

Cosmic Rays in Galaxies: From Microscales to Macroscales

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The Contributions of Professor Parker



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Professor Parker moved plasma effects in astrophysics from the fringe to the mainstream (*“What about magnetic field used to be a JOKE!”*).

He did this through a unique scientific style that combined clearly articulated physical reasoning, analytical math, & the ability to distill issues to their essence.

He had unbounded curiosity and left his mark on every problem he touched. No one has more equations and processes named after them. ***This talk focusses on cosmic rays.***

Plan of This Talk

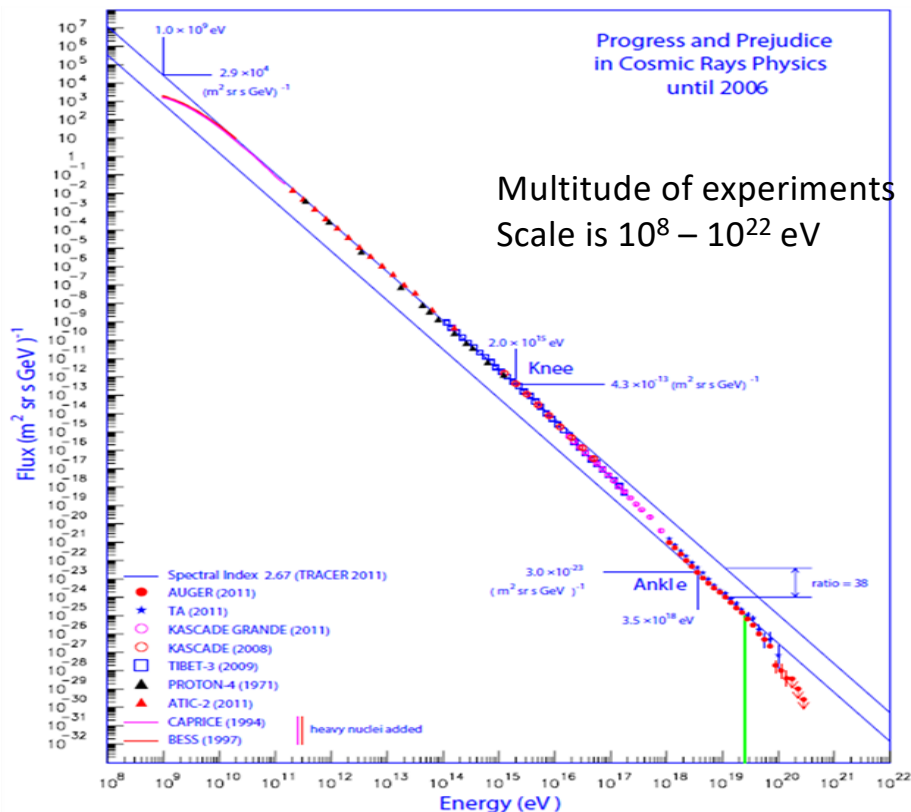
- A few salient properties of cosmic rays.
- Cosmic ray propagation in the interstellar medium.
- Parker's Instability
- Some unsolved problems

Cosmic Rays at a Glance

- Cosmic rays are the relativistic charged particle component of interstellar, circumgalactic, and galaxy cluster plasmas.
 - Broken power law distribution in energy from < 1 GeV to 10's of EeV
 - Essentially collisionless (*but radiative processes & collisions -> observational tools*).
- Exchange energy & momentum with ambient gas through scattering from kinetic scale fluctuations: *The Microscales*
- A fluid theory describes how they can:
 - drive global instabilities in galaxies & galaxy clusters,
 - heat interstellar & intracluster gas
 - quench star formation
 - launch outflows on galactic scales.
- But there are many incomplete aspects of this picture.

The Macroscales

Cosmic Ray Energy Spectrum

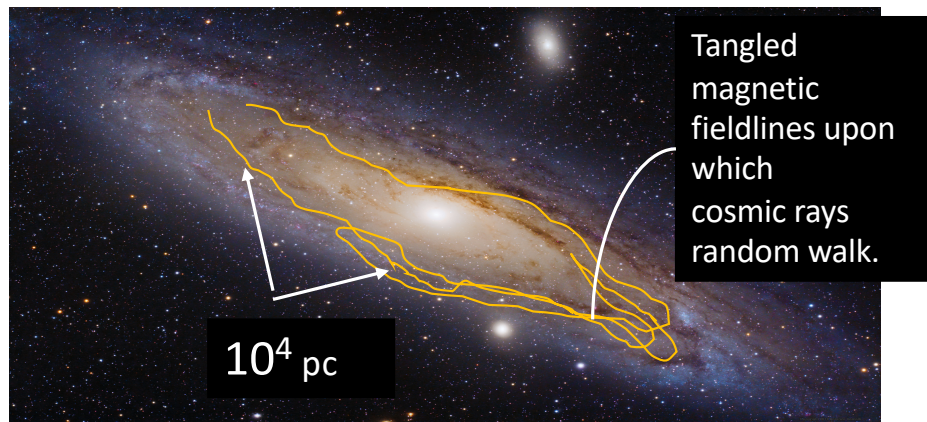


- Mostly protons ($n_i/n_e \sim 50 - 100$)
- $U_{\text{cr}} \sim 1 \text{ eV cm}^{-3}$, similar to magnetic, thermal, & radiation energy densities in galaxies
- $n_{\text{cr}} \sim 3 \cdot 10^{-10} \text{ cm}^{-3}$
- $\langle E \rangle \sim 3 \text{ GeV}$
- Above $\sim 10^9 \text{ GeV}$, not magnetically confined by Galactic magnetic field.

