

Statistical description in terms of Ensemble of Blocked states (inherent structures)

- ⊗ Basics of statistical mechanics; fundamental postulates and the ensembles of equilibrium statistical mechanics (We worked these out yesterday)
- ⊗ Consequences of the Boltzmann distribution:
 - ⊗ Fluctuation-dissipation relations
 - ⊗ All of Thermodynamics
- ⊗ Work out consequences of the Edwards' distribution
- ⊗ Write configurational partition function of a liquid in terms of inherent structures, explore consequences
- ⊗ Observations in simulations of Lennard-Jones systems
- ⊗ Introduce a toy model where we can work things out explicitly

Toy model

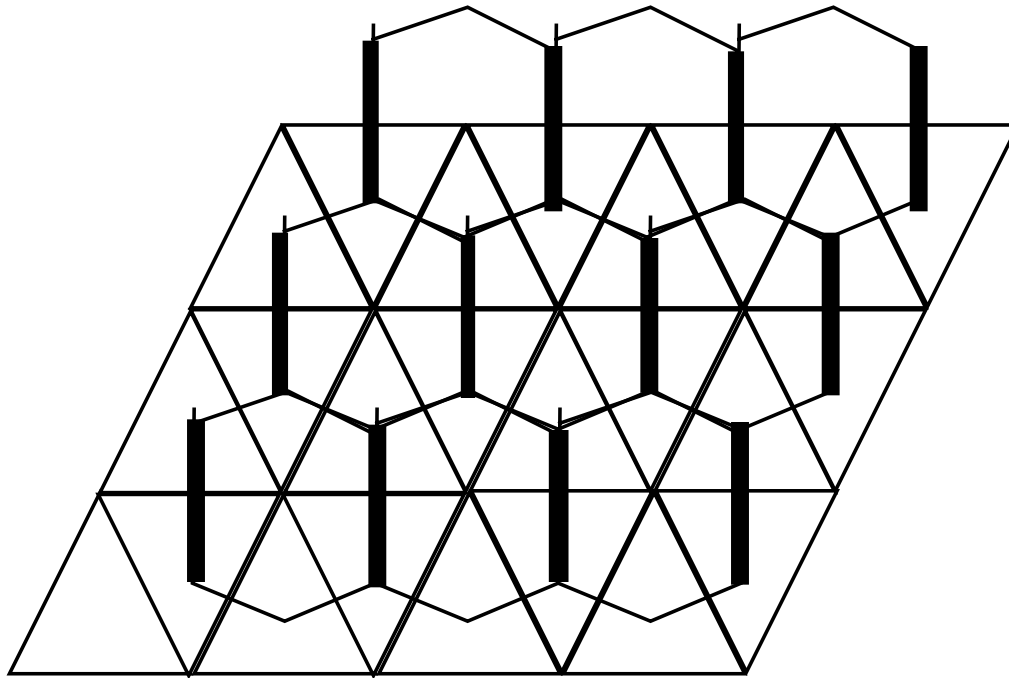
Triangular-lattice Ising antiferromagnet (interaction strength J)

Frustrated system with exponentially large # of ground states

Making the lattice compressible or introducing a staggered field lifts degeneracy and picks out a globally stable state

