

# Dynamics of Oceans: Pt 3: Small-Scales and Climate

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NOAA (NA19OAR4310366), Gulf of Mexico Research Initiative, Schmidt Futures

Boulder School for Condensed Matter and Materials Physics

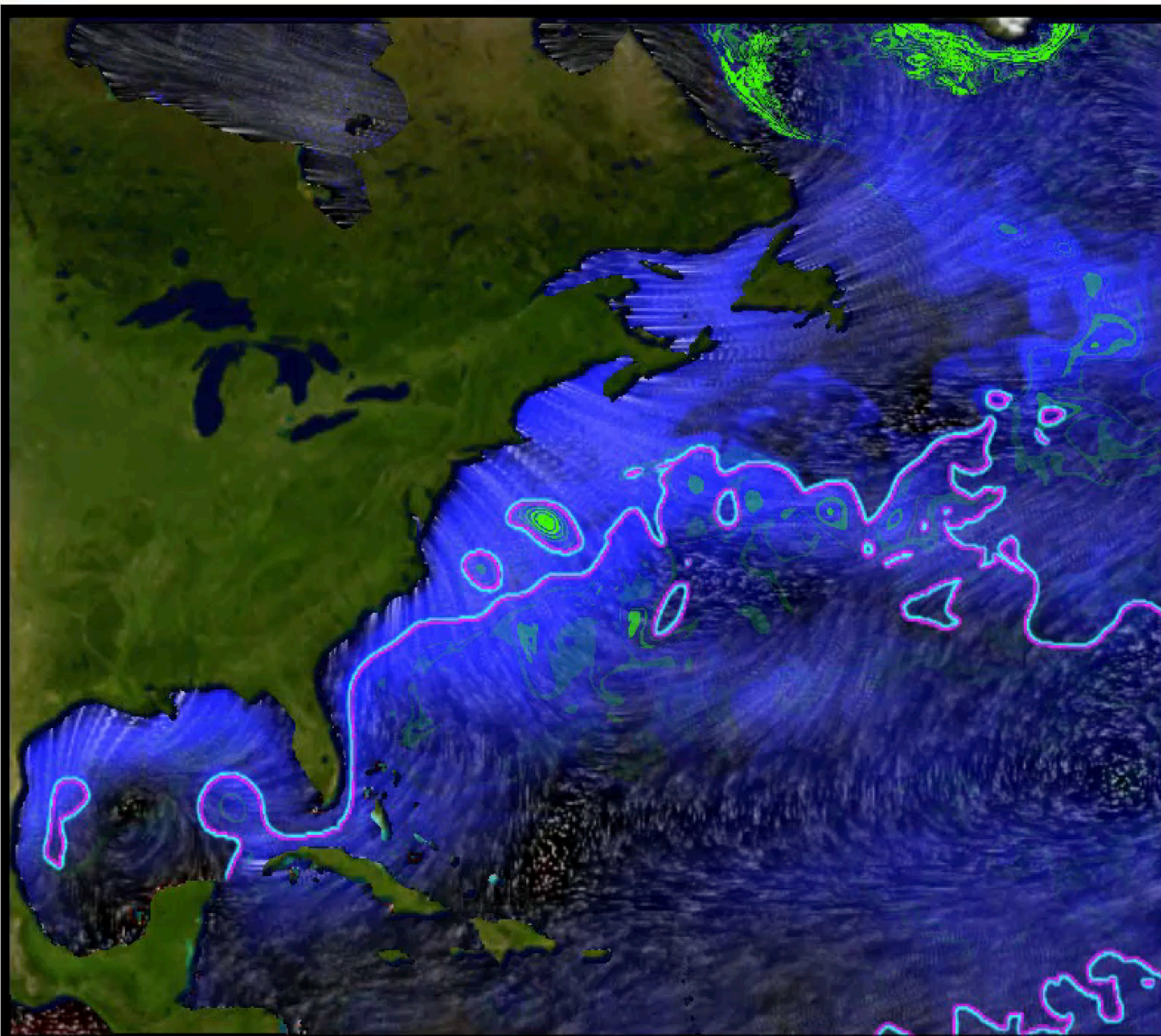
July 18, 2022



Weather,  
Atmosphere  
Fast

Ocean,  
Climate  
Slow

3.4m of ocean  
water has  
same heat  
capacity as  
the **WHOLE**  
atmosphere



ECCO Movie: Chris Henze, NASA Ames

tau / qflux / theta200m / kppMLD

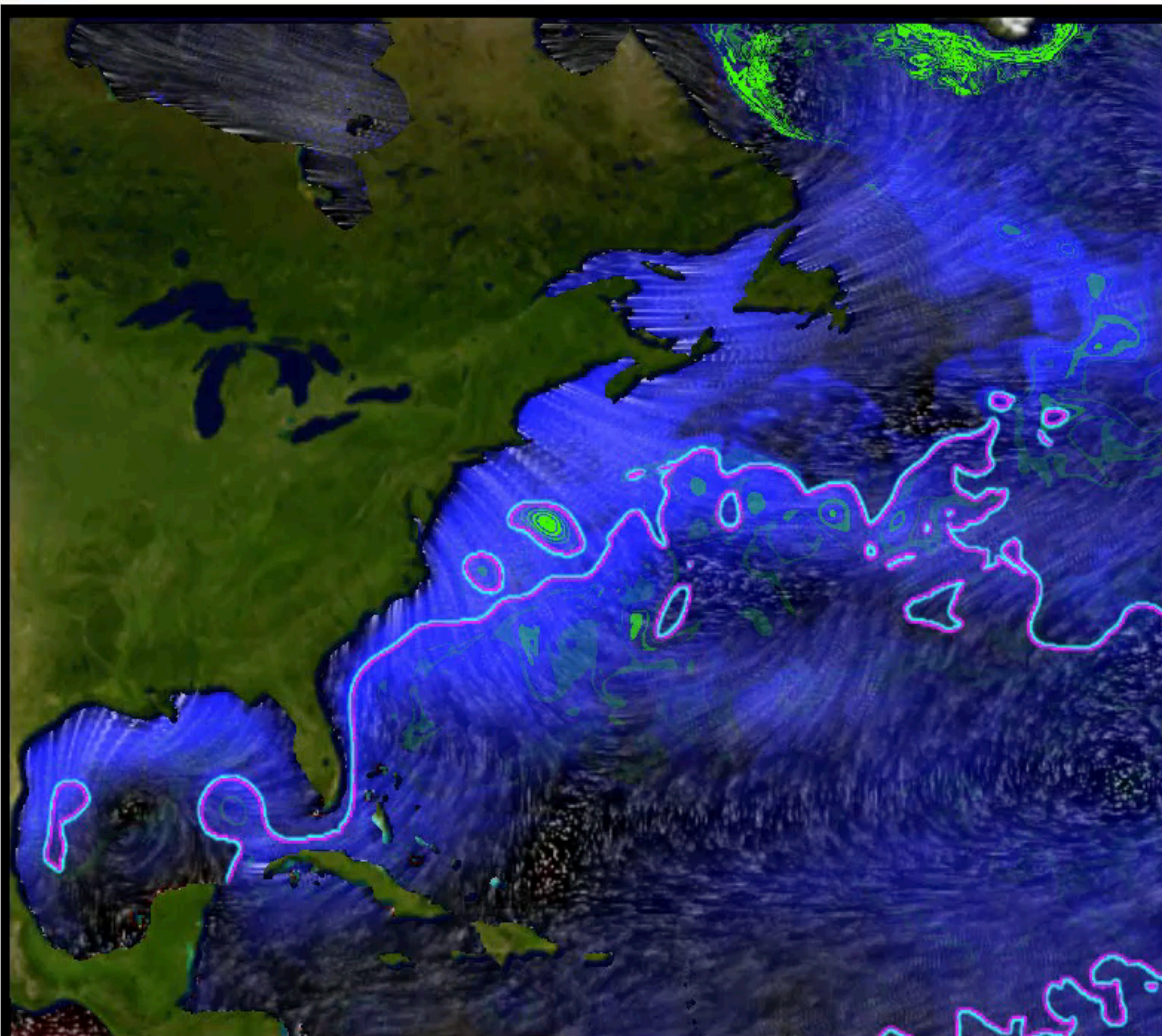
Jan 1 00:30 2001



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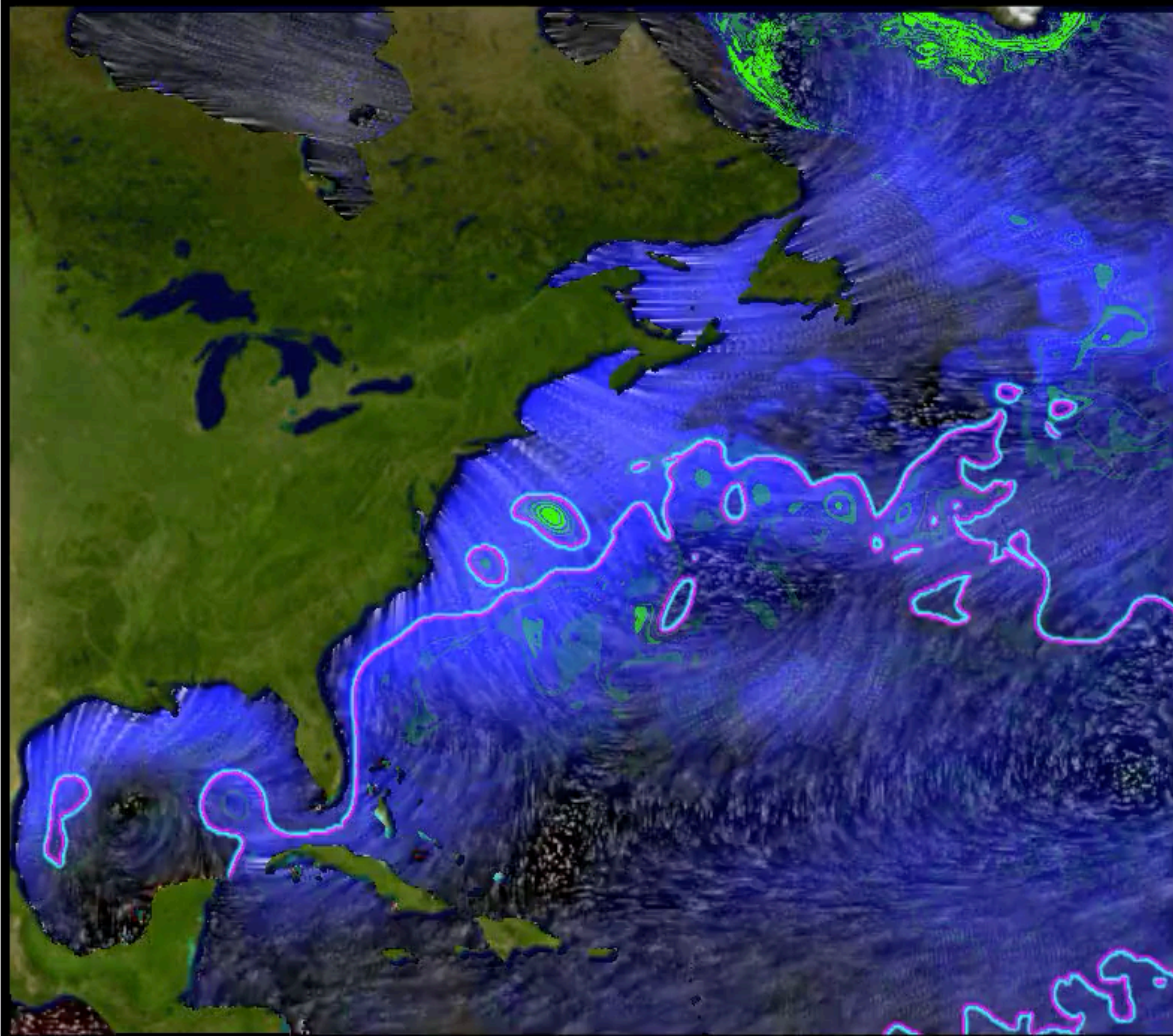


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The essence of the  
Hasselmann 2021  
Nobel in Physics!



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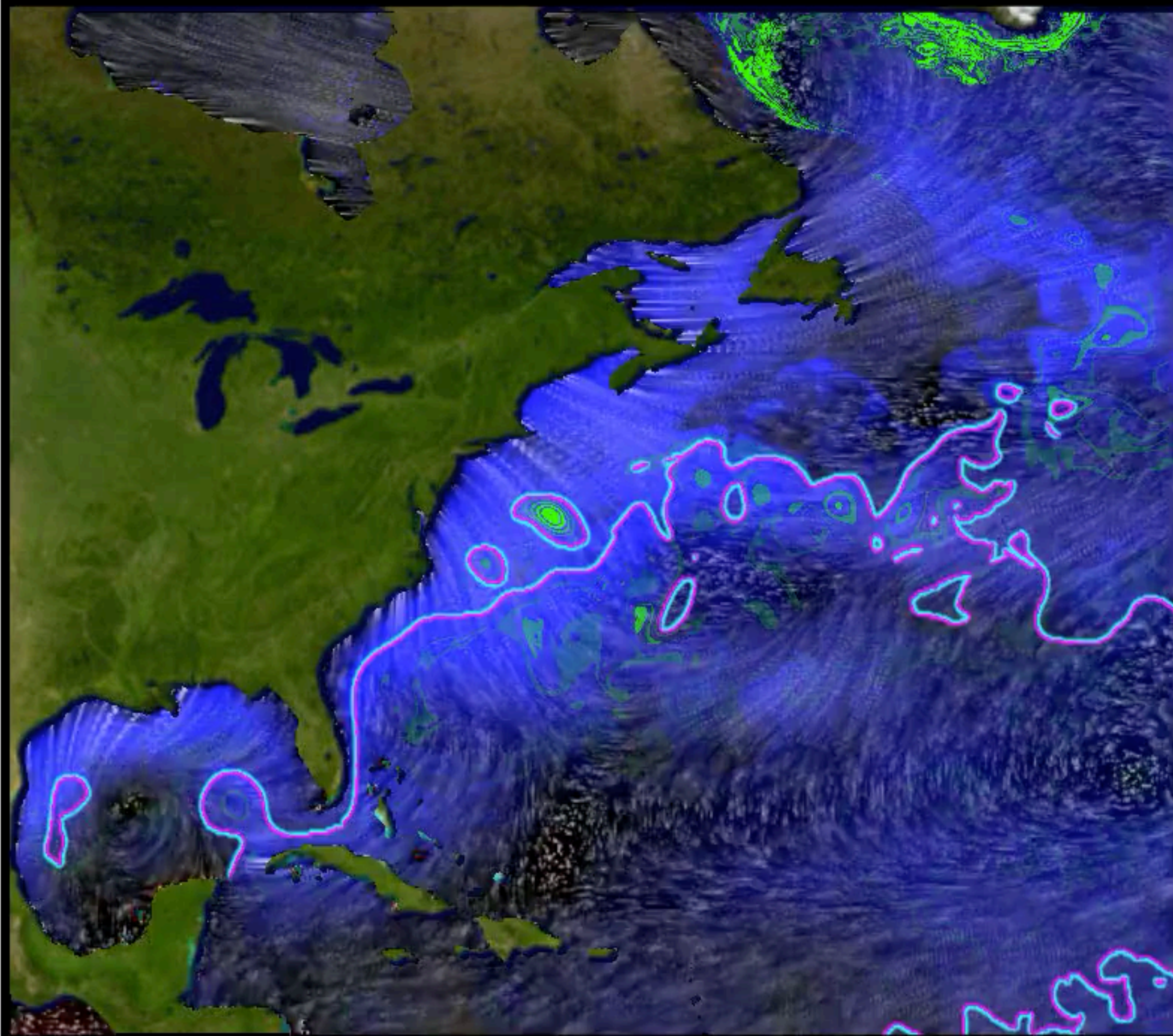


