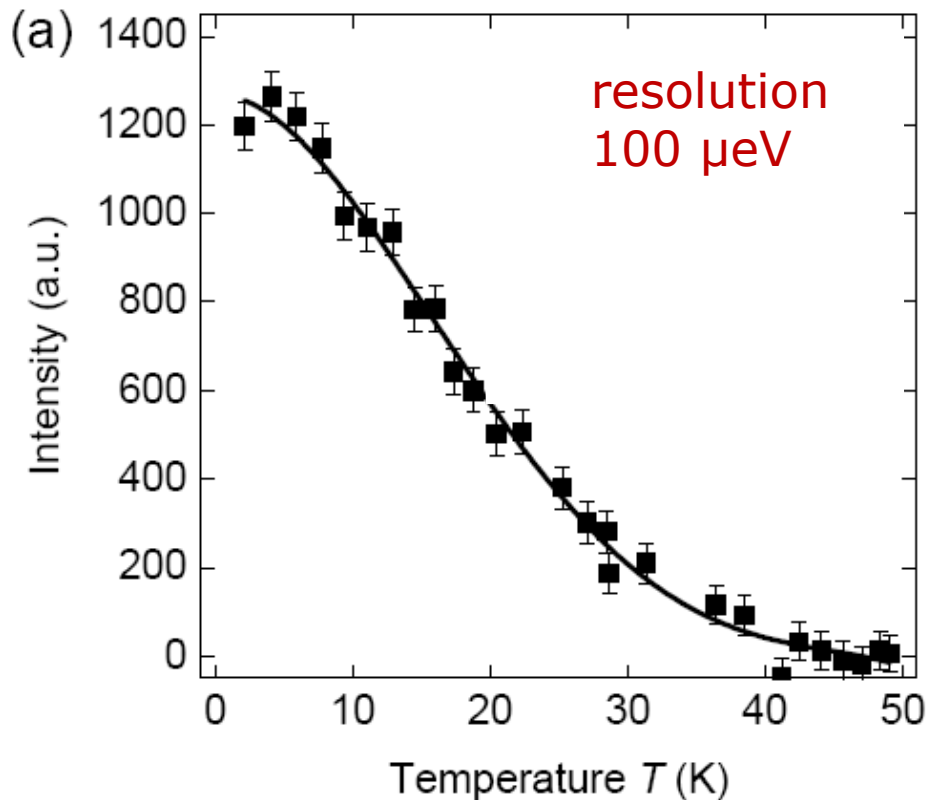
Haug et al.
NJP 2010

uniaxial incommensurate magnetic order

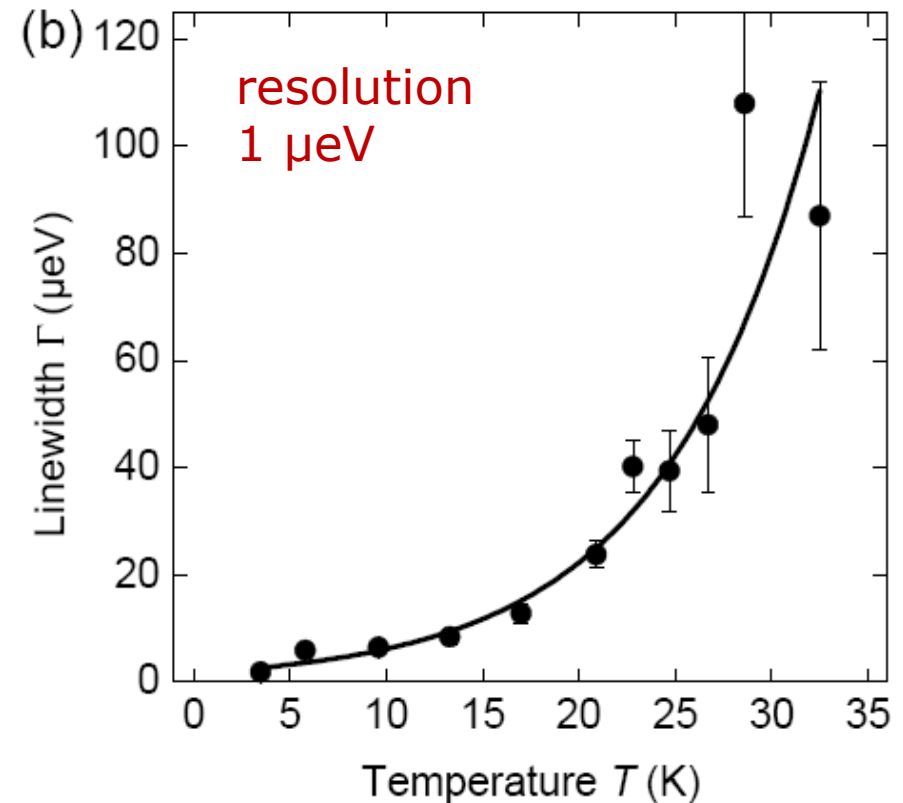
in lightly doped, weakly metallic YBCO → "spin density wave"

