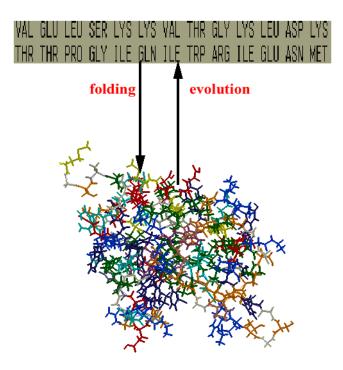
Proteins can fold in vivo and in vitro (Anfinsen' 59): protein folding problem

PROTEIN FOLDING: THE INTELLECTUAL CHALLENGE.

HOW IS BIOLOGICAL 1d ORDER (INFORMATION) TRANSLATED INTO PHYSICAL 3d ORDER (STRUCTURE)?



THEORETICAL EFFORT IN PROTEIN FOLDING: STATISTICAL-MECHANICAL STUDIES.

I Homopolymer microscopic models: ('64-'82)

```
Ptistyn'64 (Flory-type), Lifshitz'68 (rigorous
mean-field), DeGennes'75, Lifshitz, Grosberg, Khohlov'79, ES and
A.Finkelstein'82 (theory of molten globules:homopolymers with
side-chain freezing)
```

II Phenomenological Model of N.Go ('69-)

```
Go'69'72- lattice models with strong bias
''by hands'' to the native state (''Principle of maximal
consistency''), Jernigan'72, Dinner and Karplus'95, Pande nad
Rokshar'97.
```

III Exact theory (Zimm-Bragg) of local interactions:

Flory'58, Volkenstein'59, Scheraga'64-72, Finkelstein'75-82, Serrano'94....

IV. Phenomenological ''landscape'' approach

(Bryngelson and Wolynes'87,89'90; Wolynes and coworkers '92-) An assumption about density of states is made at the beginning and its implications are explored. ''Principle of minimal frustartions'' suggested.

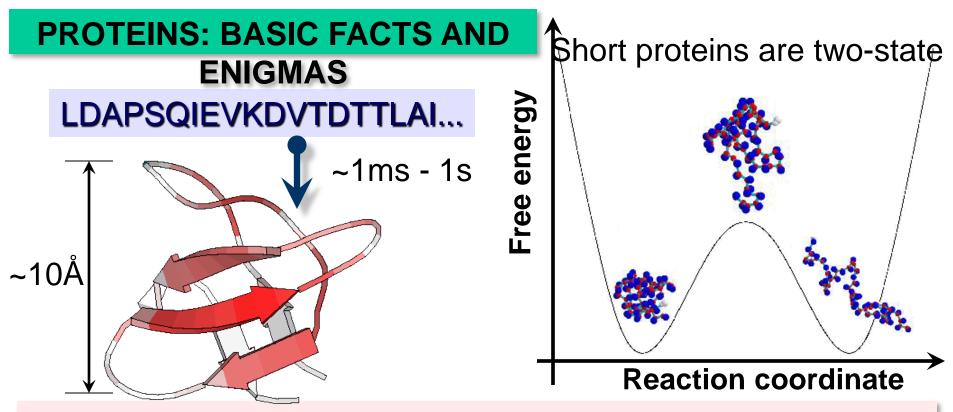
V Phenomenological heteropolymer model: Dill'85

VI.Microscopic Models of heteropolymers:analytical studies using replica approach.

```
The study starts from a microscopic hamiltonian. Density of states
(''landscape'') emerges as a result, not as assumption.
Garel and Orland'88, ES and A.Gutin'89,
Sasai and Wolynes'91, ES and Ramanathan'94, Grosberg and Pande'95,
ES and Archontes'94, ..... - still very active.
```

```
VII Lattice model simulations: Ball and Fink, Broglia,
Dill, Goldstein, Grosberg, Karplus, Onuchic, ES, Skolnick,
Tang, Thirumalai, Unger, Wolynes....
```

IN 1994 THESE APPROACHES WERE NOTED BY BIOCHEMISTS AND CALLED ''A NEW VIEW''



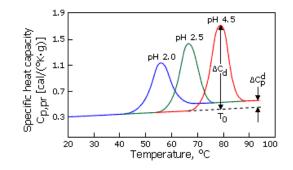
1. Protein sequence uniquely defines protein native (ground state) structure

2. Folding transition is cooperative (resembling first order transition)

3. Native state is thermodynamically stable

4. Native state is kinetically accessible --- reachable in a biologically reasonable time

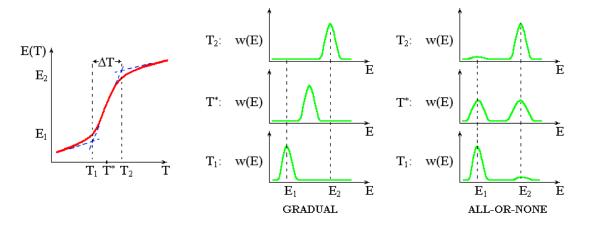
Calorimetry: an important experimental test of cooperativity



Calorimetric study of the lysozyme heat denaturation at various pH. The position of the heat capacity (C_p) peak determines transition temperature T_0 , the peak width gives the transition width ΔT , and the area under the peak determines heat ΔH absorbed by a gram of the protein. The values ΔT , $\Delta H \times \text{protein's}_M.W.$, and T_0 satisfy van't Hoff equations indicating hat the denaturation occurs as an *"all-or-none" (first order) transition*. The increased heat capacity of the denatured protein (ΔC_p) originated from the enlarged interface between its hydrophobic groups and water after denaturation.

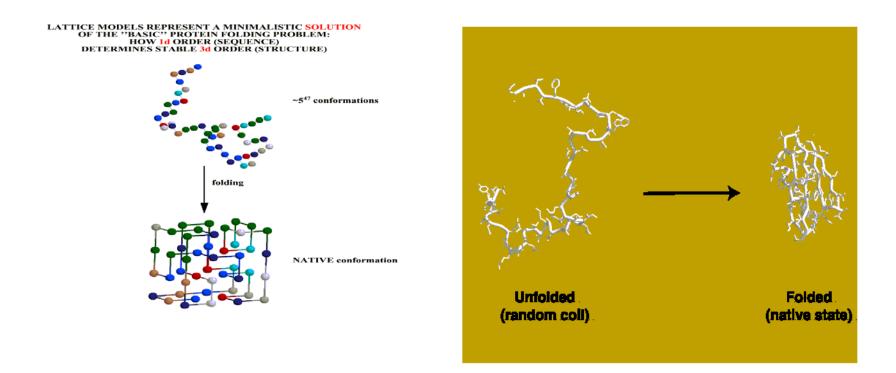
Adapted from P.L.Privalov & N.N.Khechinashvili, J. Mol. Biol. (1974) 86:665-684.