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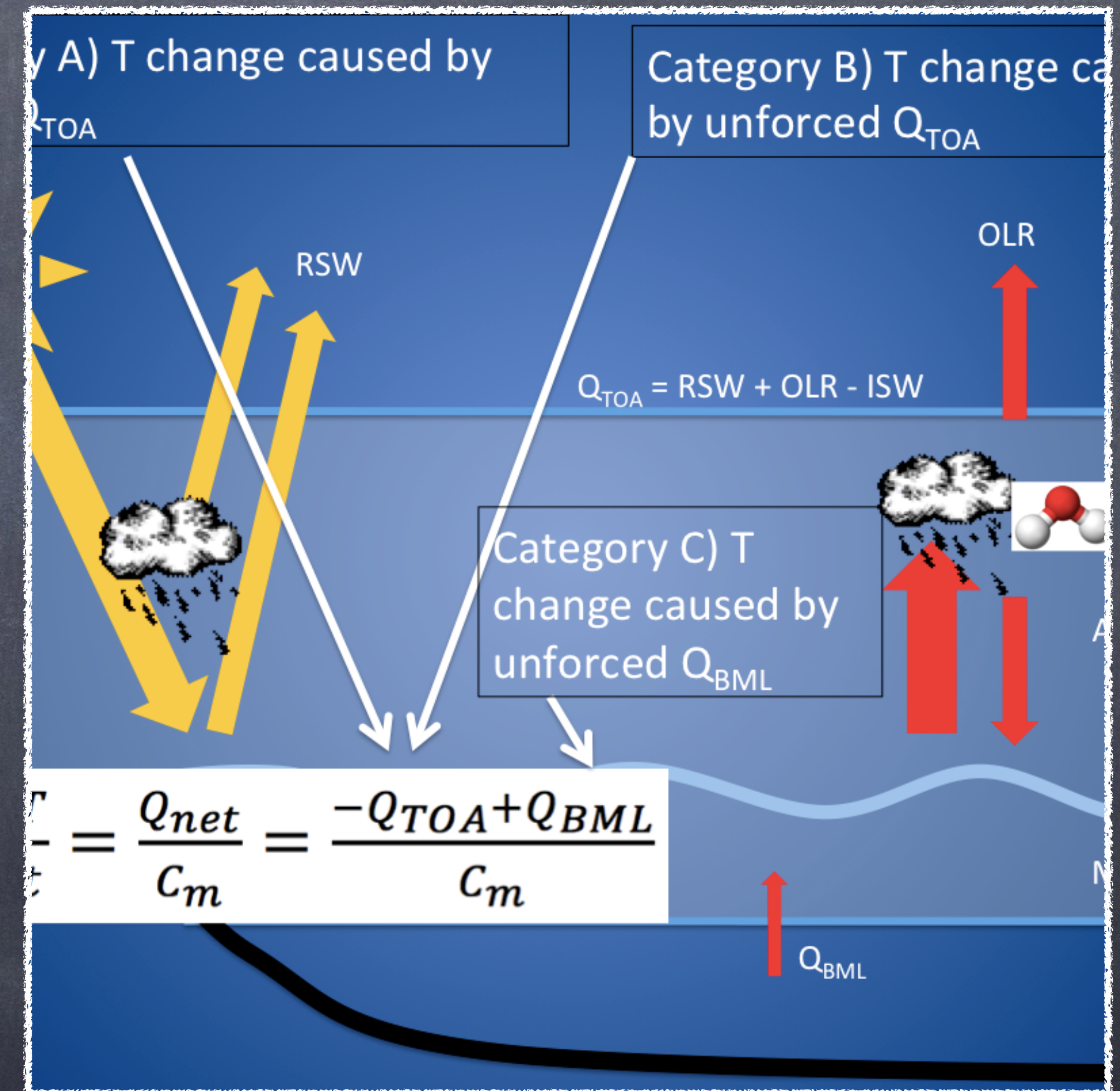
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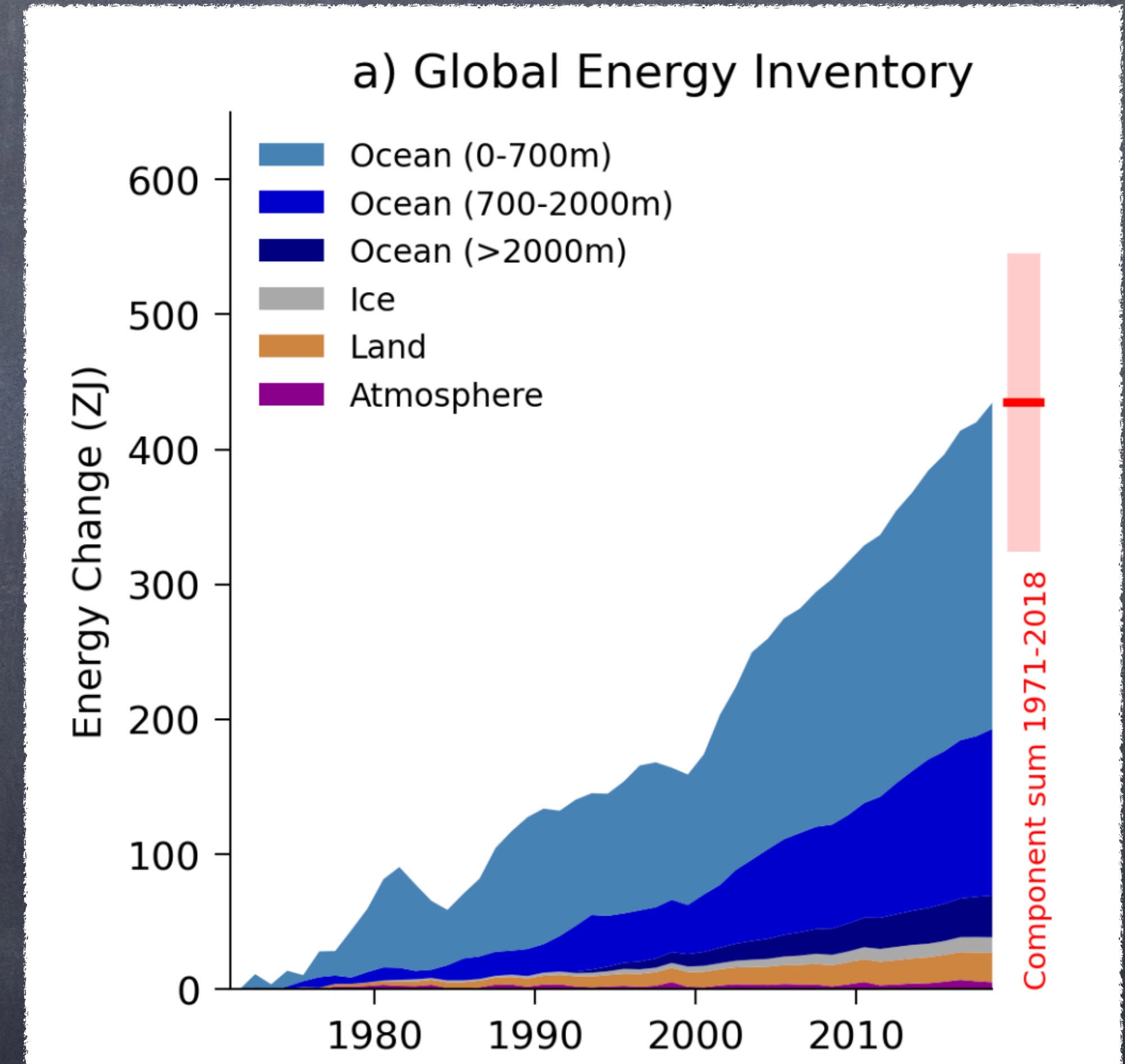
# For Today: Small-scales affect Climate

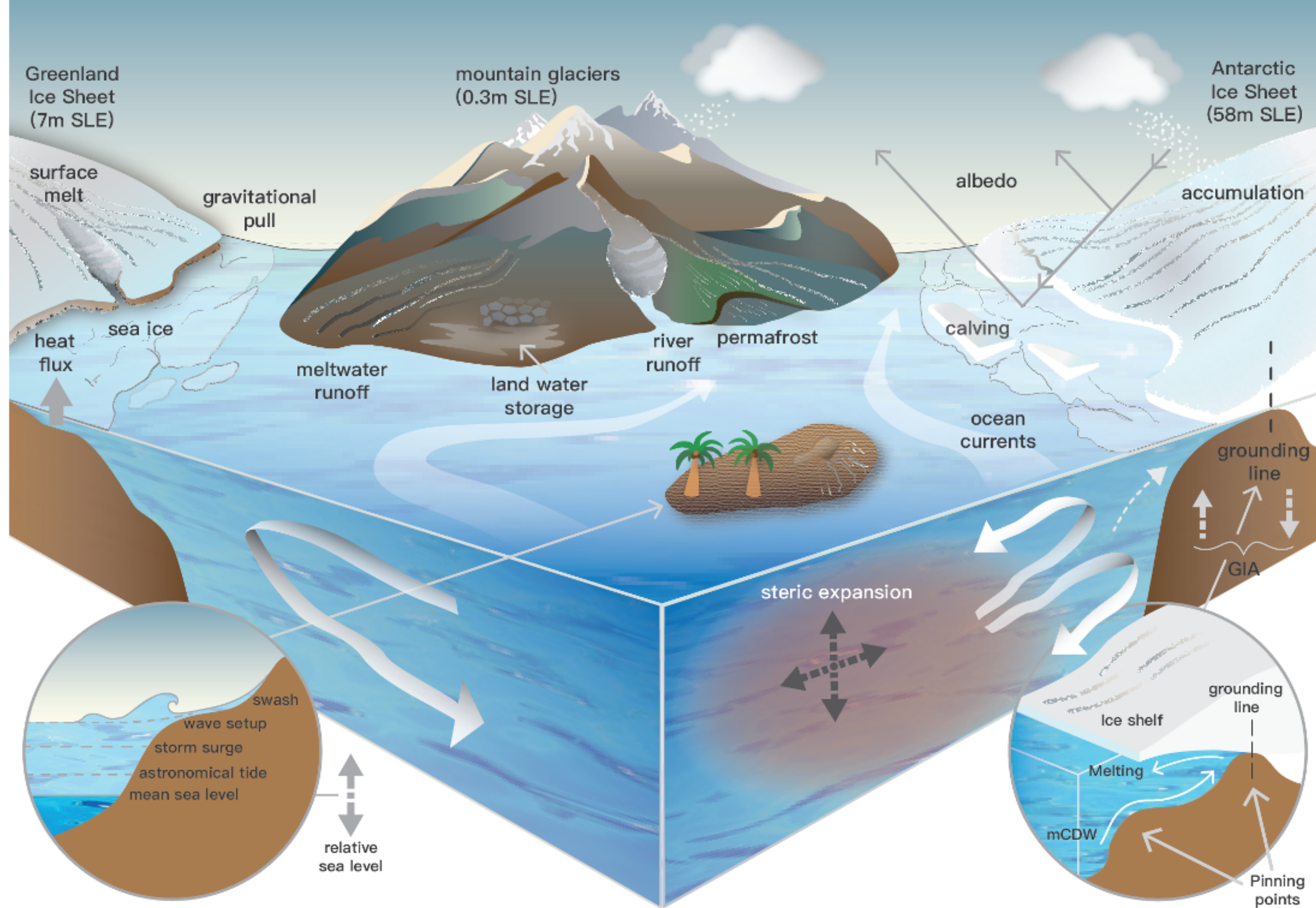
- Key Ocean Climate Questions
- Large Eddy Simulation Closures
  - Smagorinsky, Leith, QG Leith
- Effects on Global Kinetic Energy
- Submesoscale affects Mixed Layer
- Wave-Driven Turbulence affects Mixed Layer
- Regional Mixed Layer Depth affects Climate Sensitivity



# Key Ocean Climate Questions

- Sea Level Rise
- Ocean Anthropogenic Heat Uptake
  - Earth's Energy Balance
- Ocean Anthropogenic Carbon Uptake
  - Earth's Carbon Balance
  - Ocean Acidification
- Will Currents & Stratification Change?
  - Affects the above & ecosystems...