Temperature dependence

**conventional triple-axis
quasielastic neutron scattering**



neutron spin-echo spectroscopy
→ energy width Γ of quasielastic peak



energy width Γ of quasielastic peak narrows continuously as $T \rightarrow 0$

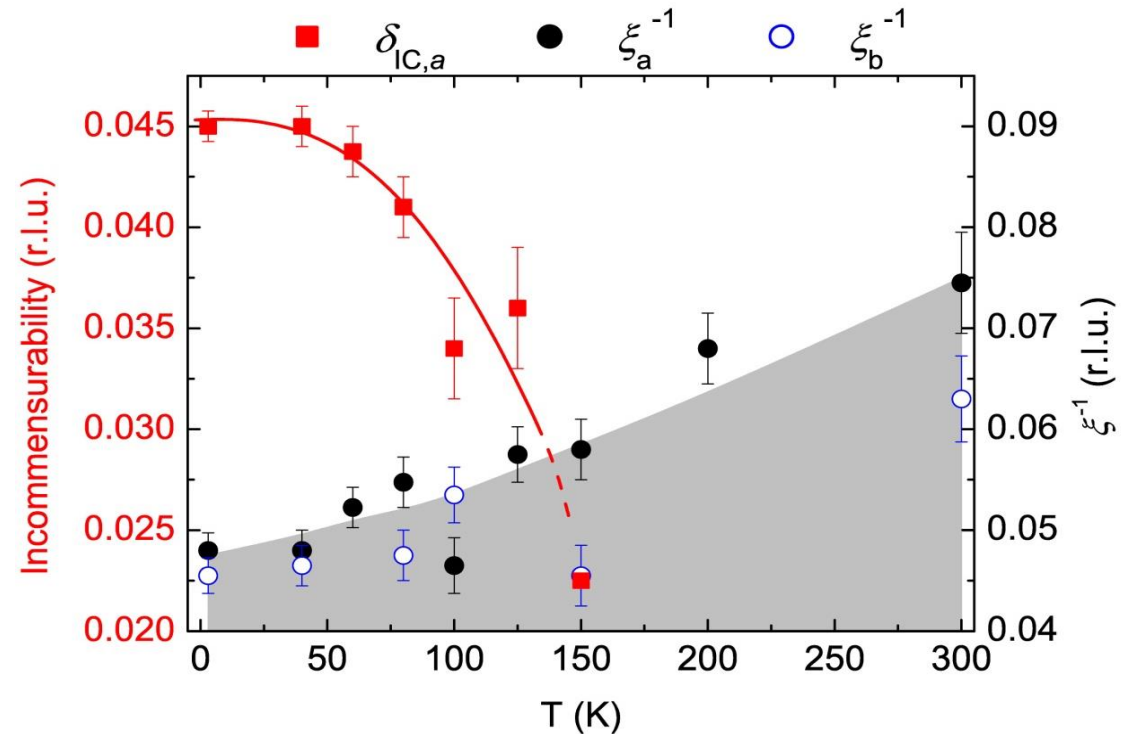
generic behavior in 2D Heisenberg systems

disorder also contributes → “cluster spin glass”

Electronic nematic state in YBCO

inelastic neutron scattering

*Hinkov et al.
Science 2008*



spontaneous onset of incommensurability upon cooling below ~ 150 K

→ orientational symmetry broken

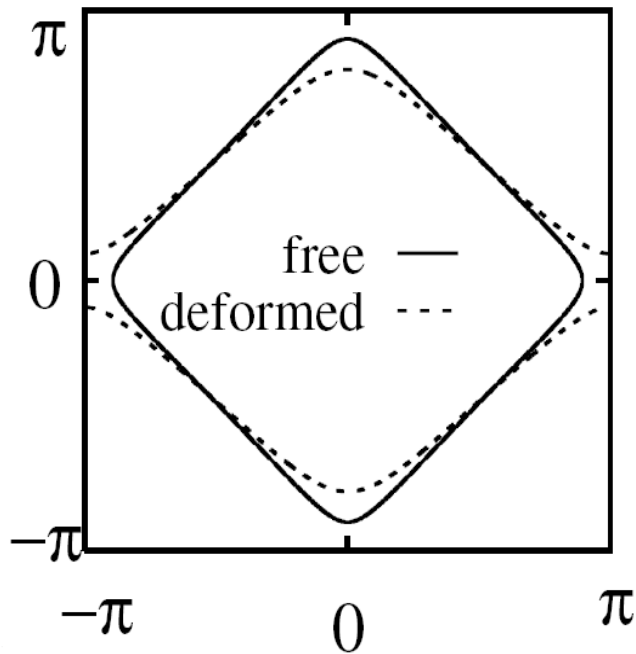
but no static magnetic order → translational symmetry unbroken

electronic analog of nematic liquid crystal

*Fradkin, Kivelson et al.
Yamase, Metzner et al.
Ando et al.
McKenzie et al. ...*