Summary of Cosmic Ray Properties

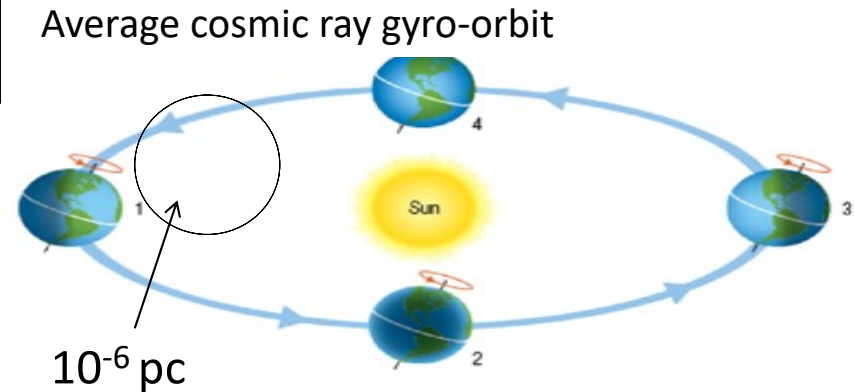
- Cosmic rays are accelerated from the interstellar medium in events, probably supernova remnant shocks, that produce an $\sim E^{-2}$ spectrum
- Spectrum is steepened by energy dependent losses & escape.
- GeV cosmic rays are confined to the Milky Way for $\sim 2 \cdot 10^7$ yr & scattered with a mean free path $\lambda \sim 1$ pc. $\sim 10\%$ of non ν supernova power required to maintain energy density in a steady state.
- Detected in star forming galaxies over a wide range of sizes, ages & activity levels.

Cosmic Ray Propagation is a Multiscale Problem



Milky Way-like galaxy

Even if we had an exact representation of the galactic magnetic field, we could never integrate enough cosmic ray orbits or account for all processes. *Must resort to a statistical description. Parker provided an early one.*



Fluid theory for global modeling must describe a collisionless, non-Maxwellian component. *Near isotropy makes this possible.*

Parker's Transport Equation For Cosmic Rays

with an Isotropic Phase Space Distribution $n(x,p,t)$

$$\frac{\partial n}{\partial t} + \mathbf{u} \cdot \nabla n = \nabla \cdot (\hat{k} \nabla n) + \frac{1}{p^2} \frac{\partial}{\partial p} \left(p^2 \kappa_{pp} \frac{\partial n}{\partial p} \right) + \frac{p}{3} (\nabla \cdot \mathbf{u}) \frac{\partial n}{\partial p} + S$$

The diagram illustrates the components of Parker's transport equation. Red arrows connect four descriptive boxes to their corresponding terms in the equation:

- Advection & compression by fluid flow \mathbf{u}** : Points to the term $\mathbf{u} \cdot \nabla n$.
- Sources & sinks**: Points to the term S .
- Diffusion along magnetic field**: Points to the term $\nabla \cdot (\hat{k} \nabla n)$.
- Diffusion In momentum**: Points to the term $\frac{1}{p^2} \frac{\partial}{\partial p} \left(p^2 \kappa_{pp} \frac{\partial n}{\partial p} \right)$.

Statistical description appears as physical & momentum space diffusion. Field line geometry appears through anisotropy of spatial diffusion.

Magnetically mediated interaction between
cosmic rays and thermal gas.

Perpendicular Dynamics is Straightforward

Cosmic ray force balance:

$$\nabla_{\perp} P_c = \frac{\mathbf{J}_c \times \mathbf{B}}{c}$$

Lorentz force on thermal gas: $\mathbf{J}_g \times \mathbf{B}/c$

$$= \frac{\mathbf{J} \times \mathbf{B}}{c} - \frac{\mathbf{J}_c \times \mathbf{B}}{c}$$

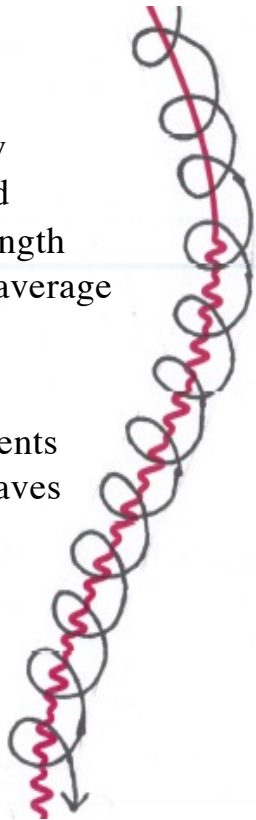
Cosmic ray pressure gradient
introduced through Lorentz force

$$= \frac{\mathbf{J} \times \mathbf{B}}{c} - \nabla_{\perp} P_c$$

Parallel Dynamics: Gyroresonant Scattering

Orbits follow fieldlines and short wavelength fluctuations average out.

Scattering agents are Alfvén waves



$$\omega - kv\mu = \pm\omega_c;$$

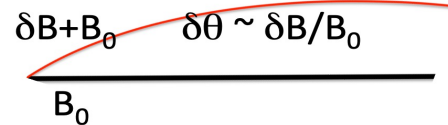
$$\mu \equiv \frac{v_{\parallel}}{v}$$

Estimate Diffusion Coefficient

Scattering frequency

$$\nu \sim \frac{\langle (\delta\theta)^2 \rangle}{\delta t}$$

$$\sim \omega_c \left(\frac{\delta B}{B_0} \right)^2$$



Scattering is nearly elastic: $v_A/c \ll 1$.

Theories Diverge over Source of Waves

Waves generated by the cosmic rays themselves through streaming instability: **Self-Confinement**

Classical Cosmic Ray Hydrodynamics

Waves are present as part of a turbulent cascade: **Confinement by Extrinsic Turbulence**

Generalized Cosmic Ray Hydrodynamics

See EZ2017 for bridging these theories.

Gyroresonant Amplification of Alfvén Waves

$$\Gamma_{cr} = \frac{\pi^2 q^2 v_A^2}{2 c^2} \sum_{\pm} \int \delta(\omega - kv\mu \pm \omega_c) v(1-\mu^2) \left[\frac{\partial f}{\partial p} + \left(\frac{kv}{\omega} - \mu \right) \frac{1}{p} \frac{\partial f}{\partial \mu} \right] p^2 dp d\mu,$$

resonance

damping

excitation by anisotropy

For drift anisotropy, a simple approximation to the growth rate:

$$\Gamma_{cr} \sim C \omega_{ci} \frac{n_{cr}}{n_i} \left(\frac{v_D}{v_A} - 1 \right)$$

$\Gamma_{cr} = 0$ means cosmic ray isotropy in the wave frame.