"All-or-none" first-order like transitions: is Sshape of a transition curve a good evidence?



With the same temperature dependence of energy E (or any other observable parameter), the cooperative ("S-shaped") transition can be either of the "all-or-none" type (example: protein denaturation), or gradual (example: helix-coil transition in polypeptides). The difference is displayed in the shape of the function W(E) showing distribution of molecules over energy (or over any other observable parameter) rather than in the shape of the curve E(T). Dashed lines in the left drawing explain a graphical determination f Δ T, the width of temperature transition.

Important remark – on protein models



Proper selection of a model is a key to success! It depends on the Question you want to ask!

PHYSICAL QUESTIONS:

A protein stays in its ground (micro) state at room temperature at which characteristic energy scale of all interactions is ~kT. (Thermodynamic uniqueness of protein structure)

Q1:

a) Is it a generic feature of any (i.e. random) sequence or some specially selected (designed) sequences only can fold?

b) What features should be optimized to achieve thermodynamic uniqueness of protein structure?

c) How many sequences can fold thermodynamically?

Q2:

Obviously, proteins can find their ground state much faster than by exhaustive search (Kinetic accessibility of the native state).

a) Which sequences can fold fast?

b) What is the kinetic mechanism(s) of folding? – e.g. hierarchical or nucleation? (i.e. Second-order like vs First-order like)

Heteropolymer Hamiltonian: a basic model for study protein folding thermodynamics and kinetics

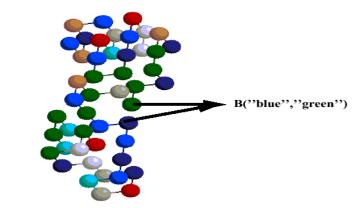
The partition function:

 $Z=\Sigma_{conf} exp(-H_v(seq,conf))\prod_i g(r_{i+1}-r_i)$

g(r_{i+1}-r_i) describe local interactions along the sequence.

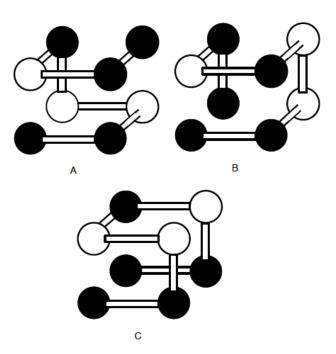
The nonlocal Hamiltonian:

 $H_{v}(seq,conf) = \sum_{ij} B(\sigma_{i},\sigma_{j})U(r_{i}-r_{j})$ $\{\sigma_{i}\} describes polymer sequence$



Frustration in Polymers

• In a 2 × 2 cube there are only three different conformations



• A quenched sequence of monomers leads to frustration even for the two-letter case.

Effect of Frustration

• In a homopolymer all different conformations with the same total number of contacts have exactly the same energy.



- On a lattice with coordination number z the number of conformations is $(z/e)^N$
- In a frustrated system all conformations have unsatisfied contacts.
- The various conformations have different energies
- Freezing. Thermodynamic preference.

 $/ \setminus /$

The Polymer Hamiltonian

• Interactions between monomers in space

$$U_{ij} = B_{ij}\delta(\mathbf{r}_i - \mathbf{r}_j)$$

• In the two-letter problem (an A, B random sequence) B_{ij} will be given by

$$B_{ij} = \chi \sigma_i \sigma_j + A \left(\frac{\sigma_i + \sigma_j}{2} \right) + B_0$$

$$B_{AA} = B_0 + A + \chi$$

$$B_{BB} = B_0 - A + \chi$$

$$\sigma_i = \pm 1$$

$$B_{AB} = B_0 - \chi$$

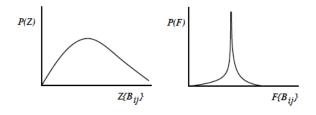
$$\chi = \frac{B_{\rm AA} + B_{\rm BB}}{2} - B_{\rm AB}$$

• Include polymeric bonds via

$$g(\mathbf{r}_{j+1} - \mathbf{r}_j) = \frac{1}{(2\pi a^2)^{3/2}} \exp\left[-\frac{(\mathbf{r}_{j+1} - \mathbf{r}_j)^2}{2a^2}\right]$$

Averaging over Randomness

- Need to average over disorder.
- Self-averaging.
- Rare sequences with low energies (e.g. homopolymer) contribute significantly to the partition function.
- The contribution of these sequences to free energy is minimal.



• Replica Trick

$$\langle \ln Z \rangle_{av} = \lim_{n \to 0} \frac{1}{n} \left(\left[Z^n \right]_{av} - 1 \right)$$

Main steps in analytical theory (just an outline: for more detail see 2006 ChemRev)

The key issue: How to average over sequences? Prescription: Consider only self-averaging (SA) quantities – whose average is representative of majority of realizations. Bad news: partition function is not SA (why???) (Sort of) Good news: Free energy (i.e. log of Z) is SA!!!!

$$\langle F(T) \rangle = -kT \sum_{\{\sigma\}} P(\{\sigma\}) \log(Z(\{\sigma\},T))$$

Where summation is taken over all 20^N sequences
and P({s}) is the probability to find a sequence {s} in the ensemble

 $P_r(\{\sigma\}) = P(\{\sigma\}) = \frac{1}{20^N}$ For Random Sequences!

However, for evolutionary selected sequences $P(\{\sigma\}) \neq P_r(\{\sigma\})$ Therefore: Evolution enters theory via $P(\{\sigma\})$!