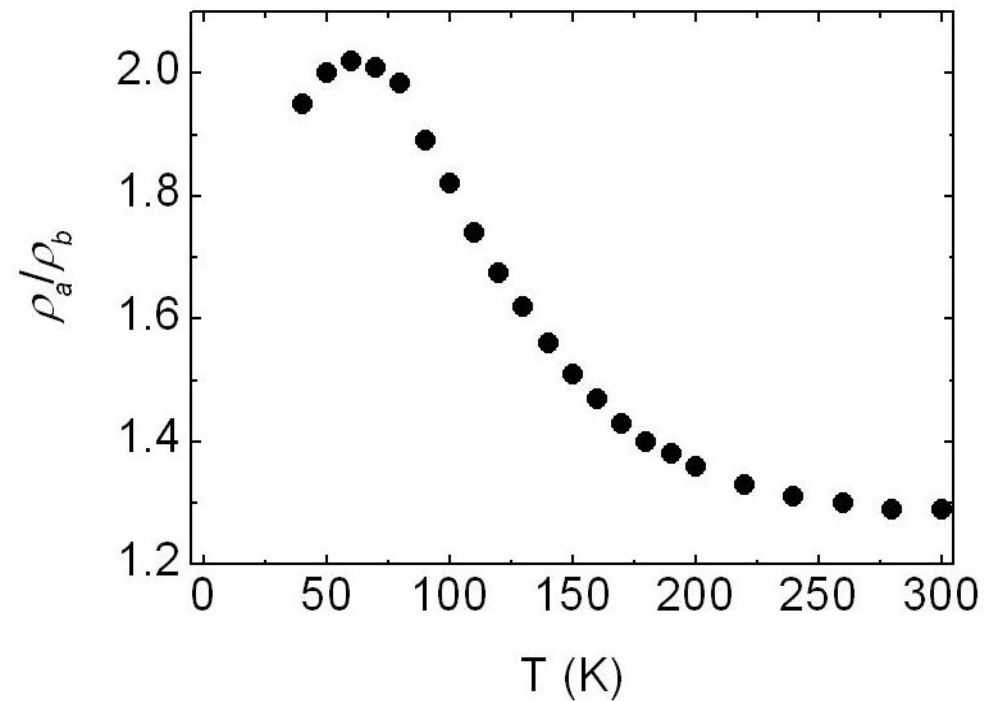


Electronic nematic state in YBCO



possible route towards nematicity:
Pomeranchuk instability

Halboth & Metzner, PRL 2000
Yamase & Kohno, JPSJ 2000

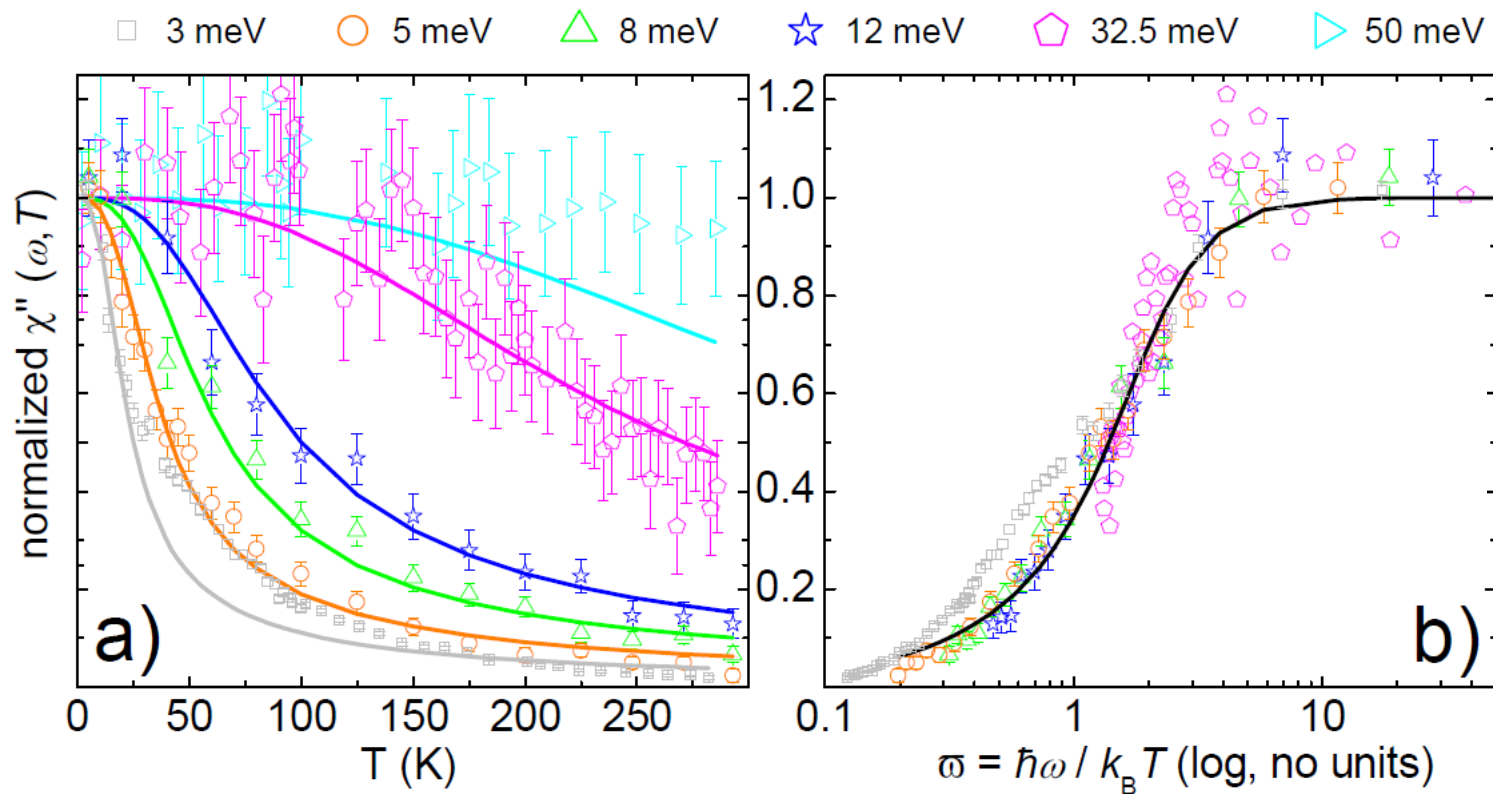


resistivity, Nernst effect anisotropies
turn up at similar temperature

Ando et al., PRL 2002
Daou et al., Nature 2010

Quantum critical point

$\text{YBa}_2\text{Cu}_3\text{O}_{6.45}$ $T = 2\text{K}$ (in IC-SDW state)