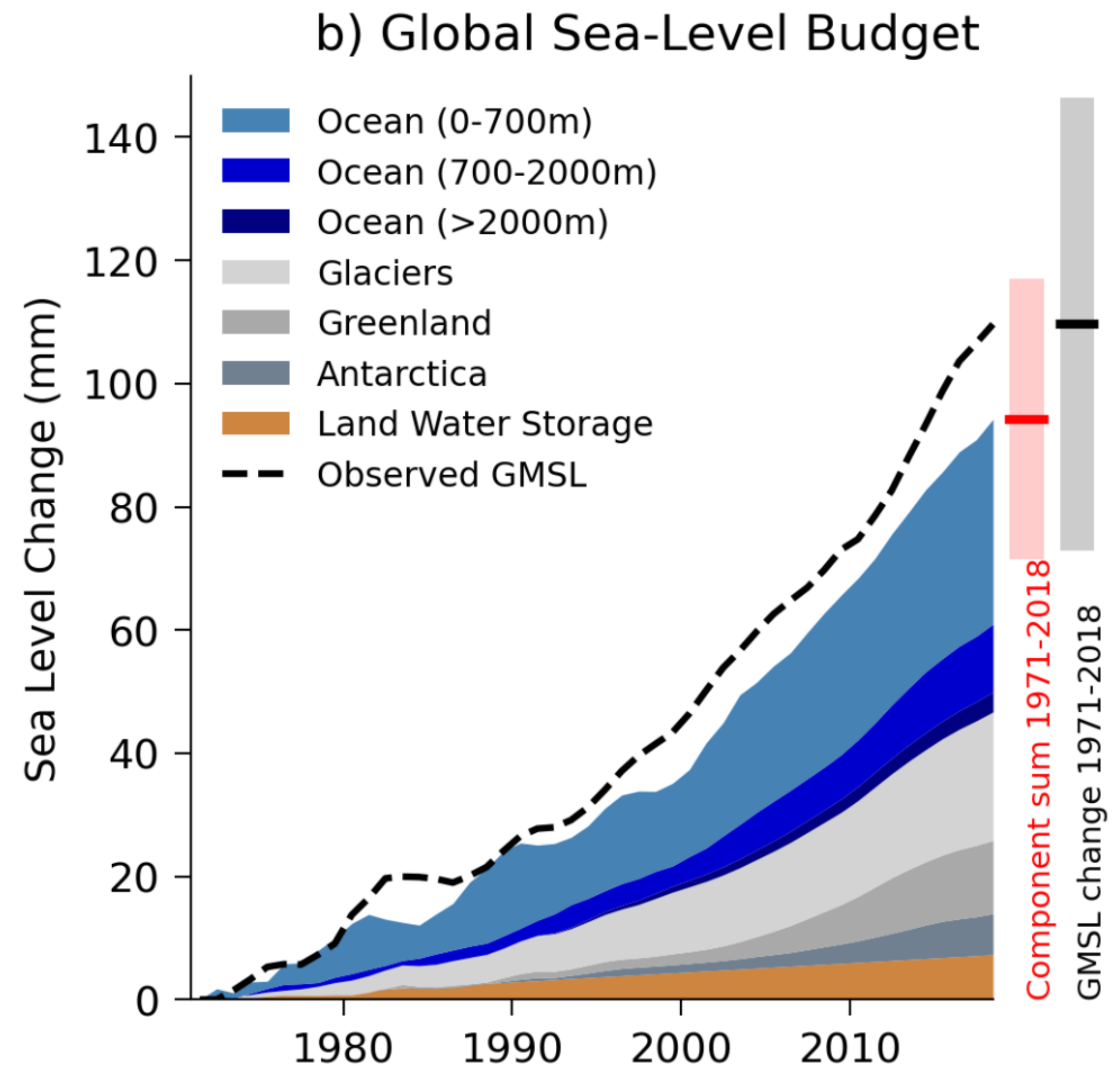
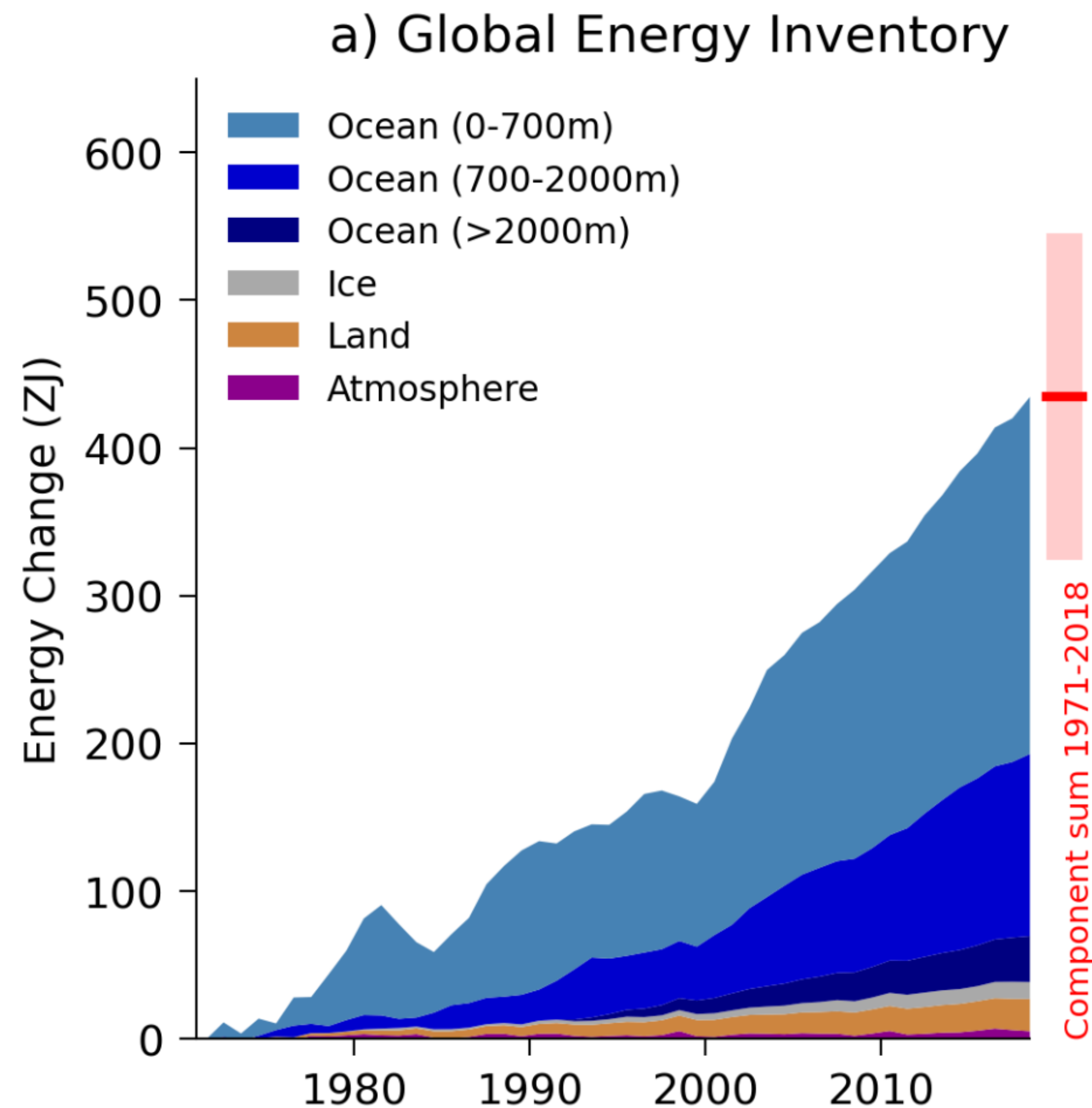
our IPCC chapter emphasizes **PROCESSES** contributing to sea level rise

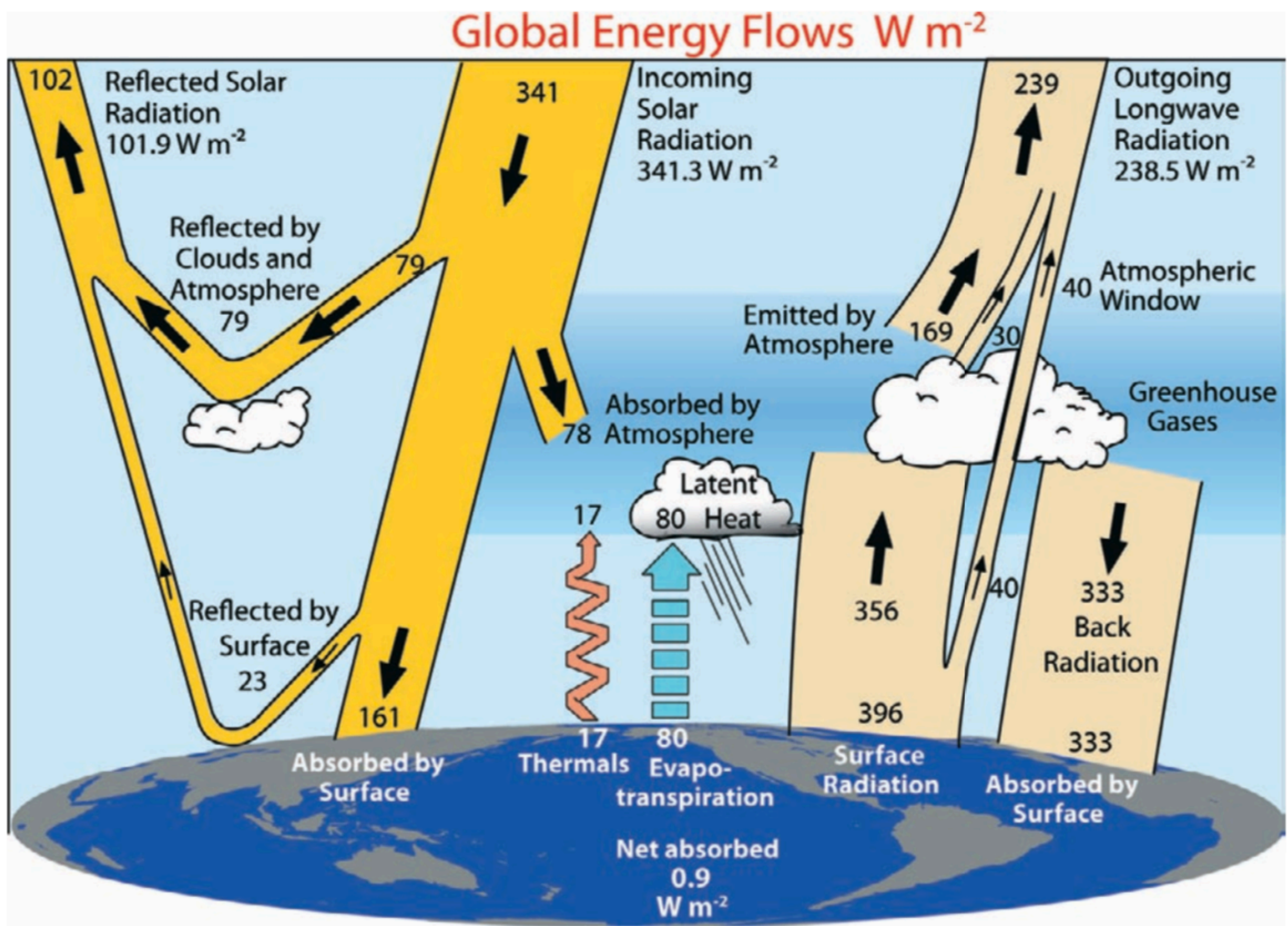
(a)

Components of ocean, cryosphere and sea level assessed in this chapter. (a) Schematic of processes (mCDW=modified Circumpolar Deep Water, GIA=Glacial Isostatic Adjustment). White arrows indicate ocean circulation. Pinning points indicate where the grounding line is most stable and ice sheet retreat will slow.



# Heating of the climate system has caused global mean sea level rise through ice loss on land and thermal expansion from ocean warming (high confidence)





Top of Atmosphere Imbalance!!

$341.3 - 101.9 - 238.5 = 0.9$

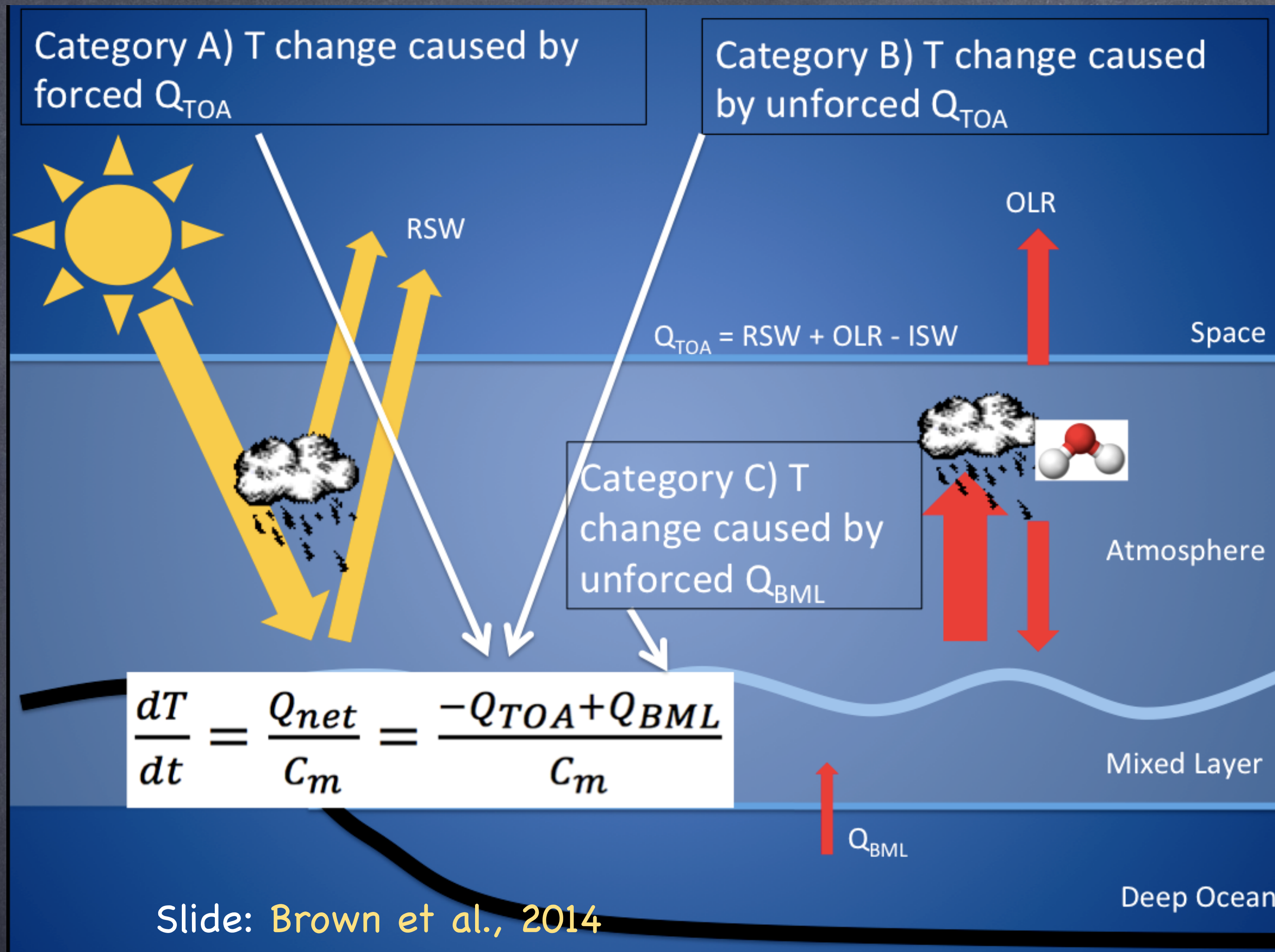
This equals net absorbed

**FIG. 1.** The global annual mean Earth's energy budget for the Mar 2000 to May 2004 period ( $W m^{-2}$ ). The broad arrows indicate the schematic flow of energy in proportion to their importance. Image: Trenberth et al. 2009

Simple: Planetary Energy Balance

$$c \frac{dT}{dt} = R_{incoming}(T) - R_{outgoing}(T)$$

# GMST: Surface Energy Budget=Ocean Heat Content Budget



Top of Atmosphere Imbalance!!

$$341.3 - 101.9 - 238.5 = 0.9$$

This equals net absorbed

Slide: Brown et al., 2014

- 3.4m of ocean has heat capacity of whole atmosphere
- Ocean Mixed Layer is about 100m deep.