ES and Gutin, Europhys Lett' 89, ibid Biophys Chem' 89, ibid JPhysA' 89 (d<=2 case)

Order Parameters

• Density of monomers

$$\rho(\mathbf{R}) = \sum_{i} \delta(\mathbf{r}_{i} - \mathbf{R})$$

• Phase separation of A from B monomers.

$$m(\mathbf{R}) = \rho_{\mathbf{A}} - \rho_{\mathbf{B}} = \sum_{i} \sigma_{i} \delta(\mathbf{r}_{i} - \mathbf{R})$$

• Overlap between conformations α and β

$$Q_{\alpha\beta}(\mathbf{R}_1, \mathbf{R}_2) = \sum_i \delta(\mathbf{r}_i^{\alpha} - \mathbf{R}_1) \delta(\mathbf{r}_i^{\beta} - \mathbf{R}_2)$$

with

$$\int d\mathbf{R}_1 Q_{\alpha\beta}(\mathbf{R}_1, \mathbf{R}_2) = \rho_\beta(\mathbf{R}_2)$$

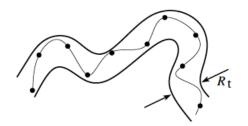
• For maximally compact state

$$Q_{\alpha\beta}(\mathbf{R}_1, \mathbf{R}_2) = Q_{\alpha\beta}(\mathbf{R}_1 - \mathbf{R}_2)$$

• $Q_{\alpha\beta}$ can be expressed as

$$Q_{\alpha\beta}(\mathbf{R}_1 - \mathbf{R}_2) = \frac{\rho}{R_t^3} \varphi_{\alpha\beta} \left(\frac{\mathbf{R}_1 - \mathbf{R}_2}{R_t}\right)$$

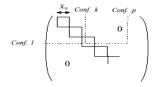
• Polymer in a tube



Solution for the Q-order parameter

Results for $Q_{\alpha\beta}$

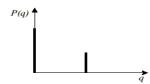
• The form of the solution for $Q_{\alpha\beta}$ gives a map of common contacts between different conformations.



The matrix elements are

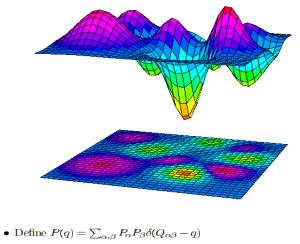
$$Q_{\alpha\beta} = \begin{cases} \rho \delta(\mathbf{R}_1 - \mathbf{R}_2) & \text{for } \alpha, \beta \text{ in the same group} \\ 0 & \text{for } \alpha, \beta \text{ in different groups} \end{cases}$$

- Conformations of different groups have no overlap
- In the frozen phase $P(q) = x_0 \delta(q) + (1 x_0) \delta(q 1)$



Potential energy surface for Random Heteropolymers

Potential Energy Surface



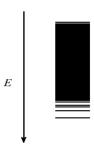
• At high temperature $P(q) = \delta(q)$

Density of states (spectrum) of the RHP

Energy spectrum

• Resulting energy spectrum and the R.E.M.

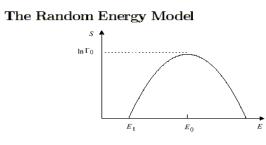
$$H = \sum_{ij} B_{ij} \delta(\mathbf{r}_i - \mathbf{r}_j)$$



 $P(q) = x_0 \delta(q) + (1 - x_0) \delta(q - 1)$

• Conformations in the bottom of the spectrum are entirely different

Random Energy Model



• The energy levels follow a random distribution

$$P(E) = \Gamma_0 \exp[-(E - E_0)^2 / \sigma^2]$$

• The entropy for each level is given by

$$S(E)=\ln\Gamma_0-\frac{(E-E_0)^2}{\sigma^2}$$

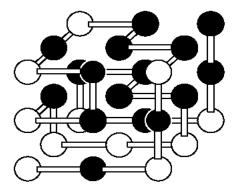
- Freezing occurs when $S\sim 0$
- The transition temperature is given by the slope

$$\frac{1}{T} = \frac{dS}{dE}$$

at E_t

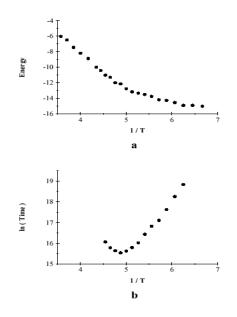
Lattice Model to test of the results of Random Heteropolymer Theory

Lattice Model



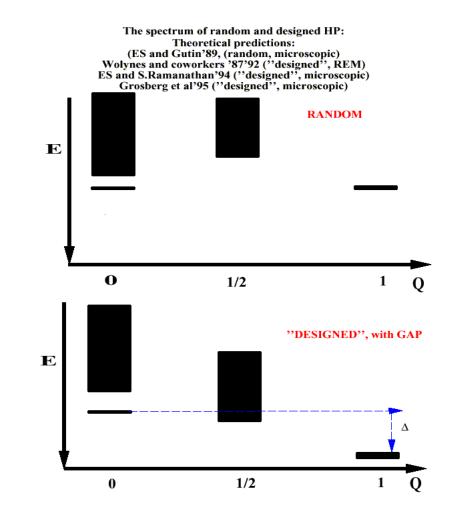
- $\Gamma_0 = 103346, K = 28, B = -1, B_0 = 1$
- $E_{\text{conf}} = -(K_{bb} + K_{ww})$

Lattice Model analysis of RHP: the results



- Experiment: Quasirandom sequence of F2 fragment of tryptophane synthase (Chafotte *et al.*, 1991) the transition is noncooperative.
- Experiment: "Binary design" (M.Hecht et al, 1995)

Why is Random HP a bad model of protein: The case for Evolution



This picture is valid only for d>2

The '' freezing'' transition in RHP is non-cooperative 22

Solution: Designed Sequences!

IS RANDOM HP A GOOD MODEL FOR PROTEIN:

- RHP Exhibit thermodynamically noncooperative "freezing" transition.
- Folding transition in proteins is cooperative (Privalov '73,'79..).