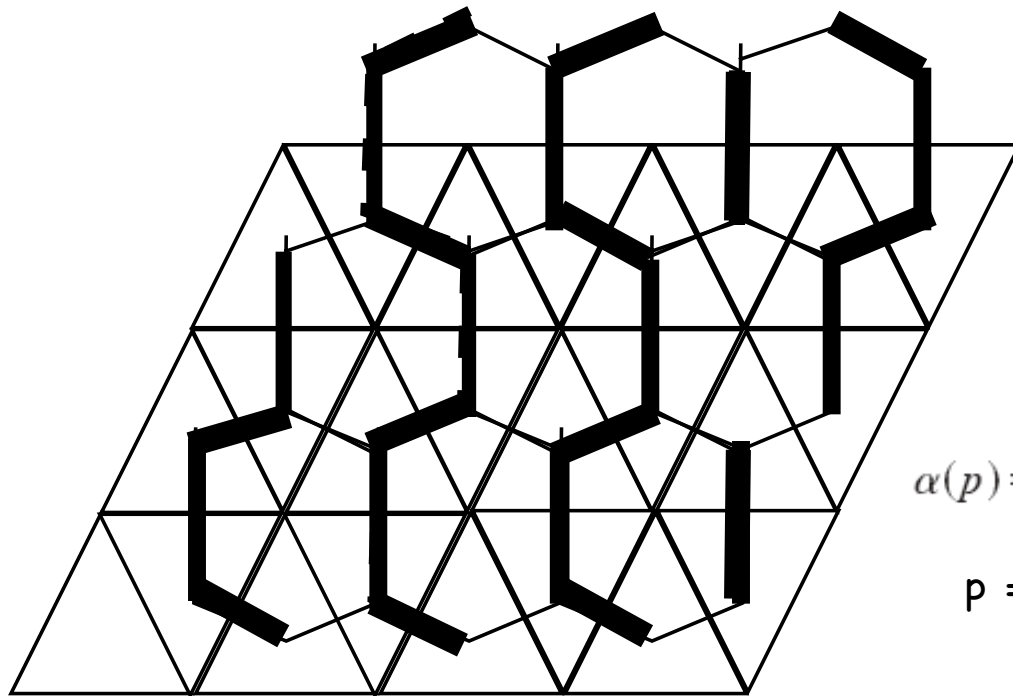
Think in terms of the inherent structures of this model (setting J to infinity analogous to hard sphere/rod constraint)

Inherent structures are dimer coverings of the honeycomb lattice



 This one is the completely ordered state and has no local excitations

Conserved quantities: strings



The "green" state has 4 strings

S goes to zero linearly !

Can divide inherent structures into string sectors.

States belonging to same string sector are connected by local dynamics

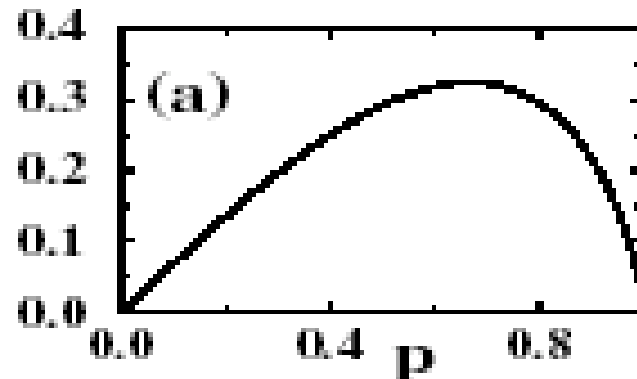
Count # of states in a string sector which gives us $S(\# \text{ of strings})$

$$S(p) = L^2 \alpha(p)$$

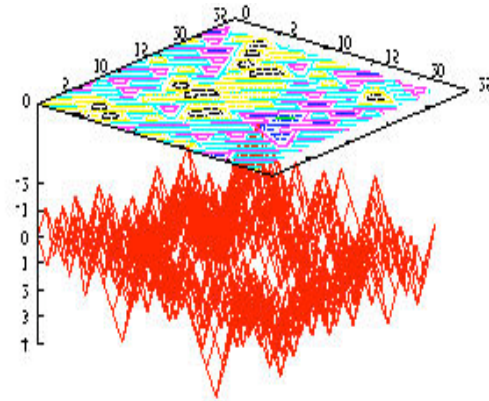
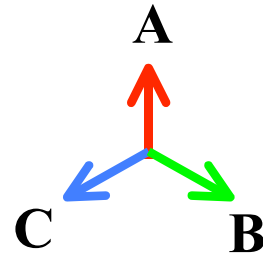
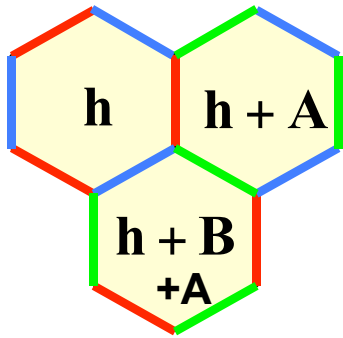
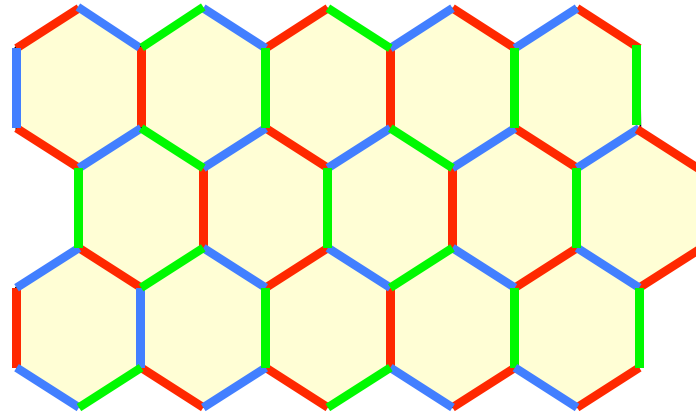
$$\alpha(p) = p \ln 2 + \frac{2}{\pi} \int_0^{\pi p/2} dx \ln[\cos(x)]$$