Here & elsewhere we're interested in the bulk cosmic rays with $\gamma \sim 1$

Fokker – Planck (F-P) Equation

Back reaction of waves (subscript 1) on zero order cosmic ray distribution function f_0

$$\frac{df_0}{dt} = - \left\langle \frac{q}{m} \left(\mathbf{E}_1 + \frac{\mathbf{v} \times \mathbf{B}_1}{c} \right) \cdot \nabla_p f_1 \right\rangle$$

$$= \nabla_p \cdot \mathbf{D} \cdot \nabla_p f_0.$$

Pitch angle scattering ($D_{\mu\mu}$) dominates:

Scattering frequency $\nu \sim \omega_c (\delta B/B)^2$

$D_{p\mu} = D_{\mu p}$ are order (v_A/c) D_{pp} is order $(v_A/c)^2$,
requires waves traveling in
both directions

Assumes small angle
scattering by nearly
periodic, randomly
phased waves.

Energy Equation

Multiply F-P eqn. by particle energy ε & integrate over momentum space:

$$\frac{\partial U_c}{\partial t} + \nabla \cdot \tilde{\mathbf{W}}_c = - \int d\omega dk 2\Gamma_c(\omega, k) I(\omega, k).$$

Energy density

Energy flux

Energy transfer to waves

Frequent Scattering Approximation

Relate anisotropy to spatial gradient:

$$D_{\mu\mu} \frac{\partial f_0}{\partial \mu} + D_{\mu p} \frac{\partial f_0}{\partial p} = -\frac{v(1-\mu^2)}{2} \frac{\partial f_0}{\partial z}$$

Left hand side related to Γ_c integrand! Energy equation simplifies to:

$$\frac{\partial U_c}{\partial t} + \nabla \cdot \tilde{\mathbf{W}}_c = \mathbf{v}_A \cdot \nabla P_c.$$

Heating looks like friction, & instability growth rate can be written in terms of pressure gradient.

Balance Wave Damping & Excitation to Ascertain Degree of Confinement

- Ion – neutral friction
 - *Important in H I, H₂ gas*
- Nonlinear energy transfer to thermal ions
 - *Important in hot gas*
- Distorted by wandering of background field
 - *Important when small scale turbulence is present, enhanced at high β by $\beta^{1/2}$.*

In steady state, damping balances growth, determining the streaming rate, wave amplitude, wave dissipation rate, & constraining the diffusion rate.

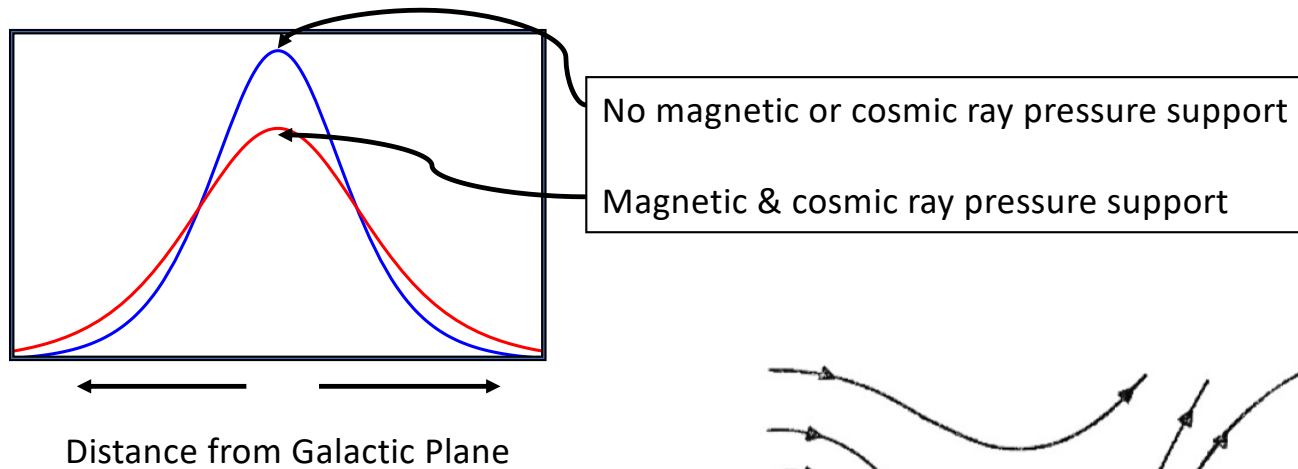
In the Milky Way, self-confinement works below 100-200 GeV (where most of the pressure is) but this is sensitive to damping rates.

Summary of the Fluid Treatment

- Classical Cosmic Ray Hydrodynamics (CCRH)
 - Cosmic rays exert a force on the thermal gas $-\nabla P_c$,
 - Cosmic rays heat the gas at the rate $|\mathbf{v}_A \cdot \nabla P_c|$
 - $P_c/\rho_g^{\gamma_c/2}$ is constant along magnetic flux tubes.
- For Generalized Cosmic Ray Hydrodynamics (GCRH), where confinement is from extrinsic turbulence, leave out streaming & heating.
- Numerical implementation methods by Sharma, Jiang & Oh, Thomas & Pfrommer.

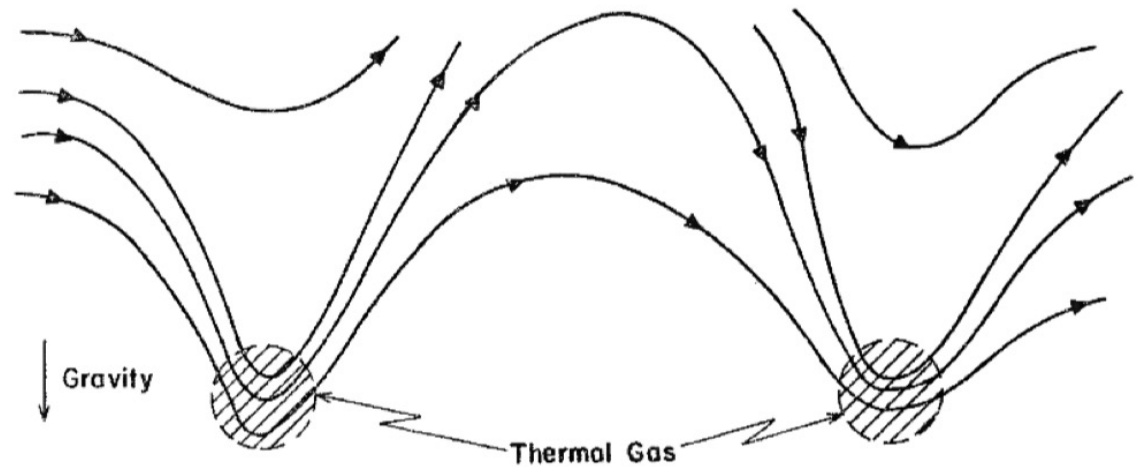
Applied to a range of gas dynamics problems in galaxies and beyond.