A POSSIBLE SOLUTION: DESIGNED SEQUENCES

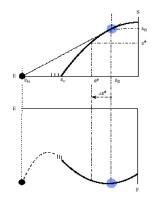


Fig.1a

WARNING: THIS IS A CARTOON. IT MAY BE VALID FOR A "TYPICAL" 3D CONFORMATION, BUT NOT FOR ALL CONFORMATIONS AND NOT FOR $D \le 2$

Statistical Mechanics of Protein-Like Sequences with Gaps

- Now we understand (or anticipate) that protein-like sequences should be selected to have large energy gap to the native state
- Thus the sequence ensemble (i.e. $P(\{\sigma\})$ should be shifted accordingly:

$$P_{T_{sel}}(\{\sigma_i\}) = \exp\left(-\frac{H(\{\sigma\},\{r^0\})}{T_{sel}}\right) = \exp\left(-\frac{\sum_{i< j} B(\sigma_i,\sigma_j)U(r_i^0 - r_j^0)}{T_{sel}}\right)$$

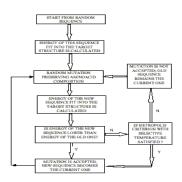
Where $\{r^0\}$ determines native conformation, T_{se} is a measure of evolutionary selection

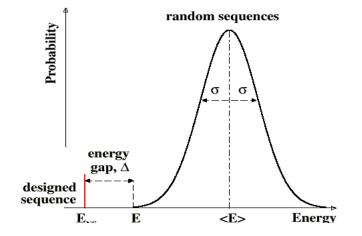
Statistical mechanics *and dynamics* in Sequence Space: Evolution (See lecture 1)²⁴

Algorithm of design: Monte-Carlo in Sequence Space

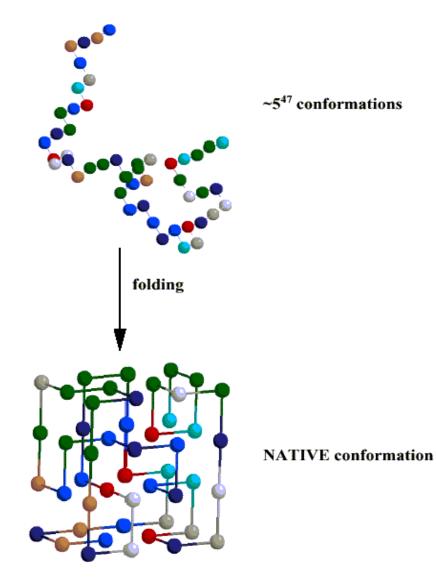
Algorithm of Selection & Design

- Microcanonical ensemble vs
- Canonical ensemble with evolutionary temperature: MC procedure in sequence space:
- $E = E_0(\{\sigma_i\}, \{r_i^0\})$
- $\{\sigma_i\}$ is a sequence $(\sigma_i = 1...20$ is a type of monomer i)

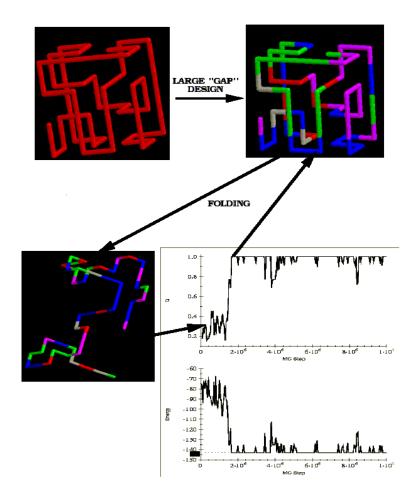




LATTICE MODELS REPRESENT A MINIMALISTIC SOLUTION OF THE "BASIC" PROTEIN FOLDING PROBLEM: HOW 1d ORDER (SEQUENCE) DETERMINES STABLE 3d ORDER (STRUCTURE)

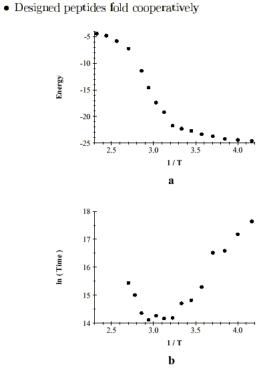


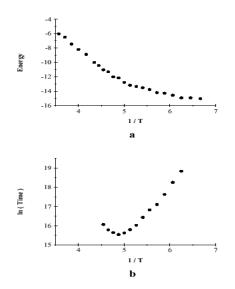
Testing the design idea in simulations: the design-folding paradigm



Design works

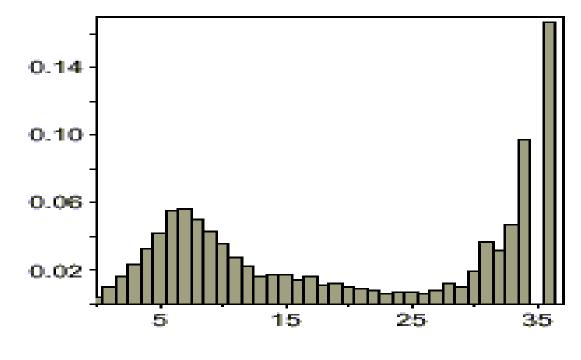
THE DESIGN WORKS:





- Experiment: Quasirandom sequence of F2 fragment of tryptophane synthase (Chafotte *et al.*, 1991) the transition is noncooperative.
- Experiment: "Binary design" (M.Hecht et al, 1995)

Cooperativity of "designed" (evolutionary selected) sequences

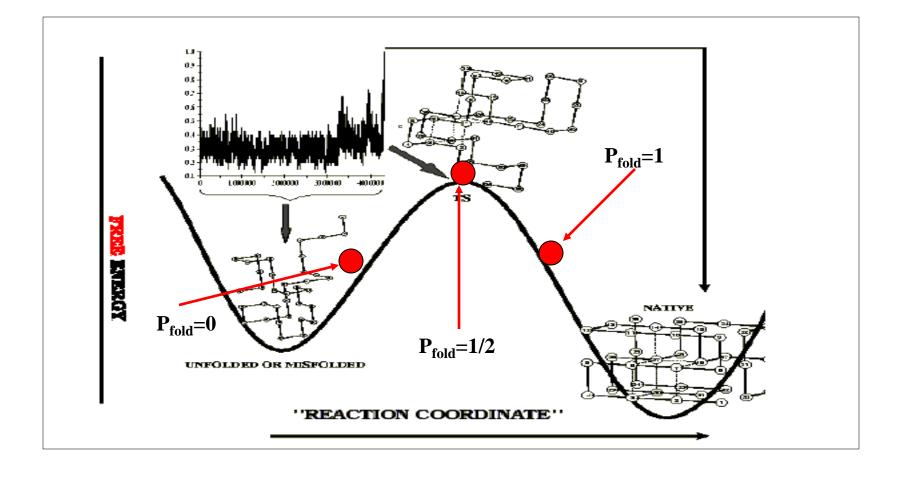


NUMBER OF NATIVE CONTACTS

Q: Why Does Large Gap Solve Kinetic Protein Folding?