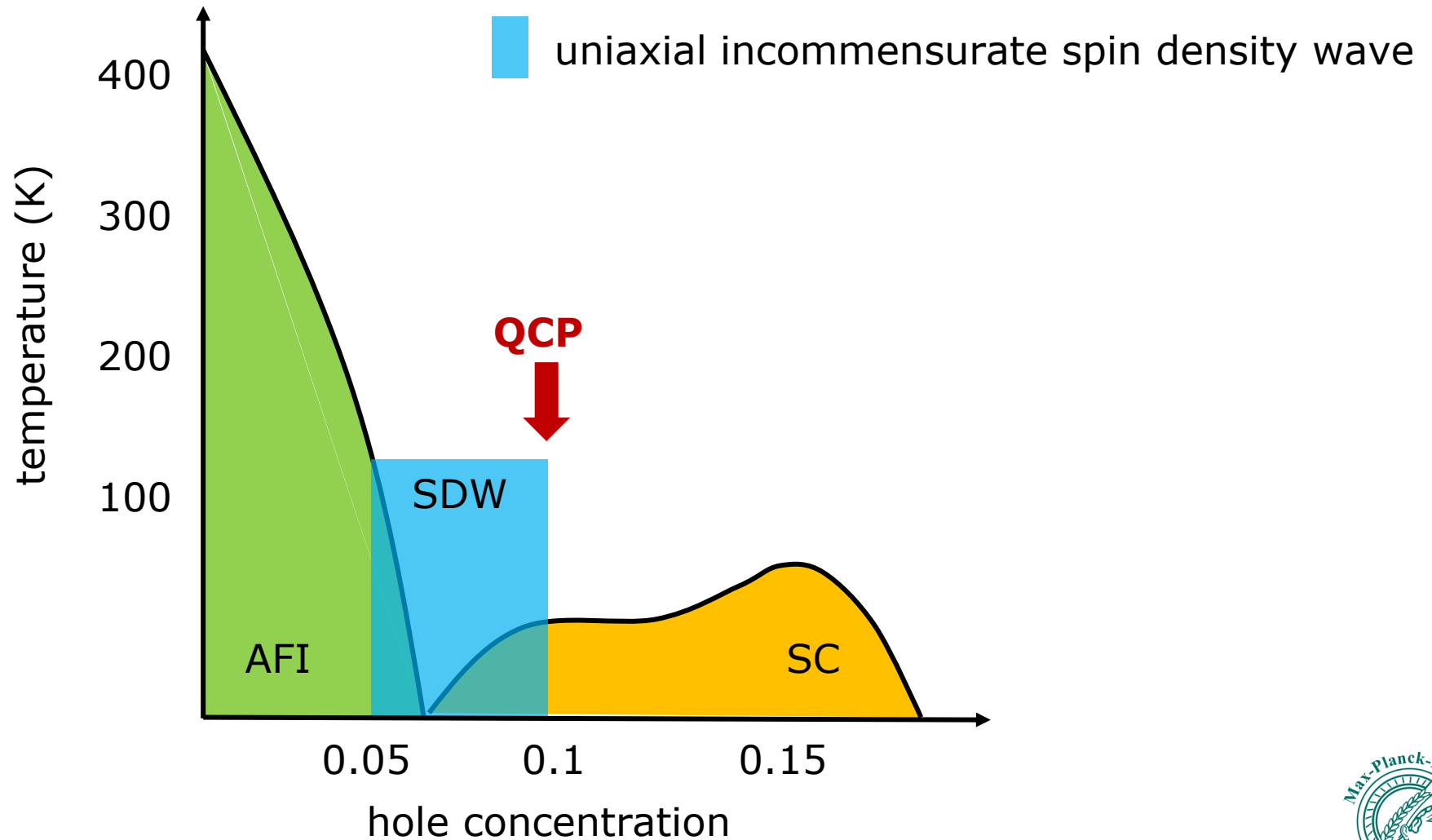


Hinkov et al.

ω/T scaling of χ''

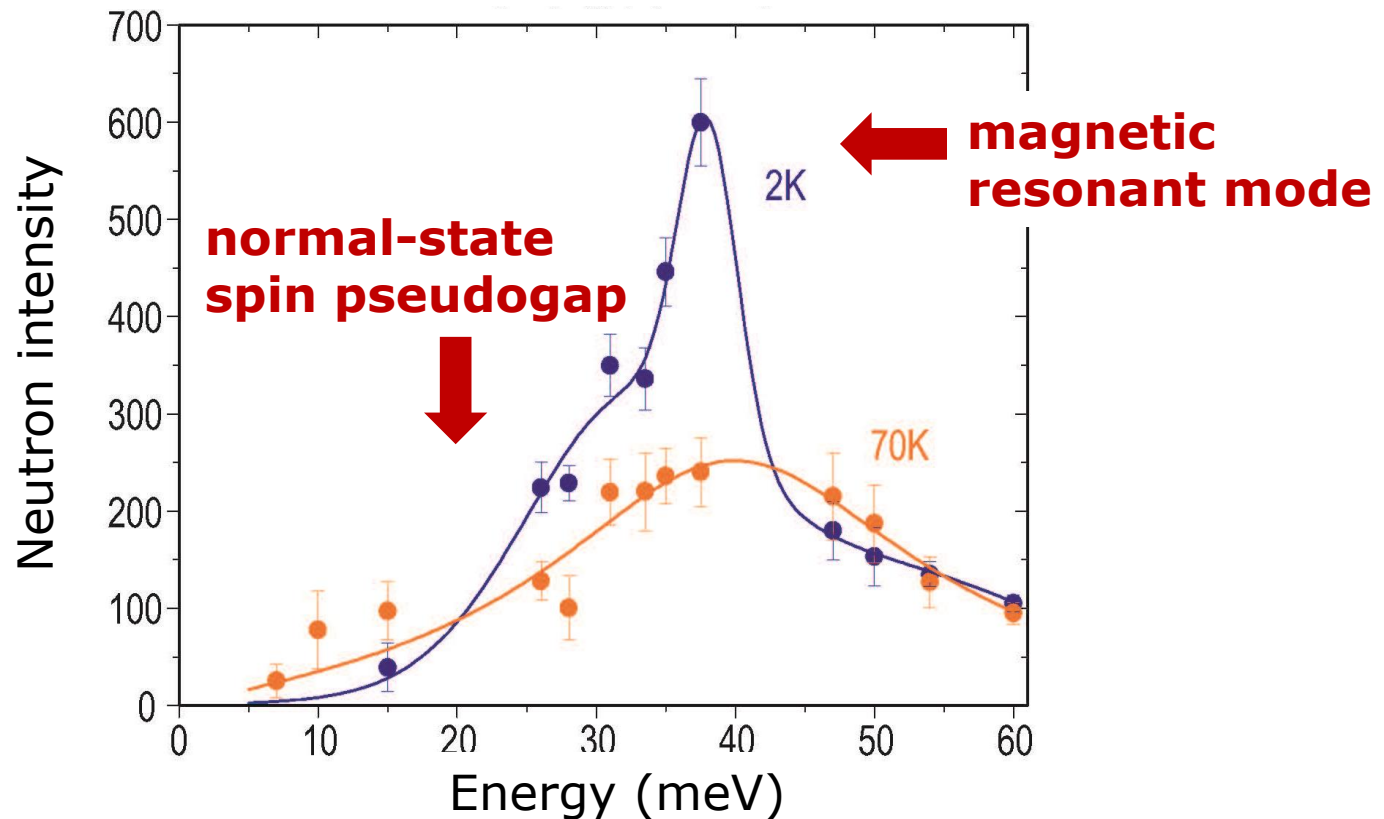
evidence of **quantum critical point** where magnetic order disappears

