

# Soft Motors: molecular to macroscopic perspectives

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# Outline: microscopic

- What is a Motor?
- Molecular Motors 1. Anomalous Photoalignment
  - ideas of Landauer, Prost & Astumian
- Molecular Motors 2. Yokoyama - Tabe Experiment
  - Lehman effect
- Molecular Motors 3. ATP synthase
  - Boyer & Walker
- Molecular Motors 4. Myosin and Actin
  - muscles
- Molecular Motors 5. LC elastomers
  - elongation/contraction
- Broken symmetry
  - broken symmetry drives motors



# Outline: macroscopic

- molecular & macroscopic length scales
- shape change
- motors based on elongation
- motors based on bend
- summary

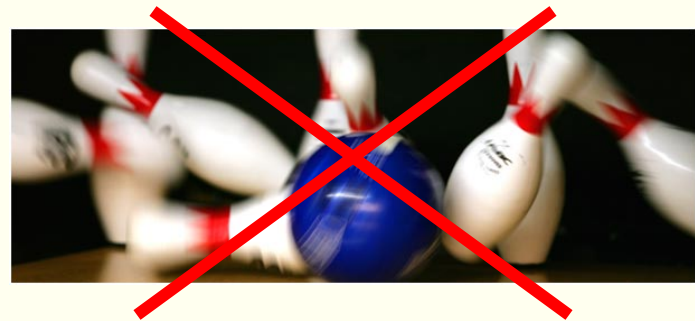


# what is a motor?



# Motor

- what is a motor?
  - anything that produces or imparts motion (*dictionary.com*)



- a machine that supplies motive power for a vehicle or other device (*O.E.D.*)



# Motor

- what is a motor?
  - a device which uses energy (but not momentum!) to cause motion



- how does motion come about?
- car convinces the road to exert a force on it, and push it forward



# Motor

- what is a motor?
  - a device which uses energy (but not momentum!) to cause motion

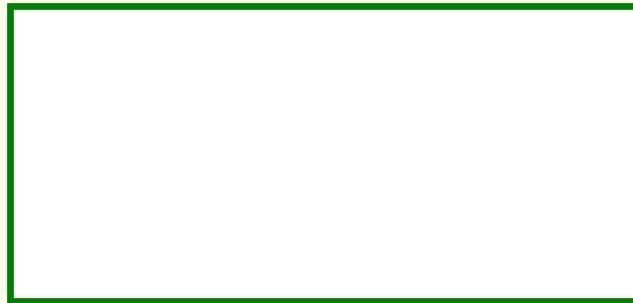


- motor causes one part of a system to exert a force/torque on another, causing motion



# Motor

- what is a motor?
  - a device which uses energy (but not momentum!) to cause motion



- motor causes one part of a system to exert a force/torque on another, causing motion

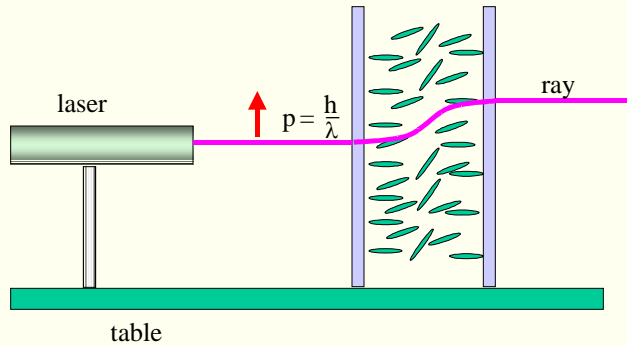




# Molecular Motors 1. Anomalous Photoalignment



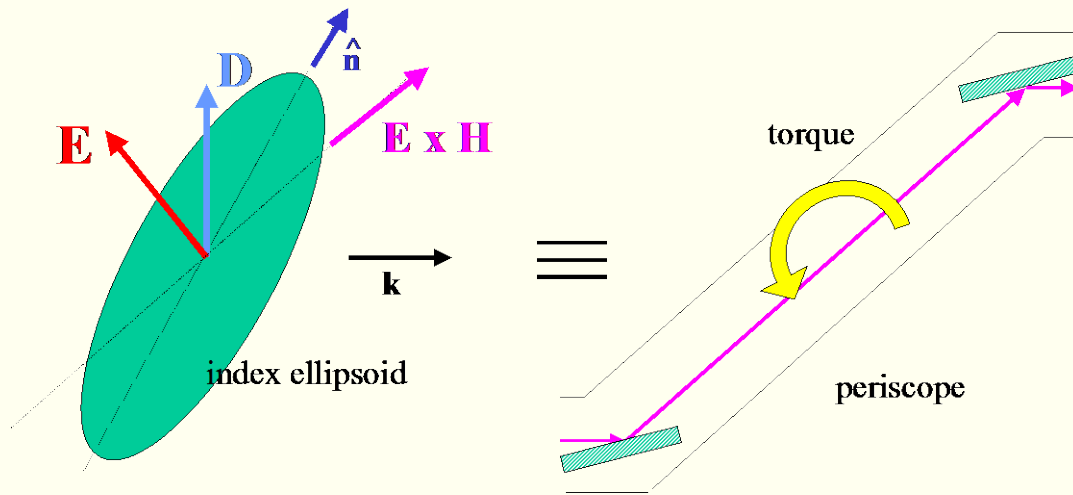
# Direct optical torque



$$\tau_{vol}^{opt} = \mathbf{D} \times \mathbf{E}$$

$$\tau_{vol}^{el} + \tau_{vol}^{opt} = 0$$

angular momentum transfer:



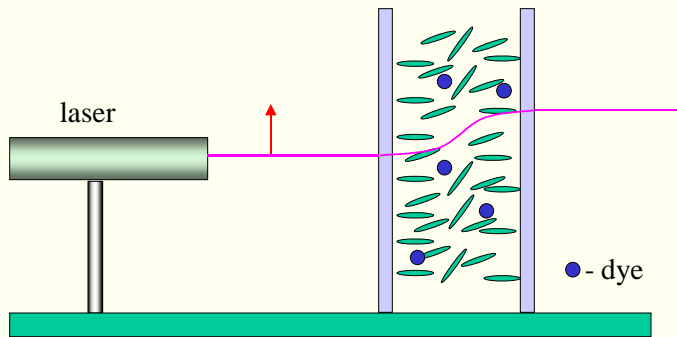
torque results from change in extrinsic angular momentum of light.



# Indirect optical torque

addition of 1% of anthroquinone dye

⇒ reduction of threshold intensity<sup>1</sup> by  $\sim \times 100!$



without dye:  $\tau_{vol}^{el} + \tau_{vol}^{opt} = 0$

with dye:  $\tau_{vol}^{el} + \frac{1}{100} \tau_{el}^{opt} + ?? = 0$

Torque on nematic causing elastic deformation  
CANNOT come from light! Source of torque??

1. I. Janossy, A.DD. Lloyd and B.S. Wherret, *Mol. Cryst.Liq. Cryst.* **179**, 1, 1990.  
I. Janossy, *Phys. Rev. E* **49**, 2957, 1994.



# Puzzle 1:

- light causes the director to reorient, against a restoring elastic torque, essentially without the transfer of angular momentum.
- light causes rotation without exerting a torque!

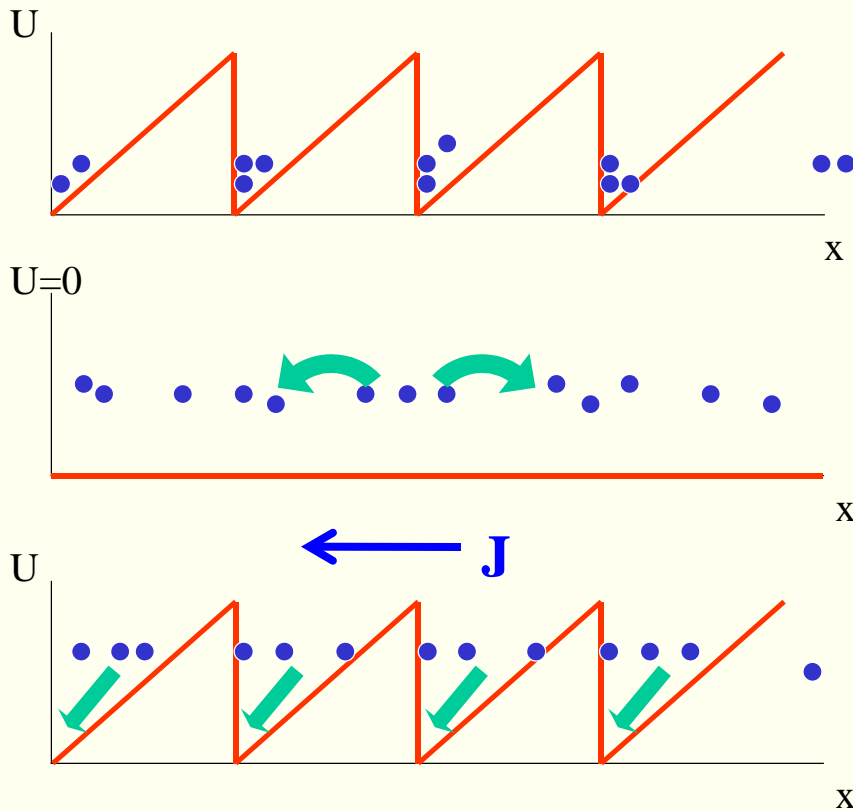
???



# Prost, Astumian



# The translational ratchet\*



Particles in asymmetric periodic potential:

When potential is turned OFF, particles diffuse, no net current

When potential is turned ON, particles move towards minimum, giving rise to current.

\* R.D. Astumian et al: *Phys.Rev.Lett.***72**, 1766, 1994.  
J. Prost et al: *Phys. Rev. Lett.* **72**, 2652, 1994.



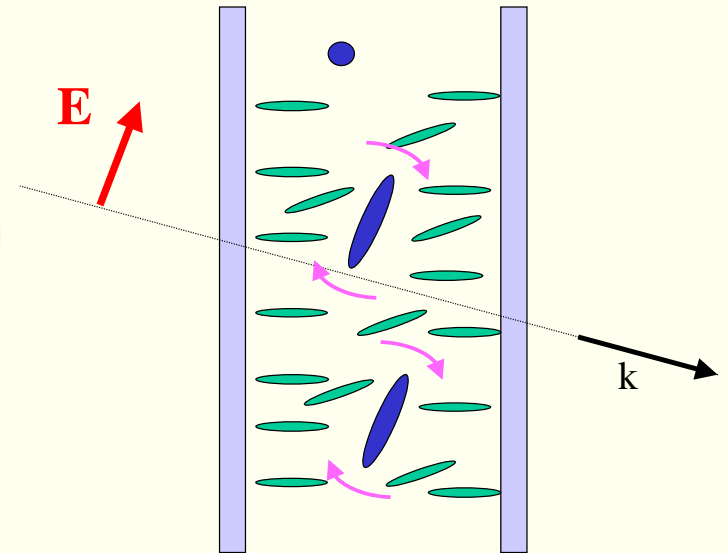
# Simple model:

- dye

- anisotropic molecule with orientation  $\hat{l}_d$
- in ground state does not interact with nematic
- in excited state becomes 'nematic' molecule
- probability of excitation depends on  $(\hat{l}_d \cdot \mathbf{E})^2$
- lifetime of excited state is  $t_o$

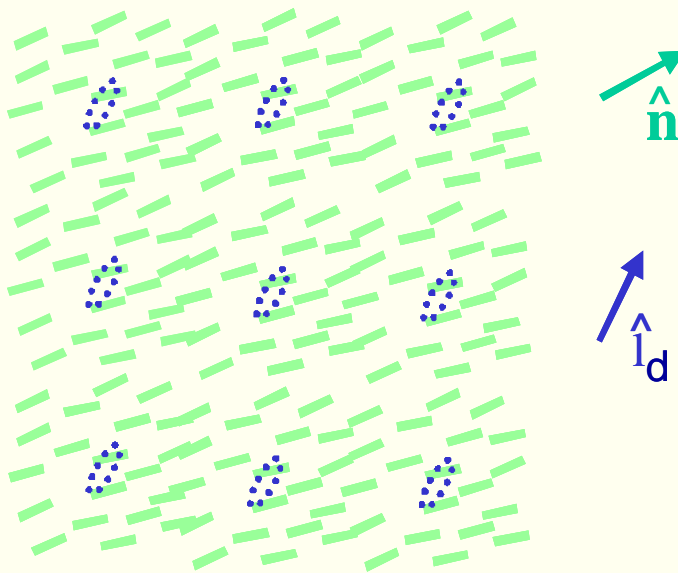
- optical field

- transfers energy, but no momentum to the sample



# Torque on liquid crystal:

Average orientation of excited dye molecule is  
NOT along nematic director



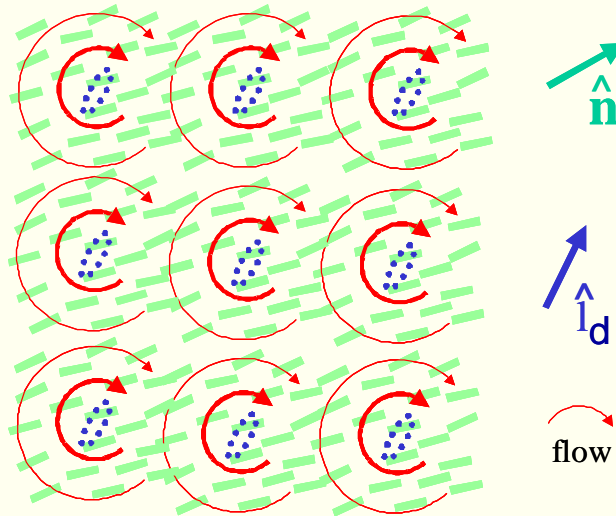
interaction between liquid crystal  
and excited dye gives rise to torque!





# Rotating dye gives rise to shear flow

- dye: - rotation  $\Rightarrow$  source of vorticity
- liquid crystal: - aligned by nematic field of dye



viscous shear carries angular momentum from cell walls to dye

angular momentum current: dye  $\xrightarrow{\text{molecular}}$  nematic  $\xrightarrow{\text{elastic}}$  cell walls  $\xrightarrow{\text{shear}}$  dye



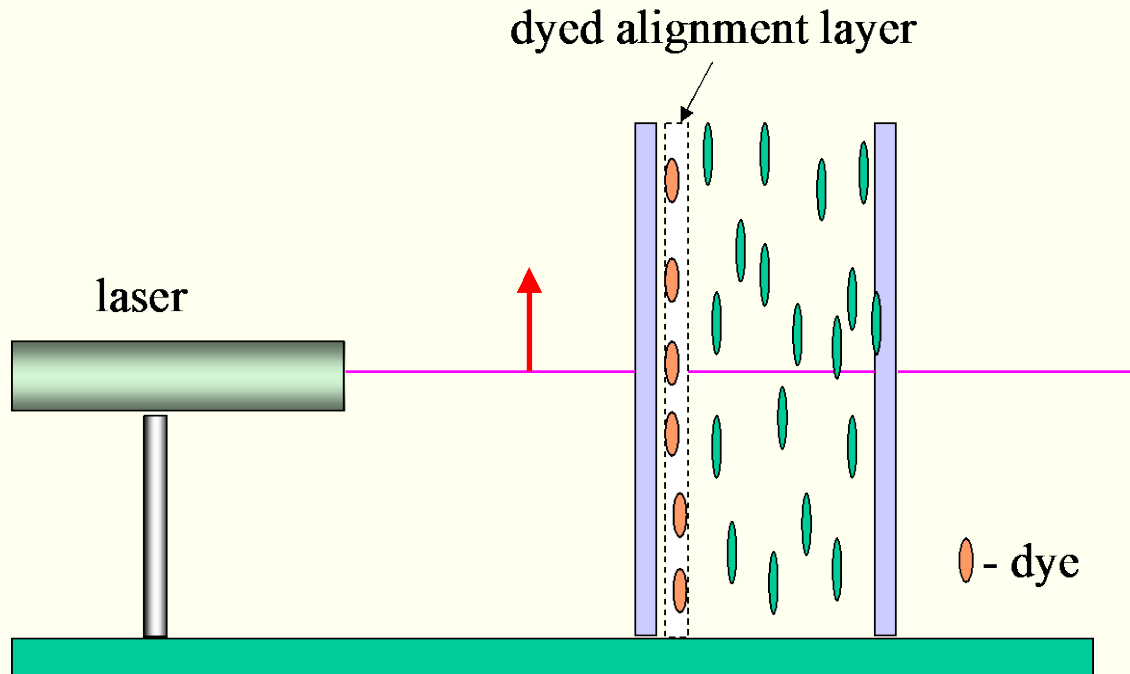
# Summary of Puzzle 1

- Janossy effect<sup>1</sup> (source provides energy but not torque yet causes rotation)  $\Rightarrow$  **orientational ratchet**<sup>2</sup>
- dye is light-driven rotor with continuous rotation
- energy from optical field drives shear flow
- viscous stress carries angular momentum from dye to cell walls

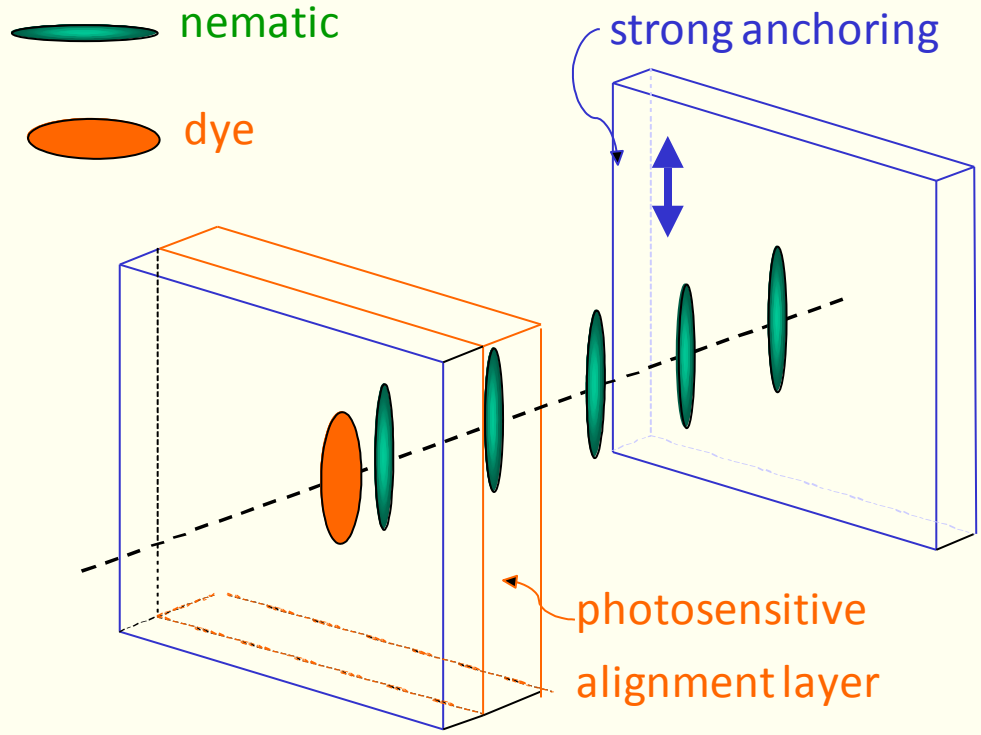
1. T. Kosa and I. Janossy, *Opt.Lett.* **17**, 1183, 1992; *Opt. Lett.* **20**, 1231, 1995.
2. W. E and P. Palffy-Muhoray, *Mol.Cryst.Liq.Cryst.* **320**, 193, 1998



# Puzzle 2. Photoinduced Twist



# Photoinduced Twist

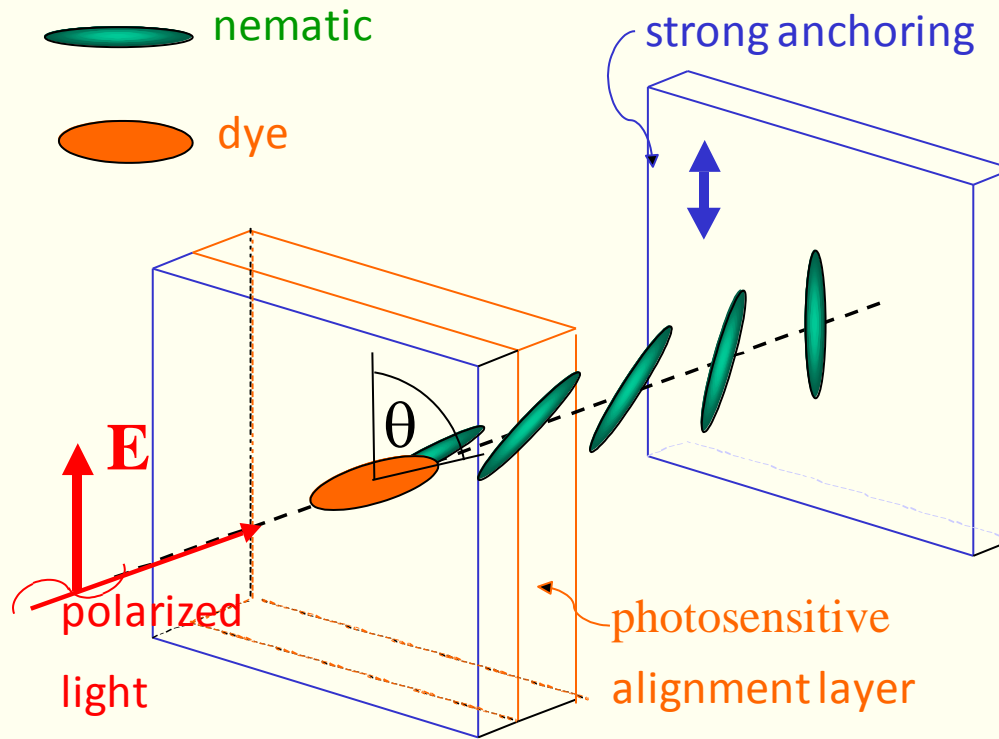


laser OFF



# Photoinduced Twist

laser ON



# Origin of torque

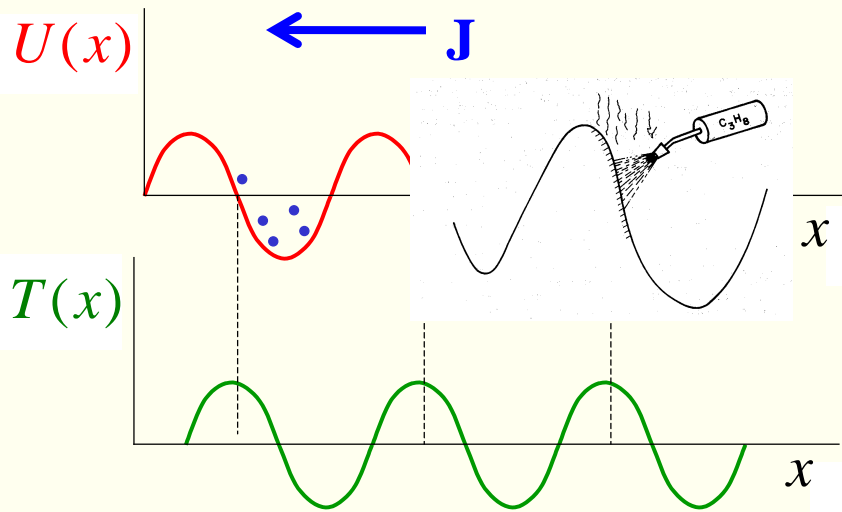
- before deformation takes place, polarization is parallel to the director, hence optical torque  $\mathbf{D} \times \mathbf{E} = 0$ .
- optical field *stabilizes* configuration!
- source of torque?