A. Dhar et al
, PRB 61,
6227 (1999)

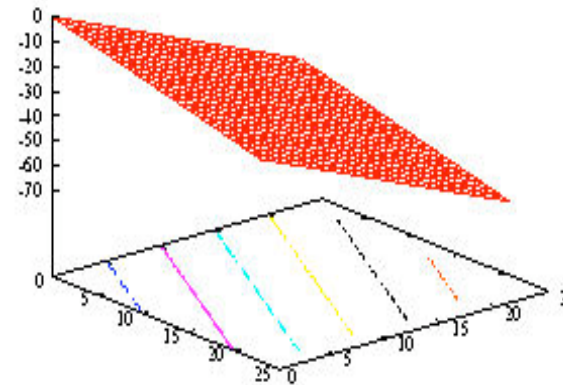
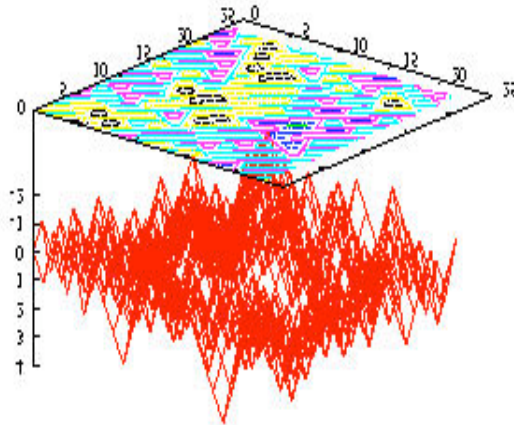
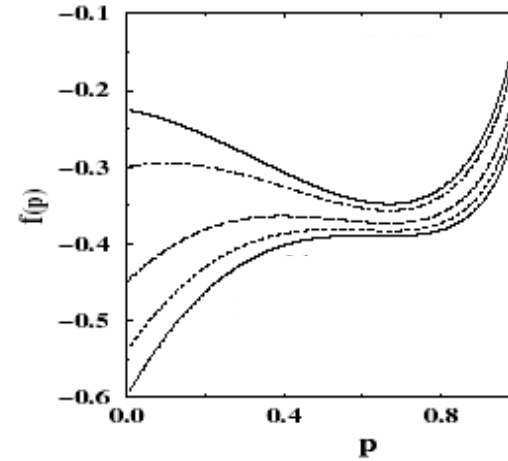
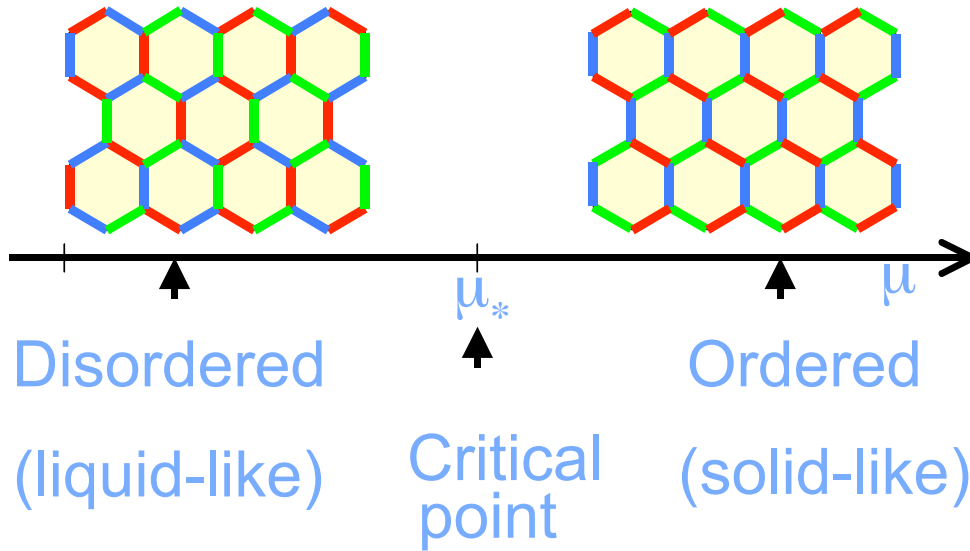
p = # of strings/unit length