Competing order in $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$



Beyond the QCP

YBCO_{6.6} T_c = 61 K p ~ 0.11
no static magnetic order

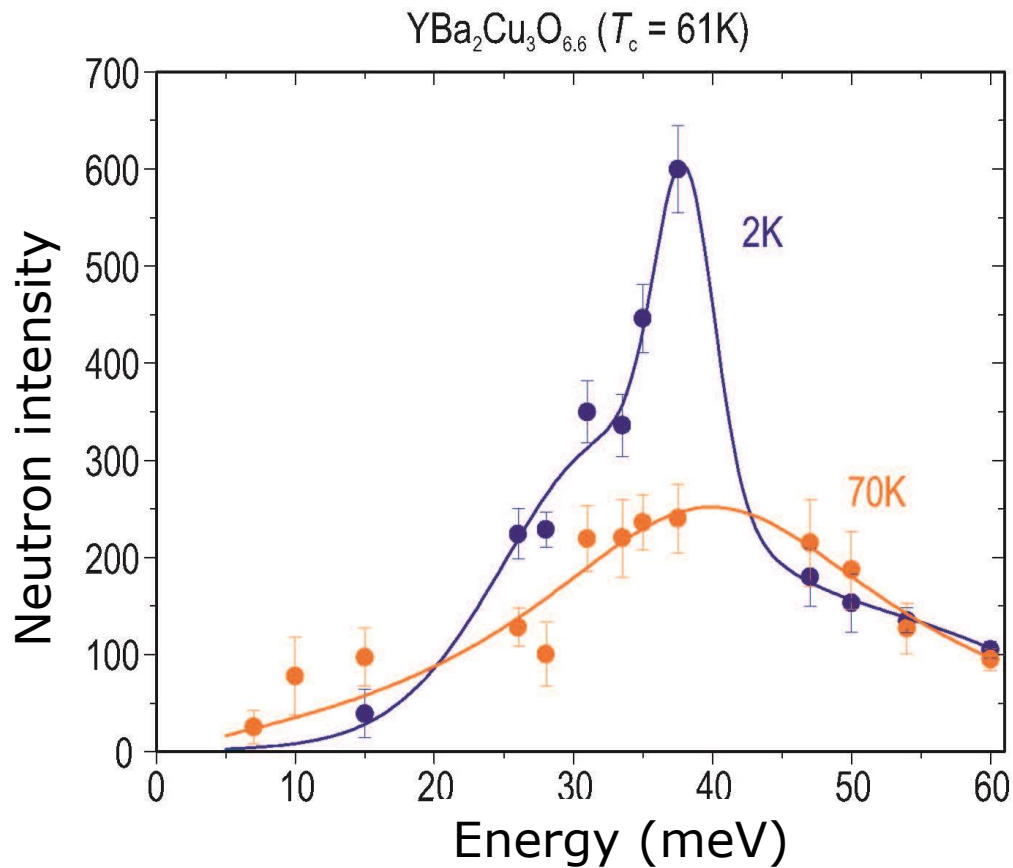


Suchaneck et al., PRL 2010

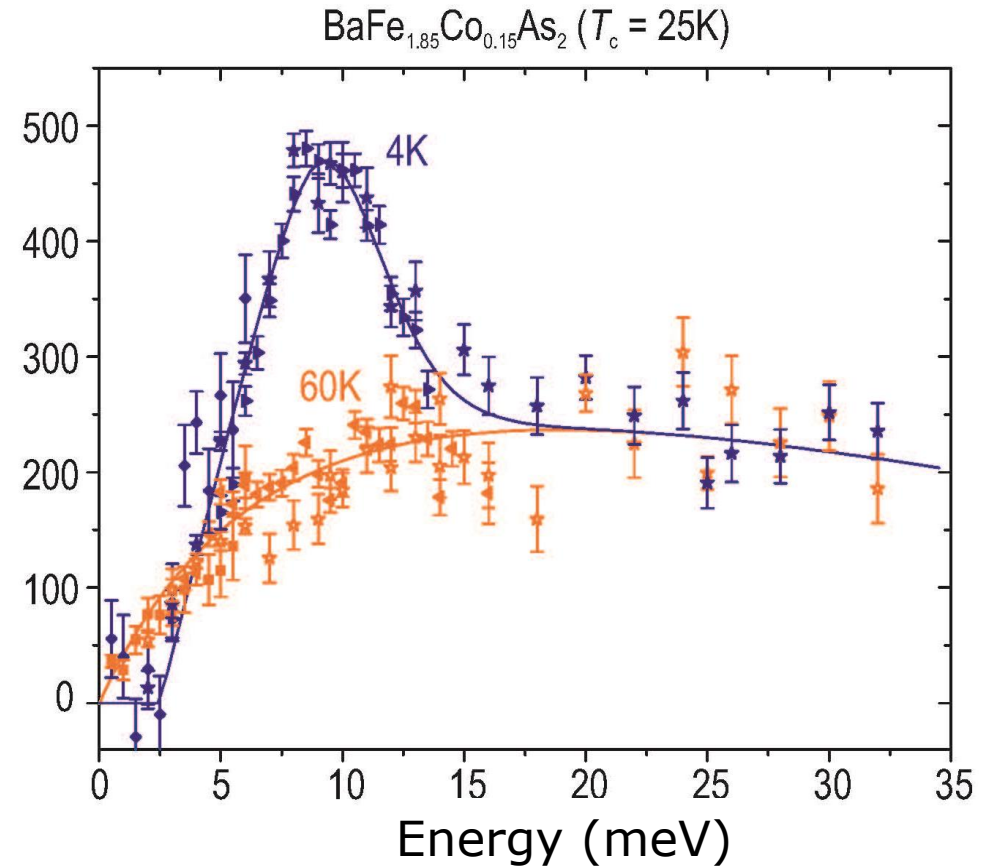
see also Fong et al., PRB 2000; Dai et al., PRB 2000



Magnetic resonant mode



Suchaneck et al., PRL 2010



Inosov et al., Nature Phys. 2010

Christianson et al., Lumsden et al ...