# Landauer



# Landauer's Blowtorch



particles in periodic potential

'hot' molecules are more excited  
& diffuse over barrier

$\therefore$  steady current: system is molecular motor (pump)

one part of system pushes on other;  
*no* momentum transfer from outside.

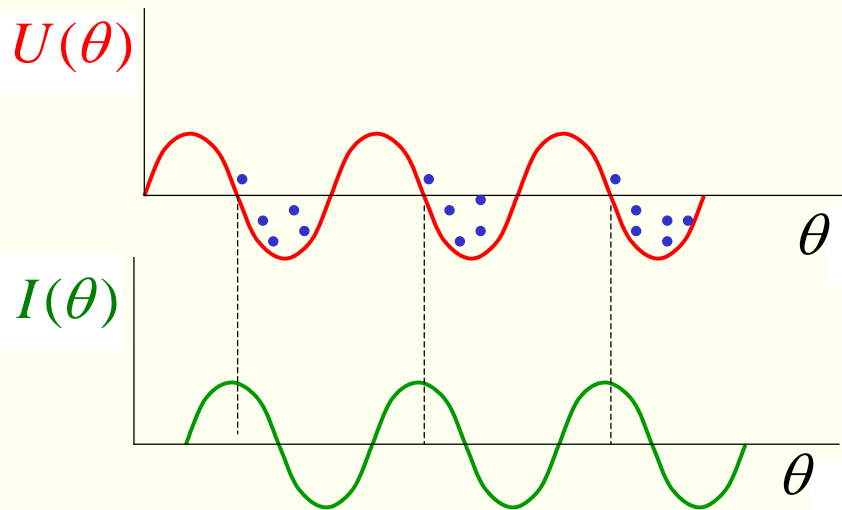
M. Büttiker, *Z. Phys. B* **68**, 161 (1987).

R. Landauer, *J. Stat. Phys.* **53**, 233 (1988)





# Indirect optical torque: Brownian ratchet



dye molecules in nematic field

dye molecules parallel to pump polarization are excited & diffuse over barrier

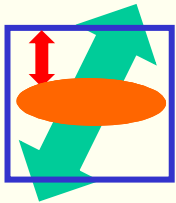
$\therefore$  steady rotation: dye is light driven molecular motor

torque results from rotation of dye;  
*no* momentum transfer from light.

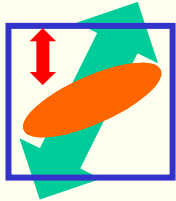
P. Palffy-Muhoray, T. Kosa and Weinan E, *Appl. Phys. A* **75**, 294 (2002)



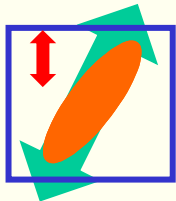
# Schematic of Brownian motor:



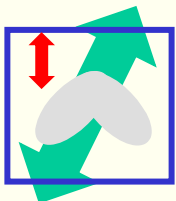
dye in *trans*-state starts to rotate towards director



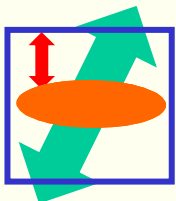
it exerts torque on director & vice versa



as it becomes parallel to pump polarization



it is excited into *cis*-state, undergoes diffusion



relaxes into *trans*-state, rotates towards director...



# Dynamics

- dye: Fokker-Planck

$$\frac{\partial \rho_t}{\partial t} = D_t \nabla^2 \rho_t + \nabla \cdot (M_t \rho_t \nabla U_t) - \rho_t f_t + \rho_c f_c$$

$$\frac{\partial \rho_c}{\partial t} = D_c \nabla^2 \rho_c + \nabla \cdot (M_c \rho_c \nabla U_c) - \rho_c f_c + \rho_t f_t$$

- transition rates

$$f_t = f_{to} e^{U_t/kT} + v e_t^2 (\hat{\mathbf{I}}_d \cdot \hat{\mathbf{E}})^2$$

$$f_c = f_{co} e^{U_c/kT} + v e_c^2 (\hat{\mathbf{I}}_d \cdot \hat{\mathbf{E}})^2$$



# Dynamics

- liquid crystal
  - bulk:

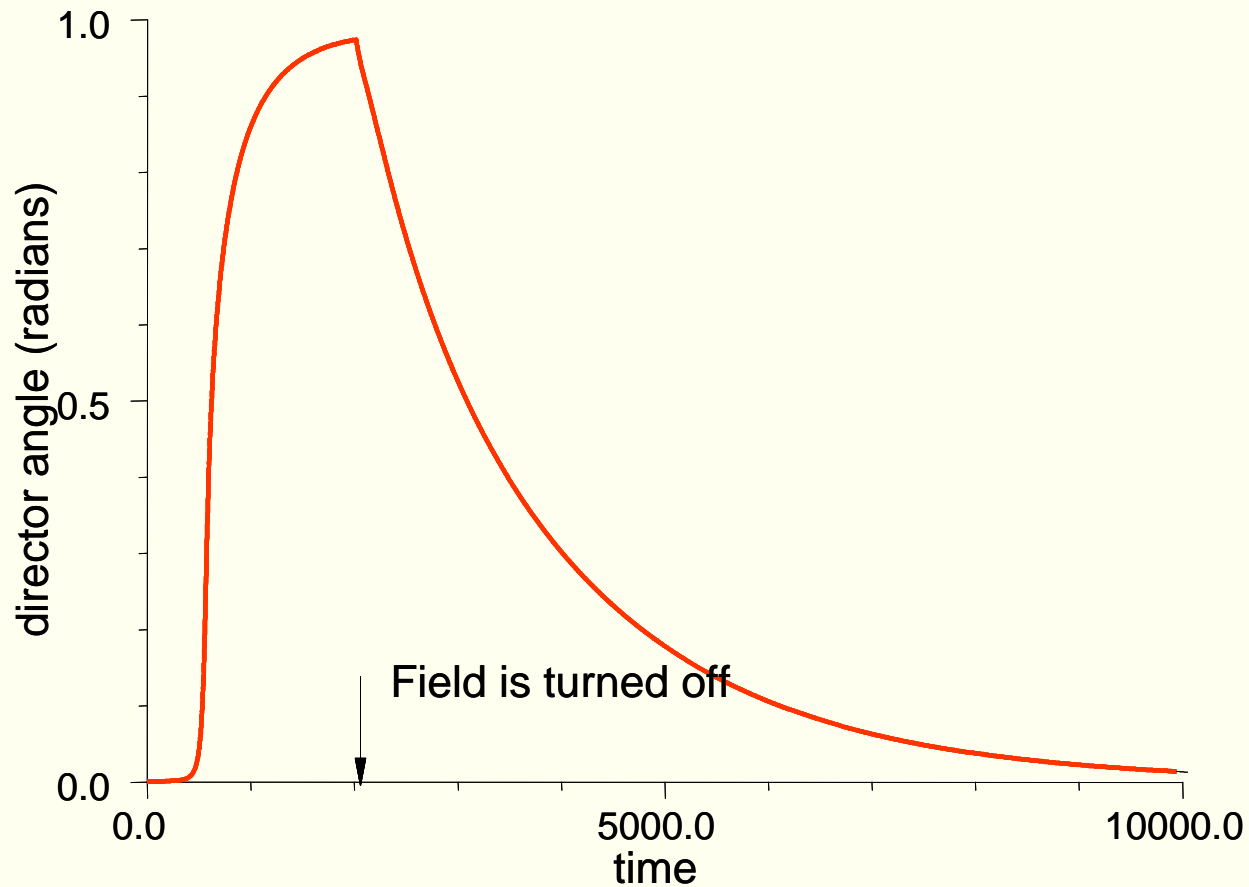
$$\gamma \frac{\partial \hat{\mathbf{n}}}{\partial t} = \mathbf{K} \nabla^2 \hat{\mathbf{n}} (I - \hat{\mathbf{n}} \hat{\mathbf{n}})$$

- surface:

$$\gamma_s \frac{\partial \hat{\mathbf{n}}}{\partial t} = \left\{ \mathbf{K} (\hat{\mathbf{N}} \nabla \cdot \hat{\mathbf{n}} + \hat{\mathbf{N}} \times \nabla \times \hat{\mathbf{n}}) - d_d \left\langle \rho_t \frac{\partial U_t}{\partial \hat{\mathbf{n}}} \right\rangle - d_d \left\langle \rho_c \frac{\partial U_c}{\partial \hat{\mathbf{n}}} \right\rangle \right\} (I - \hat{\mathbf{n}} \hat{\mathbf{n}})$$



# Simulations



- optical field drives orientational current
- dye exerts steady torque on nematic



# Summary of Anomalous Photoalignment

- light-driven molecular motors
- dye\* molecules are rotors in nematic potential
- energy (but not momentum) from light
- steady unidirectional rotation
- orientational ratchet / Landauer's blowtorch mechanism
- work: steady torque on director by viscous stress + dissipation
- achiral



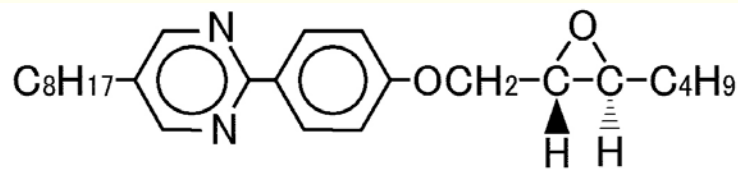
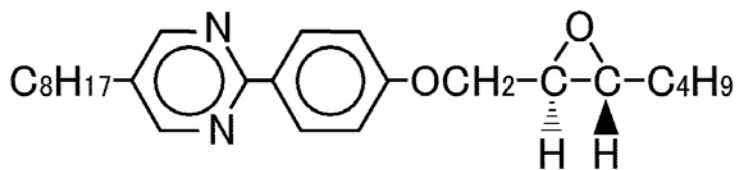
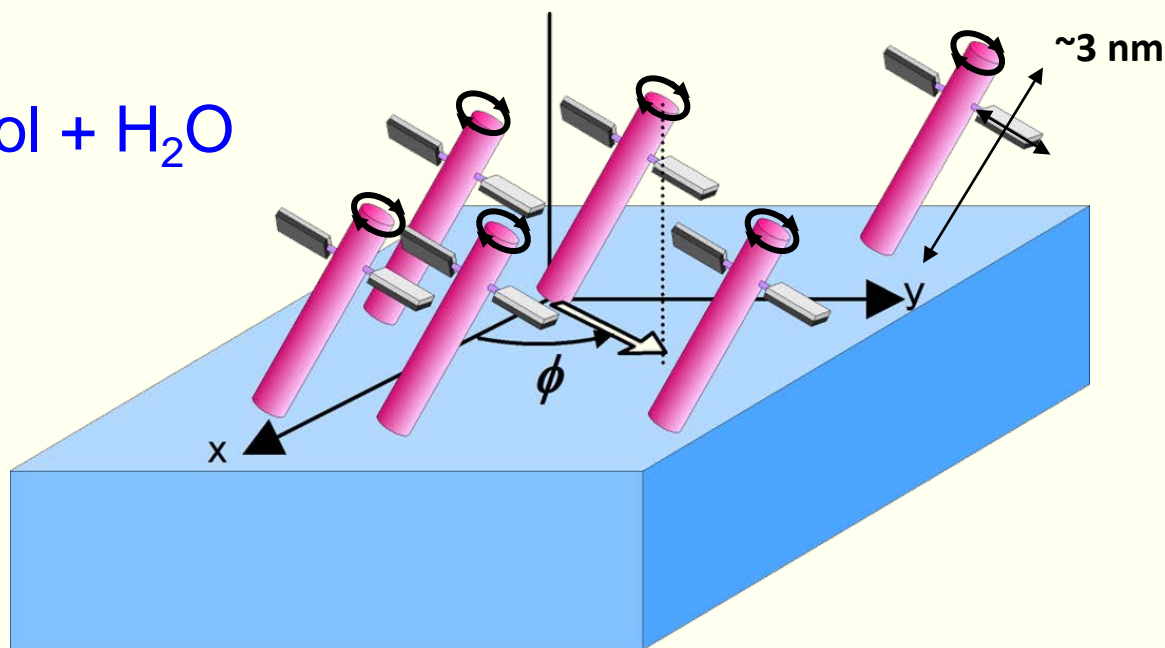
# Molecular Motors 2.

## Yokoyama - Tabe Experiment



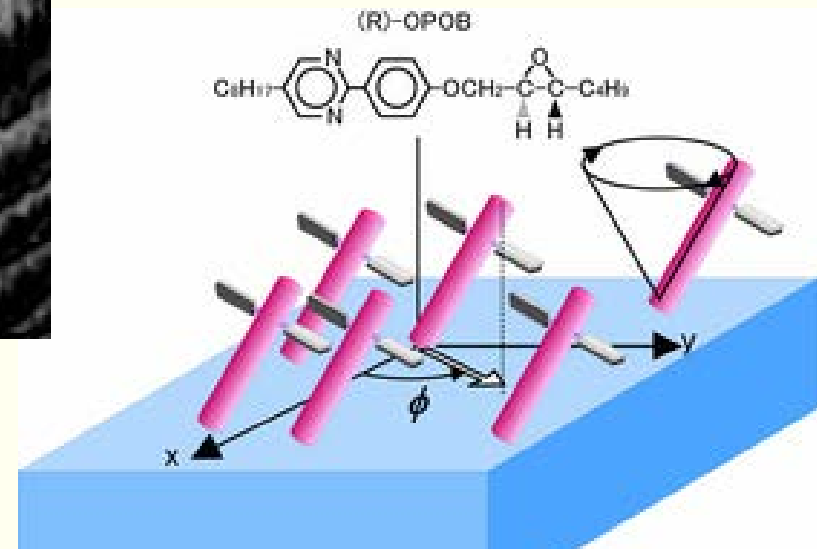
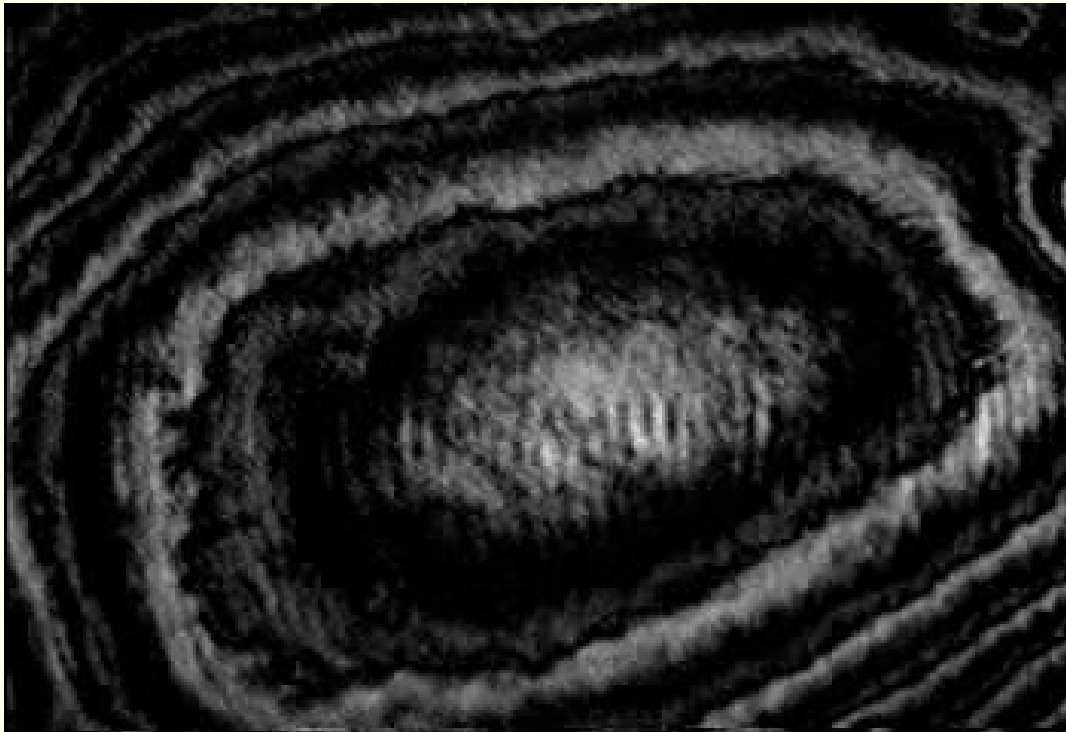
# Yokoyama – Tabe Experiment

- chiral Langmuir monolayer
- substrate: glycerol + H<sub>2</sub>O





# Collective Molecular Rotor/Pump

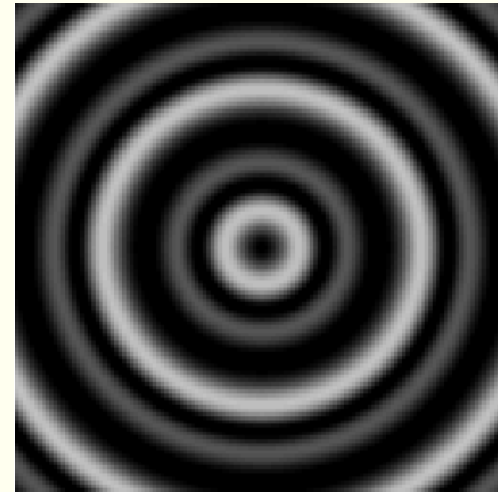
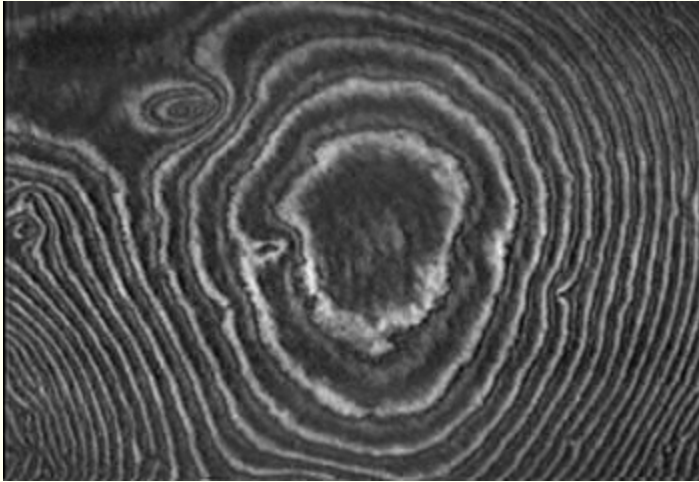


Yuka Tabe and H. Yokoyama, *Nature Materials*, 2, 806(2003).