Evaluating the Performance  
of Past Climate Model  
Projections

Z. Hausfather, H.F. Drake, T.  
Abbott, G.A. Schmidt, 2019

<https://doi.org/10.1029/2019GL085378>

Video: ZEKE HAUSFATHER, Carbon Brief, 2017

[www.carbonbrief.org/analysis-how-well-have-climate-models-projected-global-warming](http://www.carbonbrief.org/analysis-how-well-have-climate-models-projected-global-warming)

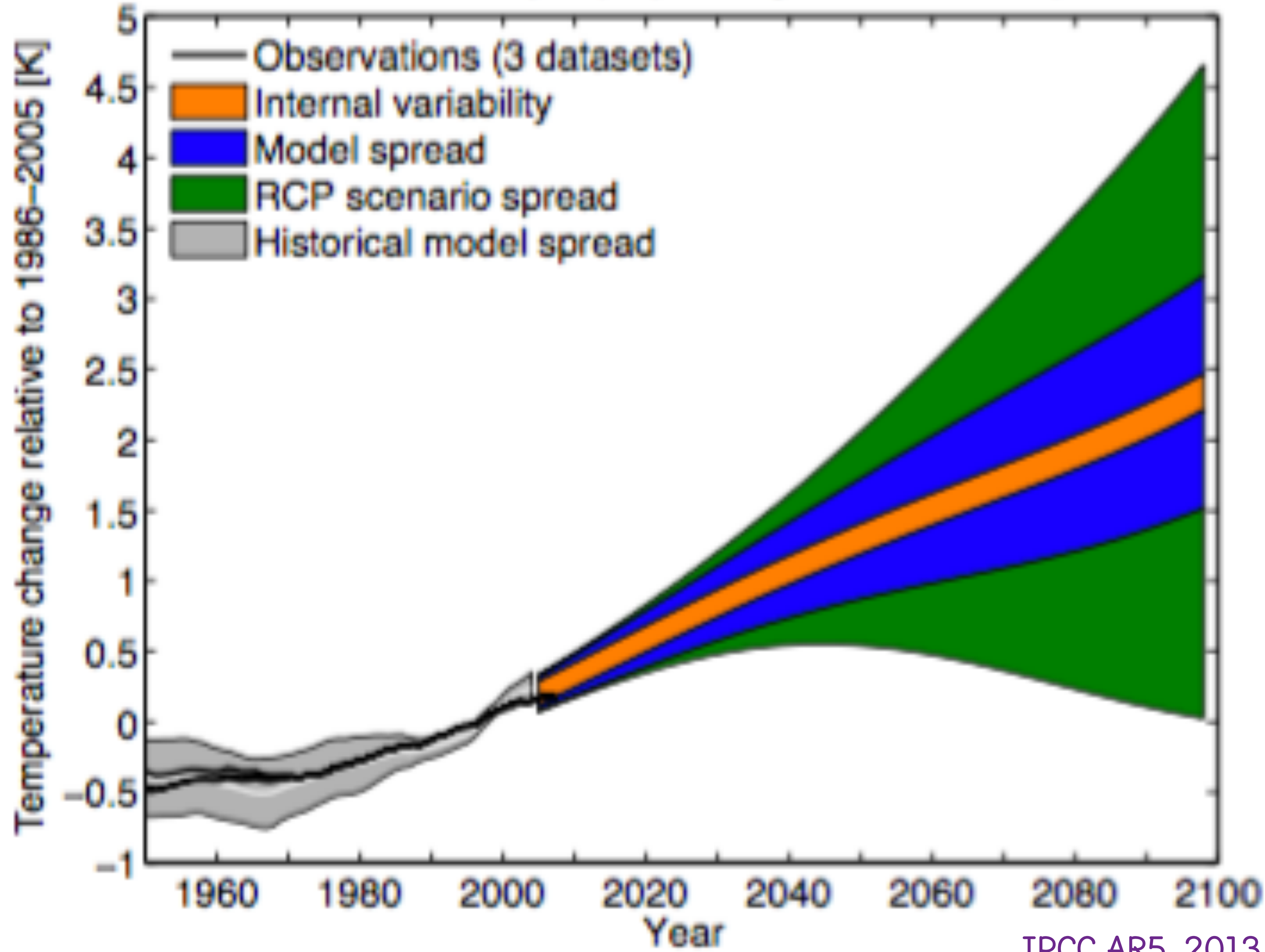


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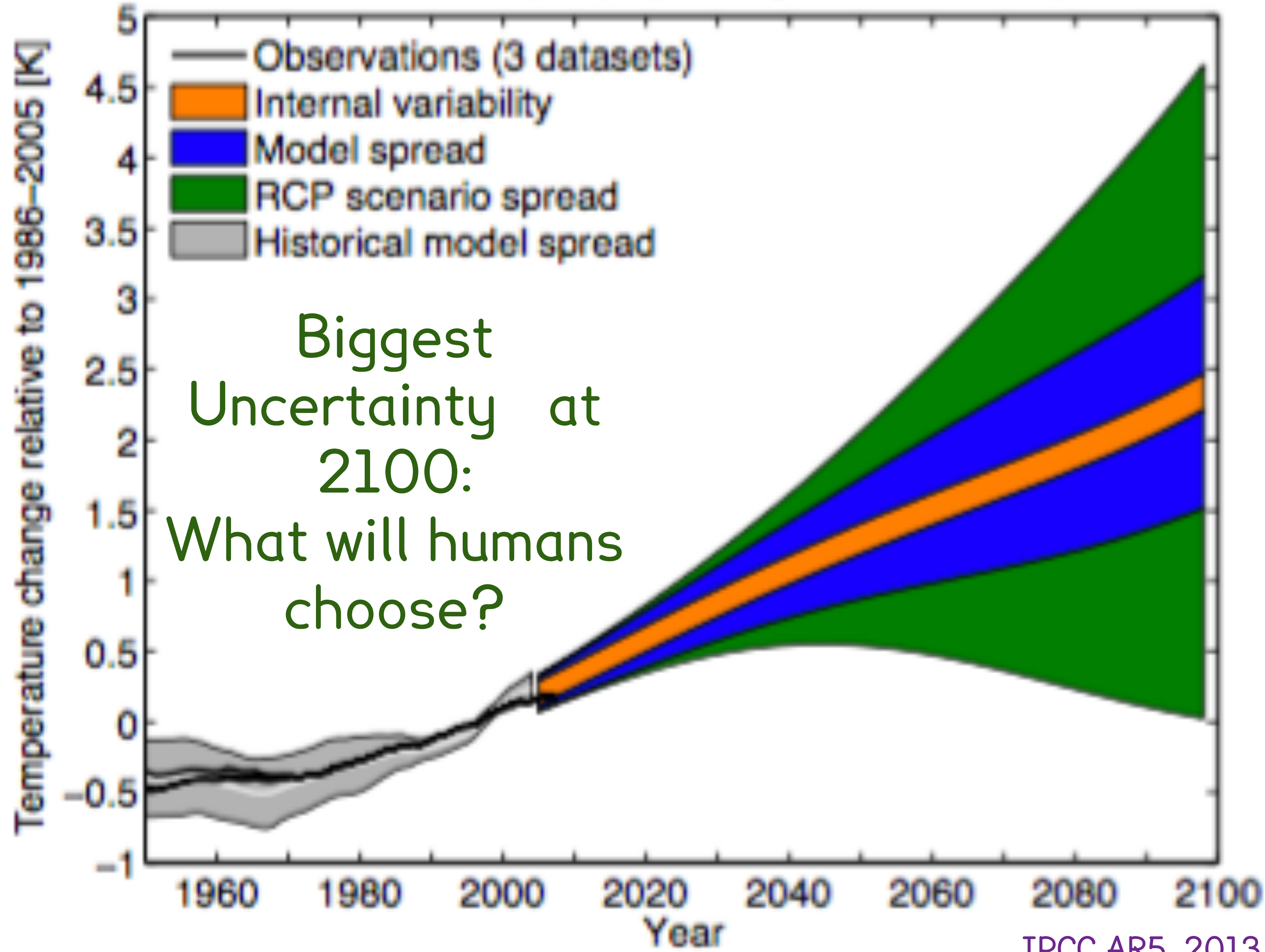
Sources of uncertainty in projected global mean temperature



IPCC AR5, 2013

Indeed, Hausfather et al. 2019 show that the early climate projections often went wrong because they assumed the wrong human emissions profile, not the wrong climate response.

Sources of uncertainty in projected global mean temperature



Biggest  
Uncertainty at  
2100:  
What will humans  
choose?

Indeed, Hausfather et al. 2019 show that the early climate projections often went wrong because they assumed the wrong human emissions profile, not the wrong climate response.

# Handling the partially-resolved mesoscale...

- In effect, the mesoscale in (highest-resolution) climate models is a Large Eddy Simulation for these modes.
- What kind of LES closure will work?



# The Character of the Mesoscale

100  
km

(Capet et al., 2008)

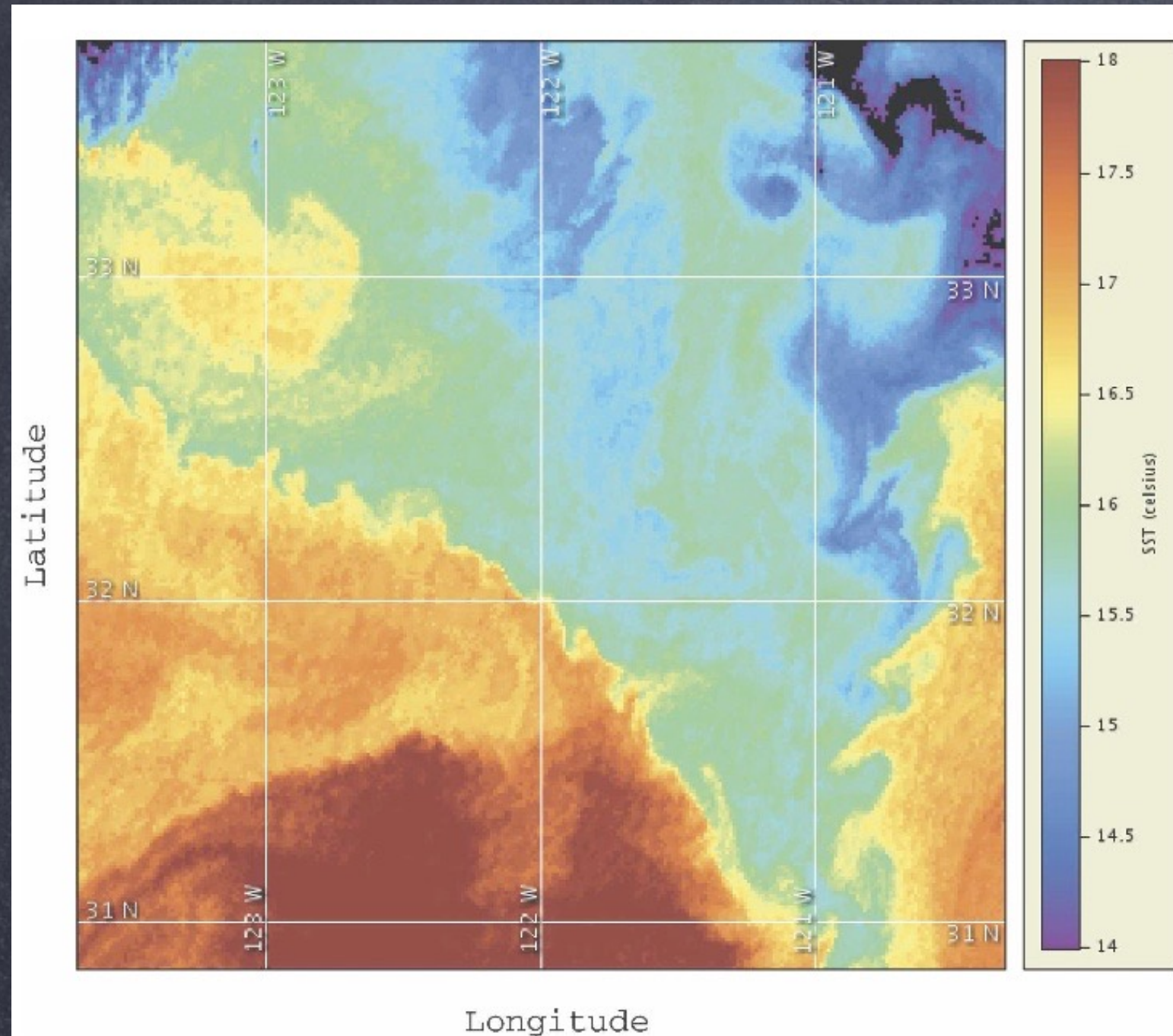
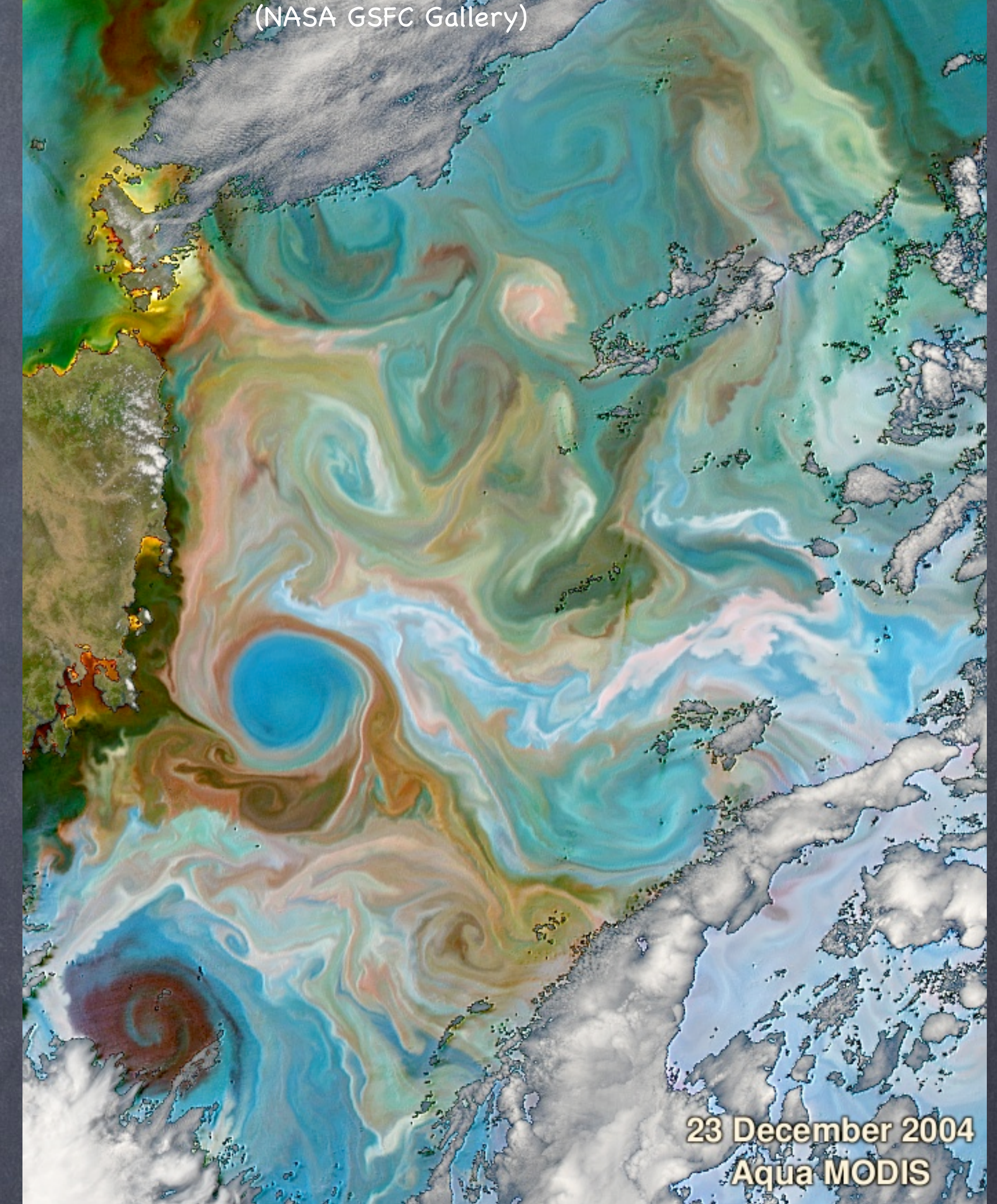