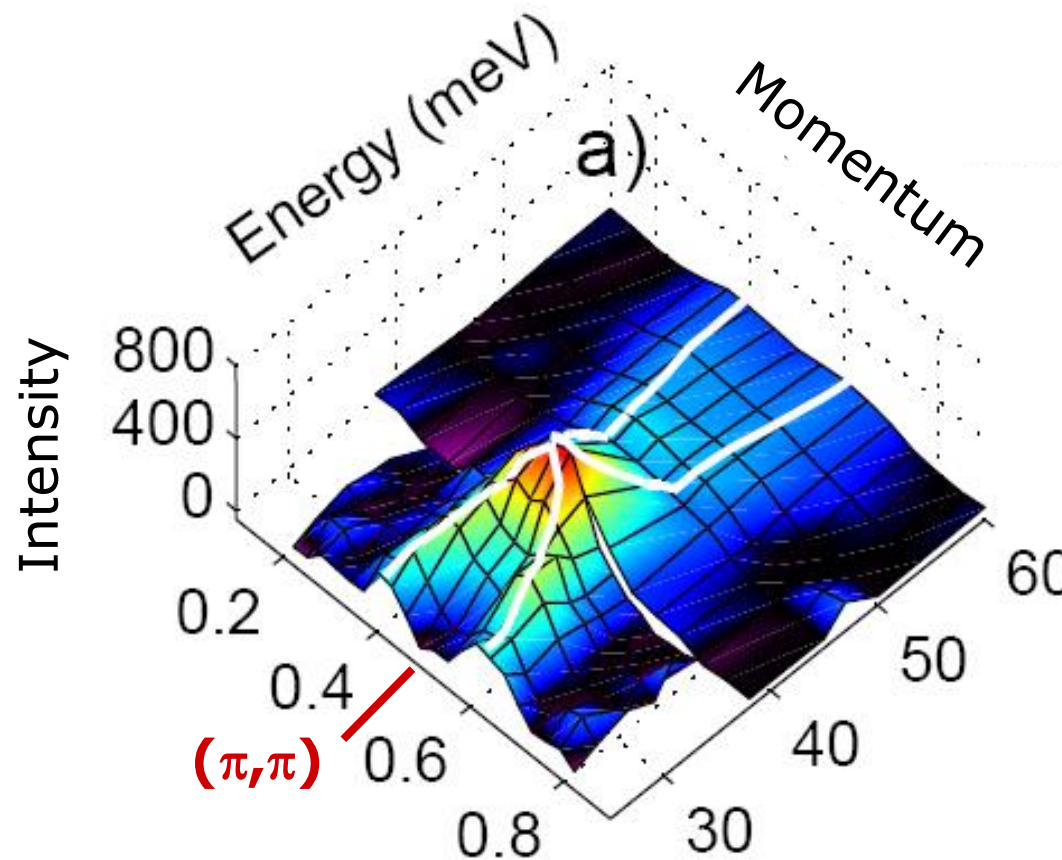
strong feedback effect of the pairing interaction on bosonic spectrum

similar amplitude, T -dependence in two families of high- T_c superconductors

Magnetic resonant mode

$\text{YBa}_2\text{Cu}_3\text{O}_{6.6}$

$T_c = 62 \text{ K}$



Hinkov et al.
Nature 2004
Nature Phys. 2007

Mook et al.
Dai et al. ...

hour-glass dispersion

- magnon-like dispersion at high energies
- “inverted” dispersion at low energies

Inelastic magnetic neutron scattering

$$\frac{d^2\sigma}{d\Omega dE} = 2(\gamma r_0)^2 \frac{k_f}{k_i} N |F(\mathbf{Q})|^2 e^{-2W} \frac{1}{\pi(g\mu_B)^2} \frac{1}{1 - e^{-\hbar\omega\beta}} \chi''(\mathbf{Q}, \omega)$$

itinerant electrons electrons → Lindhard function & RPA

$$\chi_0(q, \omega) = \sum_k \frac{f(E_{k+q\uparrow}) - f(E_{k\downarrow})}{\hbar\omega - (E_{k+q} - E_k - \Delta) + i\varepsilon} \quad \text{band dispersions}$$

$$\chi(q, \omega) = \frac{\chi_0(q, \omega)}{1 - J(q)\chi_0(q, \omega)}$$

RPA expression

$J(q)$ peaked at $q = (\pi, \pi)$

INS from superconductors

coherence factor

$$\chi(q, \omega) = \sum_k \left\{ \frac{1}{2} \left(1 + \frac{\varepsilon_k \varepsilon_{k+q} + \Delta_k \Delta_{k+q}}{E_k E_{k+q}} \right) \frac{f(E_{k+q}) - f(E_k)}{\omega - (E_{k+q} - E_k) + i\delta} \right.$$

$$+ \frac{1}{4} \left(1 - \frac{\varepsilon_k \varepsilon_{k+q} + \Delta_k \Delta_{k+q}}{E_k E_{k+q}} \right) \frac{1 - f(E_{k+q}) - f(E_k)}{\omega + (E_{k+q} + E_k) + i\delta}$$

$$\left. + \frac{1}{4} \left(1 - \frac{\varepsilon_k \varepsilon_{k+q} + \Delta_k \Delta_{k+q}}{E_k E_{k+q}} \right) \frac{f(E_{k+q}) + f(E_k) - 1}{\omega - (E_{k+q} + E_k) + i\delta} \right\}$$