Parker's Instability: A Large Scale Instability of Stratified Disks Supported by Magnetic Fields & Cosmic Rays



Free energy source: gravitational potential energy

Must exceed the work needed to compress gas in the valleys.



Parker 1966

Parker (1966) Quantified These Considerations

$$P_g \equiv \rho u^2$$

$$\alpha \equiv \frac{B^2}{8\pi P_g}$$

$$\beta \equiv \frac{P_c}{P_g}$$

$$H = \frac{u^2}{g} (1 + \alpha + \beta)$$

$$\gamma_g > \frac{(1 + \alpha + \beta)^2}{1 + \frac{3\alpha}{2} + \beta}$$

α & β are constants

Gas supported above its natural scale height by magnetic fields and cosmic rays

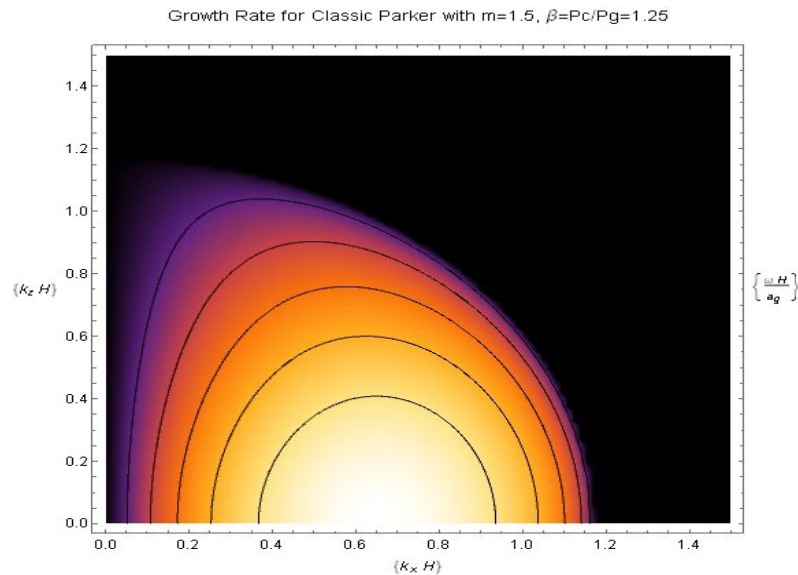
Required for stability. Parker argued $\gamma_g = 1$, $\gamma_c = 0$.

Enduring Message of Classic Parker

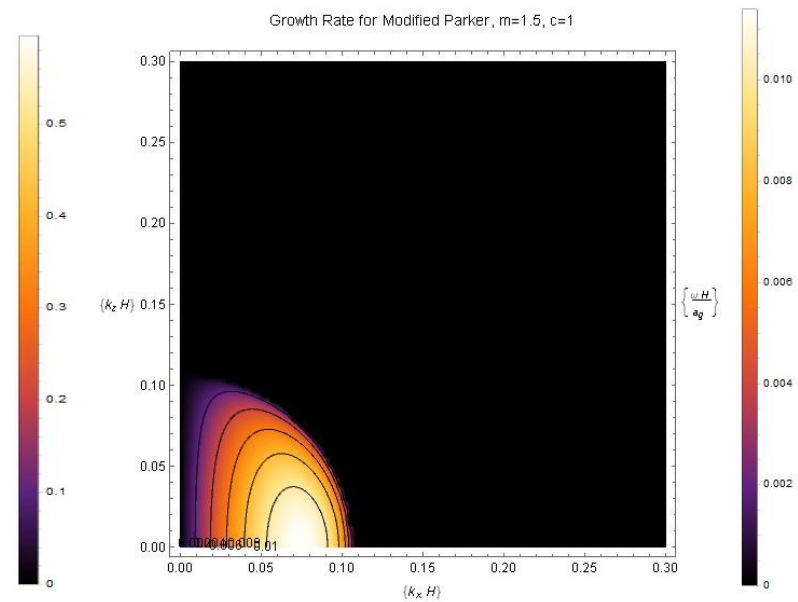
- Parker's Instability sets an upper limit to the fraction of hydrostatic support that can be provided to a gas layer by magnetic fields and cosmic rays
 - Rayleigh-Taylor analogy: light fluid/heavy fluid
 - Compressibility of ISM is critical to determining that level
 - *Thermal pressure is partly turbulent pressure (Zweibel & Kulsrud 1975)*
 - *Extremes of temperature greater than appreciated in the 1960s*
 - *Cosmic ray physics*
- Applications to accretion disks & galactic dynamo as well as the vertical structure of the disk.