A: It Provides **Nucleation** (akin to first order phase transition kinetics)

Where is folding Transition State on a trajectory and in microscopic picture?

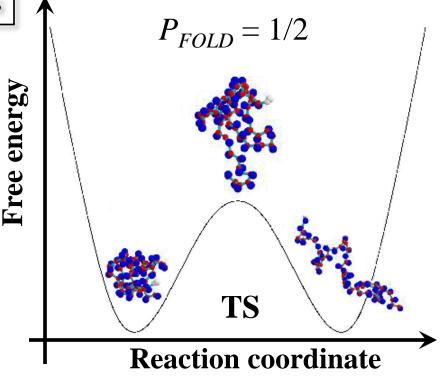




Nucleation scenario – formation of obligatory contacts between amino acids in the *transition state*.

GOAL:

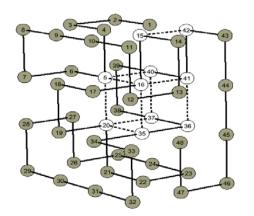
- Test nucleation scenario
- Identify transition state ensemble



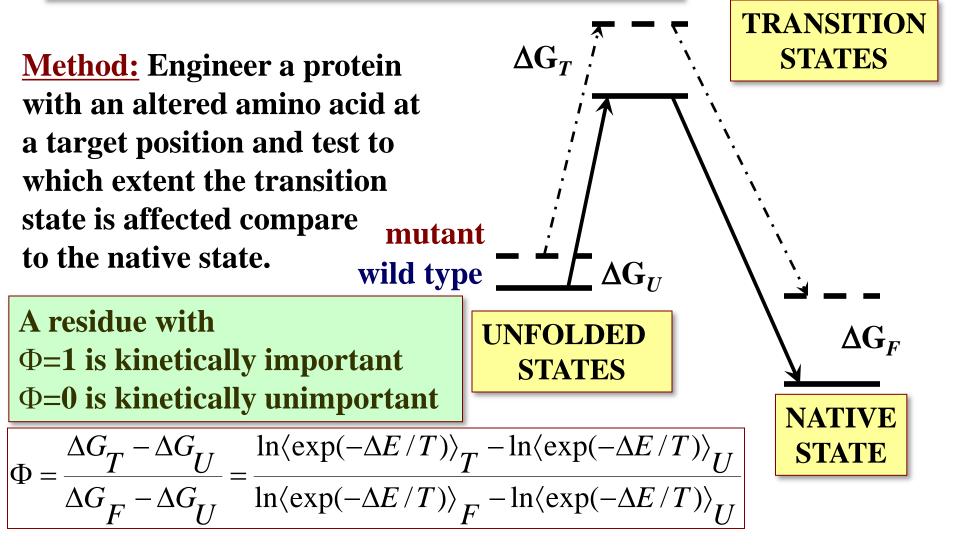
Shakhnovich et al., *Nature* **379**, 96-98 (1996) Abkevich et al., *Biochemistry* **33**, 10026-10036 (1994) Dokholyan et al., *J. Mol. Biol.* **296**, 1183--1188 (2000)

Nucleation mechanism: Formation of set of Obligatory contacts (Folding Nucleus) that define TSE

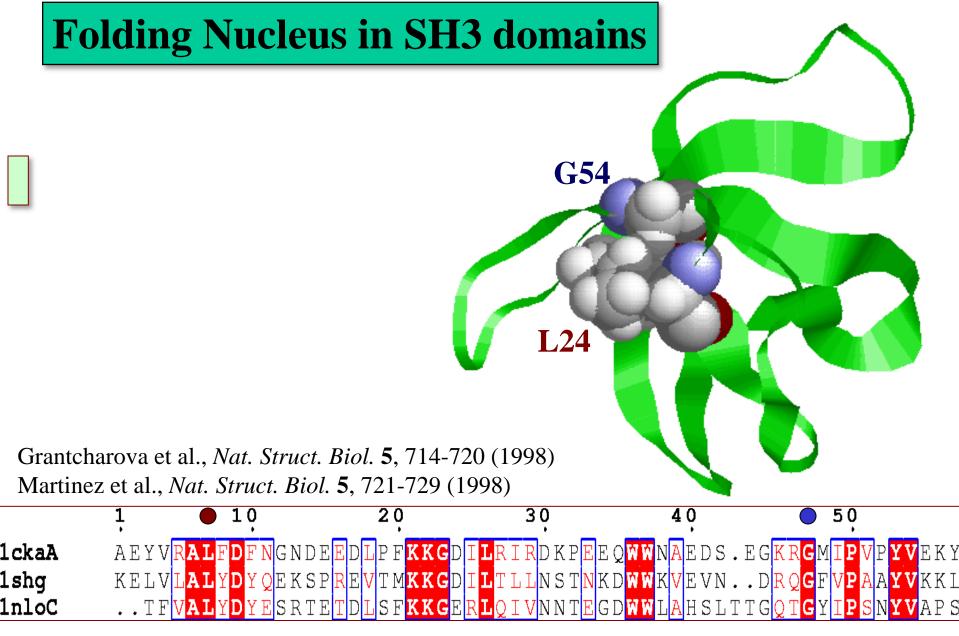
(Abkevich et al' Biochemistry'94)



PROTEIN ENGINEERING: Φ-VALUE ANALYSIS



Fersht, Curr. Opin. Struct. Biol. 7, 3-9 (1997)



Ding et al., *Biophys. J. submitted* (2002) Borreguero et al., *J. Mol. Biol. in press* (2002)

Evolutionary control of folding rates and stability

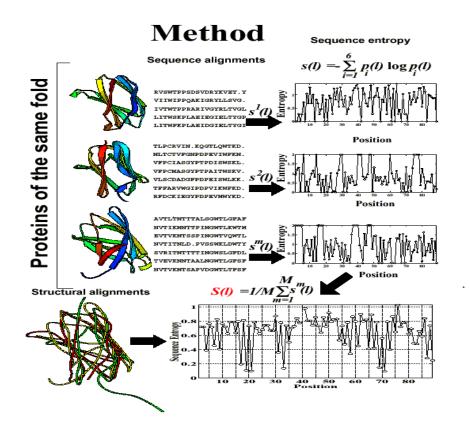
The idea: Nucleus residues may determine folding rate

>Therefore if evolution cared about folding it could have exerted extra pressure on nucleus residues.