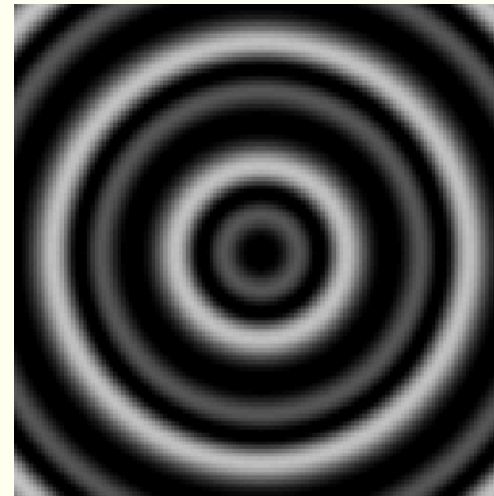
## Chiral LC molecules on glycerol

FELIX013 on  
pure glycerol  
(SmC\* phase)



counter-  
clockwise  
rotation

CS4001 on pure  
glycerol  
(SmCA\* phase)



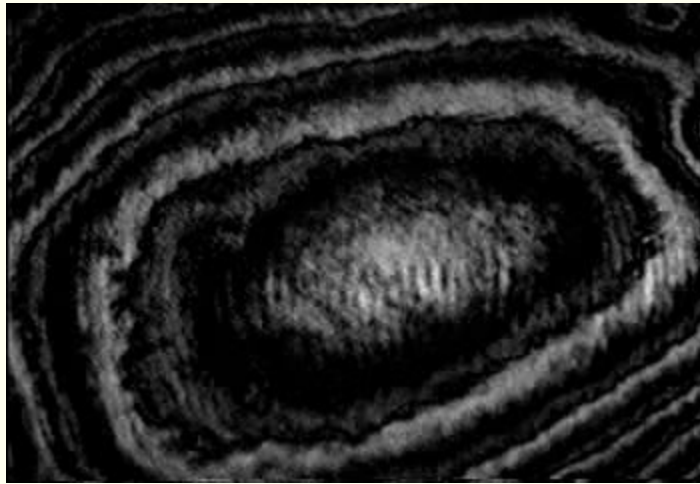
clockwise  
rotation



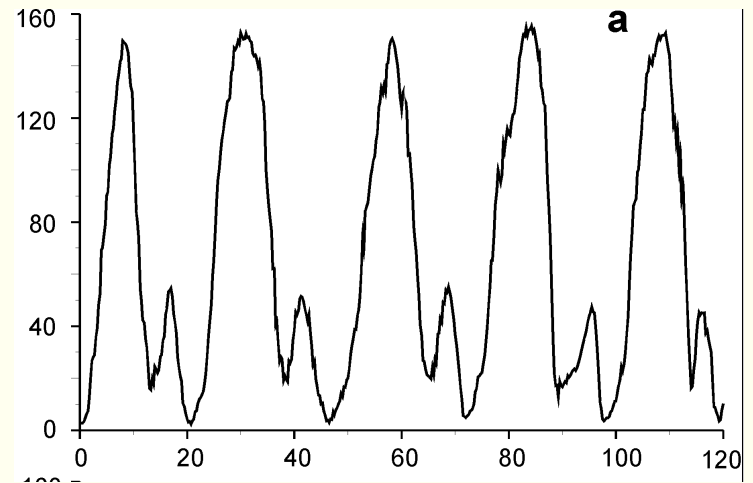
# Chirality inversion

(R)-OPOB

on pure glycerol

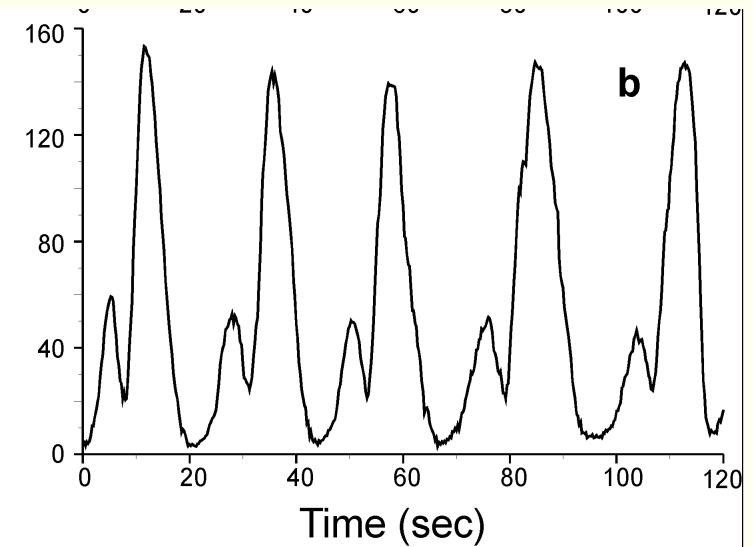
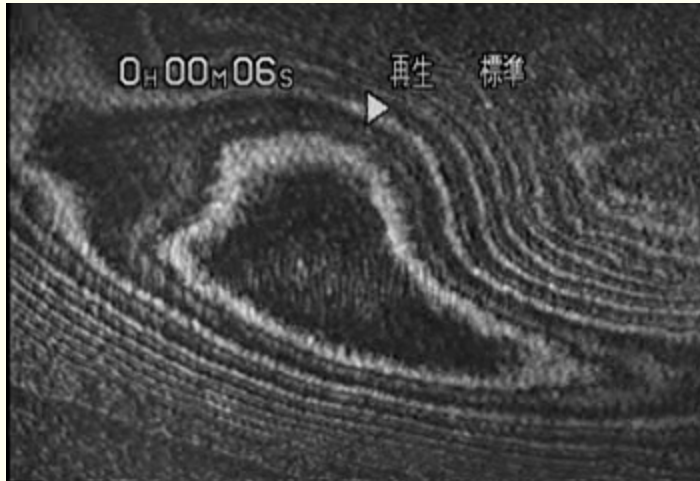


100 $\mu$ m

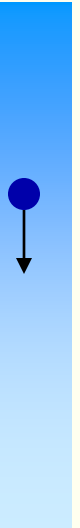
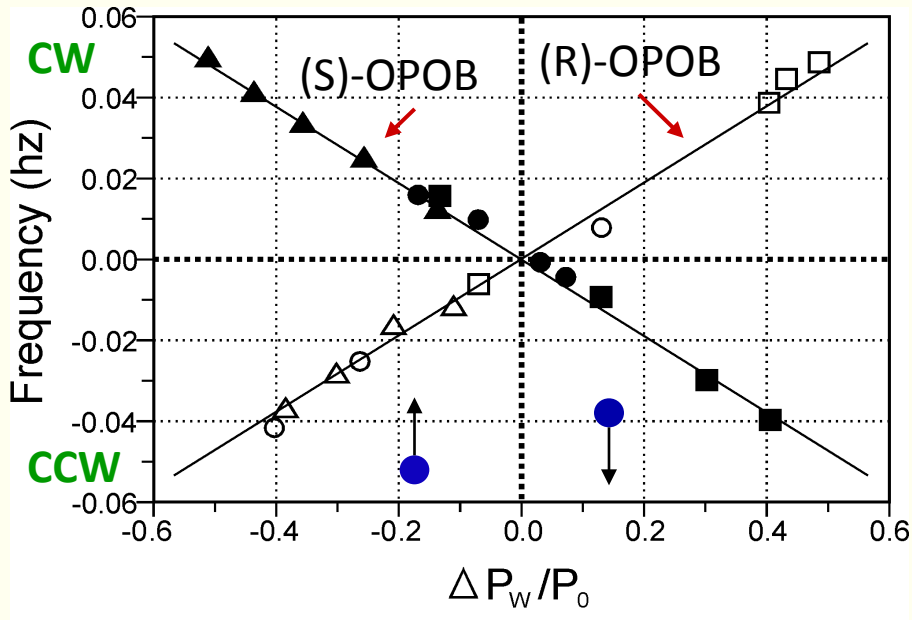


(S)-OPOB

on pure glycerol



# Precession speed linearly depends on $\Delta P_w/P_0$



$$\Delta P_w/P_0 = P_v - P_s$$

$P_v$ : actual water vapor pressure

$P_s$ : saturated water vapor pressure



# Onsager Reciprocal Relations

Entropy Production:

$$T \dot{S} = \tau \frac{\partial \phi}{\partial t} + J \cdot \nu_m \Delta P$$

Onsager Relations:

$$\frac{d\phi}{dt} = \frac{1}{\gamma} \tau + b \nu_m \Delta P$$

$$J = b \tau + \frac{1}{\eta} \nu_m \Delta P$$

$\nu_m$ : molecular volume of water in vapor

$\gamma$ : rotational viscosity of LC director

$\eta$ : water mobility through the film

$b$ : Lehmann coefficient ( $\propto$  chirality strength)



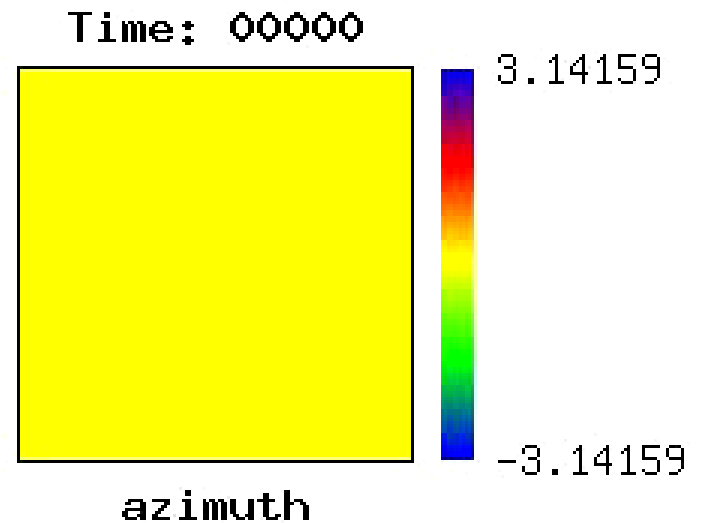
# Phenomenological Theory of Lehmann effect

- Shibata & Mikhailov, Europhys. Lett., 73 (2006)
- H. Brand et al., Phys. Rev. Lett. 96 (2006)
- Tsori & de Gennes, Euro.Phys. J. E 14 (2004)
- T. Okuzono, Kyoto Workshop (2005)

Equation of motion :

$$\gamma \frac{\partial \mathbf{c}}{\partial t} = - \frac{\delta F}{\delta \mathbf{c}} + \underbrace{v \mathbf{c} \times \frac{\partial P}{\partial \mathbf{z}}}_{\substack{\uparrow \\ \text{Lehman term}}}$$

$$F = \int dr \left[ \frac{K}{2} \sum_{i=x,y} |\nabla c_i|^2 - \frac{\tau}{2} |\mathbf{c}|^2 + \frac{u}{4} |\mathbf{c}|^4 \right]$$



# Summary of Yokoyama – Tabe experiment

- substrate & monolayer system is motor
- energy stored in chemical potential of H<sub>2</sub>O
- vapor current drives rotation of chiral molecules
- work: produces circulating flow, dissipation



# Molecular Motors 3. ATP Synthase





# F<sub>1</sub>F<sub>0</sub> ATP Synthase

- enzyme that synthesizes  
adenosine-5'-triphosphate (ATP)  
from  
adenosine diphosphate (ADP) + phosphate

- reaction:

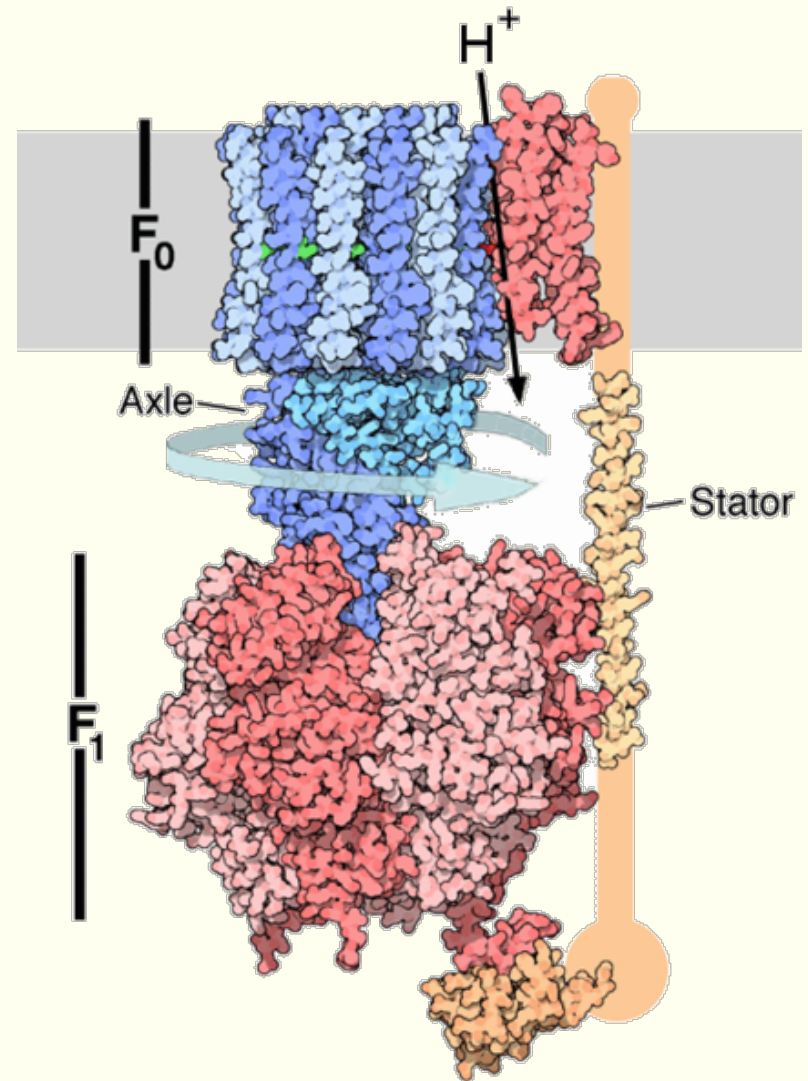
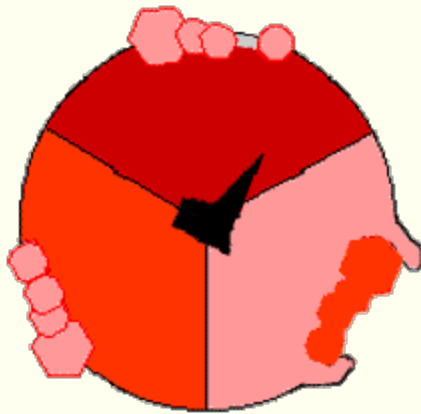


- ATP:  
"molecular unit of currency" of intracellular energy transfer

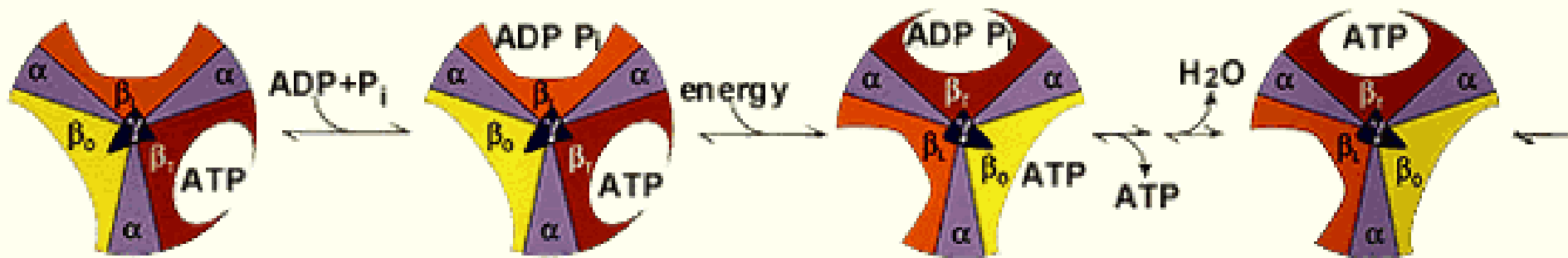


# F<sub>1</sub>F<sub>0</sub> ATP Synthase

- energy from glucose sets up H<sup>+</sup> gradient
- H<sup>+</sup> current drives motor
- stator opens & closes, producing ATP



# F<sub>1</sub>F<sub>0</sub> ATP Synthase



A complete ATP molecule is bound to  $\beta_T$ .

$\beta_L$  binds ADP and inorganic phosphate ( $P_i$ ).

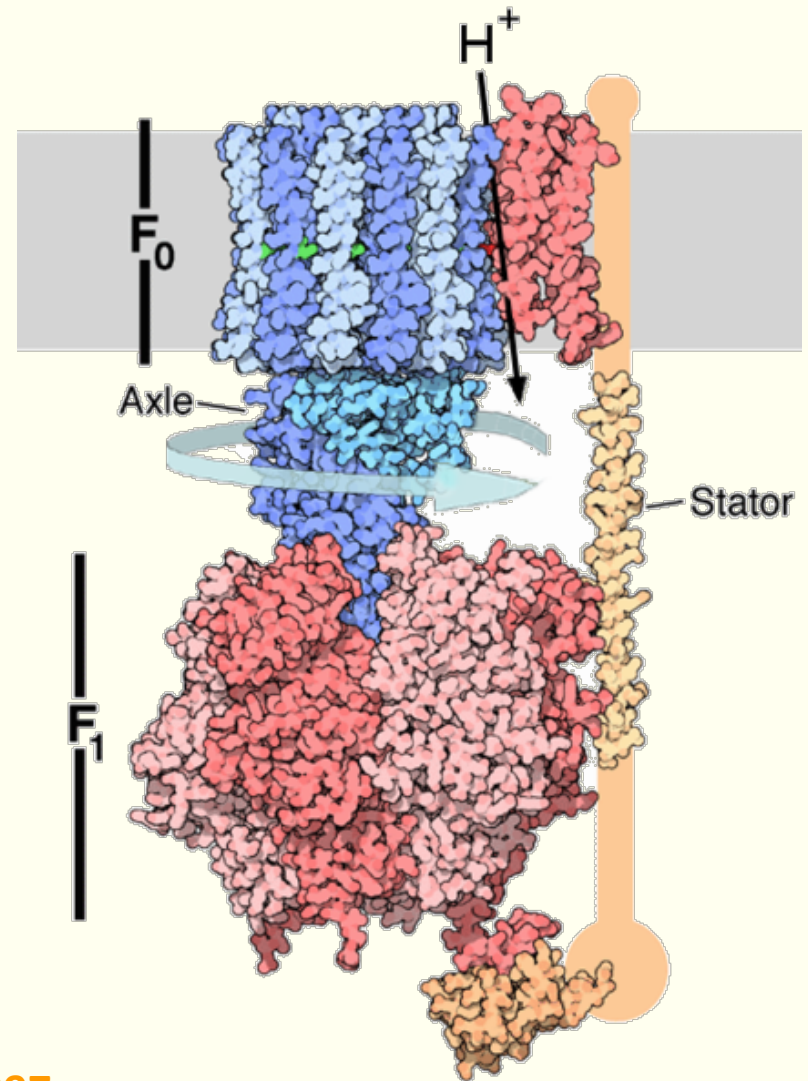
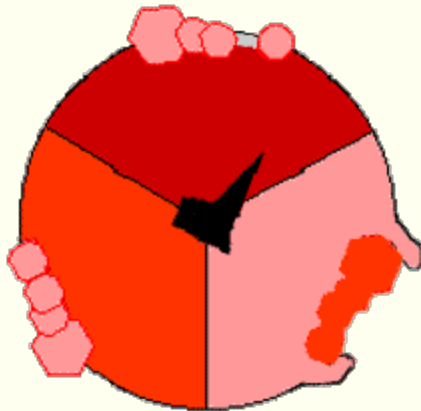
The  $\gamma$  subunit has rotated.  $\beta_T$  becomes open and ATP is released.  $\beta_L$  becomes tight and  $\beta_0$  becomes loose.

The phosphate ion reacts with the ADP molecule so that a new ATP molecule is formed. We are back to the first stage.



# F<sub>1</sub>F<sub>0</sub> ATP Synthase

- normal human produces ~100lbs of APT/day
- motor can run backwards

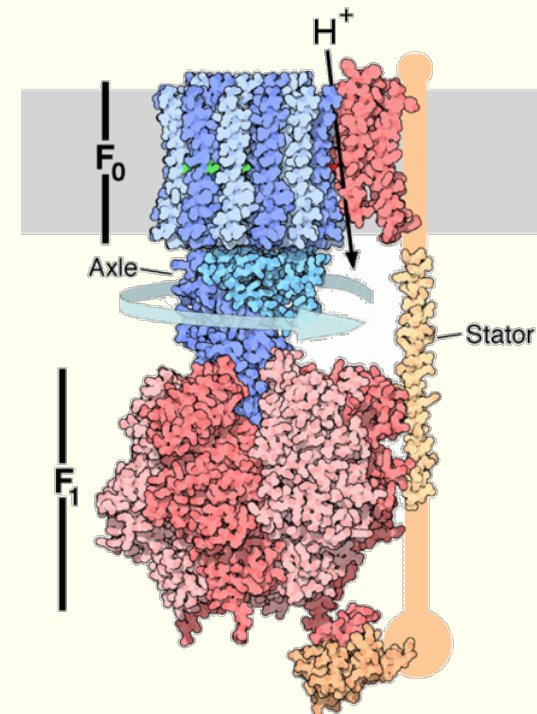
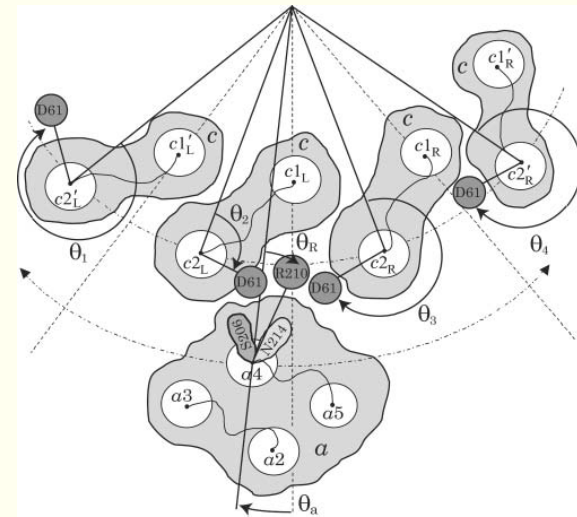


Boyer & Walker, Nobel Prize in Chemistry, 1997



# Why does rotor turn?

- sequential protonation & deprotonation of side groups coupled with rotation
- chemical potential difference of  $H^+$  drives current, which drags surface of C-rotor.
- pot. diff.  $\rightarrow$   
current  $\rightarrow$   
rotation.