Effect of Cosmic Ray Compressibility

$$\gamma_c = 0$$



$$\gamma_c = 4/3$$

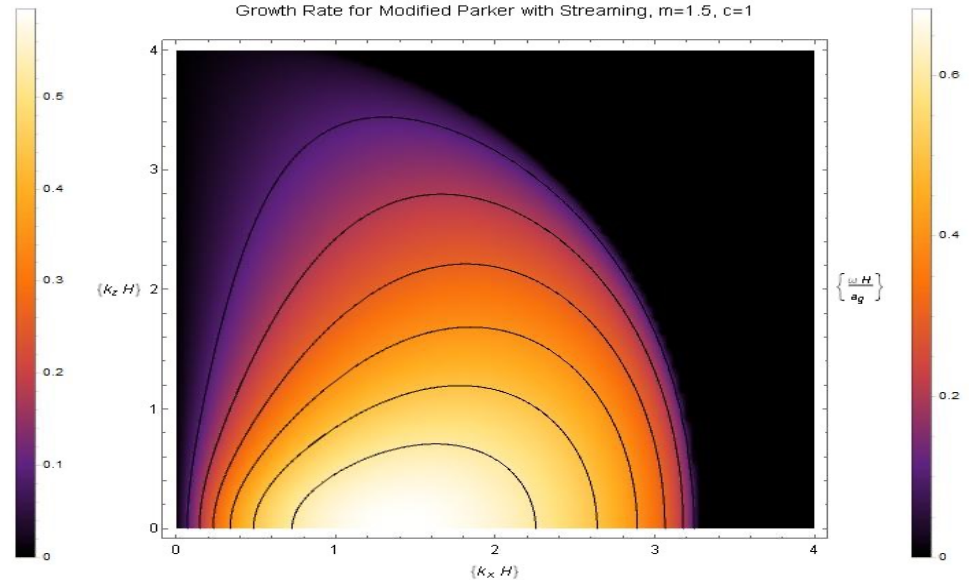
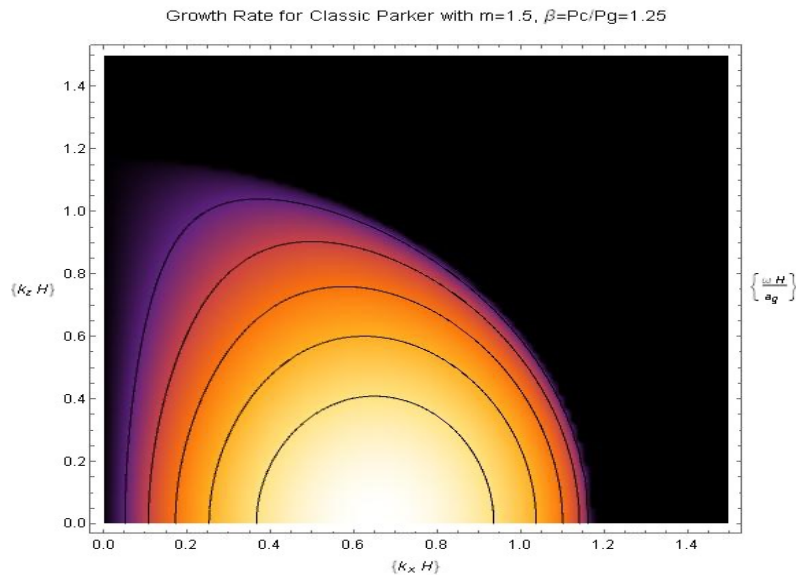


Growth rate contours in the horizontal (k_x), vertical (k_z) plane for $\gamma_c = 0$ ("Classic Parker") and $\gamma_c = 4/3$. Note different scales.
From EH & EZ 2018.

Effect of Cosmic Ray Streaming

$$\gamma_c = 0$$

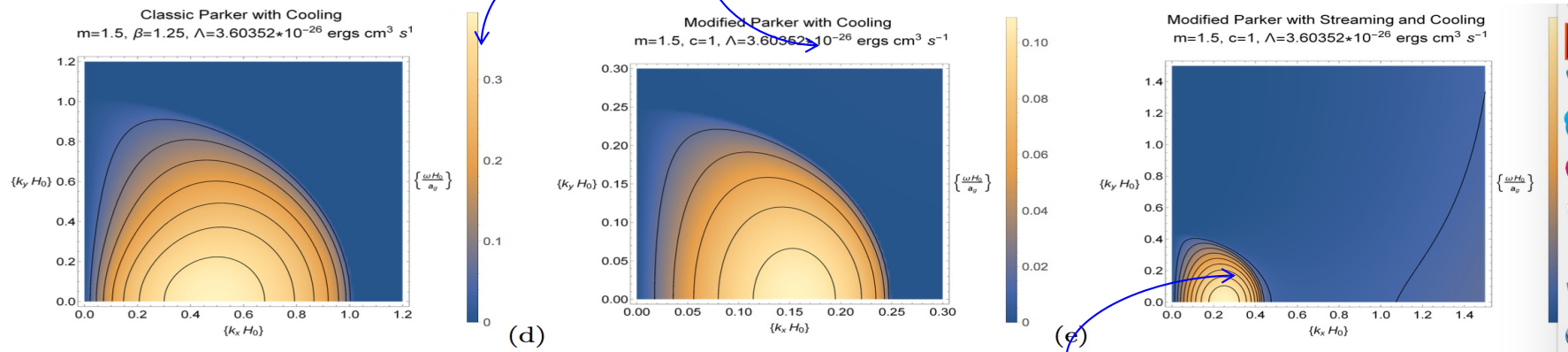
$$\text{Alfvénic streaming \& } \gamma_c = 4/3$$



Growth rate contours in the horizontal (k_x), vertical (k_z) plane for $\gamma_c = 0$ ("Classic Parker") and $\gamma_c = 4/3$. From EH & EZ 2018. Cosmic ray heating increases the growth rate & destabilizes shorter wavelengths.