FIG. 16. Sea surface temperature measured at 1832 UTC 3 Jun 2006 off Point Conception in the California Current from CoastWatch (<http://coastwatch.pfeg.noaa.gov>). The fronts between recently upwelled water (i.e., 15°–16°C) and offshore water ( $\geq 17^\circ\text{C}$ ) show submesoscale instabilities with wavelengths around 30 km (right front) or 15 km (left front). Images for 1 day earlier and 4 days later show persistence of the instability events.

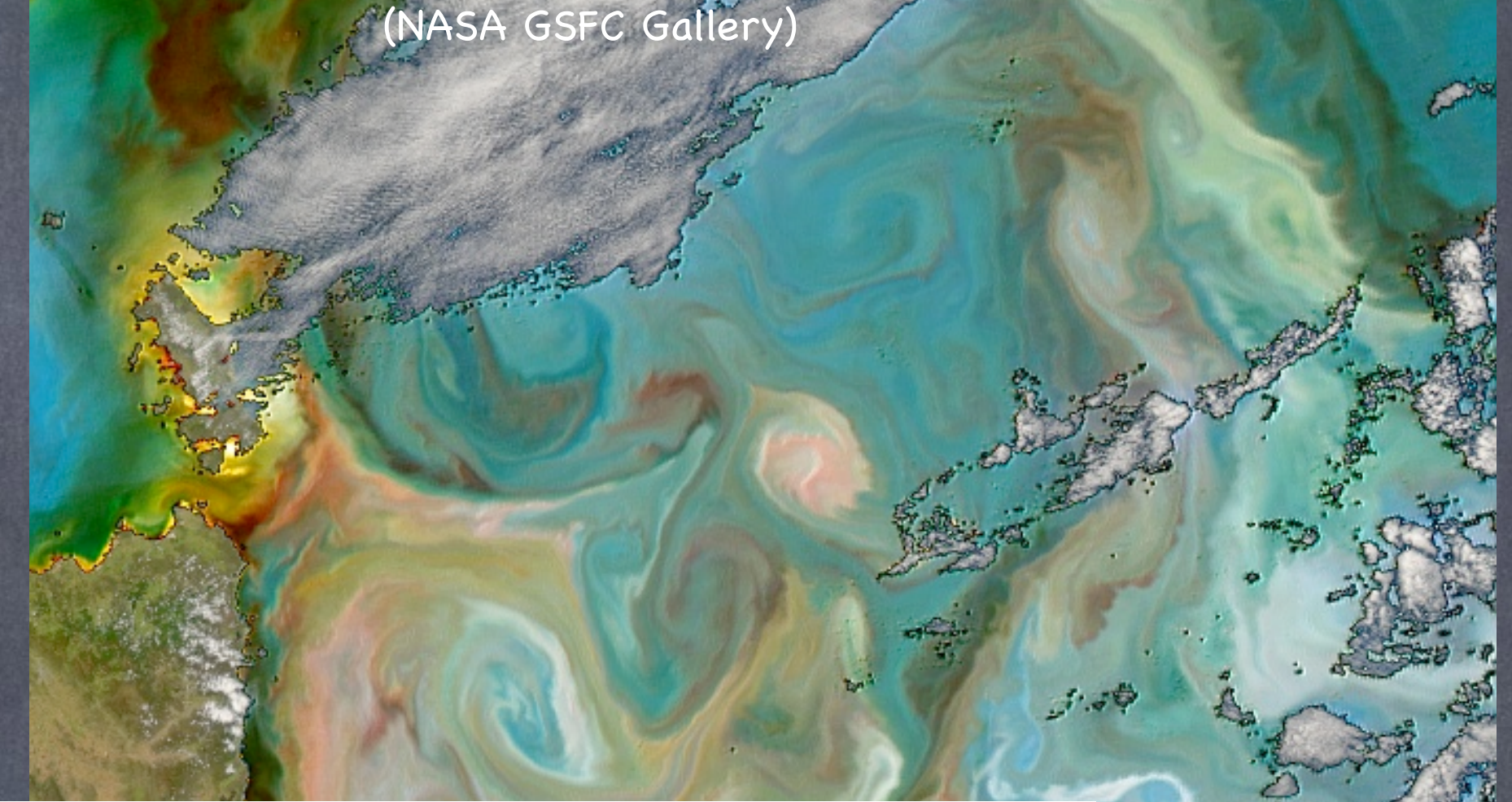
- Boundary Currents
- Eddies
- $Ro = O(0.1)$
- $Ri = O(1000)$
- Full Depth
- Eddies strain to produce Fronts
- 100km, months



Eddy processes mainly **baroclinic & barotropic instability**. Parameterizations of baroclinic instability (GM, Visbeck...).

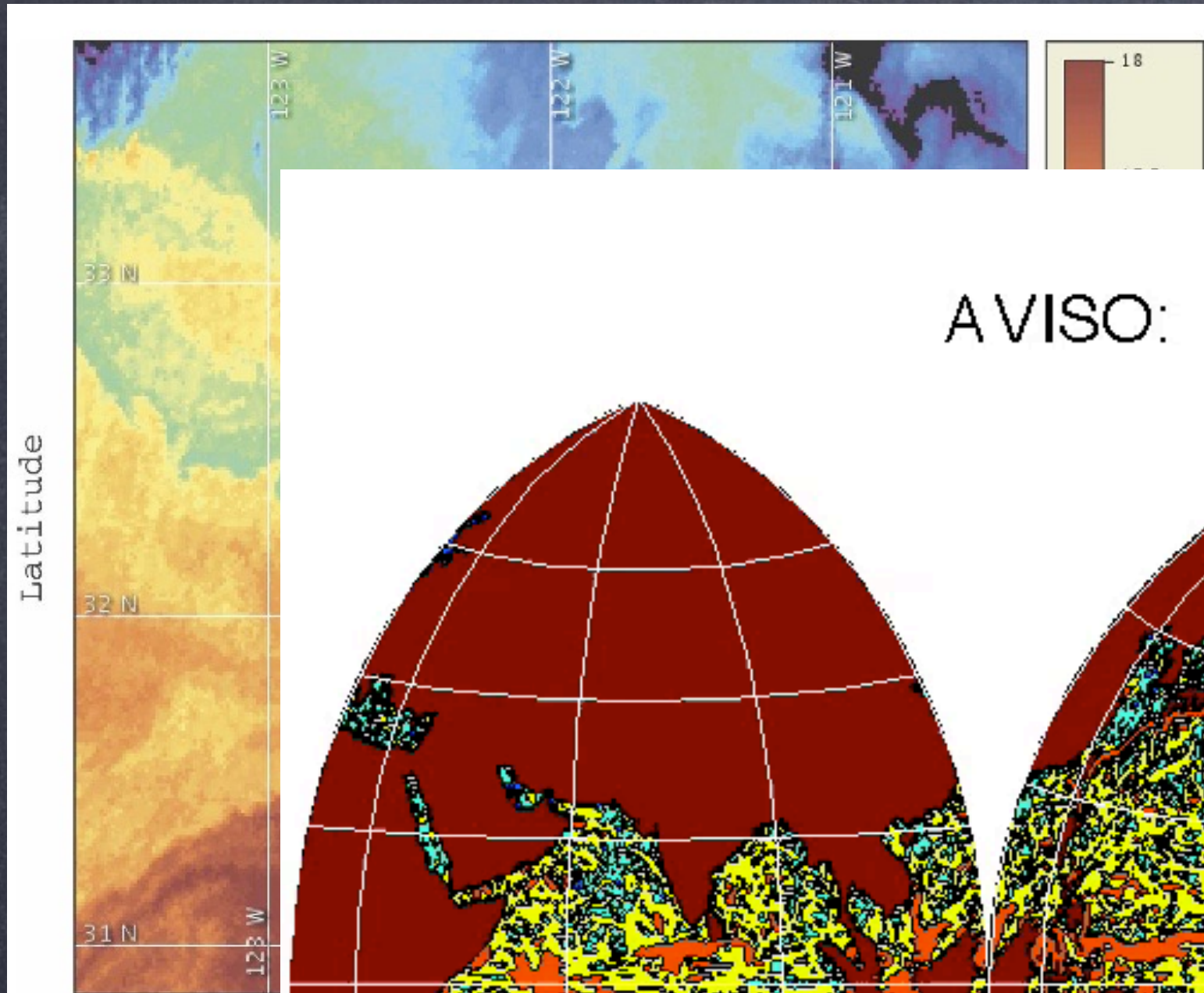
# The Character of the Mesoscale

100 km



(Capet et al., 2008)

Boundary Currents



AVISO:  $\log_{10}(0.5(u^2+v^2))$  on 19940101

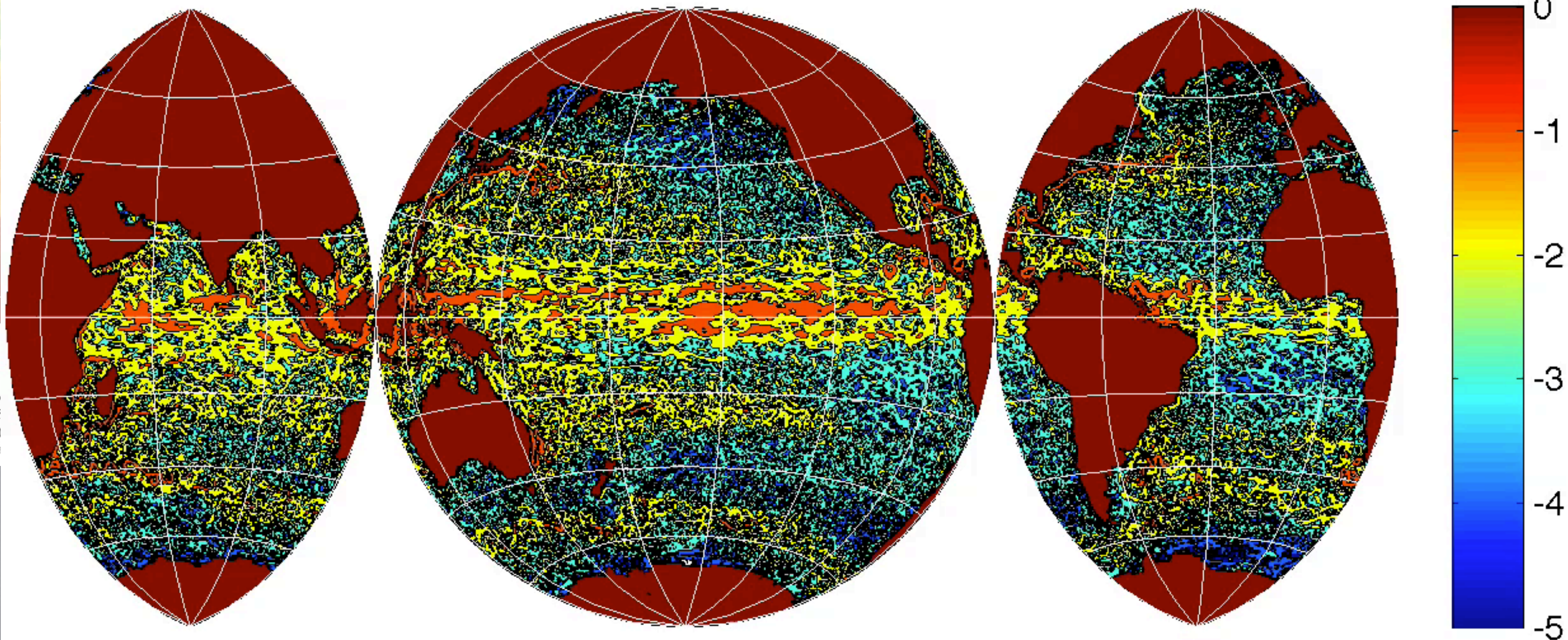


FIG. 16. Sea surface California Current from upwelled water (i.e., 15 lengths around 30 km (persistence of the insta