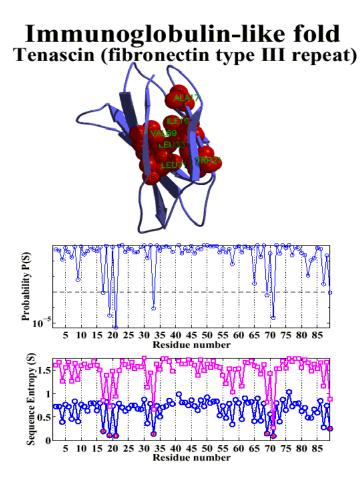
>How to select "folding pressure" on the background of many other, strong pressures (function, stability...)?

>Look at proteins featuring similar structure but diverse functions!

Conservatism-of-Conservatism: The Idea

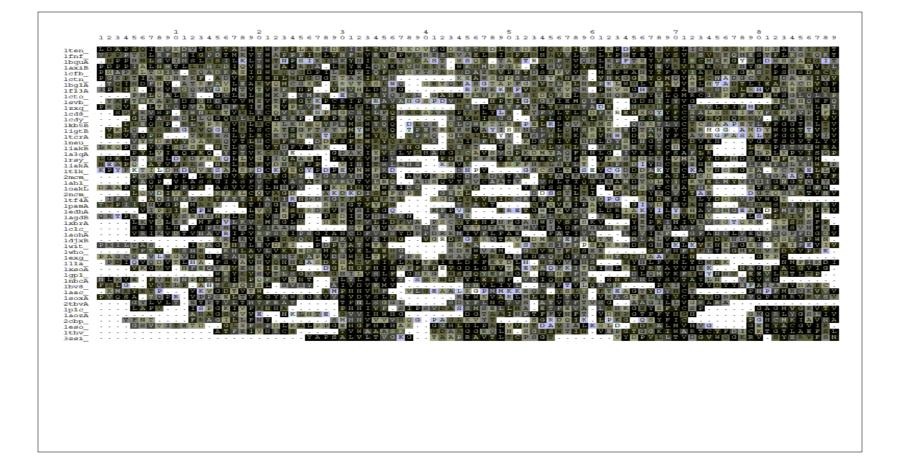


CoC Analysis correctly predicts folding nucleus in Ig-fold proteins

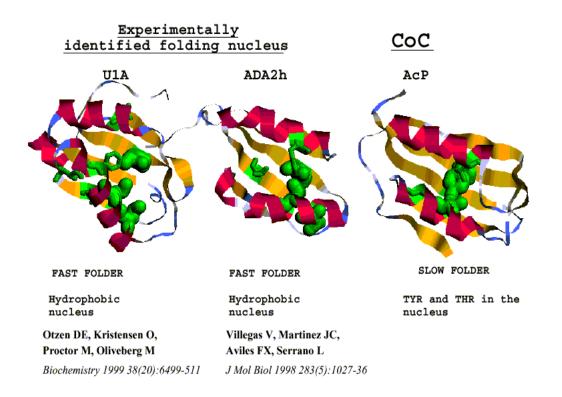


Experiment: J. Clarke and coauthors (2001)

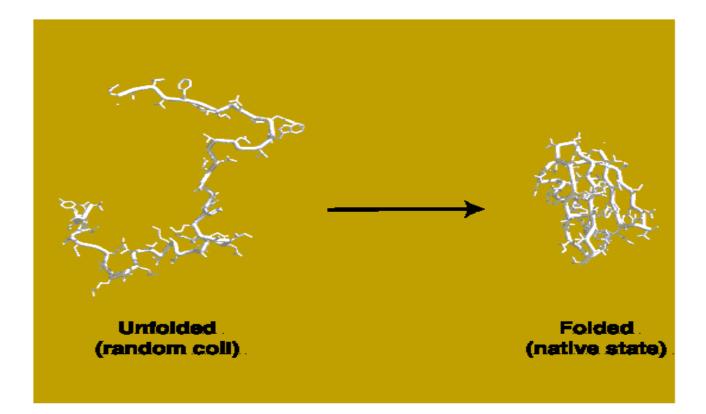
Example: Ig-Fold proteins



Understanding diversity of protein folding rates....

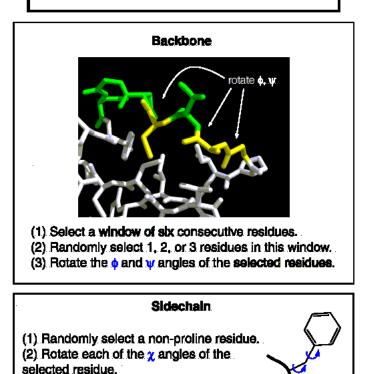


An all-atom Monte-Carlo Folding Simulation



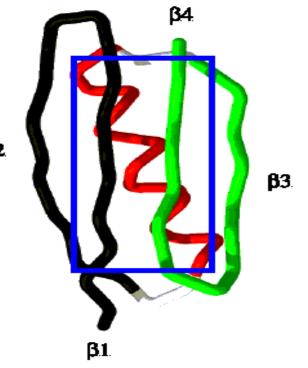
Move set

1 MC step = 1 backbone move + 10 sidechain moves.



All angle sizes are chosen from a normal distribution with mean zero, and standard deviation 2 for backbone and 10 for sidechains.

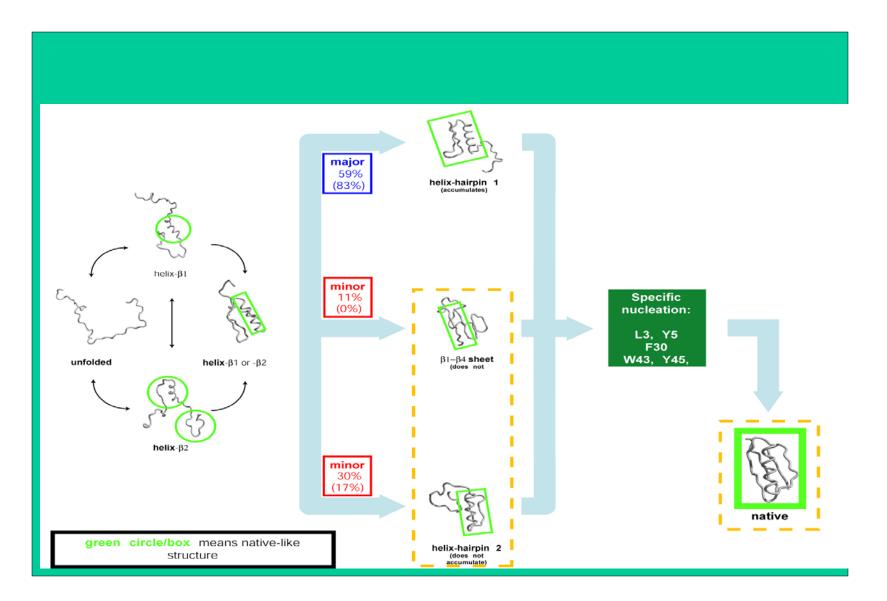
Protein G – Folding of a small protein in allatom detail – Go model