Effects of Radiative Cooling

More unstable: softens thermal gas eqn. of state

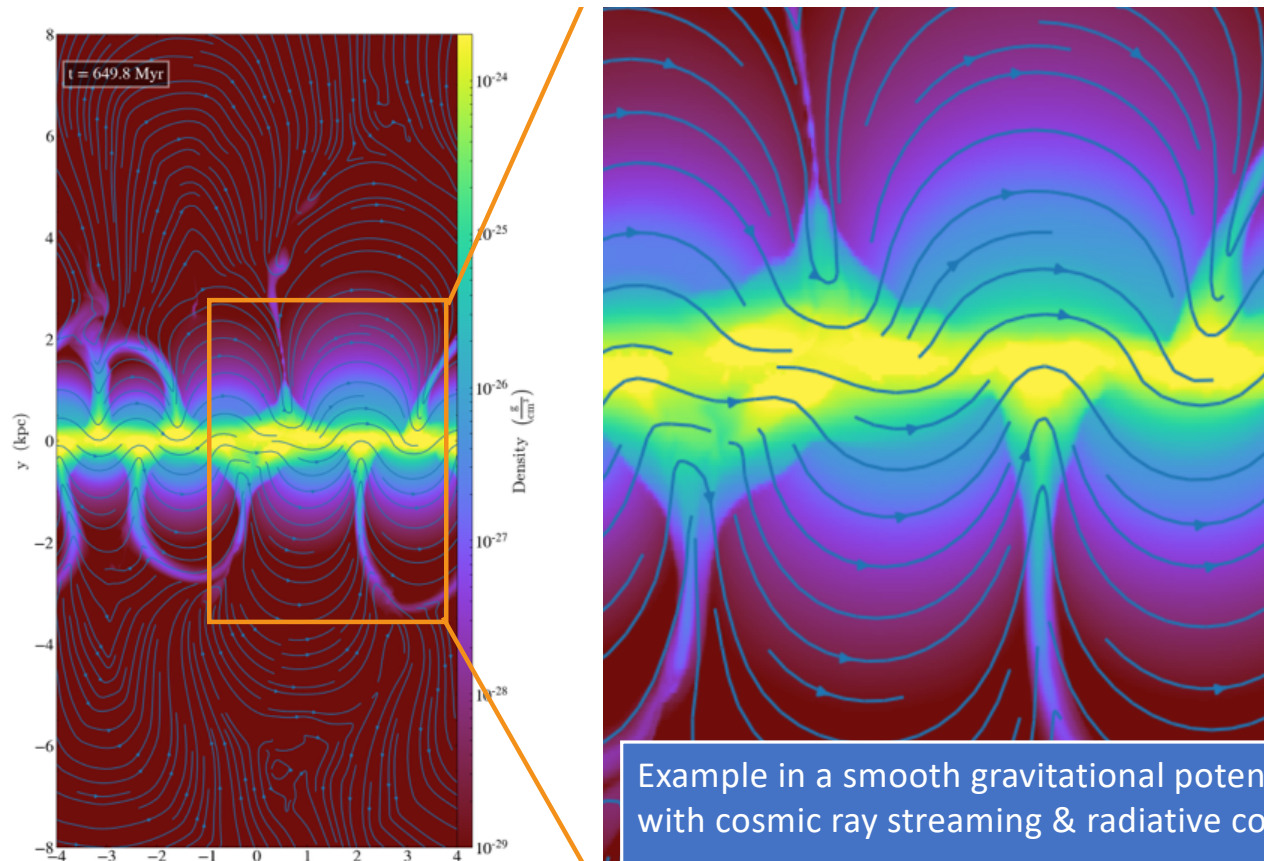


More stable: reduced effect of cosmic ray heating

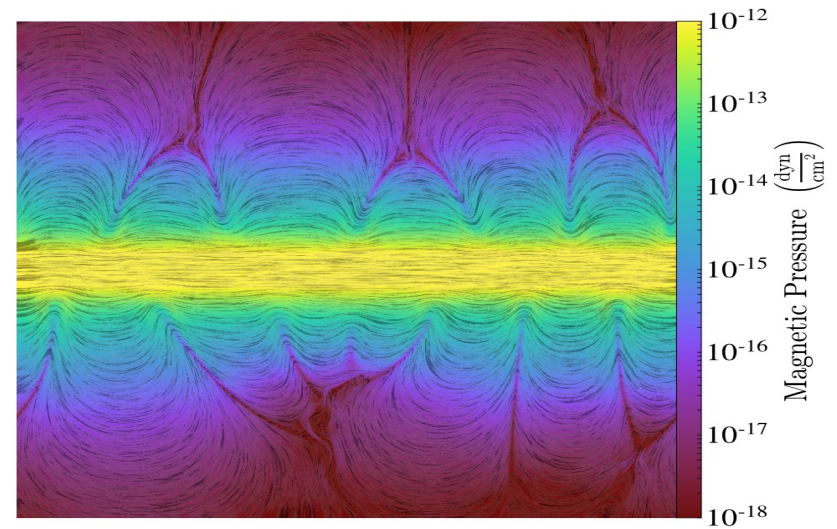
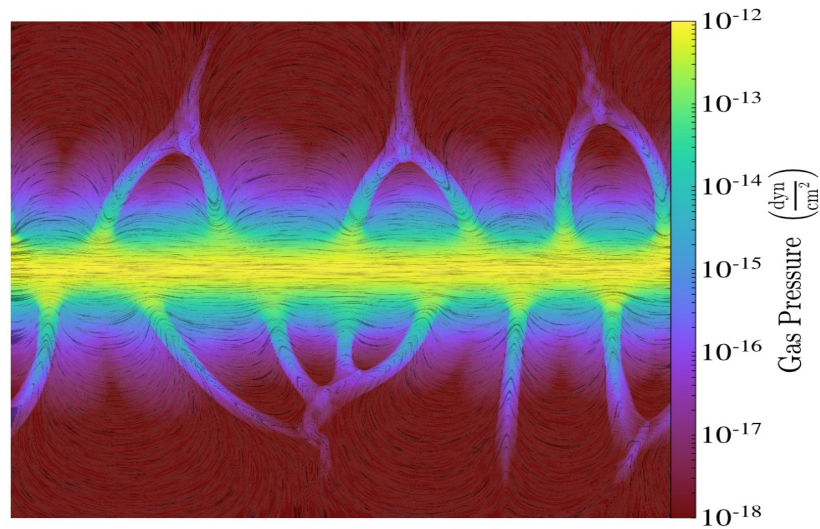
Modeling with FLASH – HBZ 2019

- Isothermal layer in a $\tanh(z/H)$ gravitational field (*smooth*).
- Radiative cooling can be implemented.
- Cosmic ray streaming & diffusion along magnetic field lines can be implemented.
- Rotation, self gravity, & stellar energy input not yet included.