December 2004  
Aqua MODIS

Th

(Co

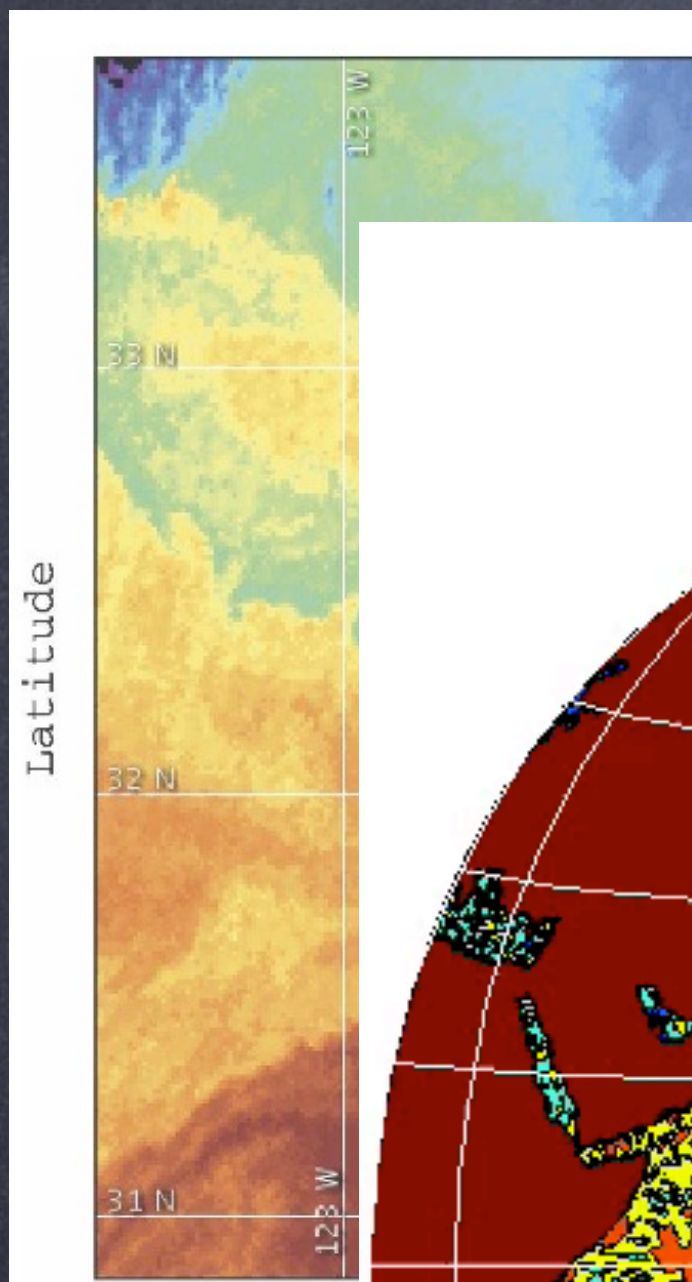
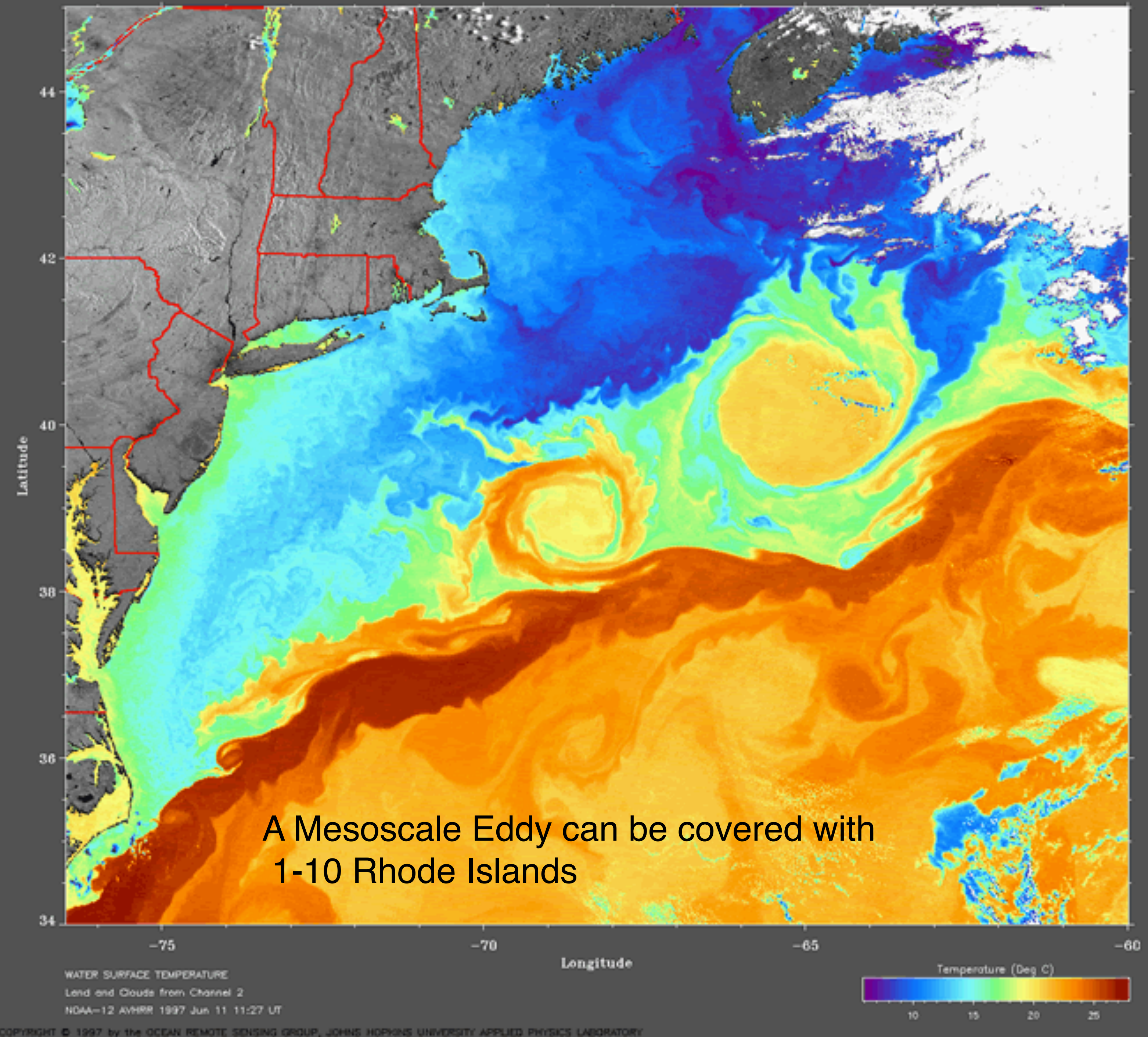
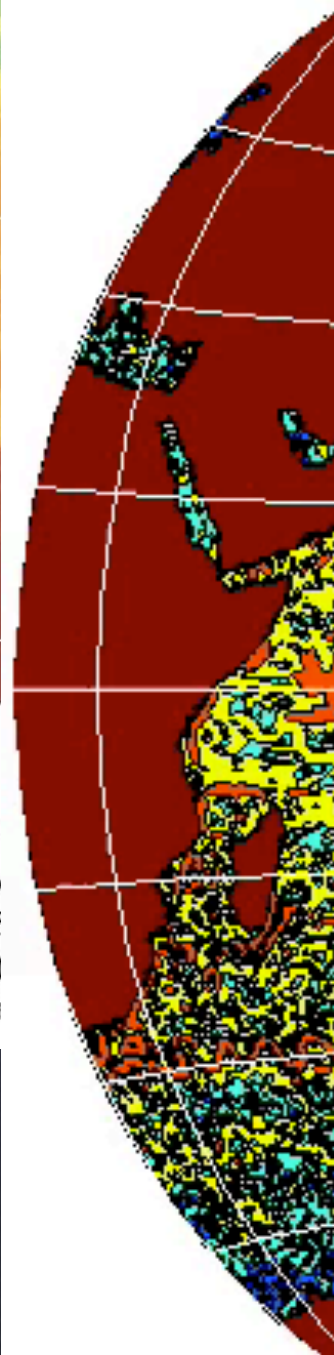
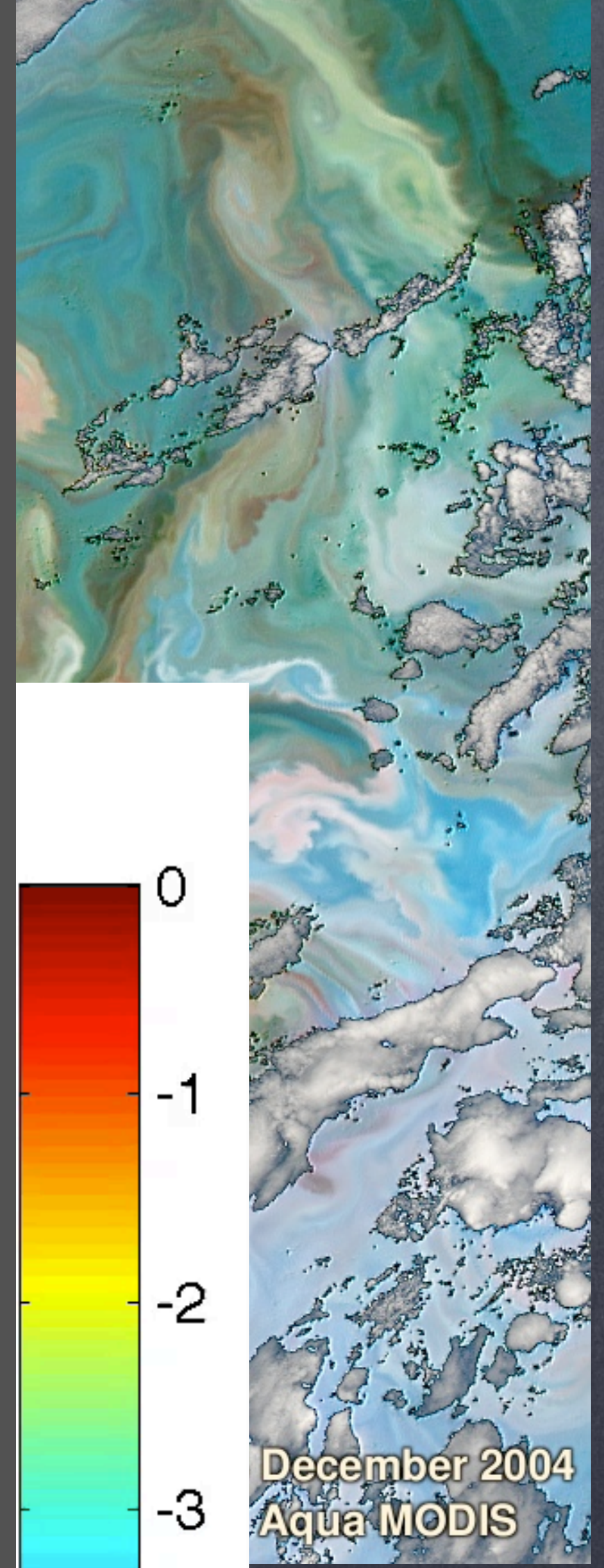


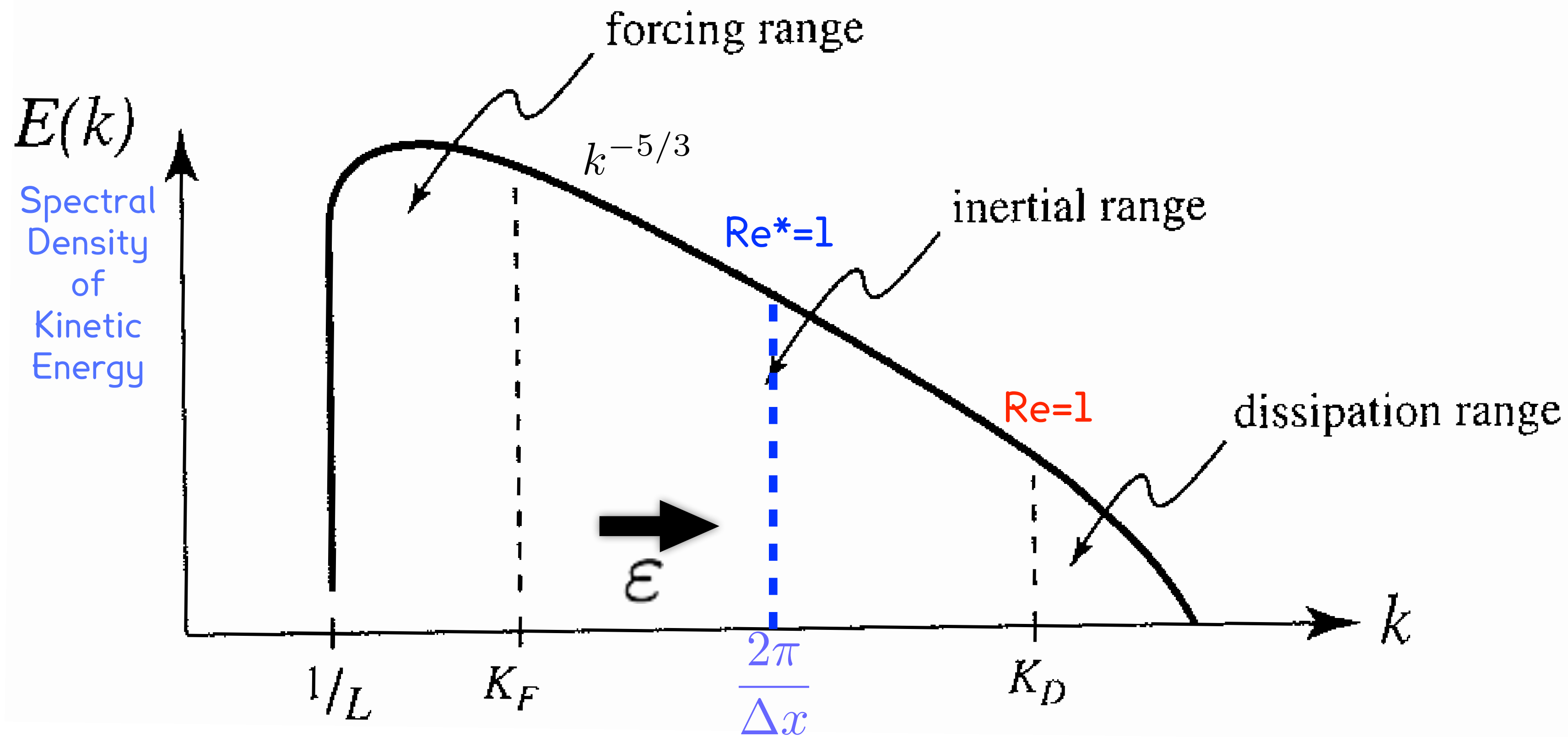
FIG. 16. Sea surface California Current from upwelled water (i.e., 15 lengths around 30 km (persistence of the insta