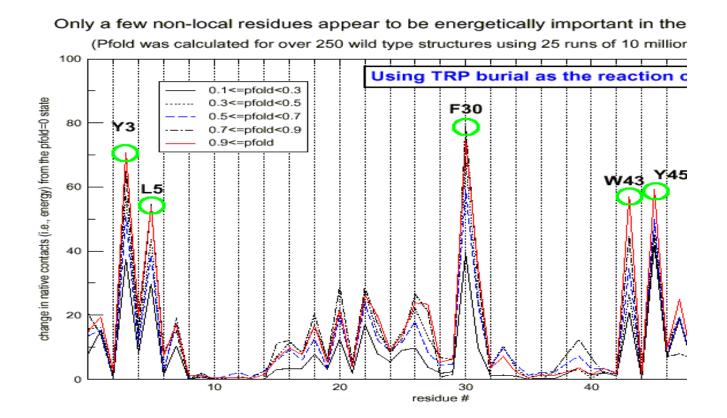
Black – first beta-hairpin
Red – alpha helix
Green – second beta hairpin

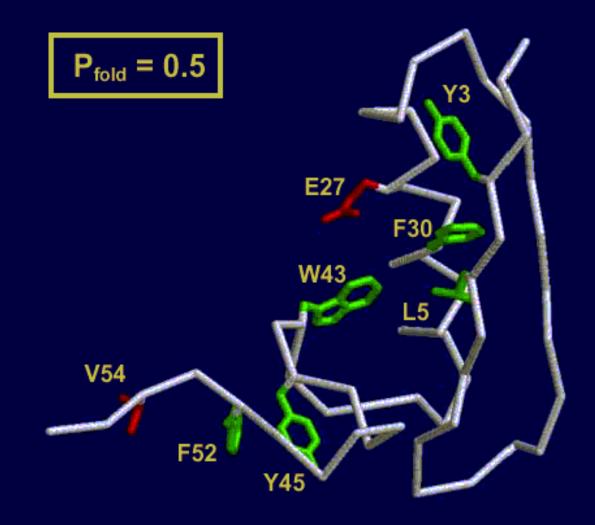
β2.

Protein G folding pathways: summary



What about TSE in protein G? Same protocol with P_{fold} identifies states that are committed to fold very fast, "downhill", in less than 10⁷ steps





green = important in WT red = important also in mutant

How to Fold A Protein?

- From Sequence to Structure (i.e. non-Go)
- All Atom
- Low RMSD....

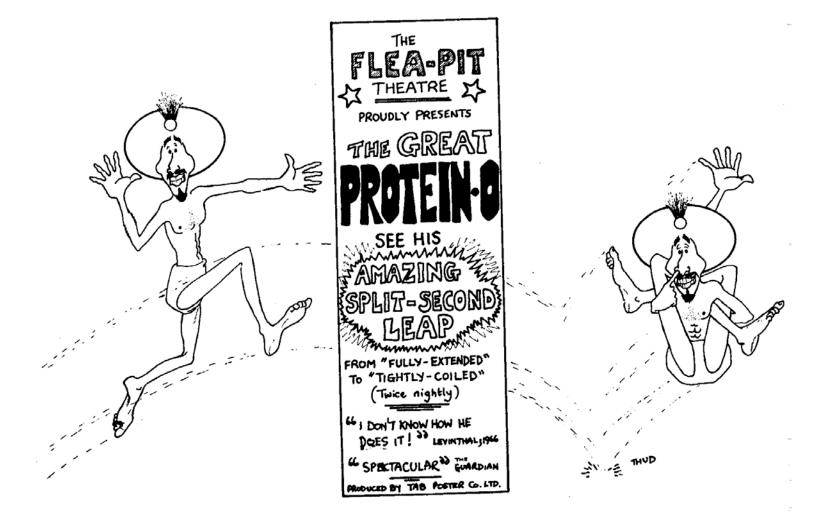
(wishlist....) Approach:

All-Atom Statistical Potentials (2-body+HB+Local) Kussell,ES PNAS' 02, Hubner, Deeds, ES, PNAS' 05

Dynamically balanced All-Atom MC: global+local BB moves + sidechain moves

Shimada, Kussell,ES JMB' 01 Shimada, ES PNAS' 02 Hubner, Shimada, ES JMB' 04

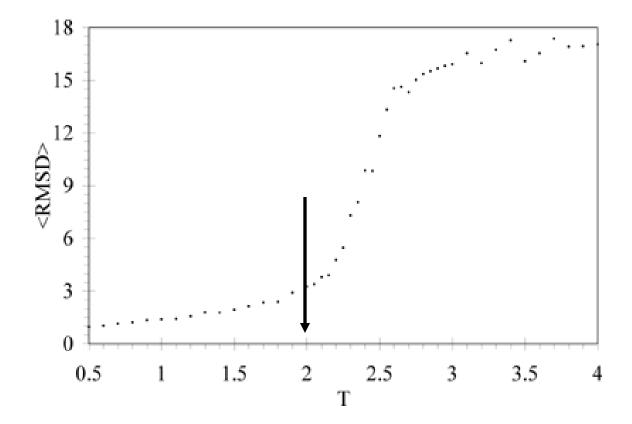
HOW DO PROTEINS FOLD? (all-atom, full detail)