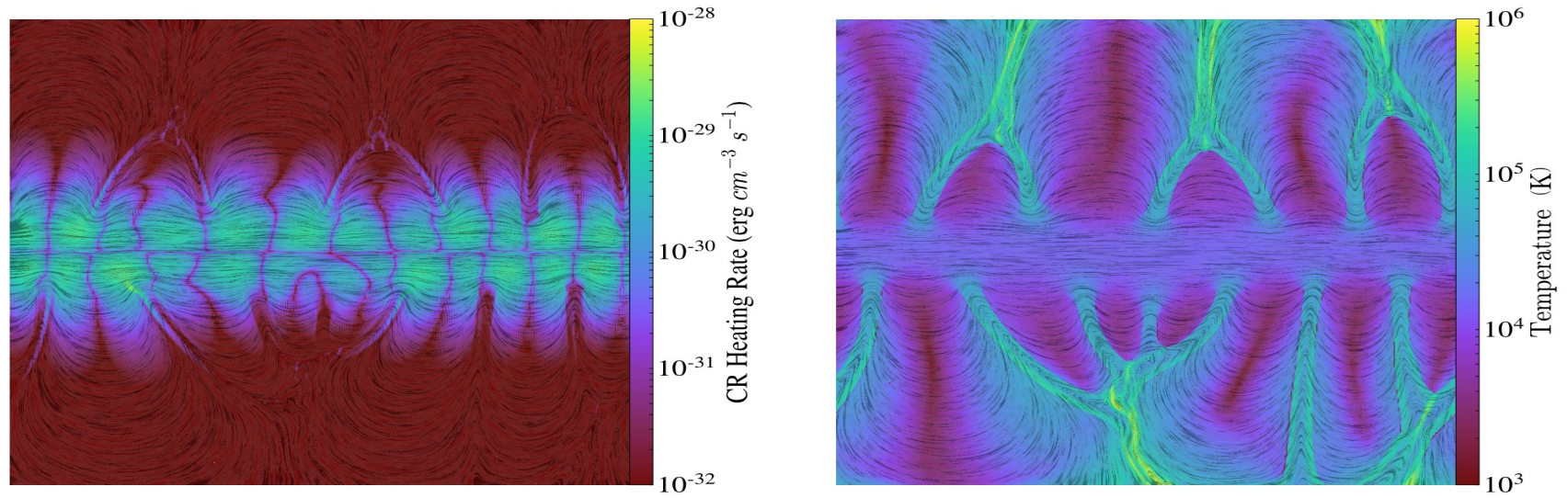
Density Evolution of an Unstable System



Pressure – Magnetic Field Lines Overplotted



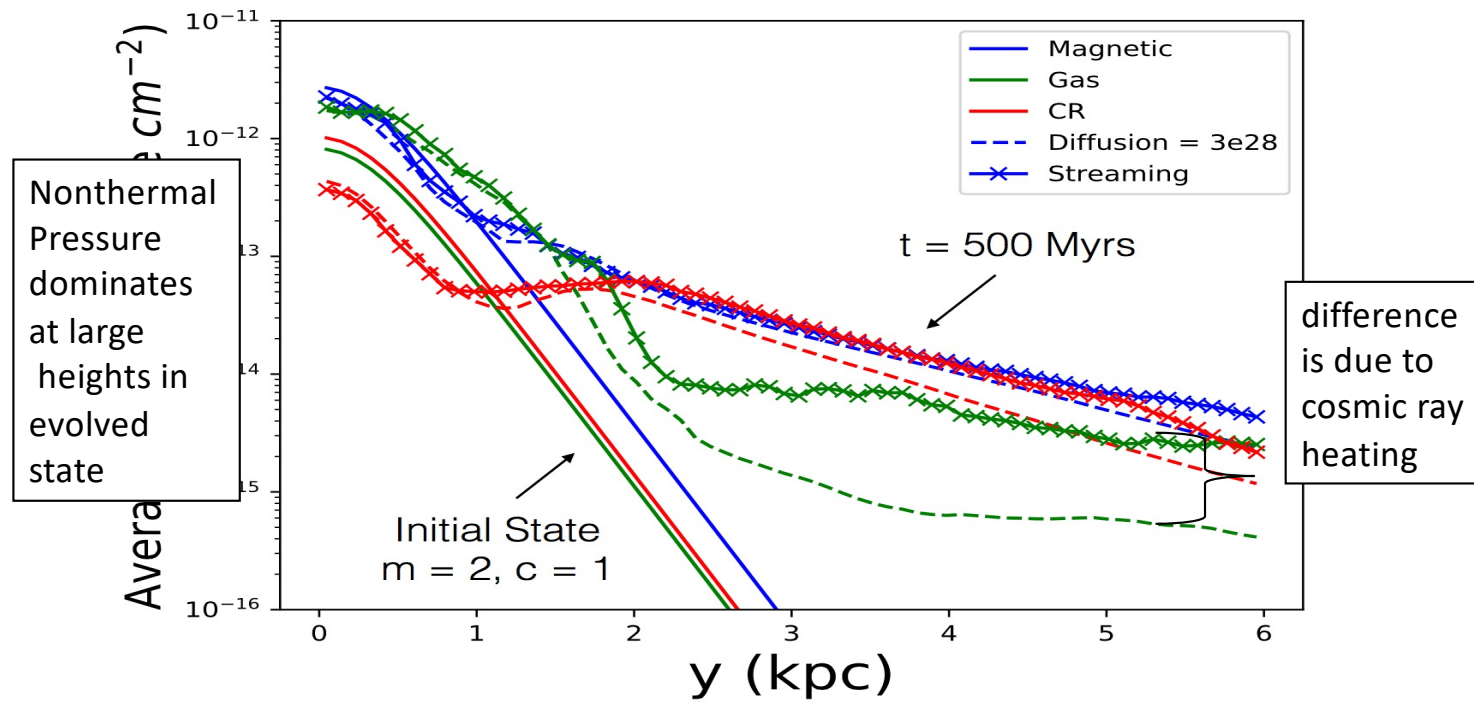
Heating & Temperature



EH,CB, EZ

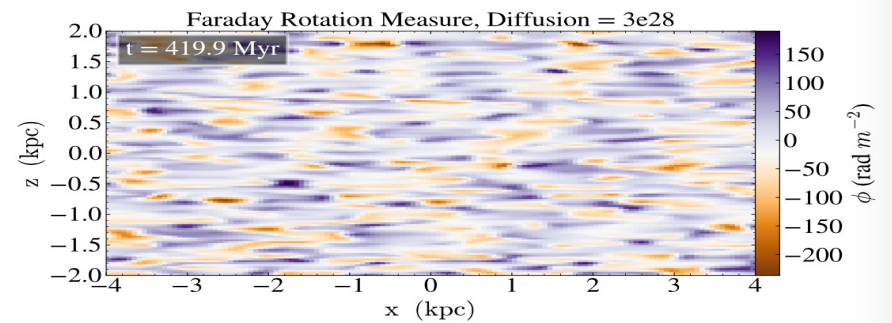
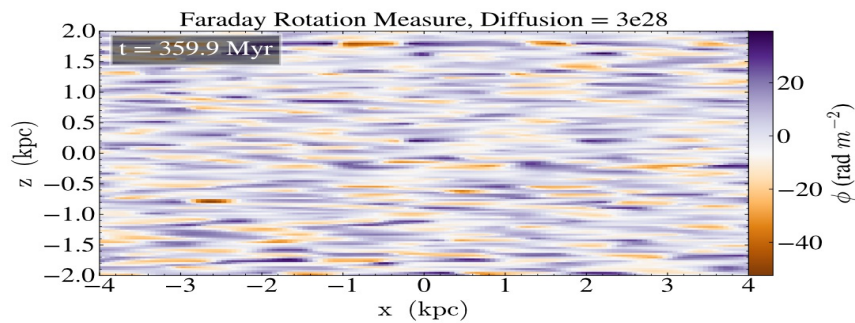
Radiative cooling is not included here.