A. Aksimentiev *et al.* Biophysical Journal **86**, 1332 (2004)



# Modeling

- system of Langevin equations

$$\xi_i \frac{d\theta_i}{dt} = - \frac{\partial \Psi(\theta_a, \theta_R, \theta_1, \theta_2, \theta_2, \theta_4)}{\partial \theta_i} + \eta_i(t),$$
$$i = a, R, 1, \dots, 4.$$

- Lennard-Jones + hydrophobic + screened Coulomb

$$U_{\text{EL}}(\vec{r}) = \frac{e^2}{4\pi\epsilon_0\epsilon} \frac{q_1 q_2}{|\vec{r}|} \exp(-\lambda|\vec{r}|);$$

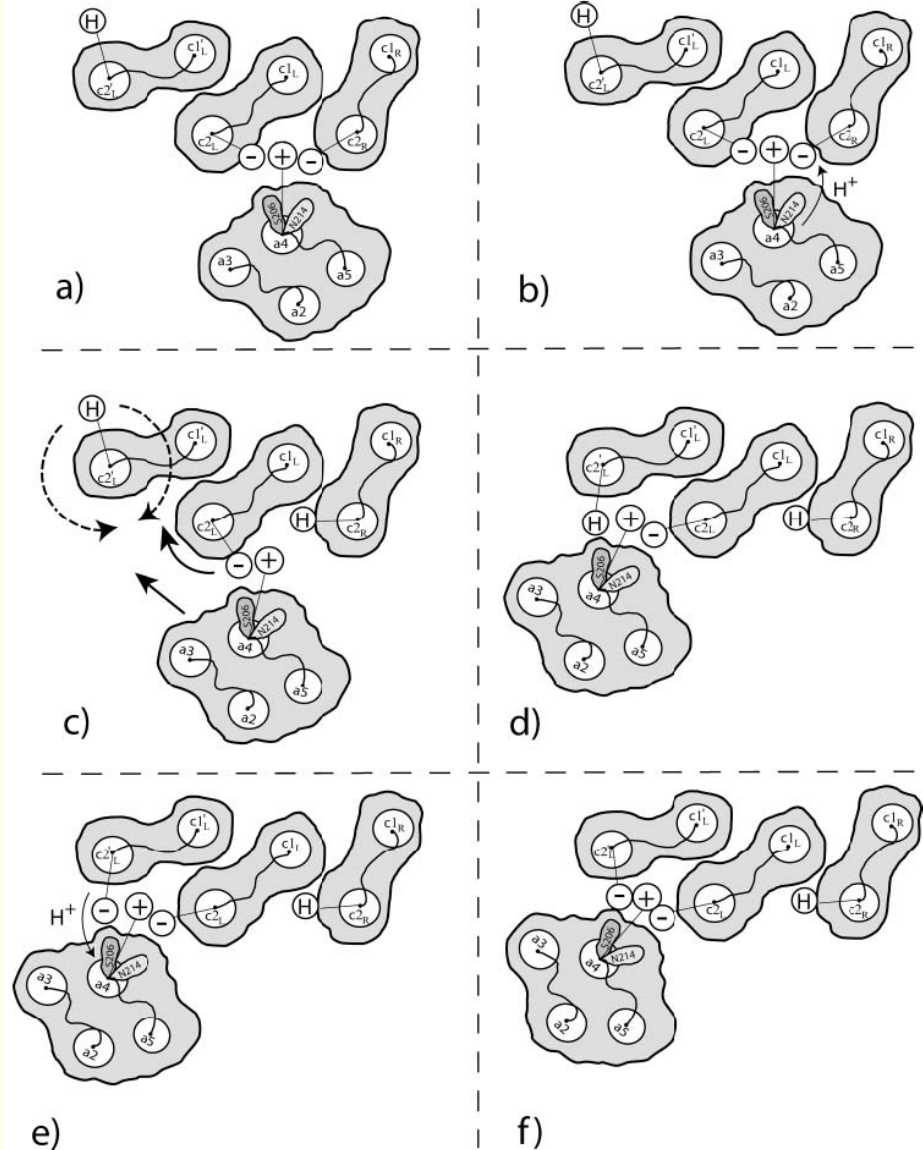
A. Aksimentiev *et al.* Biophysical Journal **86**, 1332 (2004)



# Modeling

- proposed sequence

based on simulations



A. Aksimentiev *et al.*  
Biophysical Journal **86**, 1332 (2004)



# Summary of ATP Synthase

- rotary molecular motor
- rotation driven by proton current
- energy from stored electrochemical potential
- work: creation of ATP + dissipation





# Molecular Motors 4. Actin & Myosin

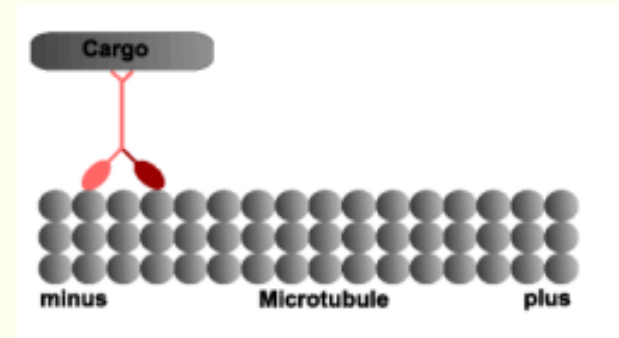


# Biological transport

- all organisms (eukaryotic cells of yeasts, plants, animals) contain 'motor proteins'
- two well known examples:

## – kinesins

- two active heads
- hydrolyses ATP
- moves along microtubules

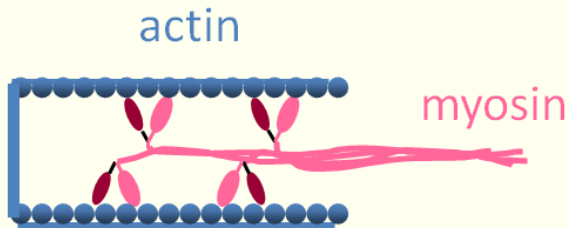


## – myosins

- two active heads
- hydrolyses ATP
- moves along actin fibers



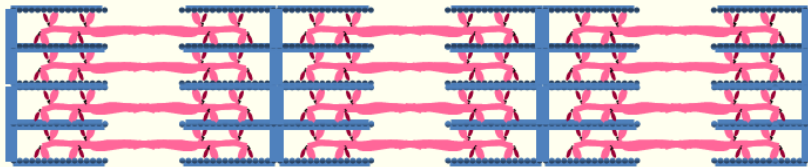
# Muscle force and movement



$F=5.3\text{pN}$  /cross-bridge  
step-size=5.5nm  
1 step/ATPase reaction



myosin filament length=1.55 $\mu\text{m}$



C.J. Pennycuick, *Newton Rules Biology* (Oxford, 1995)



# Summary of Actin & Myosin

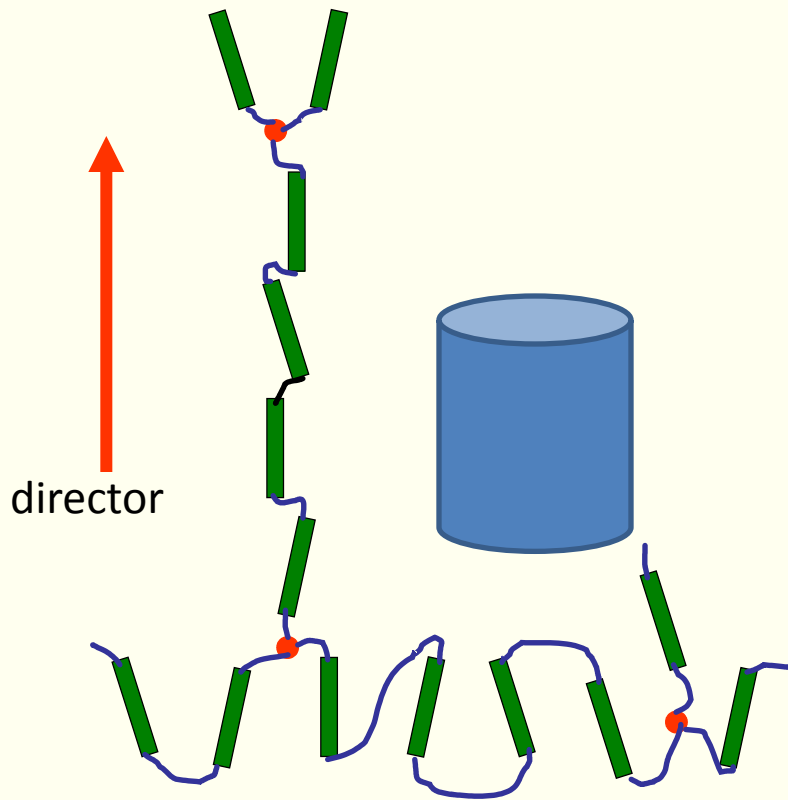
- processive motor (unidirectional translation)
- hand-over-hand motion
- driven by chemical energy of ATP
- work: translation of cargo against force, dissipation



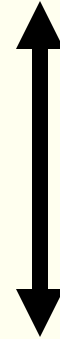
# Molecular Motors 5. LC elastomers



# LC elastomer network changes shape



If orientational order increases,



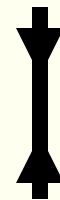
expansion



contraction



If orientational order decreases,



contraction



expansion



# Summary of LC elastomers

- elongational motor:
  - change in orientational order produces stress
  - stress produces change in shape
- driven by heat, light, or impurities
- work: motion against force, dissipation



# Broken Symmetry





# Proper- and pseudo-tensors

tensor rank	0	1	2	3	4
proper	$\pi, x$	<b>E, P, D</b>	$\delta_{\alpha\beta}, Q_{\alpha\beta}$ $\epsilon_{\alpha\beta}, \mu_{\alpha\beta}$	$d_{\alpha\beta\gamma}$	etc.
pseudo	<b>A•(B × C)</b>	<b>H, M, B</b>	$\Gamma_{\alpha\beta}$	$\epsilon_{ijk}$	

scalar

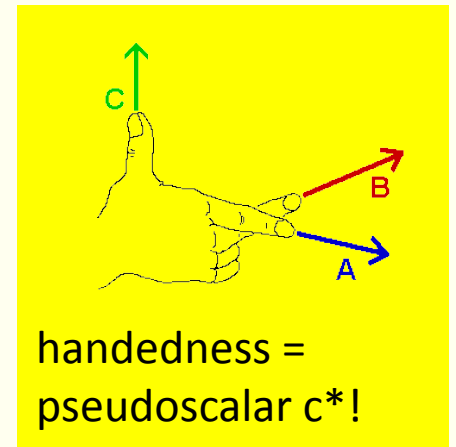
vector



changes sign on inversion



does not change sign on inversion



# Symmetry and motors

- symmetry of the system determines motion of motor!



# Translational motion

**YAHOO!** ANSWERS

## **My rugs are alive and moving!?**

I have wall to wall carpets in my apartment. I have placed large rugs on top of the carpet. Without damaging the carpet or the rug, how can I stop them from wrinkling and moving? Is it caused by static electricity?



# Translational motion

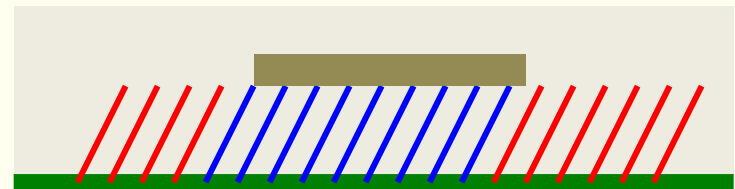
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tilted strands



# Translation

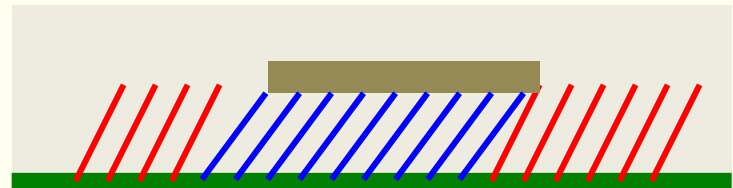
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# Translation

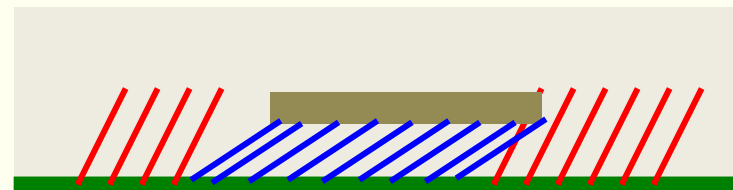
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# Translation

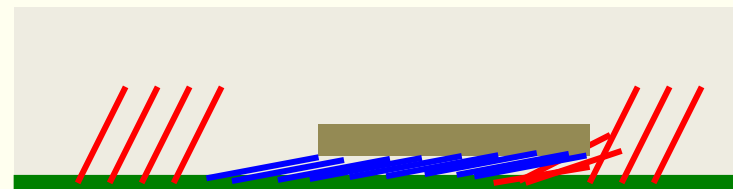
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tilted strands



# Translation

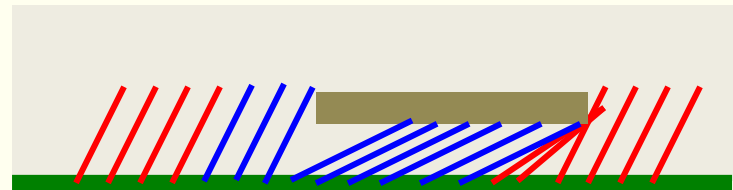
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tilted strands





# Translation

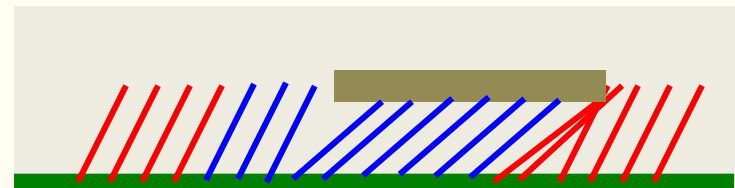
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tilted strands



# Translation

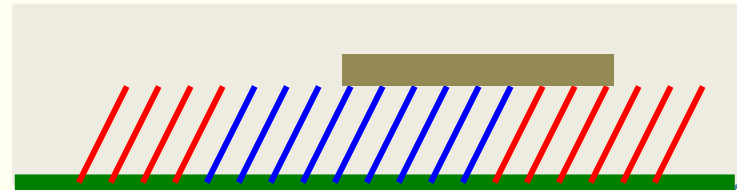
## YAHOO! ANSWERS

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tilted strands



# Symmetry and motors

- symmetry of the system determines motion of motor!
- broken left-right symmetry:
  - proper vector exists in system
  - translation in direction of vector (carpet, kinesin, myosin)



# Symmetry and motors

- symmetry of the system determines motion of motor!
- broken translational symmetry:
  - proper vector exists in system
  - translation in direction of vector (carpet, kinesin, myosin)  
[Q: is there a vector in Landauer's blow torch?]
- broken chiral symmetry
  - pseudoscalar  $c^*$  exists in system
- broken chiral + broken translational symmetry
  - pseudoscalar + proper vector = pseudovector
  - rotation (Yokoyama – Tabe, ATP Synthase, etc.)


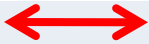


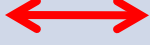
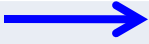
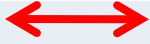


# Symmetry and motors

- if system has proper dyad (director)
  - strain (elongation)
- if system has proper dyad and gradient (vector),
  - proper vector (curvature, bend)
- if system has proper dyad, gradient & pseudoscalar
  - pseudovector (twist)



# Symmetry & deformation/motion

Symmetry	can construct	deformation/motion	material
	vector	translation	in kinesin, actin
	dyad	uniaxial strain	nematic networks
$c^*$	-	-	
$c^*$ 	pseudovector	rotation	Lehman, ATP, Tabe
  $c^*$	pseudovector	twist	LCE twist
 	proper vector	bend	LCE photoactuation



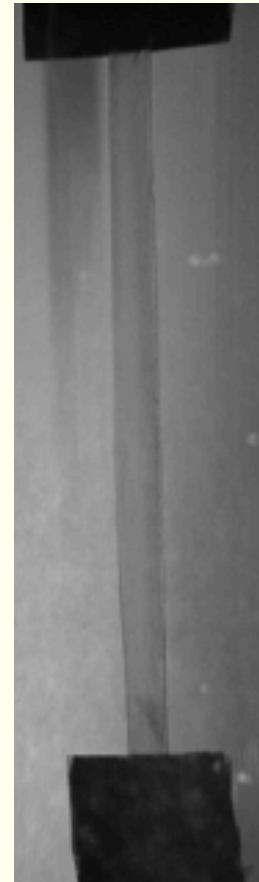
# Connecting molecular & macroscopic scales

- cumulative effect of many molecular motors → macroscopic response
  - rotating dye system, Yokoyama-Tabe expt:
    - distributed local vorticity → bulk torque & bulk flow
  - ATP synthase:
    - produces bulk ATP
  - myosin & actin
    - muscle fibers form long muscles, these contract
  - effective shape change
    - contraction, elongation and bend and twist



# Smooth shape change: elongation

- liquid crystal elastomers:
  - coupling of orientational order and strain
  - large mechanical response to excitations
  
- elongation due to temperature change
  
- motor based on elongation



H. Finkelmann, nematic LCE





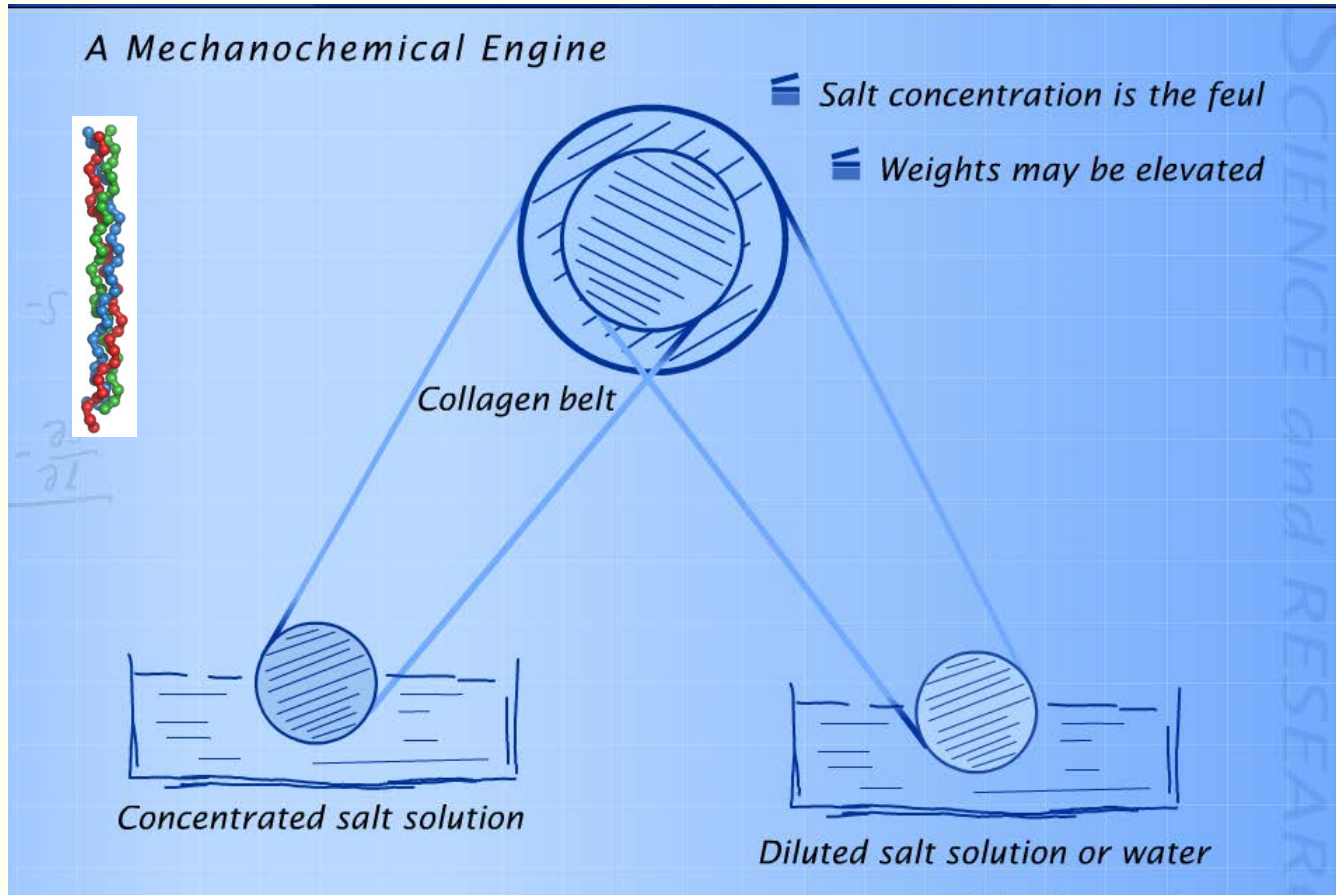
# Another motor based on elongation

- rubber band heat engine



# Another motor based on elongation

- collagen contracts in salt water

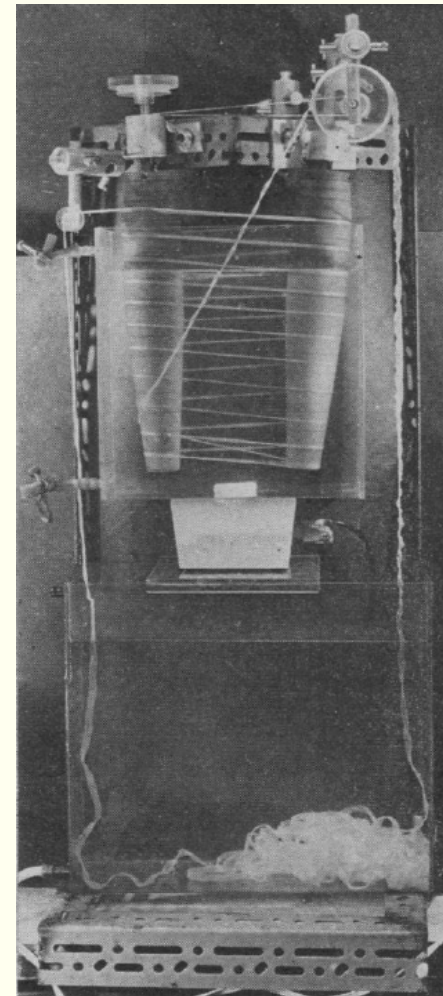
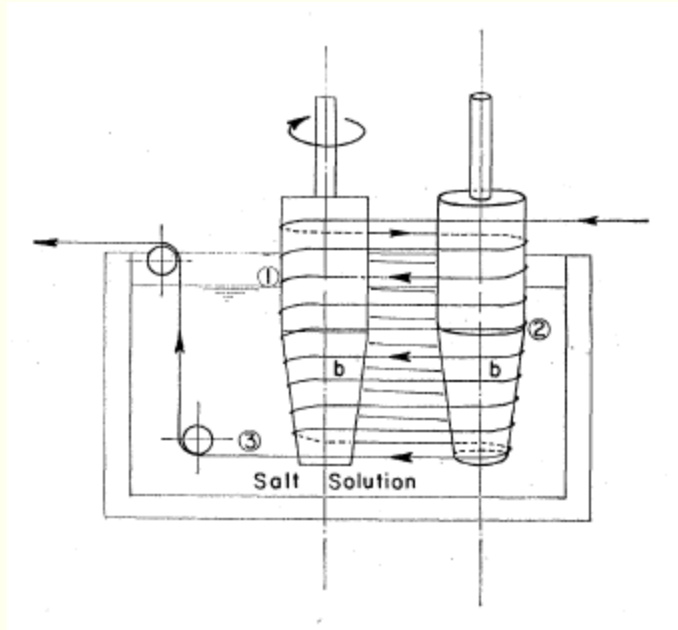