scattering of
thermally excited pairs

pair annihilation

pair creation

$$E_k = \sqrt{\varepsilon_k^2 + \Delta_k^2}$$

Fong et al., PRL 1995

Monthoux & Scalapino, PRL 1994

$\chi'' \rightarrow 0$ at $q = (\pi, \pi)$ in s-wave superconductor

resonant mode implies **sign change** in superconducting gap function

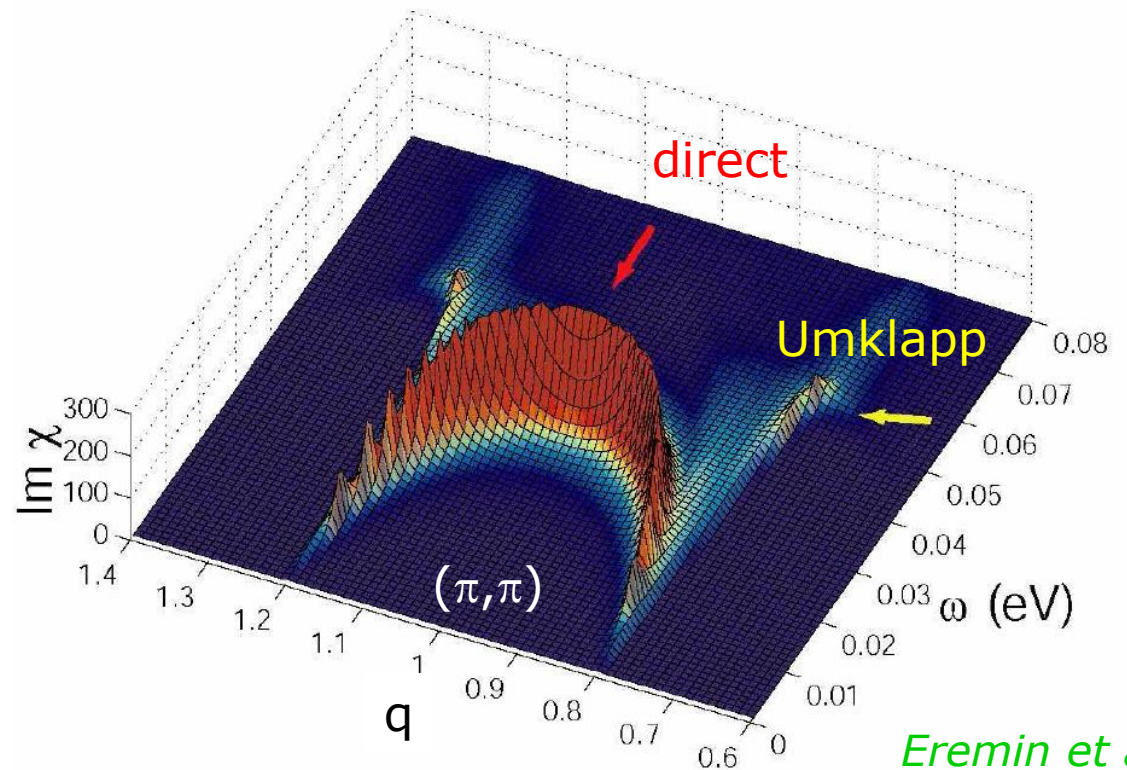
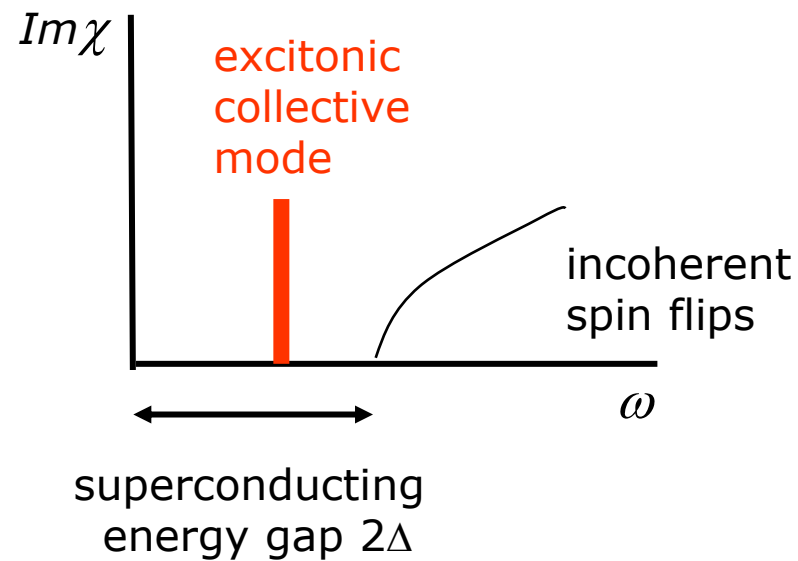
d-wave in cuprates, s_{\pm} in iron pnictides



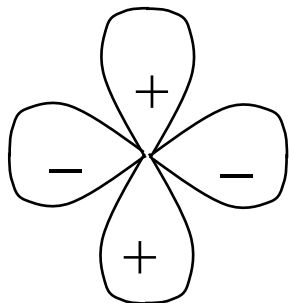
Magnetic resonant mode

spin excitations of a d-wave superconductor

RPA reproduces lower branch of hour-glass dispersion



*Eremin et al.
PRL 2005*



dispersion of resonant mode

momentum-space signature of Cooper-pair wave function

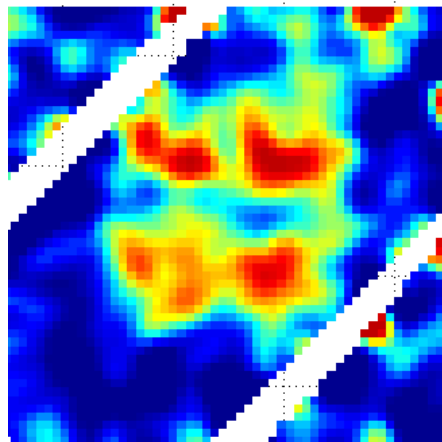
Universal high-energy spin excitations

underdoped cuprates

high-energy excitations weakly dependent on temperature, doping, structural details

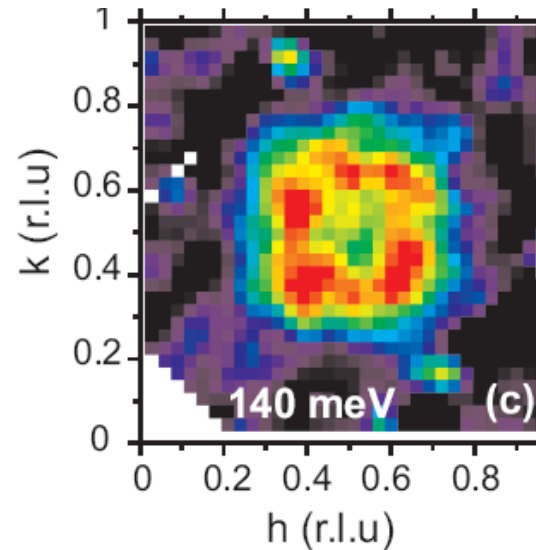


63 meV



0.2 0.4 0.6 0.8
H (r.l.u.)

Hinkov et al.
Nature Phys. 2007

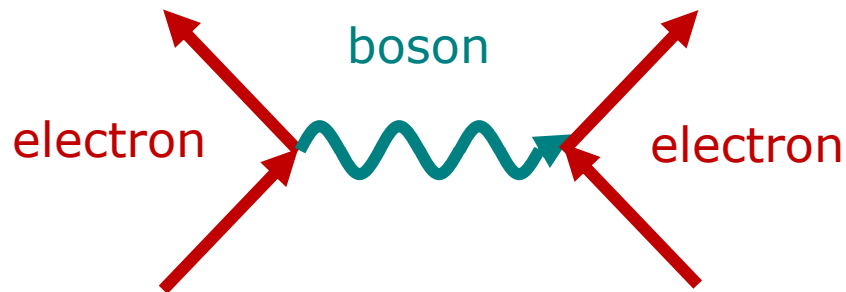


Lipscombe et al.
PRL 2009

Understanding unconventional superconductors

experimental observations

- gapped, non-critical spin fluctuations
- feedback effect of superconductivity on spin fluctuation spectra
- fermionic quasiparticles, at least at low temperature