A Mesoscale Eddy can be covered with 1-10 Rhode Islands



# 3D Turbulence Cascade



Suitable  
For  
Nonhydrostatic  
Boussinesq;  
Wave-averaged

1963: Smagorinsky Scale & Flow Aware Viscosity Scaling,  
So the Energy Cascade is Preserved,  
but order-1 gridscale Reynolds #:  $Re^* = UL/\nu_*$

$$\mathbf{v}_{*h} = \left( \frac{\Upsilon_h \Delta x}{\pi} \right)^2 \sqrt{\left( \frac{\partial u_*}{\partial x} - \frac{\partial v_*}{\partial y} \right)^2 + \left( \frac{\partial u_*}{\partial y} + \frac{\partial v_*}{\partial x} \right)^2}$$

# Smagorinsky Viscosity (Cited by 17445) in 2 min

$$\frac{d \langle E_* \rangle}{dt} = -F_E|_0^{k_*} - \int_0^{k_*} \nu k^2 E(k) dk + \int_0^{k_*} S_E(k, t) dk.$$

Make  $k_*$   
Kolmogorov Scale

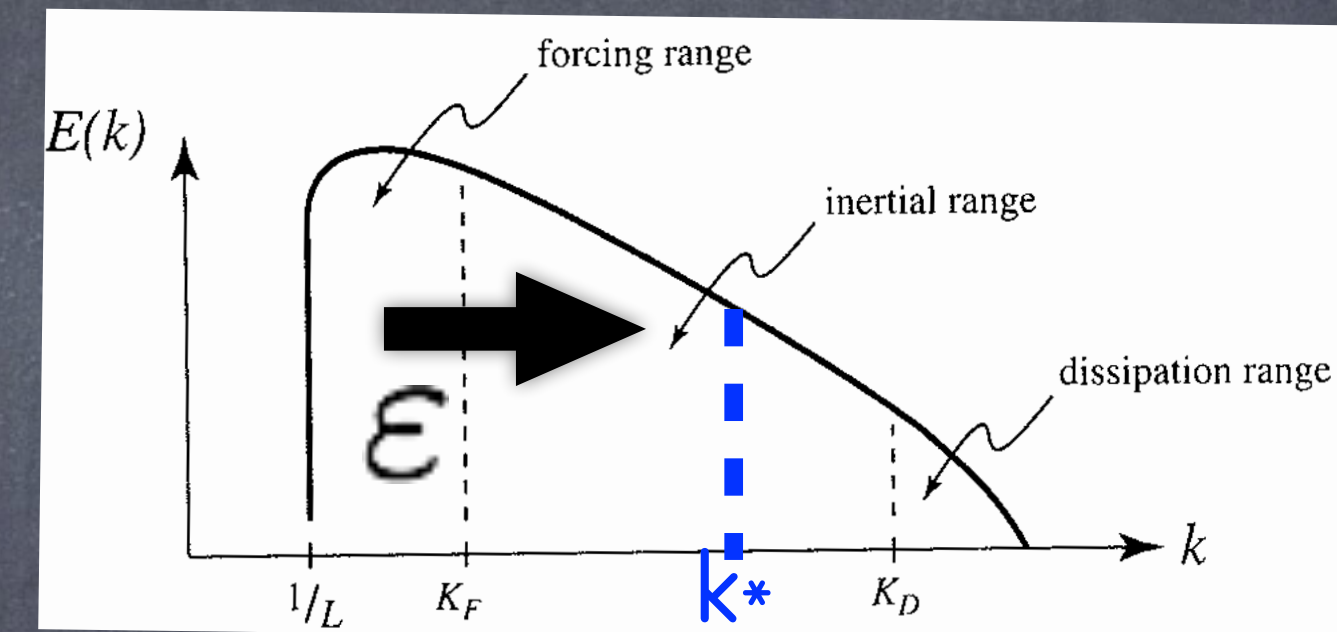
$$k_* = \Upsilon \varepsilon^{1/4} \nu_*^{-3/4},$$

$$\varepsilon = F_E(k_*) + \int_0^{k_*} \nu k^2 E(k) dk \equiv \int_0^{k_*} \nu_* k^2 E(k) dk$$

If the viscous term is evaluated in real space rather than wavenumber space, then

$$\int_0^{k_*} \nu_* k^2 E(k) dk = \langle \nu_* S_*^{ik} S_{*ik} \rangle, \quad (20)$$

Finally, because the turbulence is assumed to be homogeneous, the domain-averaged friction is replaced with a local value, and the result for  $\nu_*$  follows.



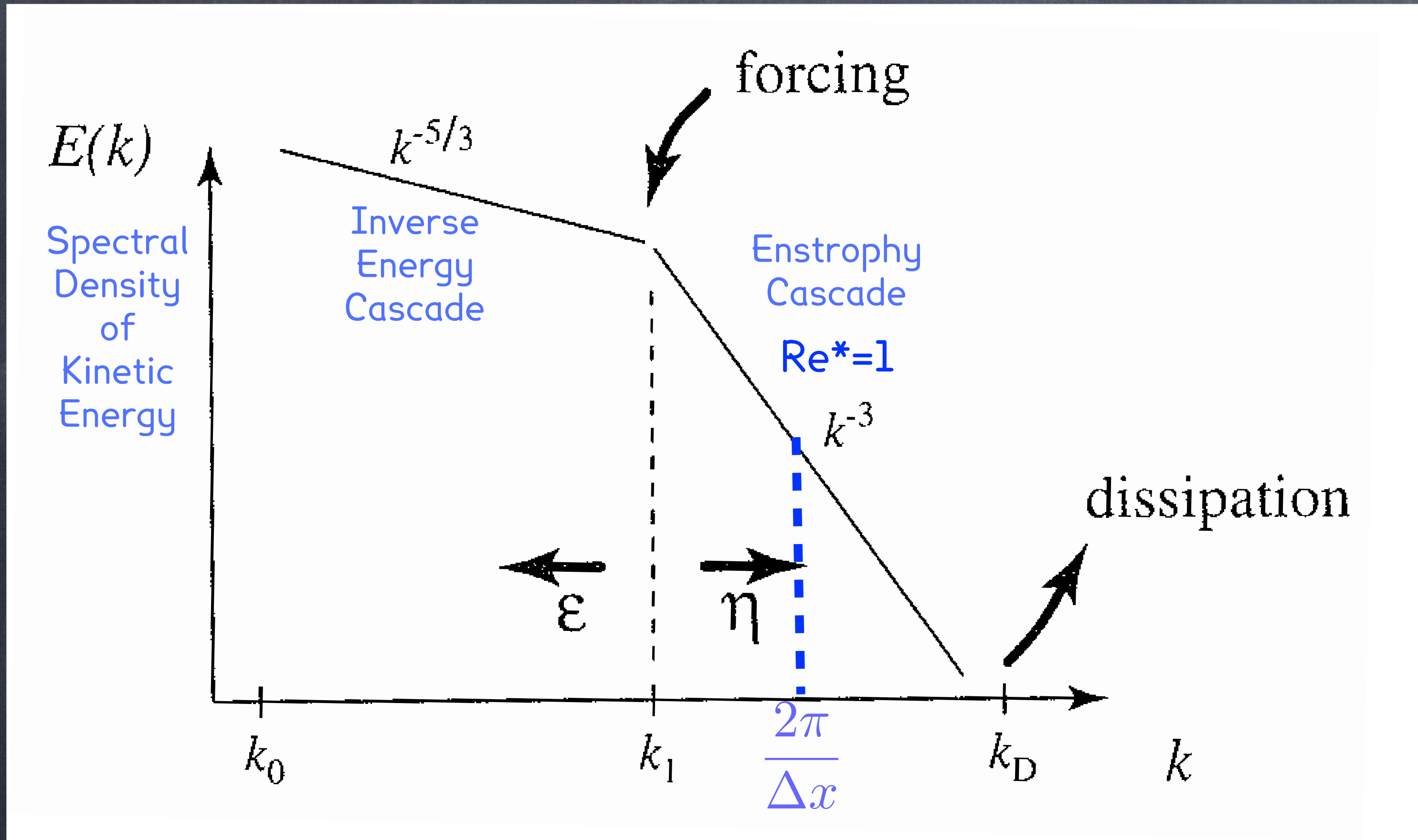
$$\nu_* = \left( \frac{\Upsilon}{k_*} \right)^2 |D_*|,$$

$$\nu_* = \left( \frac{\Upsilon \Delta x}{\pi} \right)^2 |D_*|,$$

$$|D_*| \equiv \sqrt{S_*^{ik} S_{*ik}}.$$

All flux at  $k_*$   
"viscous"

# 2D Turbulence Differs



Suitable  
For  
2D Oceans,  
E.g., Stommel  
& Munk Gyres

Best of: Graham &  
Ringler (2013)

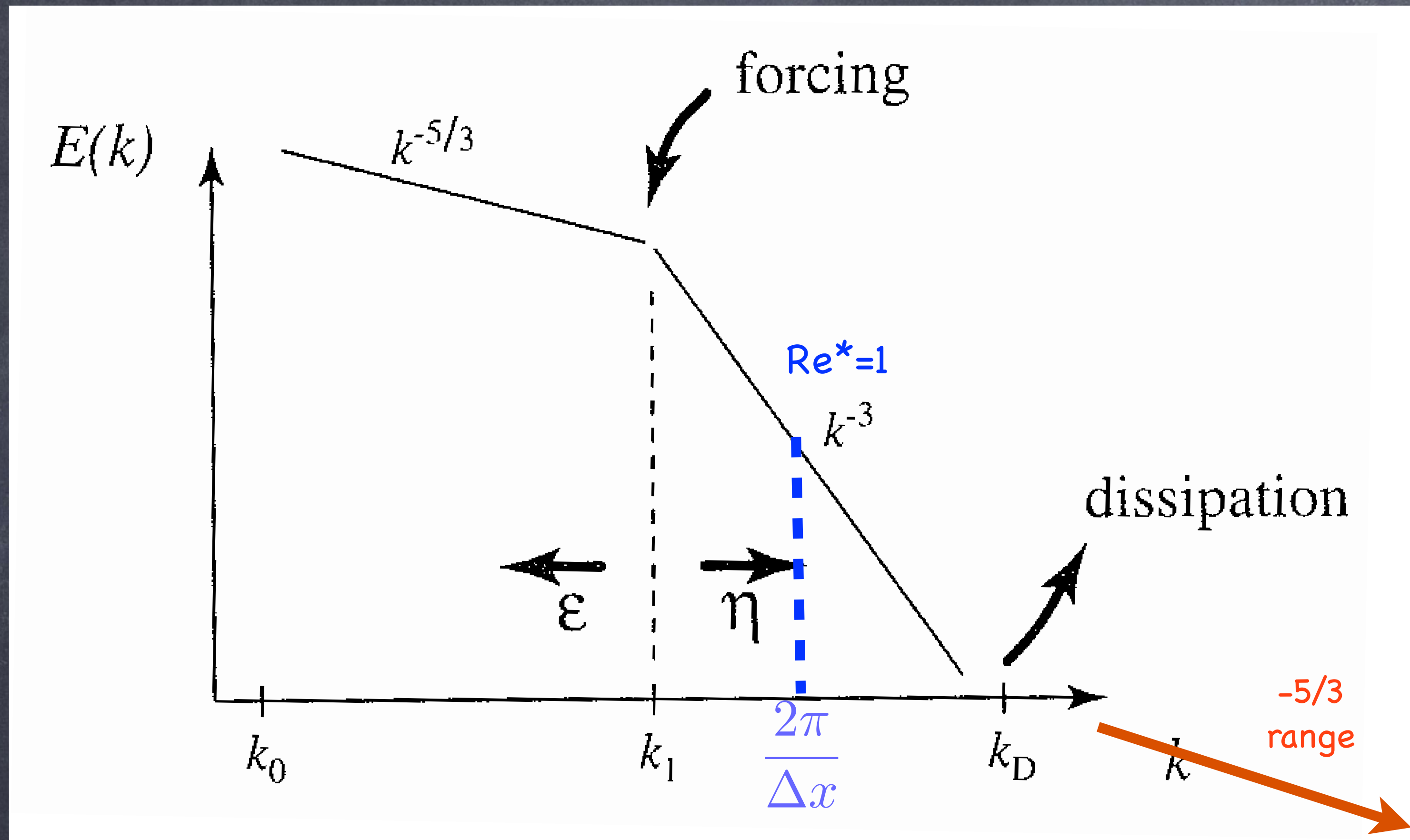
1996: Leith Devises Viscosity Scaling,  
So that the Enstrophy (vorticity<sup>2</sup>) Cascade is Preserved

R. Kraichnan, 1967 JFM  
C.E. Leith, 1996 Physica D

$$\mathbf{v}_* = \left( \frac{\Lambda \Delta x}{\pi} \right)^3 \left| \nabla_h \left( \frac{\partial u_*}{\partial y} - \frac{\partial v_*}{\partial x} \right) \right|$$