Aharon Katzir-Katchalsky 1913-1972



# Another motor based on elongation

- a modified version



M. V. Sussman and A. Katchalsky, *Science*, **167**, 45 (1970)



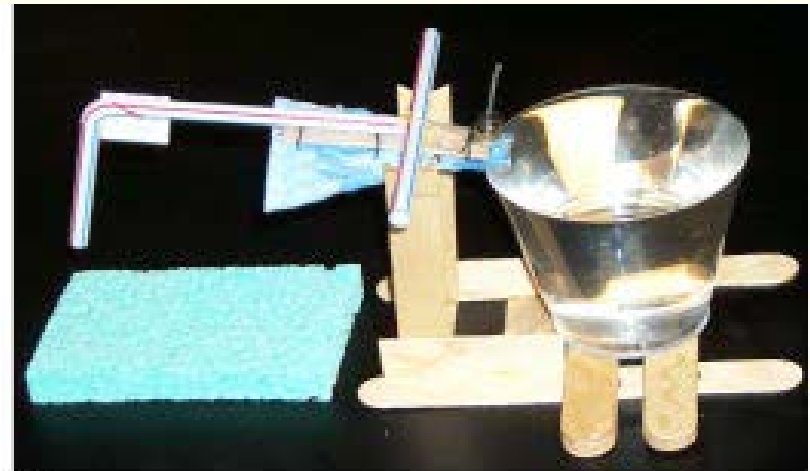
# Dunking bird of the first kind

- heat engine
  - Carnot efficiency limit

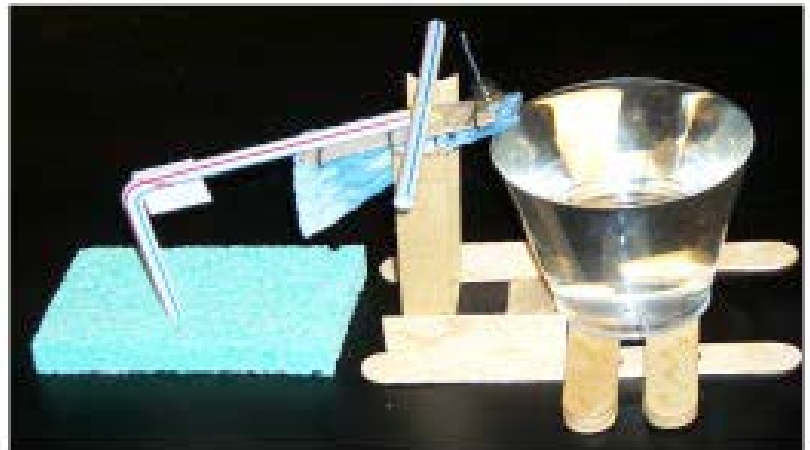


# Dunking bird of the second kind

- not a heat engine
- isothermal
- involves shape change



(a)



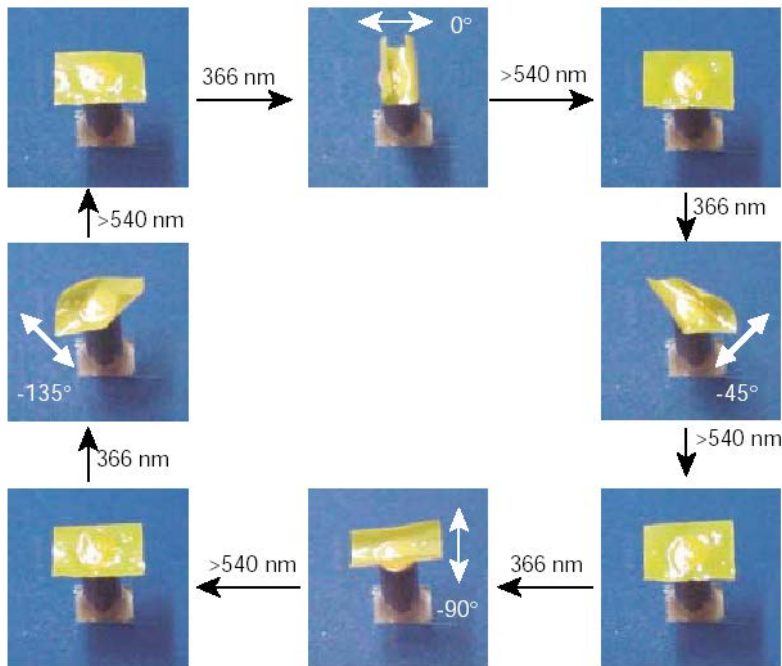
(b)

N. Abraham, P. Palffy-Muhoray., *Am.J. Phys.* **72**, 782 (2004)



# Bend Motors

- samples of liquid crystal elastomers with azo dye bend on exposure to light



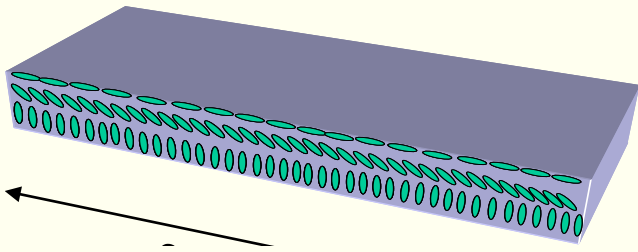
LC + diacrylate network  
+ functionalized  
azo-chromophore

timescale: 10 s

Yanlei Yu, Makoto Nakano, Tomiki Ikeda, *Nature* **425**, 125 (2003)



# Water vapor sensitive network

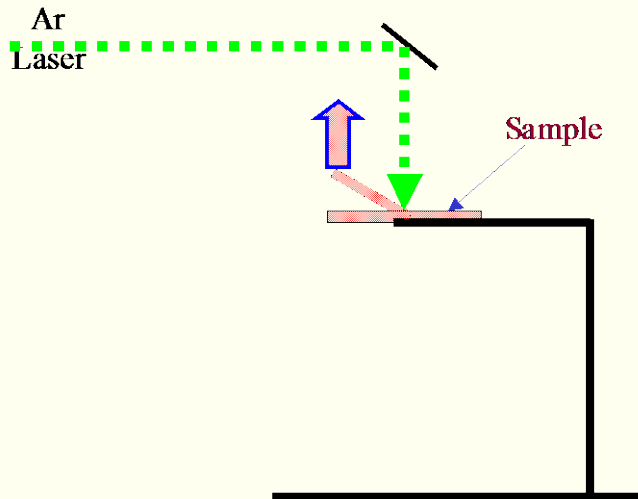
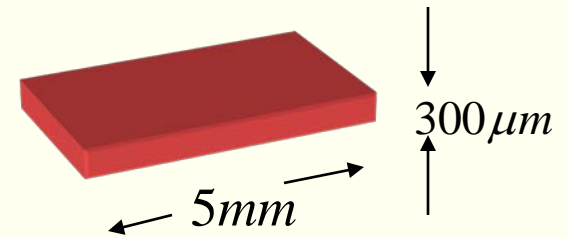


K. Harris, C. Bastiaansen, D. Broer, *J. MEM Syst.*, **16**, 480 (2007)



# Bend Motors

sample: nematic elastomer EC4OCH3  
+ 0.1% dissolved  
Disperse Orange 1 azo dye



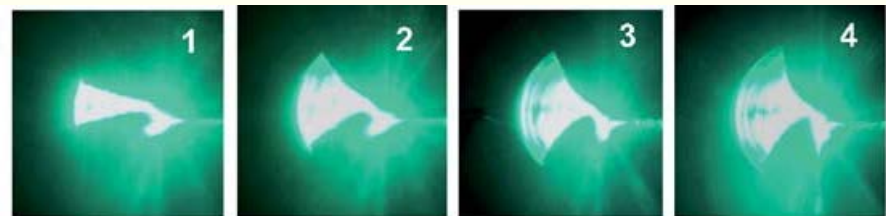
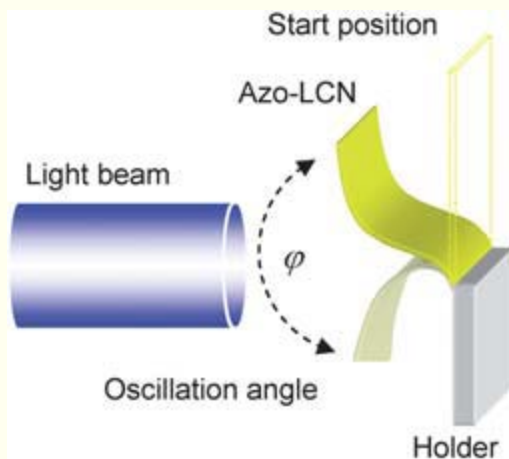
Response time: 70ms





# Photoinduced oscillations

- if sample bends  $> 90^\circ$  both sides are illuminated, producing oscillations



nematic azo-elastomer

sample size: 5 mm 0.8 mm 0.05 mm

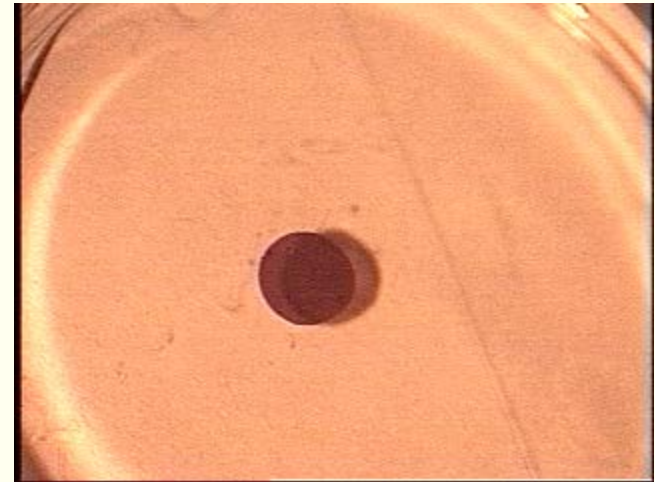
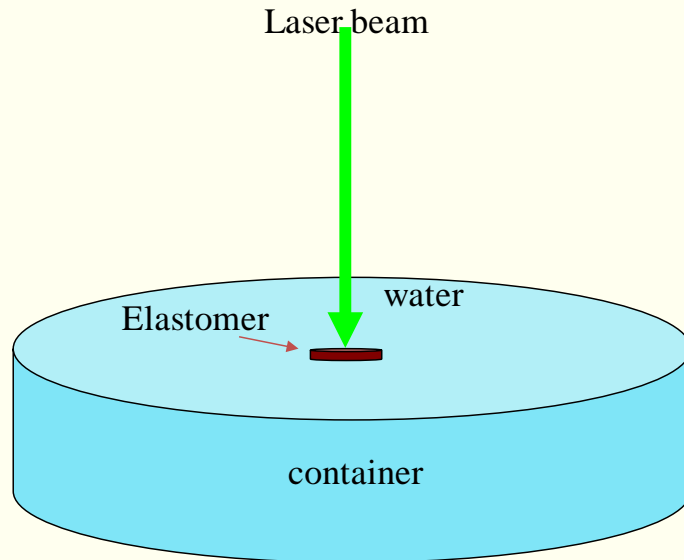
max. frequency = 270Hz

S. Serak, N. Tabiryan, R. Vergara, T. White, R. Vaia, T. Bunning, *Soft Matter* **6**, 779–783 (2010)



# Bend Motors

- floating nematic LCE sample illuminated from above

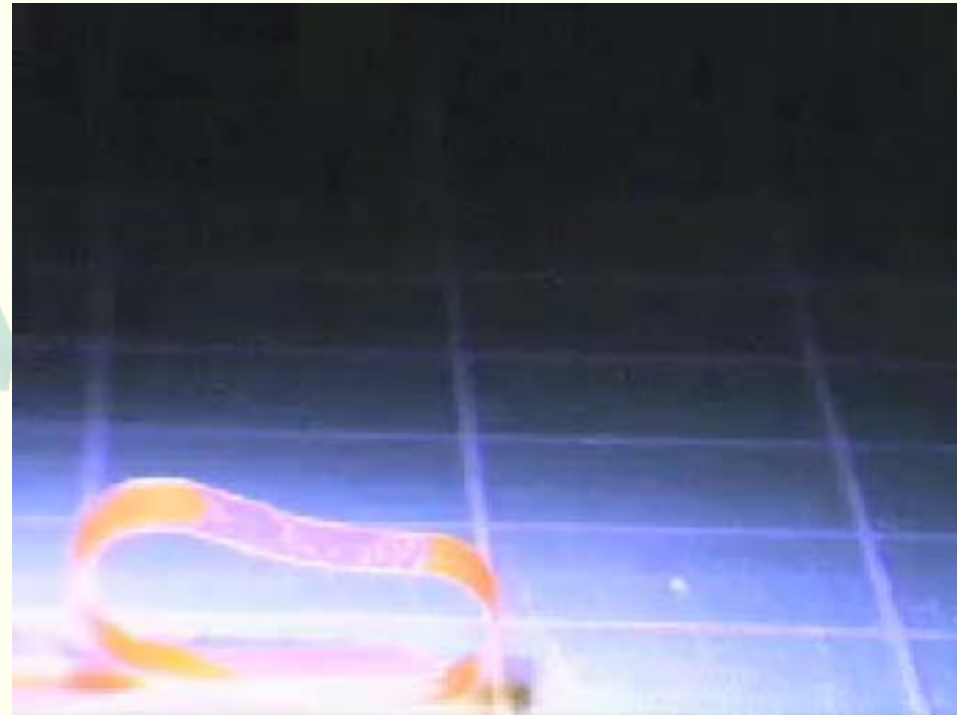
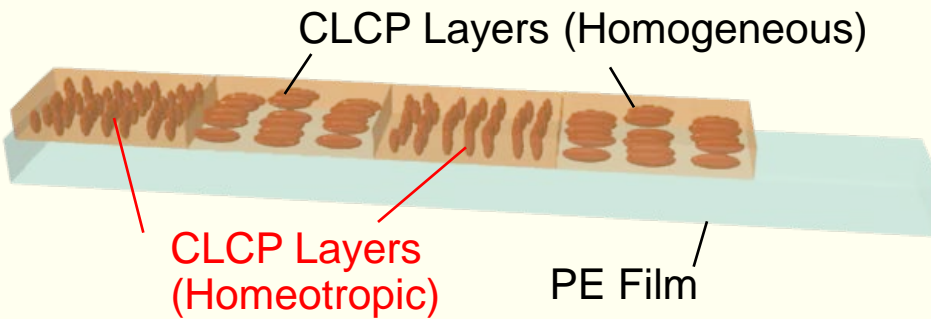


M. Chamacho-Lopez, H. Finkelmann, P. Palffy-Muhoray, M. Shelley, *Nature Mat.* **3**, 307, (2004)



# Bend Motors

- robotic arm



T. Ikeda, Chuo University (priv. comm.)



# Bend Motor: focus of current interest

- azo-dye doped liquid crystal elastomer

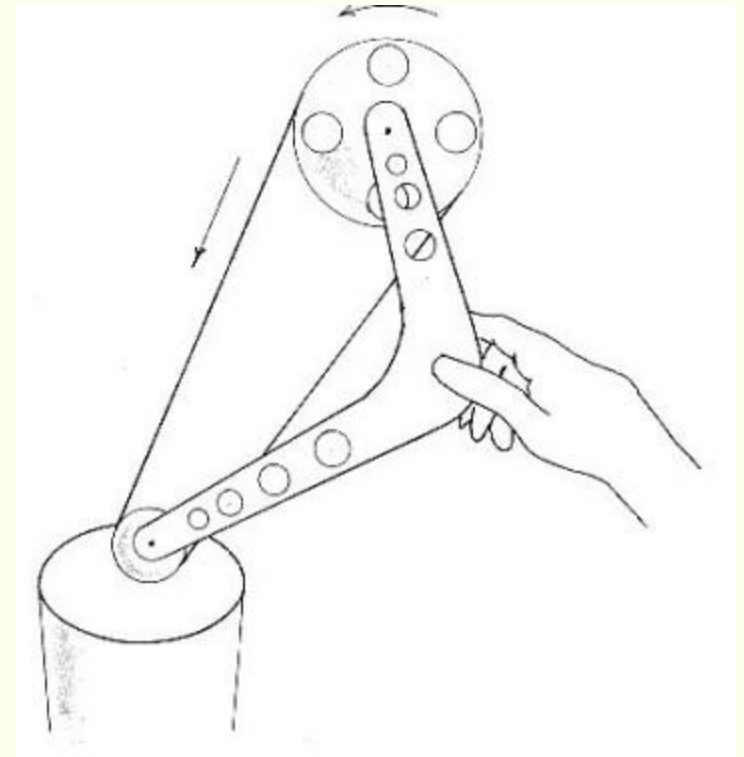


M. Yamada, M. Kondo, J. Mamiya, Y. Yu, M. Kinoshita, C. Barrett, T. Ikeda, *Angew. Chem.* **47**, 4986 (2008)

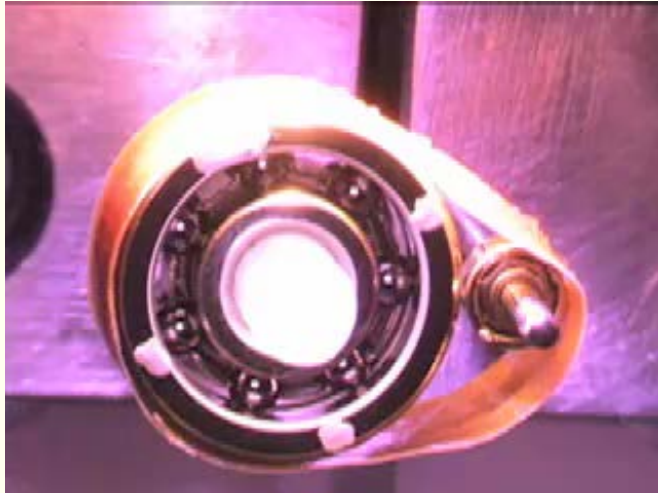


# Bend Motor: focus of current interest

- shape-memory Ni-Ti alloy,
  - two phases:
    - orthorhombic martensite (cold)
    - cubic austenite (hot)



# Existing rotary bend motors



LCE bends on exposure to light

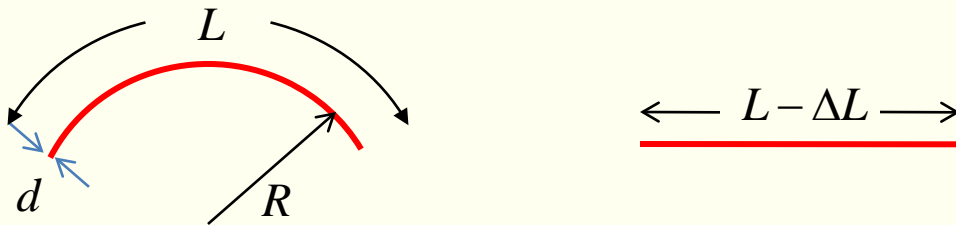


wire becomes stiff when heated



# Model

- bend vs. compression:
  - if ends are brought closer, will filament compress or bend?



$$\frac{\Delta L}{L} \approx \frac{1}{24} \left(\frac{L}{R}\right)^2$$

$$\mathcal{E}_b / l = \frac{1}{2} E \left(\frac{d}{R}\right)^2 \frac{wd}{12}$$

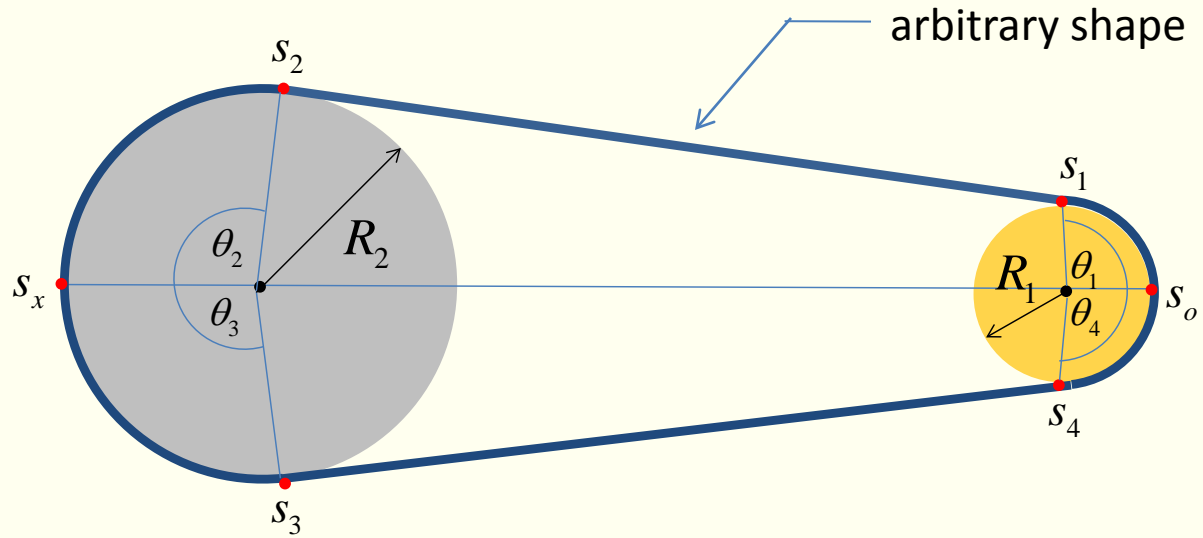
$$\mathcal{E}_c / l = \frac{1}{2} E \left(\frac{\Delta L}{L}\right)^2 wd$$

$$\boxed{\frac{\mathcal{E}_c}{\mathcal{E}_b} \sim \left(\frac{L}{d}\right)^2 \left(\frac{L}{R}\right)^2}$$

Bend is *much* cheaper energetically than stretch;  
→ length is constant.