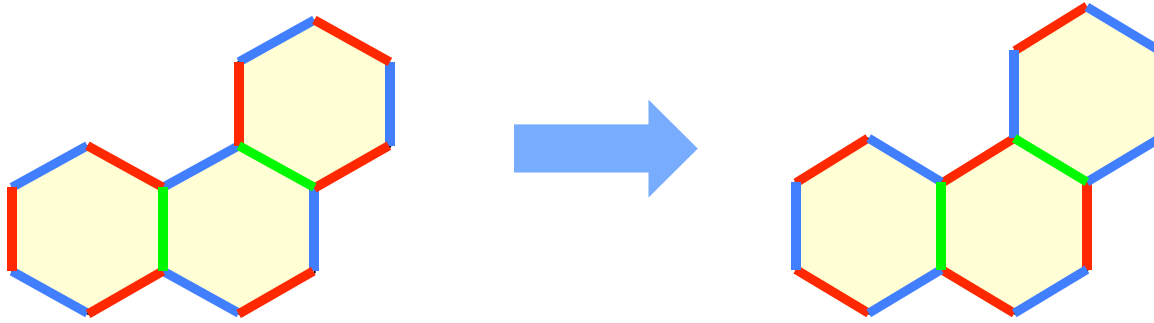
Three-coloring model



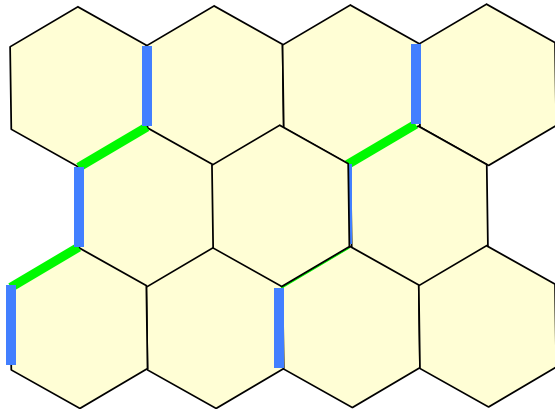
Equilibrium transition



Constraints \rightarrow Loop dynamics



Loops of all sizes get updated



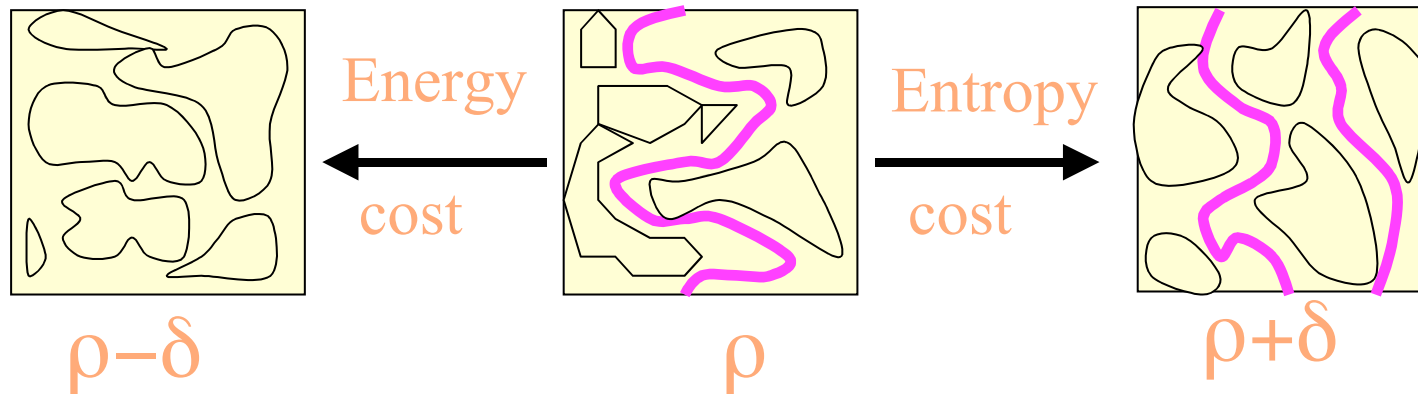
Loops with non-zero winding numbers

Necessary to update these to change the tilt

Transition Rates

$$W_{L_\rho \rightarrow L_{\rho-1}} \sim \exp[-\Delta E]$$

$$W_{L_{\rho-1} \rightarrow L_\rho} \sim \exp[-\Delta S]$$



Master Equation

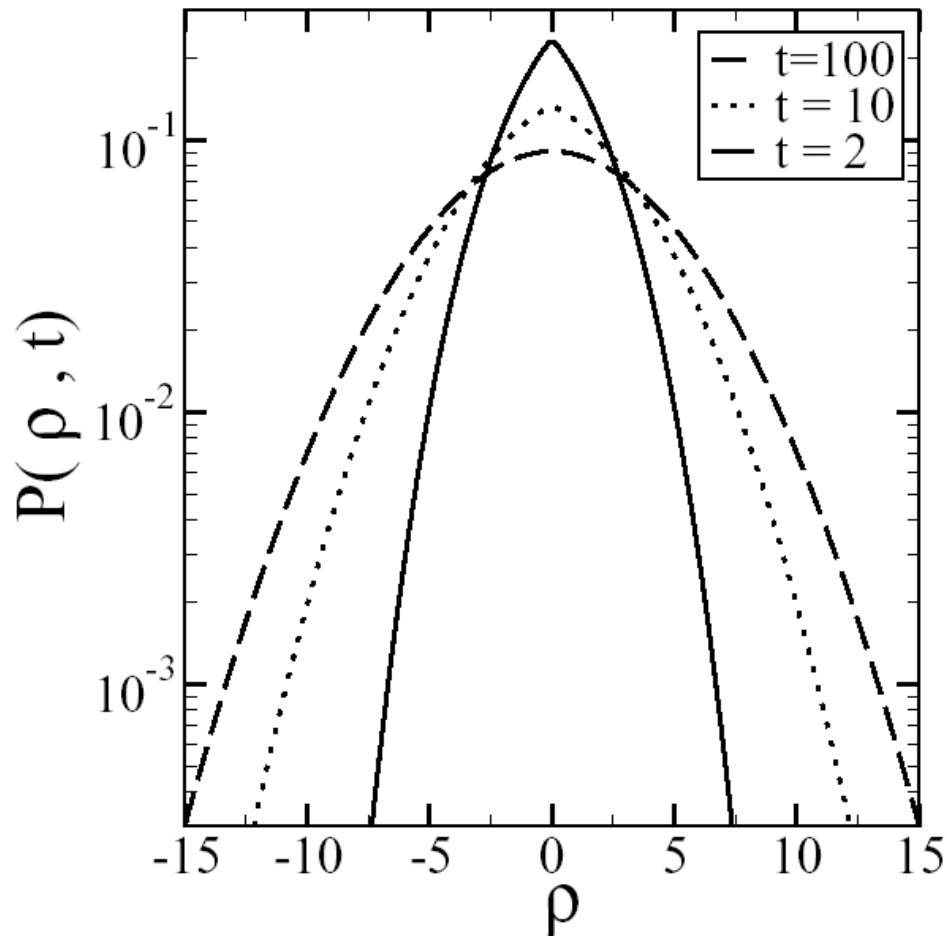
$$\frac{\partial P(\rho)}{\partial t} = -(W_{\rho \rightarrow \rho-1} + W_{\rho \rightarrow \rho+1})P(\rho) + W_{\rho-1, \rho}P(\rho-1) + W_{\rho+1 \rightarrow \rho}P(\rho+1)$$

Continuum limit

$$\frac{\partial P}{\partial t} = \frac{\partial}{\partial \rho} [\epsilon \rho e^{-\mu^* |\rho|} P + 2D e^{-\mu^* |\rho|} \frac{\partial P}{\partial \rho}]$$

Without these factors, this equation would lead to usual Langevin equation for ρ

$P(\rho, t)$ away from critical point



$l_1 \simeq \log t / \mu^*$ distance from critical point

$$l_2 = \sqrt{\frac{2D}{\lambda}}$$

Crossover time:

$$\tau_c = e^{2\mu^*} \sqrt{D/\lambda}$$

Also shows log-normal distribution of hopping times