Methods

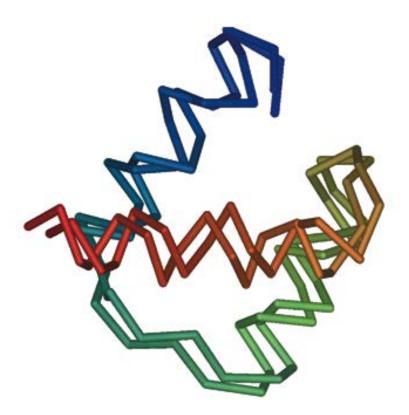
- 4000 folding runs from random coil chain
- At Constant T (NOT Replica Exchange!)
- TSE by p_{fold}
- Graph analysis of massive data (like PDUG etc.) instead of projecting on "reaction coordinates"
 - E_{min} and folding graphs
 - Clustering (PDUGs) in multiple Order
 Parameters: a Multidimensional view

Folding at physiological T~25C



Identifying the native state

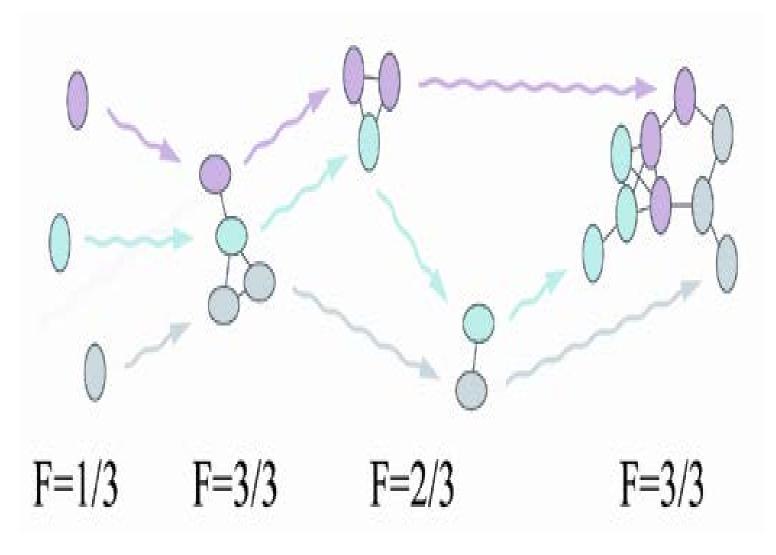
- Lowest E predition is 2.44Å (best of 4000)
- a < 3Å prediction is statistically significant with p-value ~10⁻⁷
- Of 4000, 44 trajectories sampled the ≤ 2Å range, 523 ≤ 3Å range, 1685 ≤ 4Å range, 2700 ≤ 5Å range, and 3331 ≤ 6Å range.
- This is consistent with usual exponential distributions of FPT



A network ensemble view folding

- Construct a *structural* (not kinetic, cf Caflisch) graph
- Allows combination of multiple trajectories
- Multidumensional view: Cluster Conformations based on various properties: RMSD, Rg, dRMS
- That will allow to fully characterize the folding mechanism, *while any single Order Parameter may be misleading*
- How to introduce Ensemble Kinetics into the graph description: *idea of flux!*

Flux:



Ensemble kinetics

- MFPT = $63 \cdot 10^6$ MC steps, so ~15% of runs would not be expected to fold (~ number of 4000 runs that did not sample < 6Å)
- relaxation from the initial random, swollen-coil conformation to a semi-collapsed state with some hydrogen bonding and non-specific hydrophobic interactions in ~8.10⁶ MC steps
- laser T-jump experiments also reveal two-step folding kinetics at 25° C, with fast ($t_{1/2} \approx$ 1.5microseconds) and slow ($t_{1/2} \approx$ 15microseconds) phases. (Fersht et al)

What is an intermediate?

- Accumulative intermediate vs intermediary state
- A structural graph can identify "parametrically" cohesive intermediates
- A non-GC cluster that is "visited" by multiple trajectories, i.e. having high flux=1

Is there an intermediate?

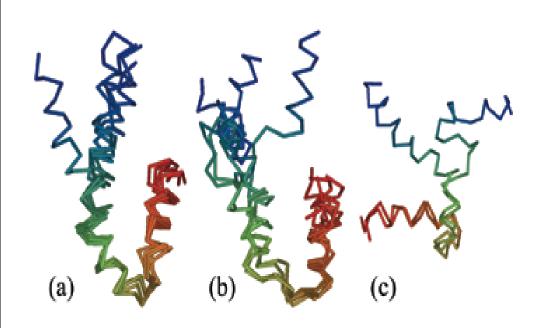
- Measure cluster size and flux (fraction of trajectories that "visit" cluster)
- Measure at all cutoffs from fully connected to fully unconnected
- Flux Analysis reveals one accumulating intermediate (Rg clustering) and one transient one (dRMS clustering)

Atomistically resolved structural intermediate

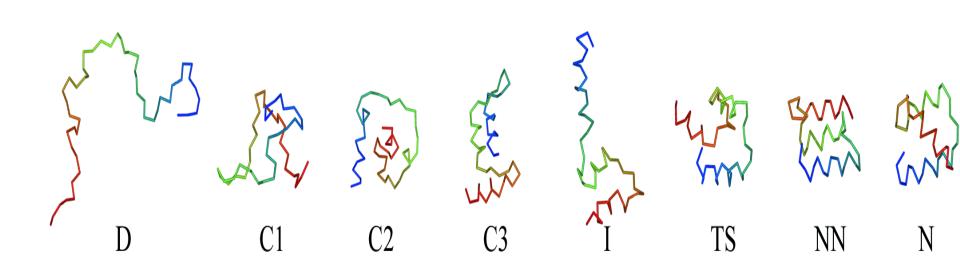
 Conformations of the dRMS intermediary cluster perfectly with NMR-derived structures of L16A model of the intermediate

(25 structures 1UZC)

• Average RMSD between dRMS cluster and 25 1UZC is 4.6A, some conformations as low as 1.5A



Folding Scenario: Summary



Important new approaches to simulations of folding kinetics

Markov Chain States:

Distributed computing generates many relatively short trajectories, the data are then collected and analyzed using markov state models: analogous to Monte-Carlo (Pande et al)

Simulations using specialized Anton computer (DE Shaw and colleagues)

Very long trajectories for several proteins (hundred microseconds) were obtained and multiple folding-unfolding events were observed and analyzed (Amber FF04 force-field). The details of folding mechanisms for villin headpiece depend on the force-field. (Science 2009,2011, Biophys Journal 2011)