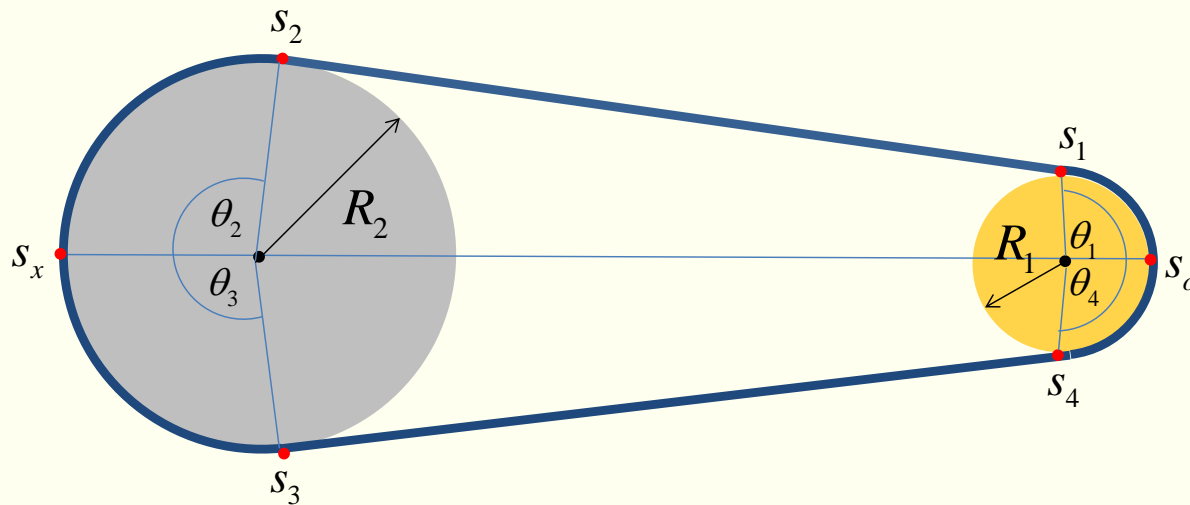


# Model



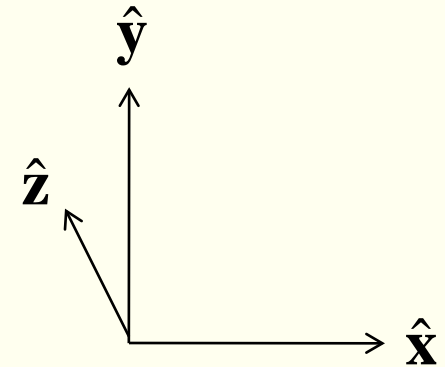


# Model

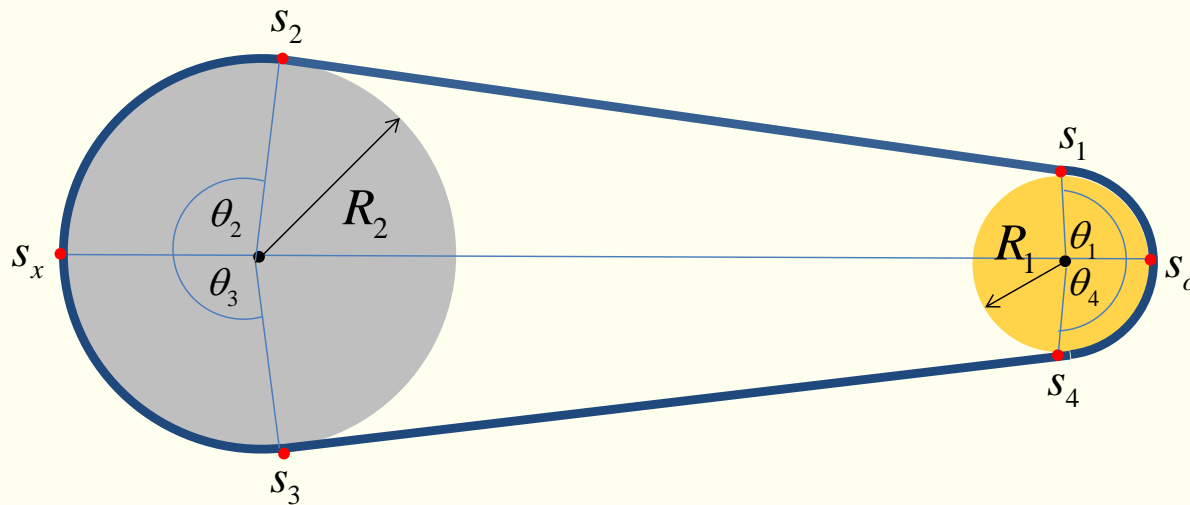


$s$  – arc length along filament

$\mathbf{R}(s)$  – position vector of point on filament



# Model



– bend energy density:

$$\mathcal{E}_b / l = \frac{1}{2} EI \frac{1}{R^2} = \frac{c}{R^2}$$

stiffness  
fn. of position

natural curvature  
fn. of position

$$\mathcal{E} = \int_{s_o}^{s_1} c \left( \frac{1}{R_1} - \kappa \right)^2 ds + \int_{s_1}^{s_2} c \left( \frac{\partial^2 \mathbf{R}}{\partial s^2} - \kappa \left( \frac{\partial \mathbf{R}}{\partial s} \times \hat{\mathbf{z}} \right) \right)^2 ds +$$

$$\int_{s_2}^{s_3} c \left( \frac{1}{R_2} - \kappa \right)^2 ds + \int_{s_3}^{s_4} c \left( \frac{\partial^2 \mathbf{R}}{\partial s^2} - \kappa \left( \frac{\partial \mathbf{R}}{\partial s} \times \hat{\mathbf{z}} \right) \right)^2 ds + \int_{s_4}^{s_o} c \left( \frac{1}{R_1} - \kappa \right)^2 ds,$$



# Model



– constraints:

- filament length is constant
- displacement in x-dir. is fixed

$$R_1(1 - \cos \theta_1) - \int_{s_1}^{s_2} \hat{\mathbf{t}} \cdot \hat{\mathbf{x}} ds + R_2(1 - \cos \theta_2) = L + R_1 + R_2,$$

$$R_1(1 - \cos \theta_4) + \int_{s_3}^{s_4} \hat{\mathbf{t}} \cdot \hat{\mathbf{x}} ds + R_2(1 - \cos \theta_3) = L + R_1 + R_2.$$

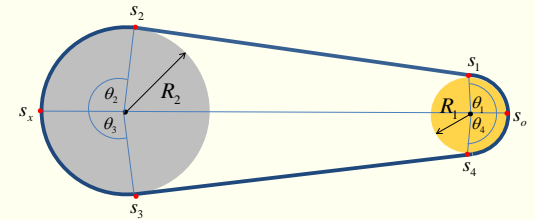
- displacement in y-dir. is fixed

$$R_1 \sin \theta_1 + \int_{s_1}^{s_2} \hat{\mathbf{t}} \cdot \hat{\mathbf{y}} ds - R_2 \sin \theta_2 = 0,$$

$$R_1 \sin \theta_4 + \int_{s_3}^{s_4} \hat{\mathbf{t}} \cdot \hat{\mathbf{y}} ds - R_2 \sin \theta_3 = 0$$



# Model: Energy with constraints



$$\mathcal{E} = \int_0^{s_1} c \left( \frac{1}{R_1} - \kappa \right)^2 ds + \int_{s_1}^{s_2} c (\hat{\mathbf{t}}' - \kappa (\hat{\mathbf{t}} \times \hat{\mathbf{z}}))^2 ds + \int_{s_2}^{s_3} c \left( \frac{1}{R_2} - \kappa \right)^2 ds$$

$$+ \int_{s_3}^{s_4} c (\hat{\mathbf{t}}' - \kappa (\hat{\mathbf{t}} \times \hat{\mathbf{z}}))^2 ds + \int_{s_4}^1 c \left( \frac{1}{R_1} - \kappa \right)^2 ds$$

$$+ f_{1x} \left\{ R_1 \left( 1 - \cos \frac{s_1}{R_1} \right) - \int_{s_1}^{s_2} \hat{\mathbf{t}} \cdot \hat{\mathbf{x}} ds + R_2 \left( 1 - \cos \left( \frac{s_x - s_2}{R_2} \right) \right) - (L + R_1 + R_2) \right\}$$

$$+ f_{1y} \left\{ R_1 \sin \frac{s_1}{R_1} + \int_{s_1}^{s_2} \hat{\mathbf{t}} \cdot \hat{\mathbf{y}} ds - R_2 \sin \left( \frac{s_x - s_2}{R_2} \right) \right\}$$

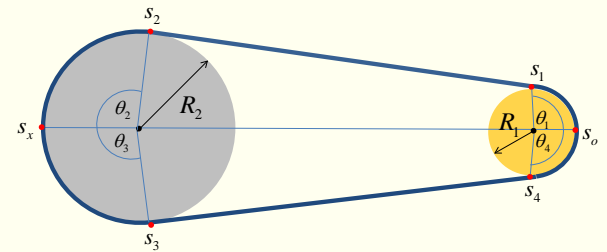
$$+ f_{2x} \left\{ R_1 \left( 1 - \cos \left( \frac{1 - s_4}{R_1} \right) \right) + \int_{s_3}^{s_4} \hat{\mathbf{t}} \cdot \hat{\mathbf{x}} ds + R_2 \left( 1 - \cos \left( \frac{s_3 - s_x}{R_2} \right) \right) - (L + R_1 + R_2) \right\}$$

$$+ f_{2y} \left\{ R_1 \sin \left( \frac{1 - s_4}{R_1} \right) + \int_{s_3}^{s_4} \hat{\mathbf{t}} \cdot \hat{\mathbf{y}} ds - R_2 \sin \left( \frac{s_3 - s_x}{R_2} \right) \right\}$$

Lagrange multipliers are components of forces in top and bottom filaments.



# Model: Energy minimization



- ODEs are obtained by minimizing  $\mathcal{E}$  w.r.t.  $\hat{\mathbf{t}} = (\cos \theta, \sin \theta)$ 
  - $-2(c\theta'' + c'\theta' + (c\kappa)') + f_{1x} \sin \theta + f_{1y} \cos \theta = 0$ , if  $s \in (s_1, s_2)$
  - $-2(c\theta'' + c'\theta' + (c\kappa)') - f_{2x} \sin \theta + f_{2y} \cos \theta = 0$ , if  $s \in (s_3, s_4)$ .
- these govern the shape of the filaments on top and bottom
  - BC for  $\theta$  is obtained by minimizing  $\mathcal{E}$  w.r.t.  $s_1, s_2, s_3, s_4$
  - this gives continuous tangent and curvature at contact pts.



# Numerical Solution

- 1. Assign initial values for  $s_1, s_2, s_x$  and  $f_{1x}, f_{1y}$  and solve ODE

$$-2(c\theta'' + c'\theta' + (c\kappa)') + f_{1x} \sin \theta + f_{1y} \cos \theta = 0$$

with BCs.

2. Assign initial values for  $s_3, s_4$  and  $f_{2x}, f_{2y}$  and solve ODE

$$-2(c\theta'' + c'\theta' + (c\kappa)') - f_{2x} \sin \theta + f_{2y} \cos \theta = 0$$

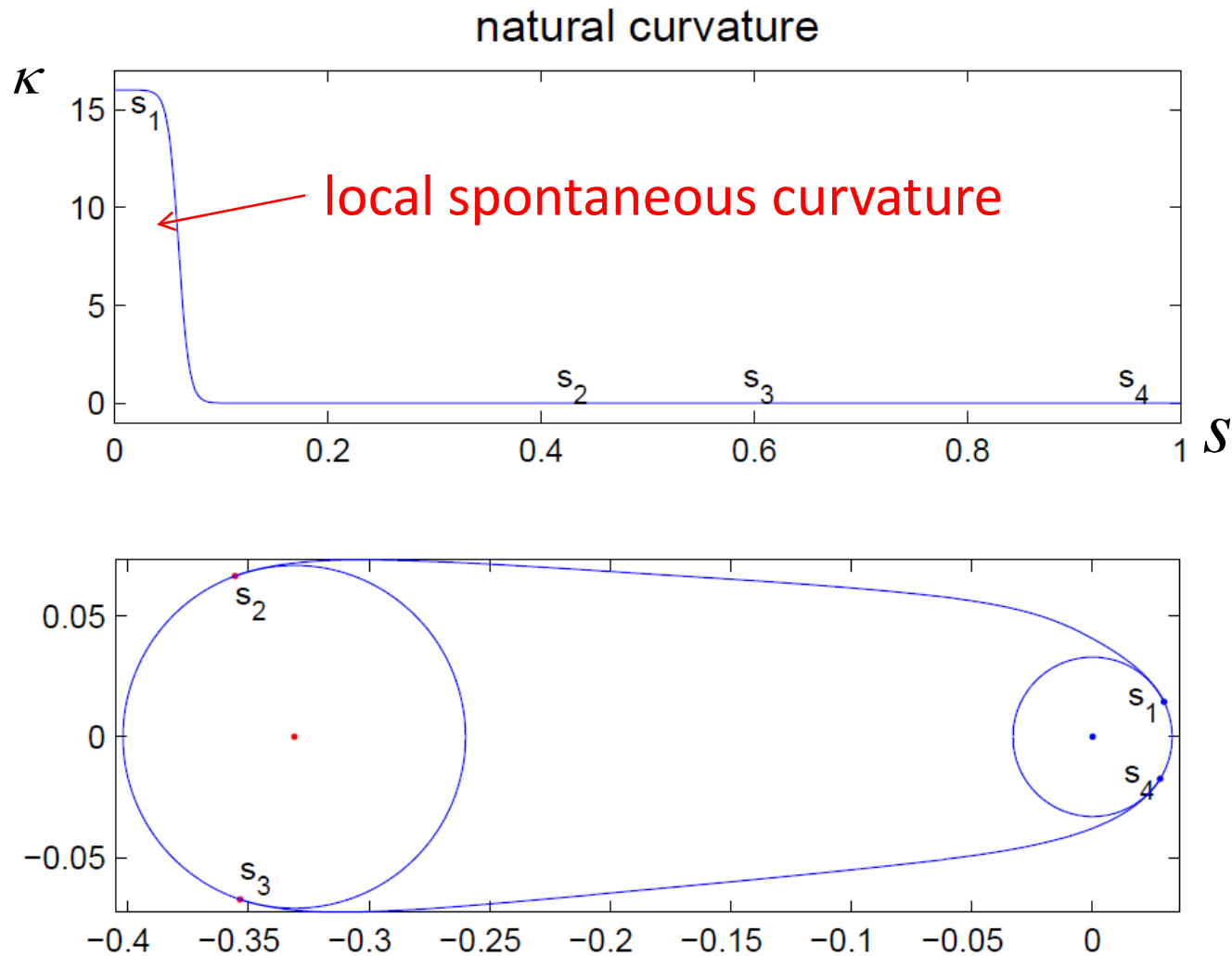
with BCs.

3. Vary the forces  $f_{1x}, f_{1y}, f_{2x}, f_{2y}$  to satisfy constraints (root finding)

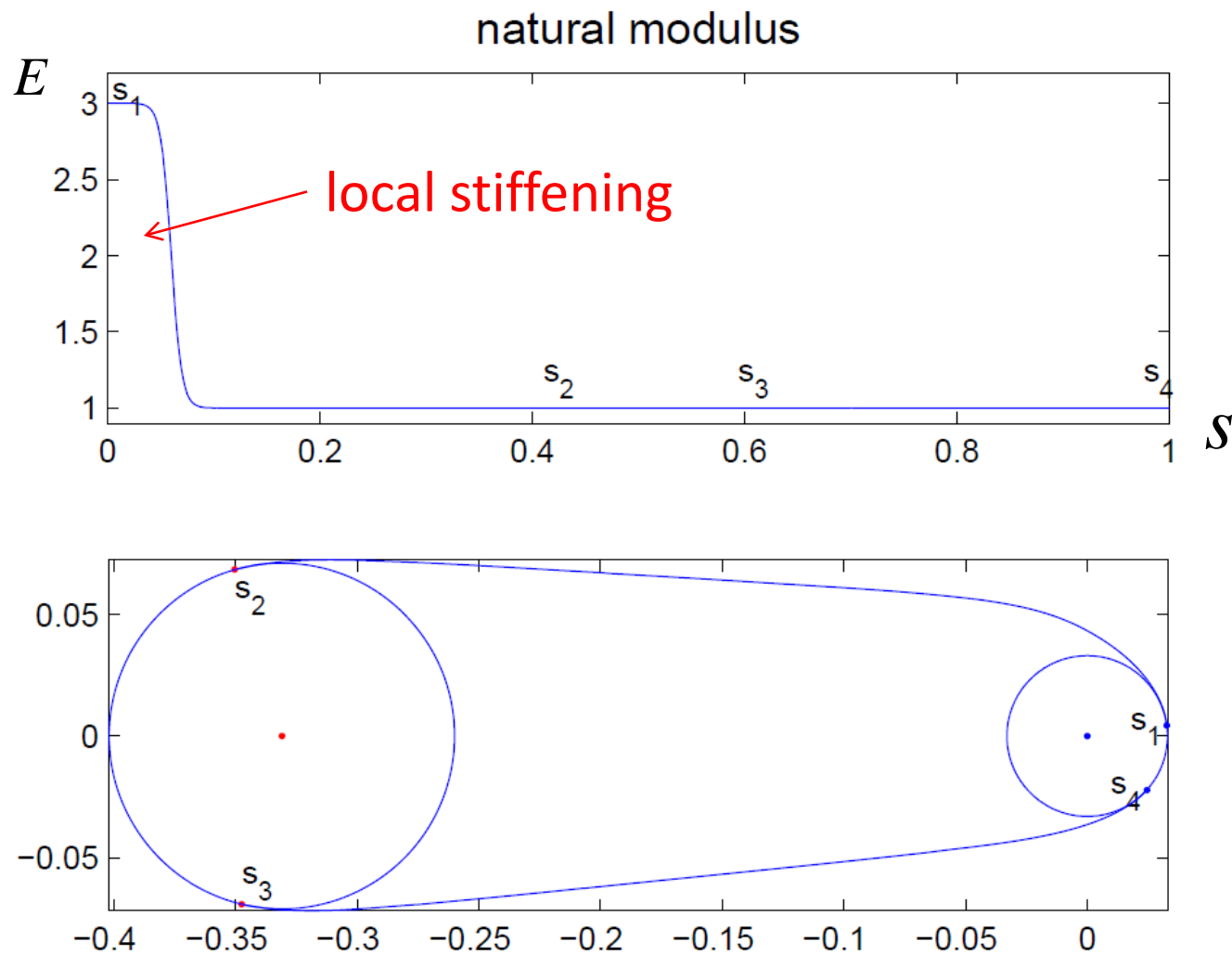
4. Vary  $s_1, s_2, s_3, s_4, s_x$  and go to 1. until  $\frac{\partial \mathcal{E}}{\partial s_i} = 0, \quad i = 1, \dots, 4, x.$



# Numerical Solution: varying the curvature



# Numerical Solution: varying the stiffness





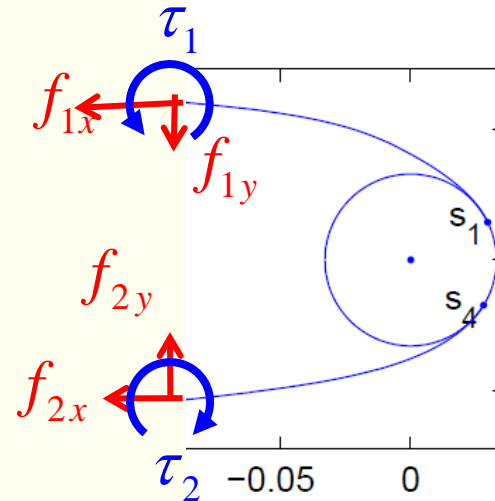
# Calculating torques

– to calculate the net torque, we consider part of system

– look at forces

– look at point torques

$$\tau = 2c(\hat{\mathbf{t}}' - \kappa(\hat{\mathbf{t}} \times \hat{\mathbf{z}}))$$



Note: integrating the ODE w.r.t. to  $s$  gives

$$2c(\theta' + \kappa) - f_{1x} \int_0^s \sin \theta(\tilde{s}) d\tilde{s} - f_{1y} \int_0^s \cos \theta(\tilde{s}) d\tilde{s} = K_1, (\text{upper part})$$

ODE is *precisely* the condition that the torque on wheel remain constant, regardless of where the filament is cut!