Pressure Profiles



Parker in 3D

Top down view of rotation measures that would be seen in face on disk with fully developed Parker



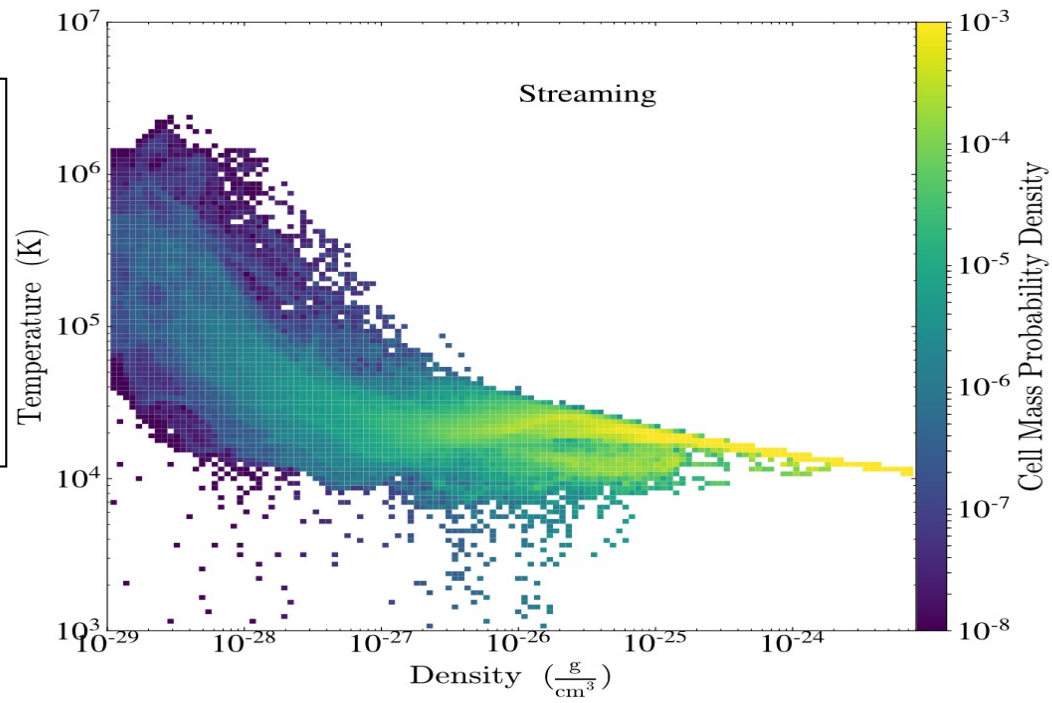
HBZ 2019

Short in-plane wavelength perpendicular to B means structure would be difficult to see edge on (*Talk by Rainer Beck!*).

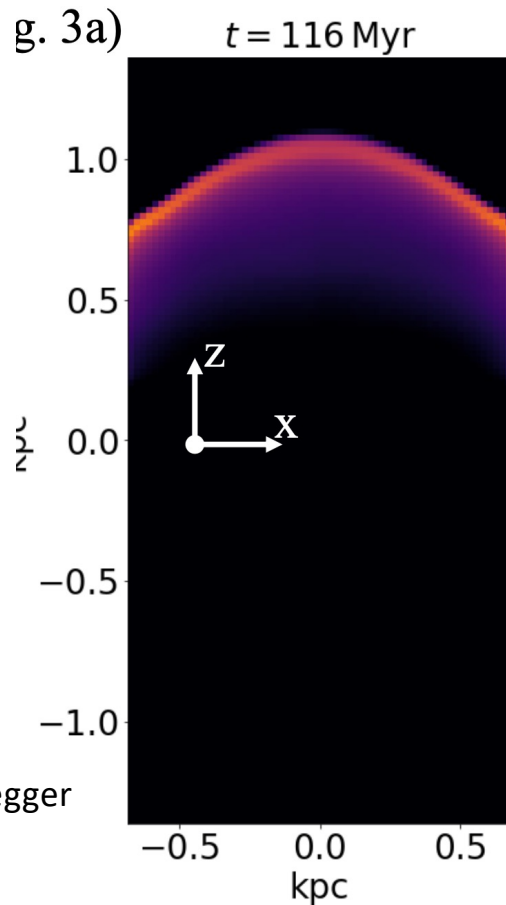
Thermal Phase Plot

Radiative cooling
is included
but unstable
T gas is present.

HBZ



Recent Work on Buoyancy



- Localized cosmic ray injection by single or clustered supernovae.
- Cosmic ray pressure pushes & heats thermal gas, driving a flow & evacuating the tube
- Buoyant flux tube rises
 - Is this faster than the Parker Instability itself?
 - Does it matter if the background medium is Parker unstable?

Summary

- Parker's instability exists when modern cosmic ray transport is included.
 - Stability of puffed up disks in simulations with no streaming or diffusion but $\gamma_c = 4/3$ is probably real.
 - Both streaming and diffusion enhance the instability, but the heating due to streaming is notable.
 - Instability produces dense pockets of cold gas and a halo dominated by nonthermal pressure on timescales of 0.5 Gyr.

Some Important Unsolved Problems

- Growth and damping of anisotropy instabilities in an inhomogeneous, e.g. turbulent background. *Important for estimating whether cosmic rays are well confined & how much they heat the thermal gas.*
- Full theory of diffusive propagation. *Important for developing a theory for the energy dependence of the diffusion coefficient & how cosmic rays cross fieldlines.*
- Propagation in a clumpy medium. Bottlenecks and time dependent propagation. Full implementation of confinement by pressure anisotropy in fluid models. *Important for understanding whether cosmic ray coupling pervades space & time or is intermittent.*