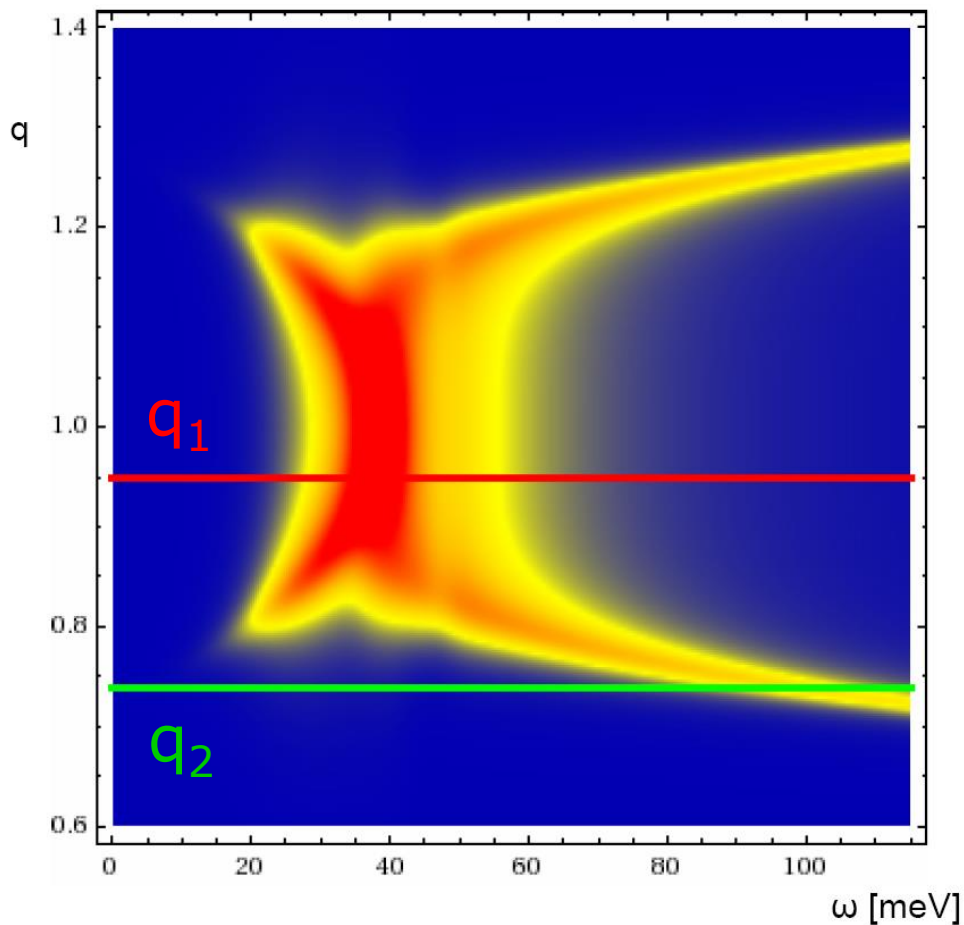
- quantify strength of spin fluctuation mediated pairing interaction
- calculate T_c , energy gap, ... → guideline for materials design

Eliashberg theory is the only method that is currently available.

Joint analysis ARPES – INS

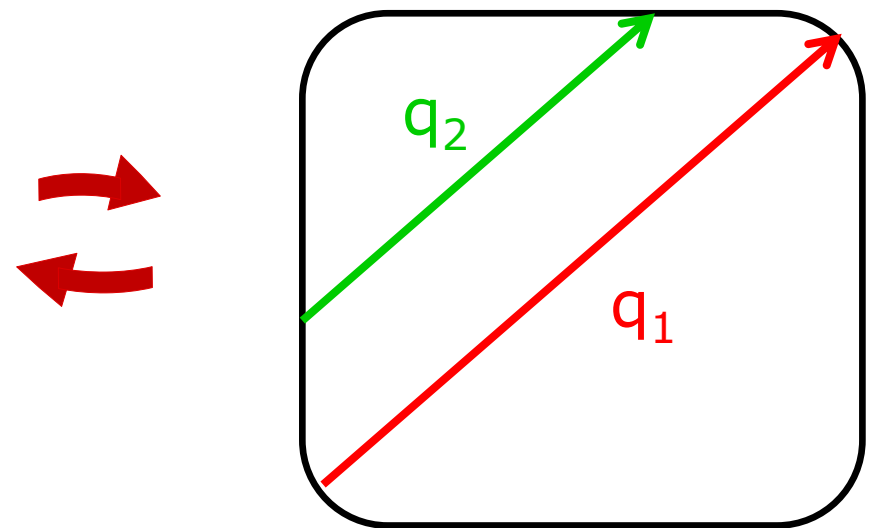
anomalous spin excitations

from neutron scattering

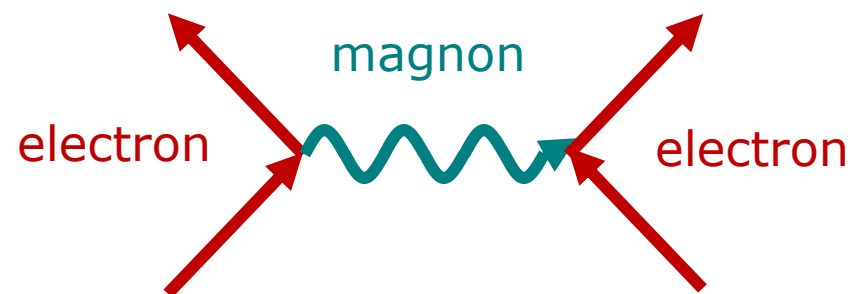


electronic band dispersions

from photoemission



quantitative cross-correlation



Eliashberg theory

- compare INS & ARPES spectra of underdoped $\text{YBa}_2\text{Cu}_3\text{O}_{6.6}$ ($T_c = 62$ K)
→ **spin-fermion coupling constant**
- mean-field superconducting gap equation
→ **$T_c \sim 170$ K**

Dahm et al., Nature Phys. 2009

spin fluctuations have sufficient strength to mediate high-temperature superconductivity

... but what about optimal doping?

ARPES dispersion