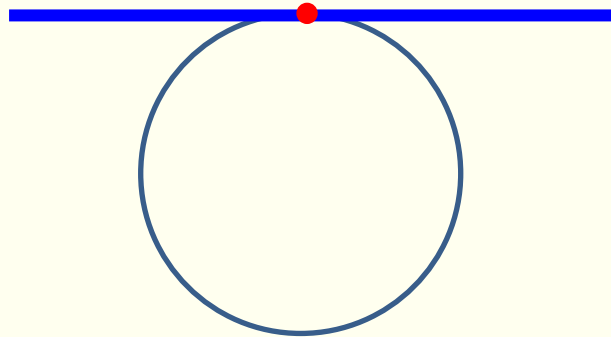
# What happens on the pulleys?

- Problem: elastic beam on rigid cylinder

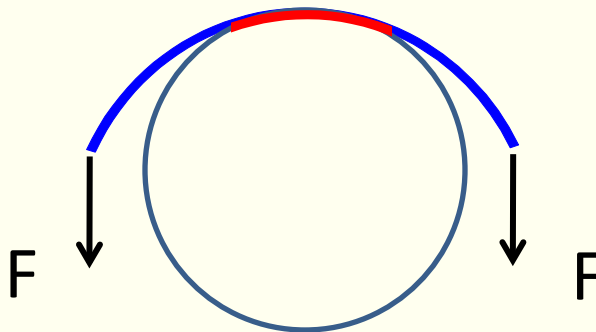


# Interesting observation

- straight beam touches cylinder along one line



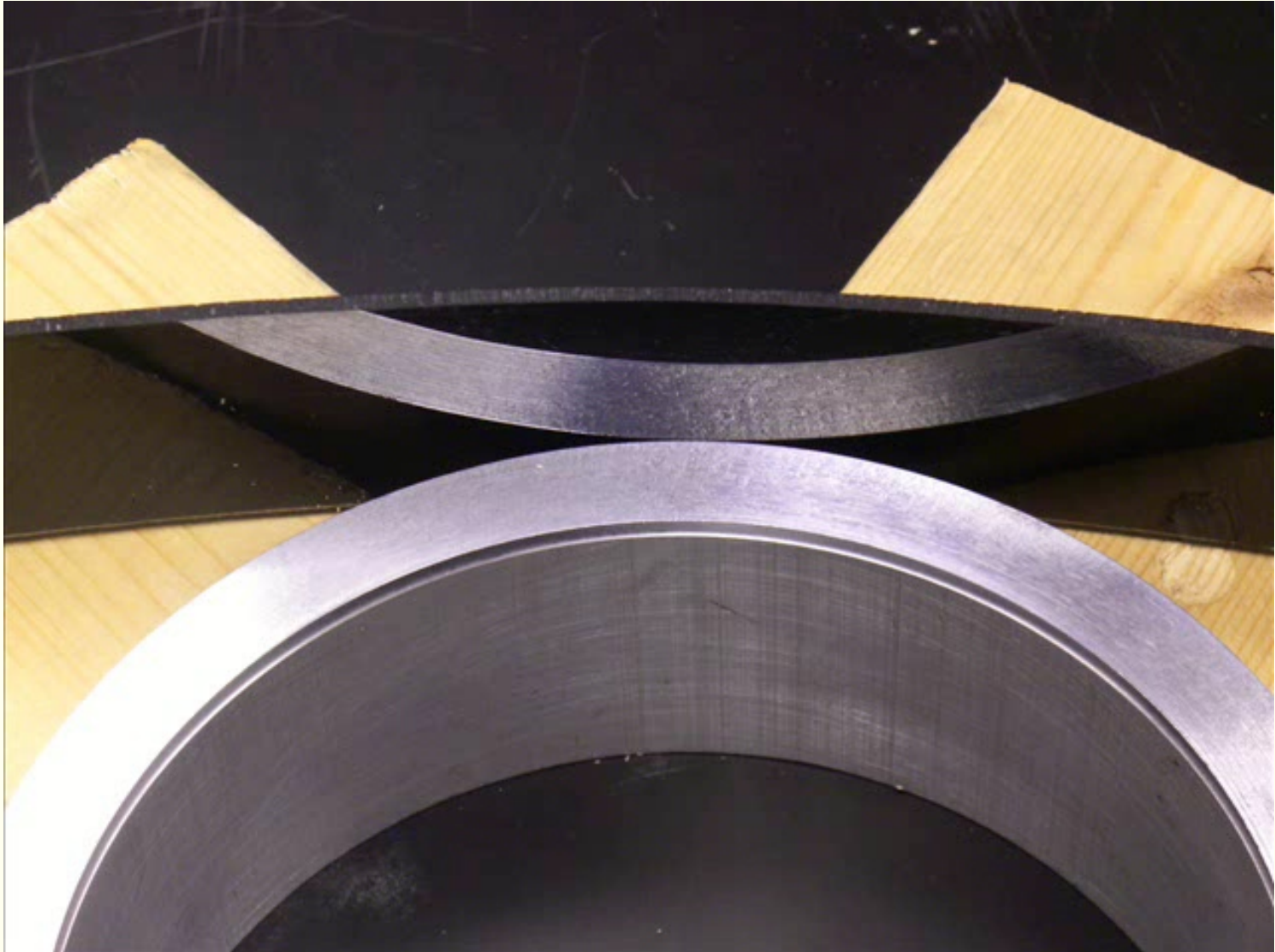
- curved beam touches cylinder over extended area



Transition?



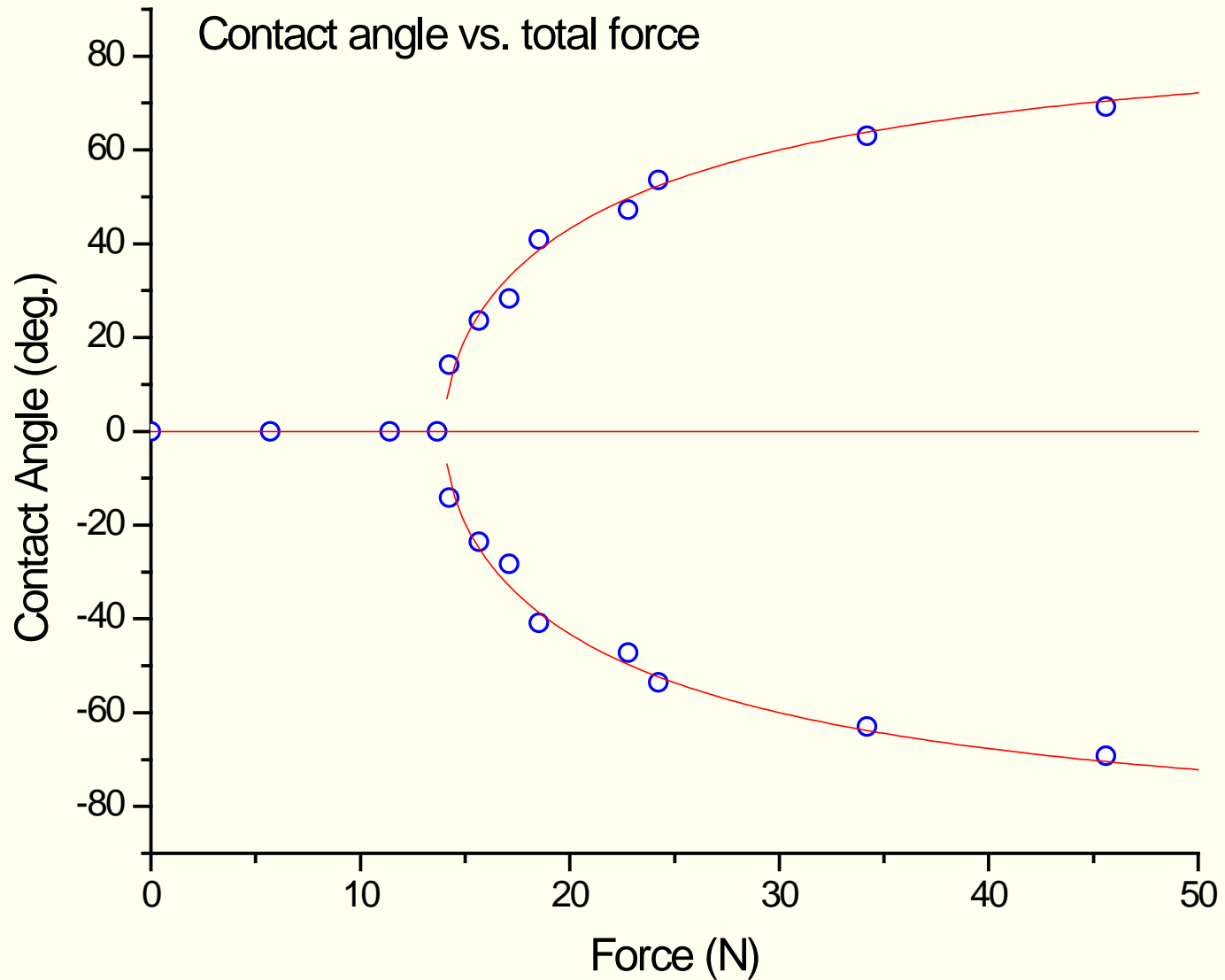
# Experiment



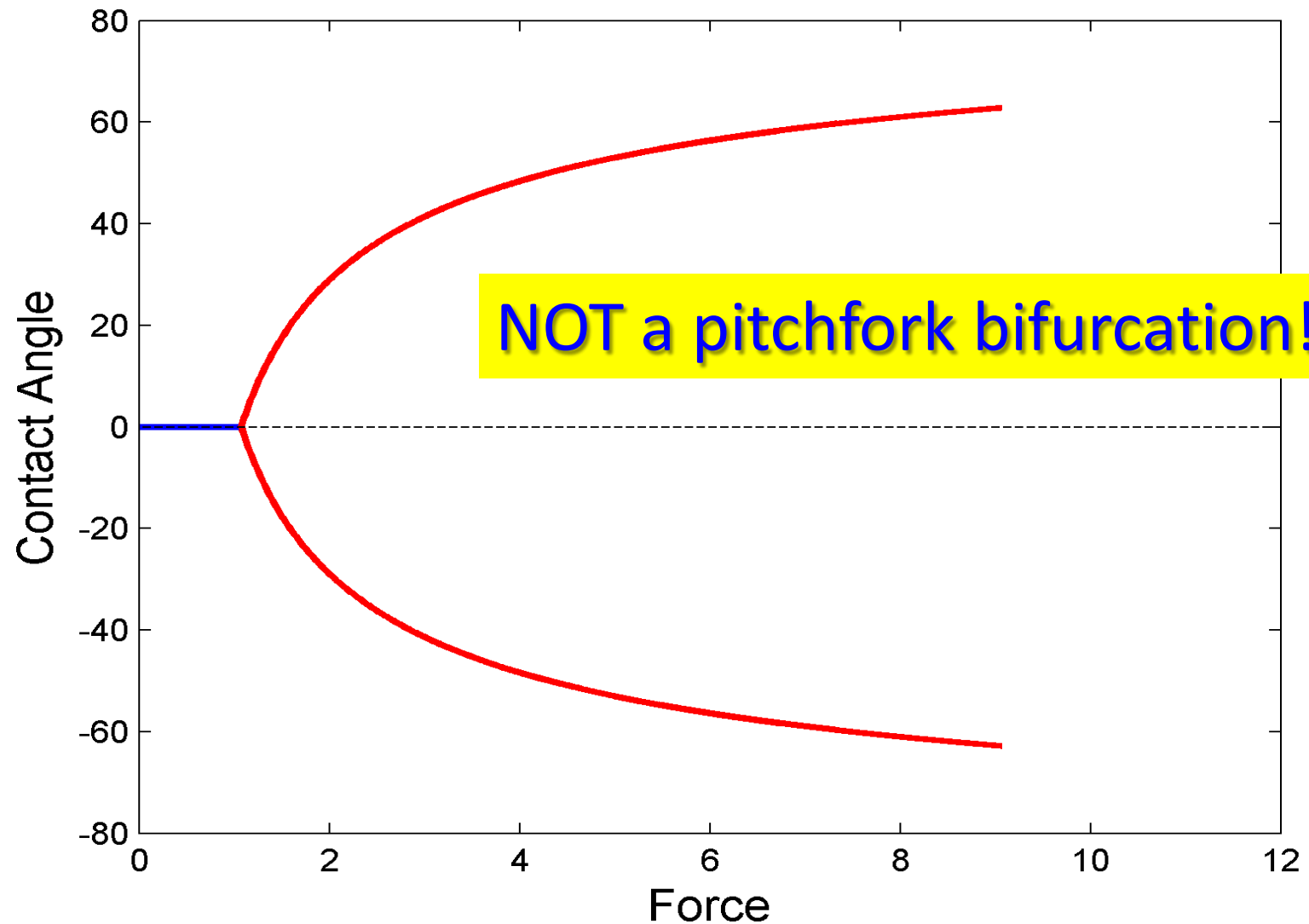
# Determination of point of contact



# Data



# Modelling: as before



# Force distribution on pulleys

- when force  $\mathbf{F}$  is applied to filament,

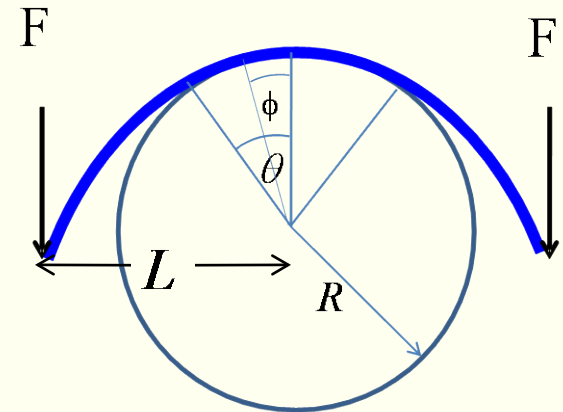
- point contact if

$$F \leq F_c$$

- extended contact if

$$F > F_c$$

$$F_c = \frac{EI}{LR}$$



- torque balance:

$$F(L - R \sin \theta) - \int_{\theta}^{\theta_c} R \sin(\phi - \theta) f(\phi) R d\phi = \frac{EI}{R}$$

force density on pulley

- Volterra integral equation of first kind.

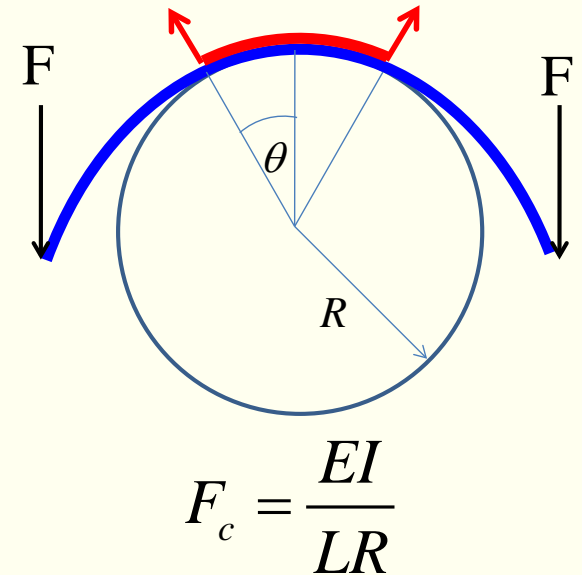




# Force distribution on pulleys

- When force  $\mathbf{F}$  is applied to filament,

- point contact if  $F \leq F_c$
- extended contact if  $F > F_c$



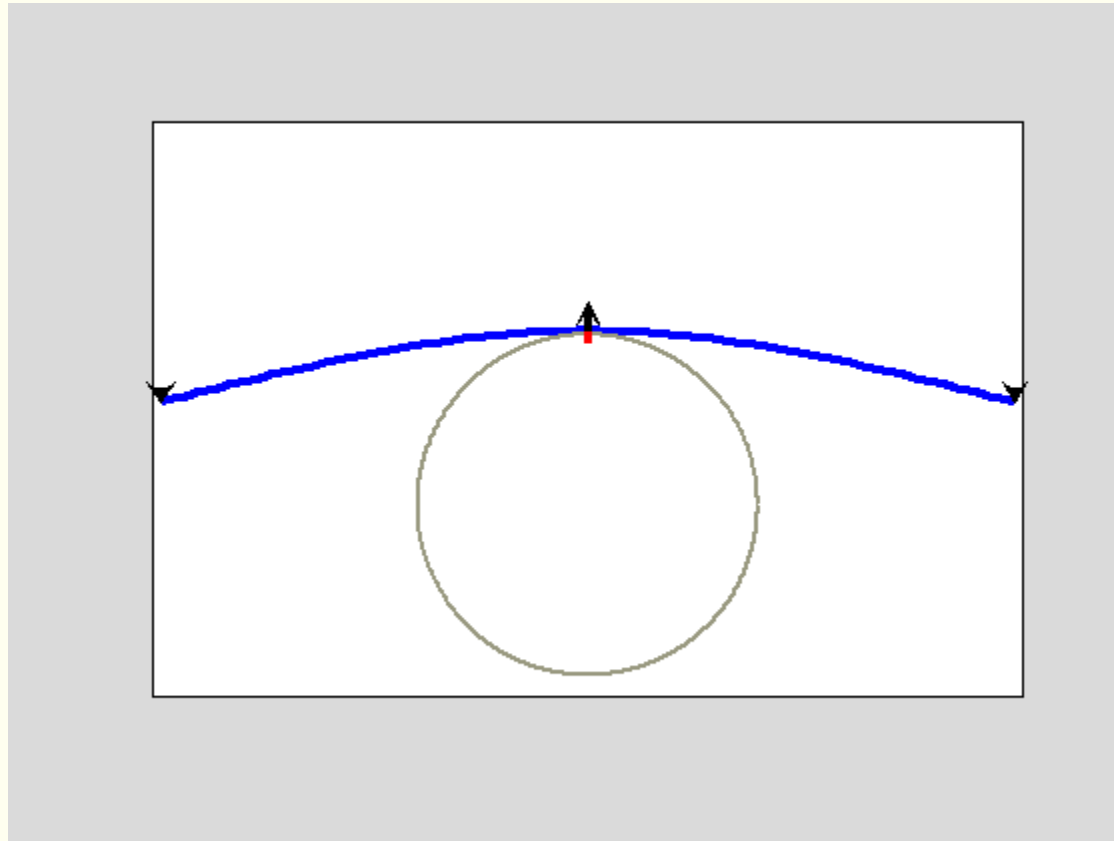
- Solution for force density:

$$f(\theta) = \frac{F}{R} (\delta(\theta - \theta_c) \cos \theta_c + \sin \theta_c)$$

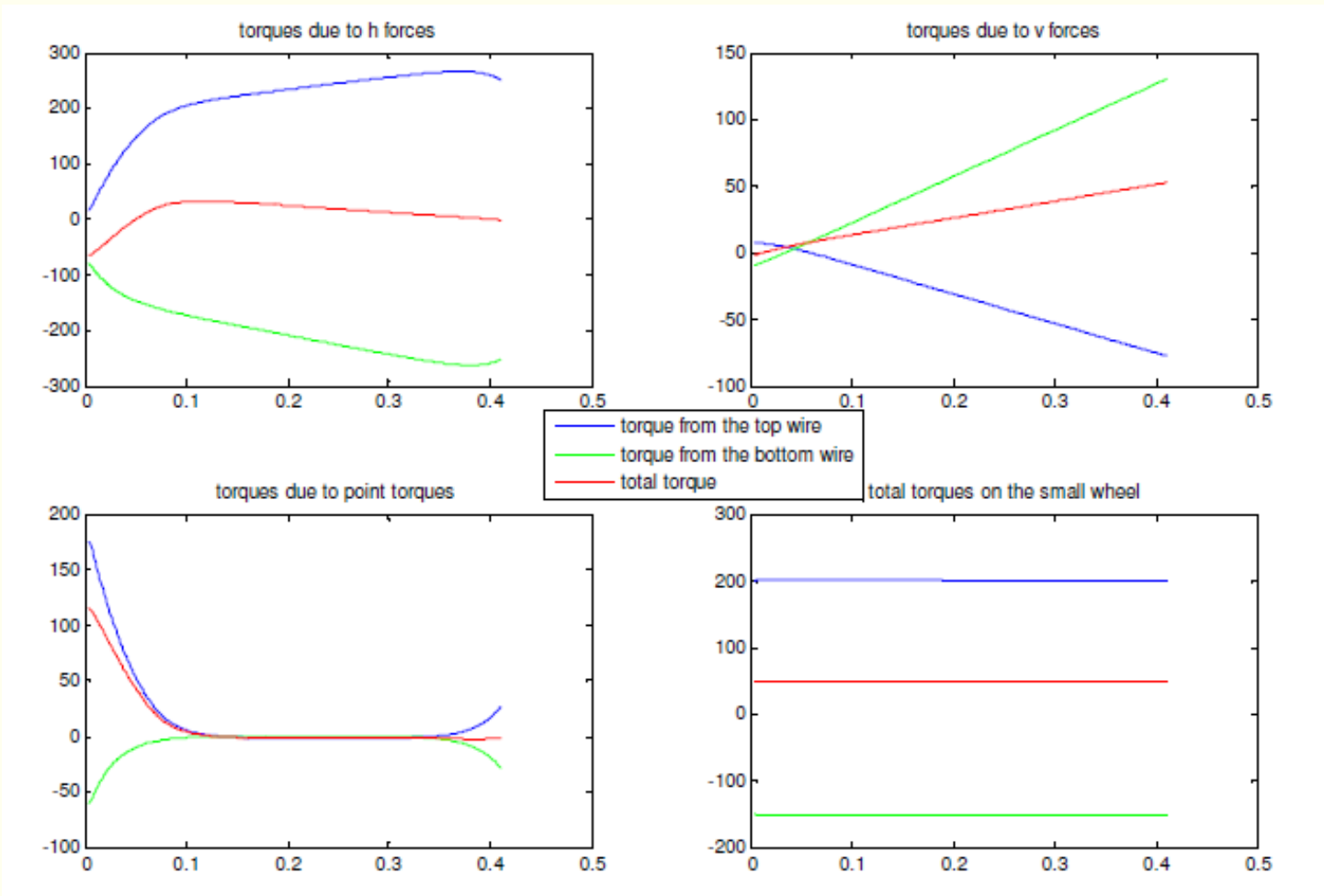
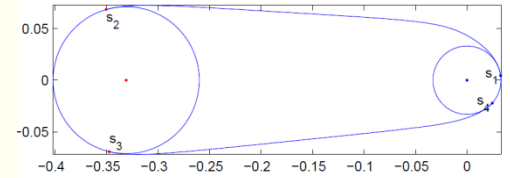
where  $\sin \theta_c = \left( \frac{F - F_c}{F} \right) \frac{L}{R}$



# Force distribution

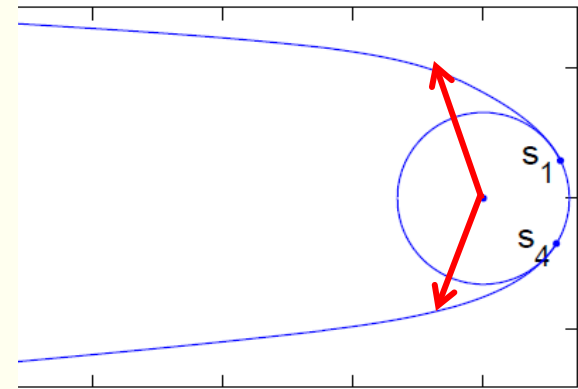


# How does rotation come about?



# How does rotation come about?

- away from the small wheel,
  - longer lever arm on one side

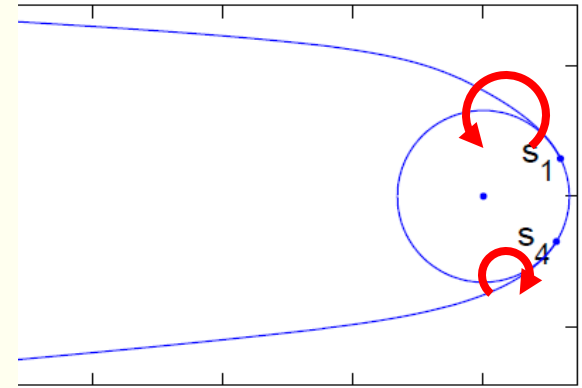


- near small wheel
  - greater point torque on one side
    - curvature is the same,
      - but stiffer on one side
      - have pontaneous curvature on one side



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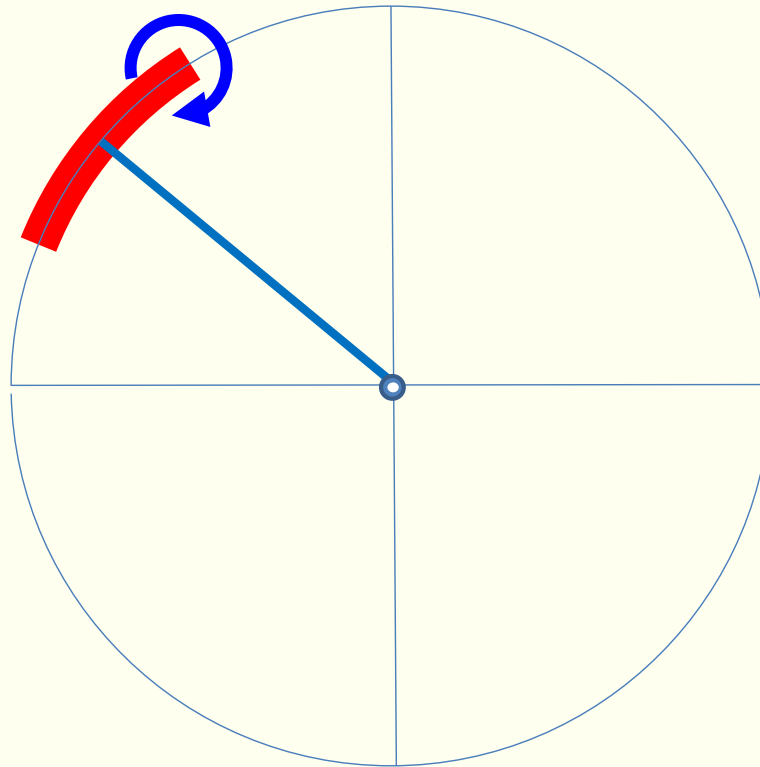


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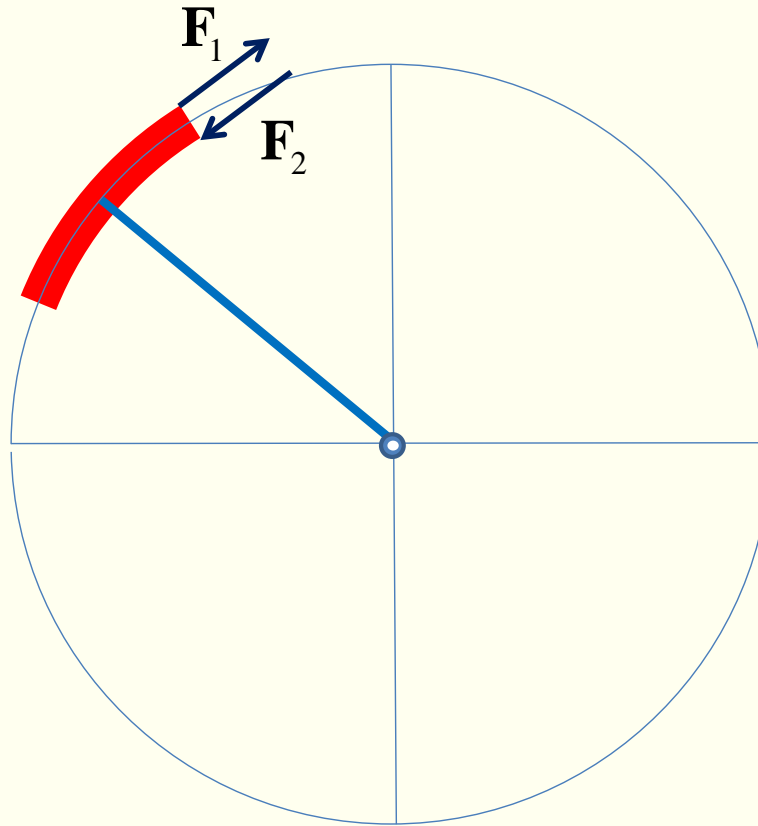
# Motion due to point torques

- how can a point torque in filament cause translation?



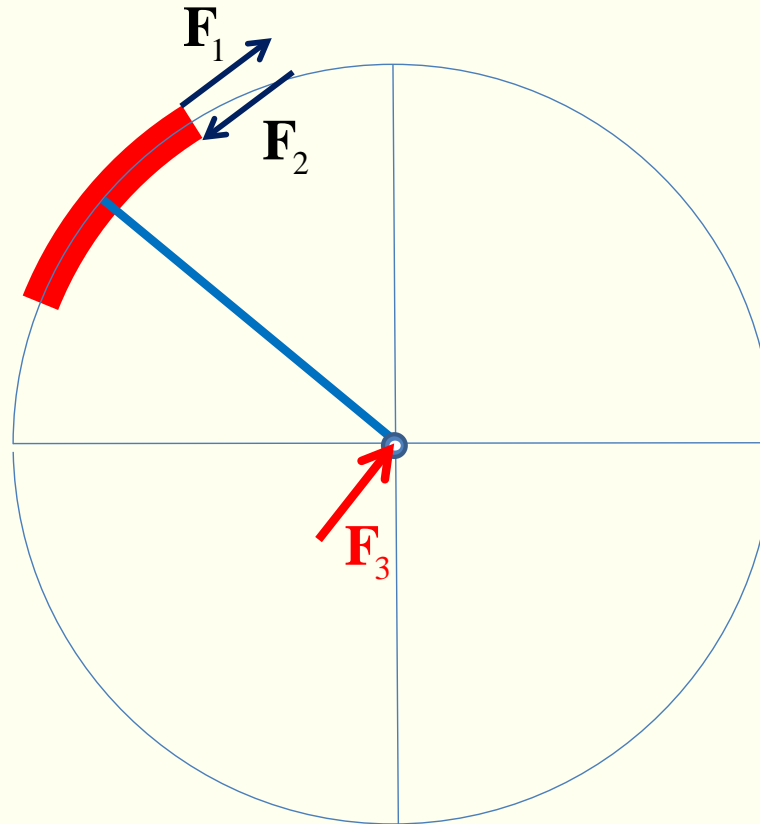
# Motion due to point torques

- how can a point torque in filament cause bulk translation?
  - net force is zero,



# Motion due to point torques

- how can a point torque in filament cause translation?
  - net force is zero, but torque causes constraint force to appear

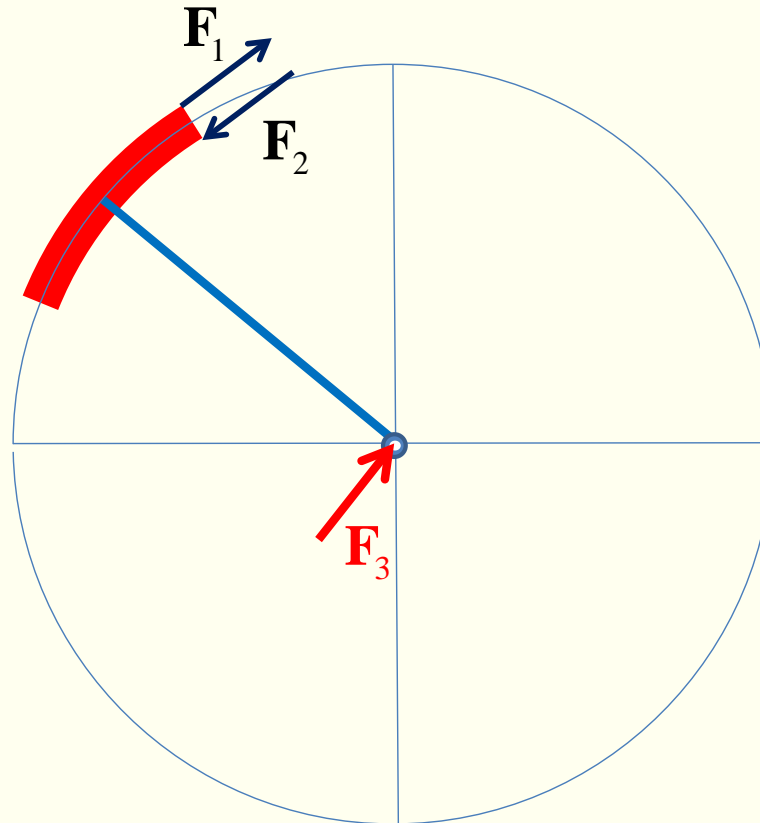




# Motion due to point torques

- how can a point torque in filament cause translation?
  - net force is zero, but torque causes constraint force to appear

- net force along  $\mathbf{F}_1$
- translation



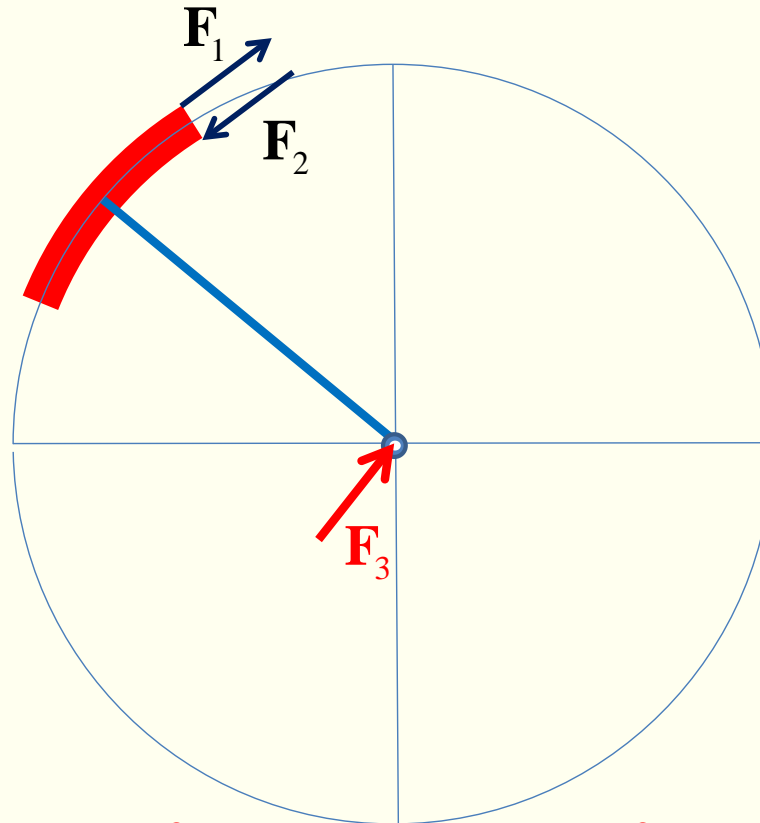
- angular momentum transport!



# Motion due to point torques

- how can a point torque in filament cause translation?
  - net force is zero, but torque causes constraint force to appear

- net force along  $\mathbf{F}_1$
- translation

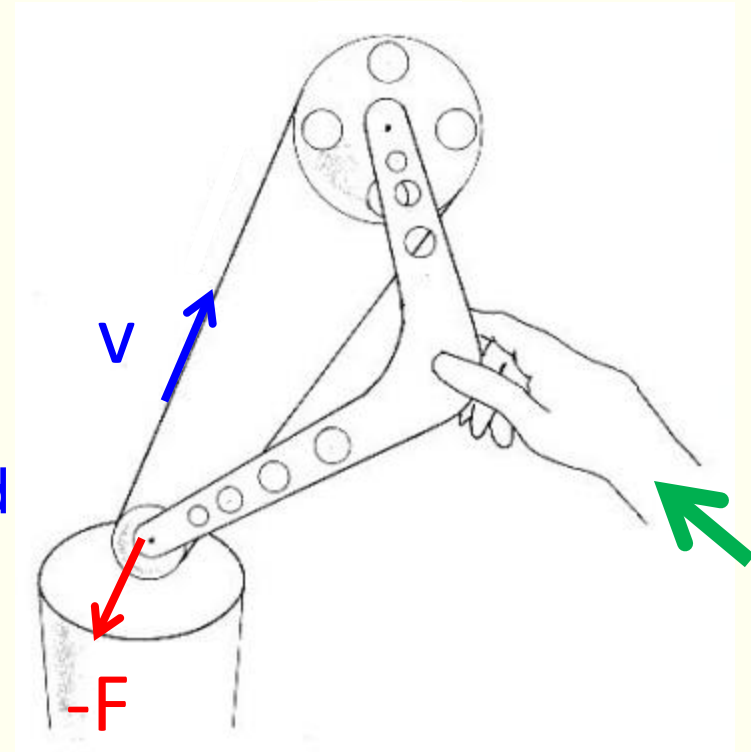


- no rotation without friction between filament & wheel!



# Angular momentum current

- point torque in filament causes constraint force  $F$  to appear\*
- reaction force  $-F$  appears; exerts torque on hand
- angular momentum current flows into bend motor from hand



- \* via stress transport: 
$$\nabla(\nabla \cdot \boldsymbol{\sigma}) = \frac{\rho}{E} \frac{\partial^2 \boldsymbol{\sigma}}{\partial t^2}$$



# Summary

- detailed model for bend motors driven by
  - local spontaneous curvature caused by light (LCE-Ikeda)
  - local stiffening caused by heat (Nitinol - Wang)
- corrects/completes existing qualitative explanations
- prediction:
  - LCE motor with perp. alignment runs the other way:  
confirmed!
- understand angular momentum transport
- efficiency calculation: remains to be done



# Angular momentum

stimulus is applied

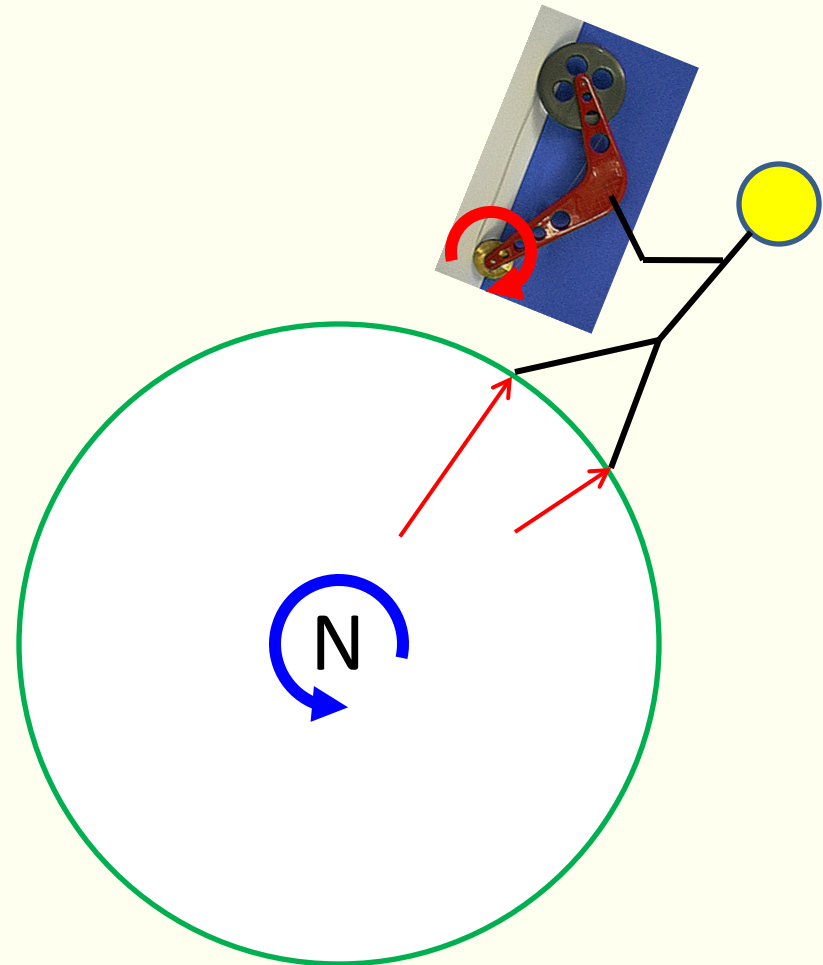
filament bends

torque appears

angular momentum flows into motor

rotation starts

angular momentum is conserved &  
earth speeds up



# Conclusions

- variety of molecular -scale mechanisms can be exploited
  - rotary (dyes, Yokoyama-Tabe)
  - need better molec. motors with processive translation
  - elongation (LC elastomers/networks)
- symmetry plays important role in determining motion
- macroscopic phenomena are cumulative effects of molecular motors
- understand bend motors well
- many exciting possibilities to explore!