Barotropic or  
stacked layers

# Mesoscale Turbulence Like 2D cascade, but a little divergent



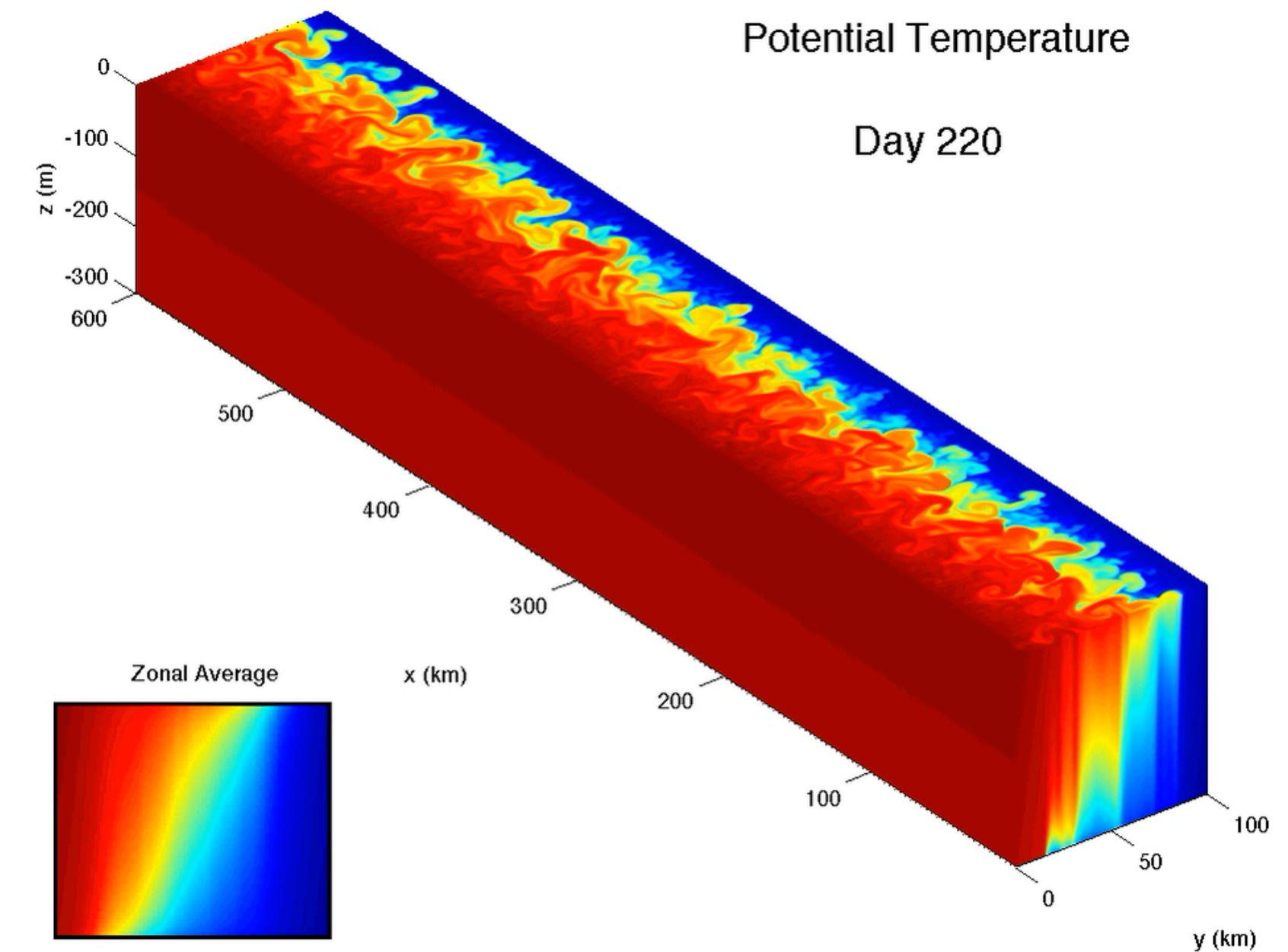
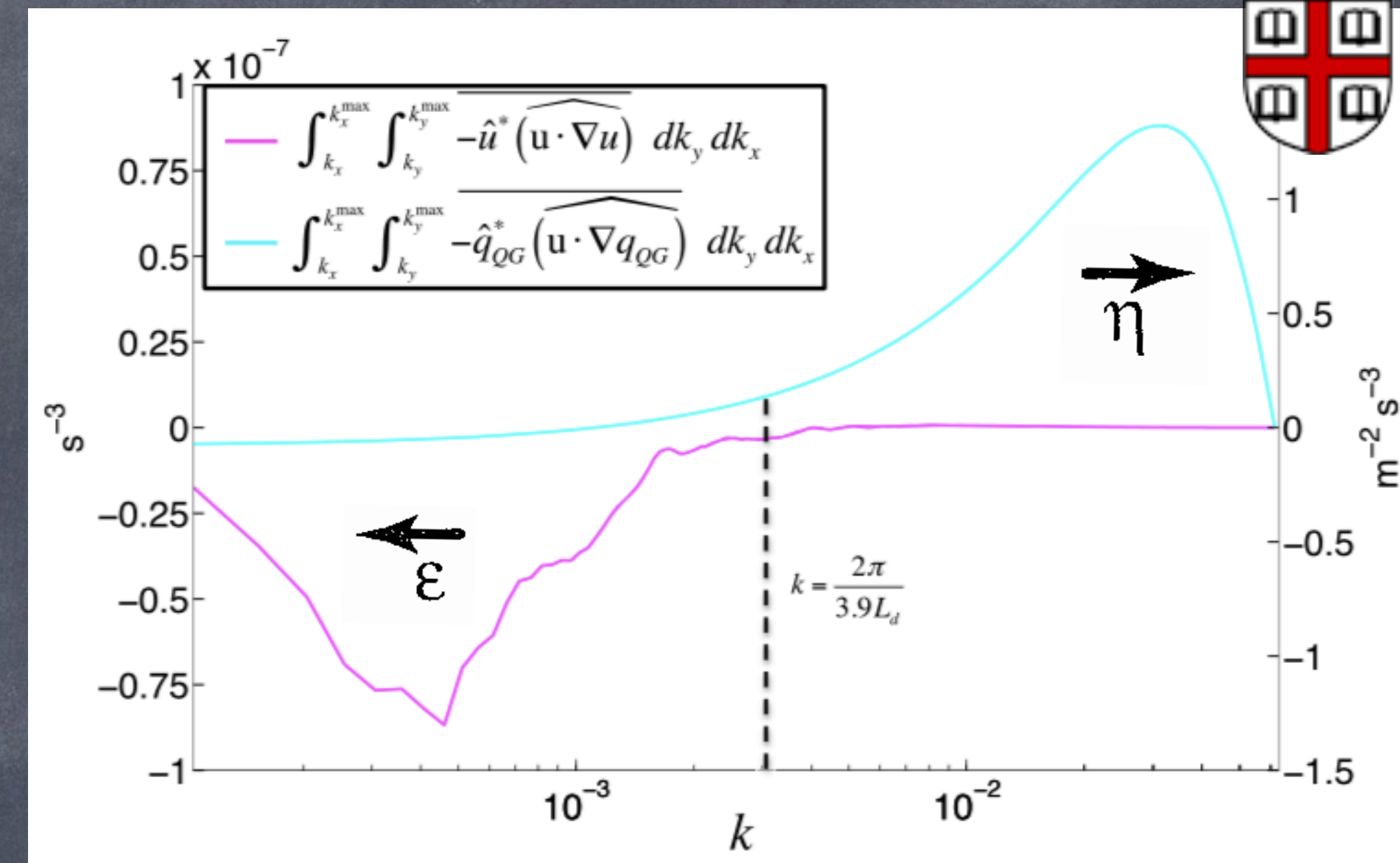
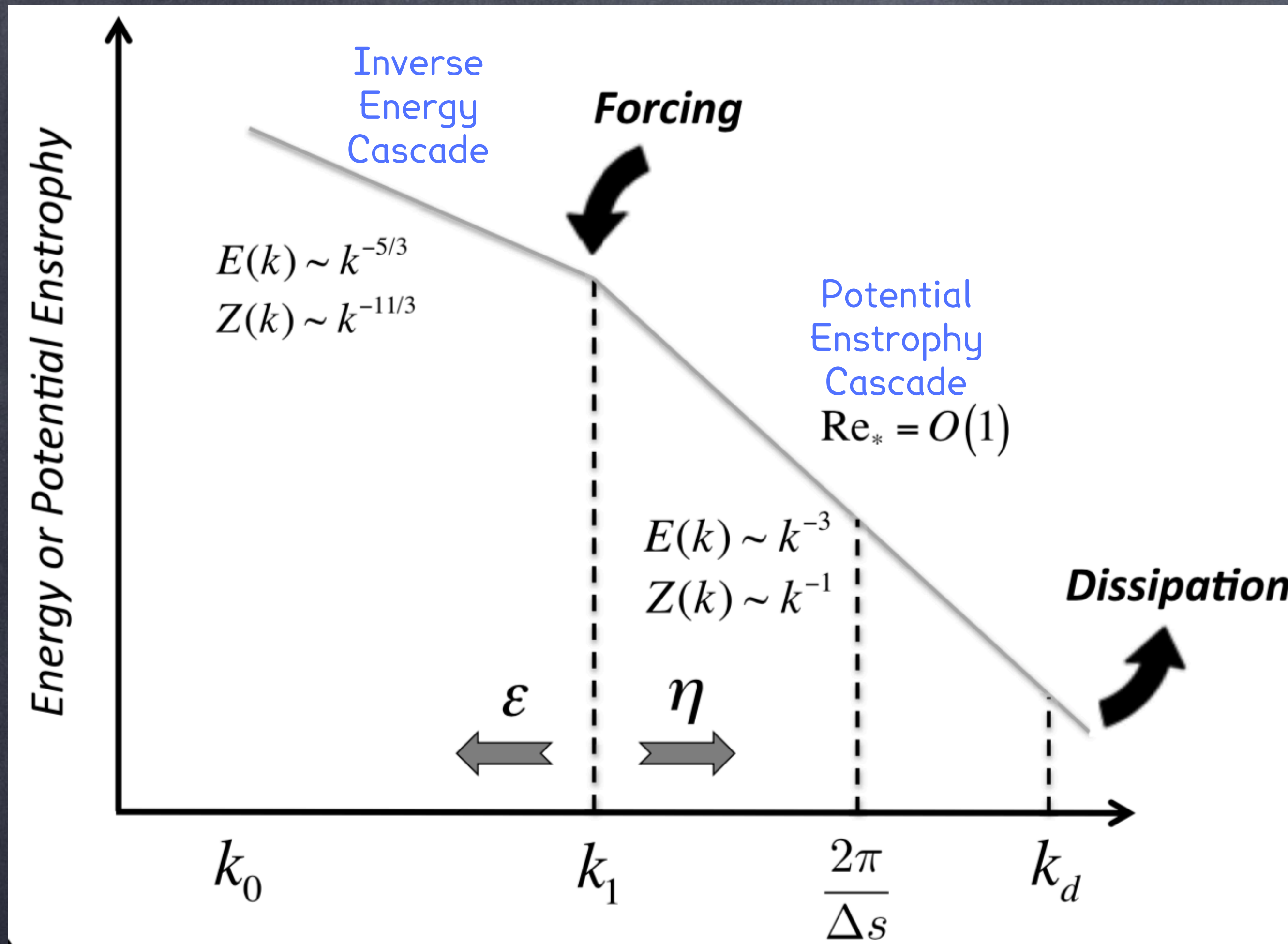
Suitable  
For  
2D Oceans,  
Or  
Hydrostatic  
Boussinesq!

F-K & Menemenlis Revise Leith Viscosity Scaling,  
So that diverging, vorticity-free, modes are also damped

$$\nu_* = \left(\frac{\Delta x}{\pi}\right)^3 \sqrt{\Lambda^6 |\nabla_h q_{2d}|^2 + \Lambda_d^6 |\nabla_h (\nabla_h \cdot \mathbf{u}_*)|^2}$$

BFK and D. Menemenlis. Can large eddy simulation techniques improve mesoscale-rich ocean models? In M. Hecht and H. Hasumi, editors, Ocean Modeling in an Eddying Regime, volume 177, pages 319-338. AGU Geophysical Monograph Series, 2008.

# QGLeith: Pot'l Enstrophy cascade (efficient versions ready for Hydrostatic Boussinesq at 10km and fine)



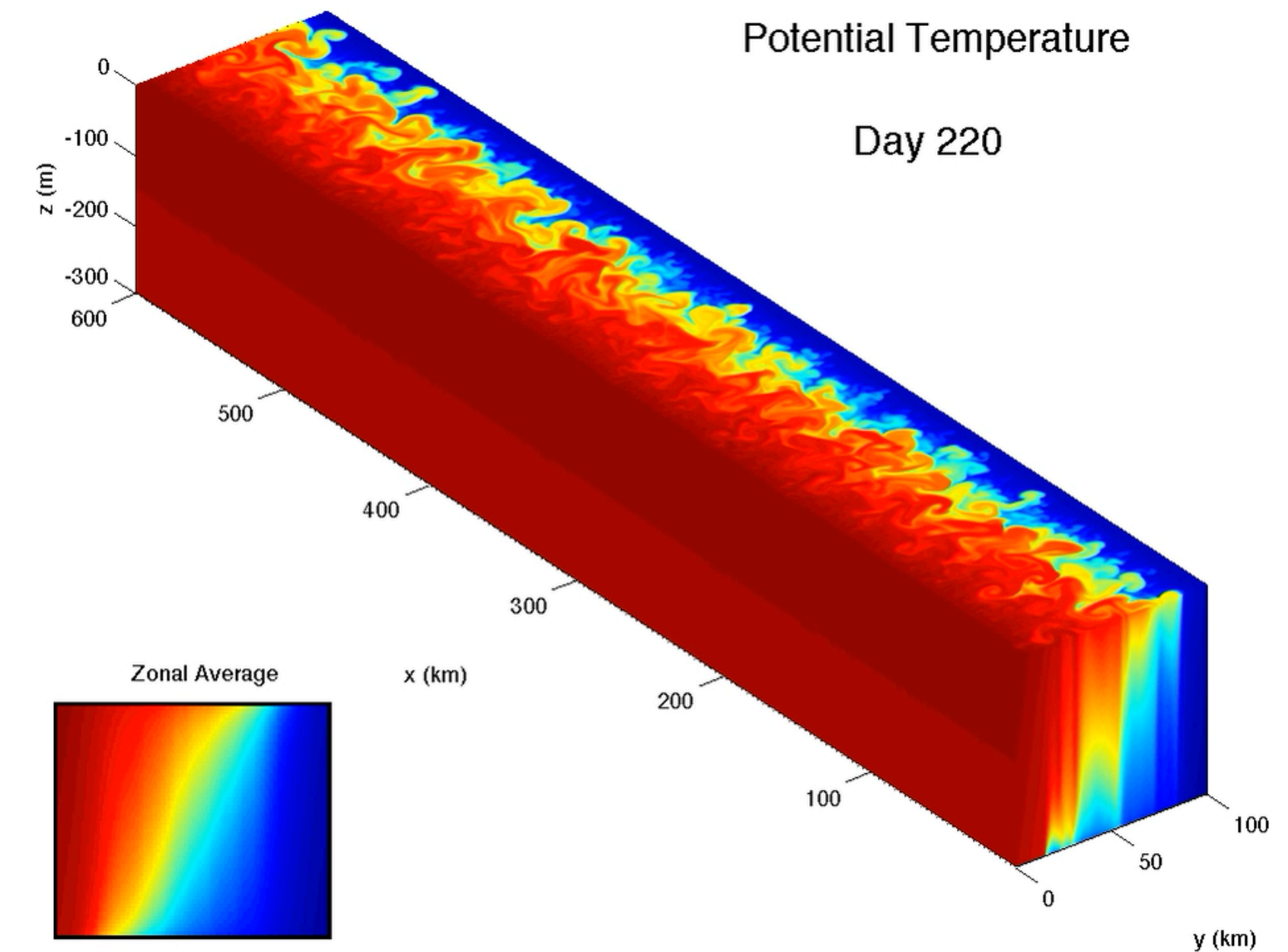
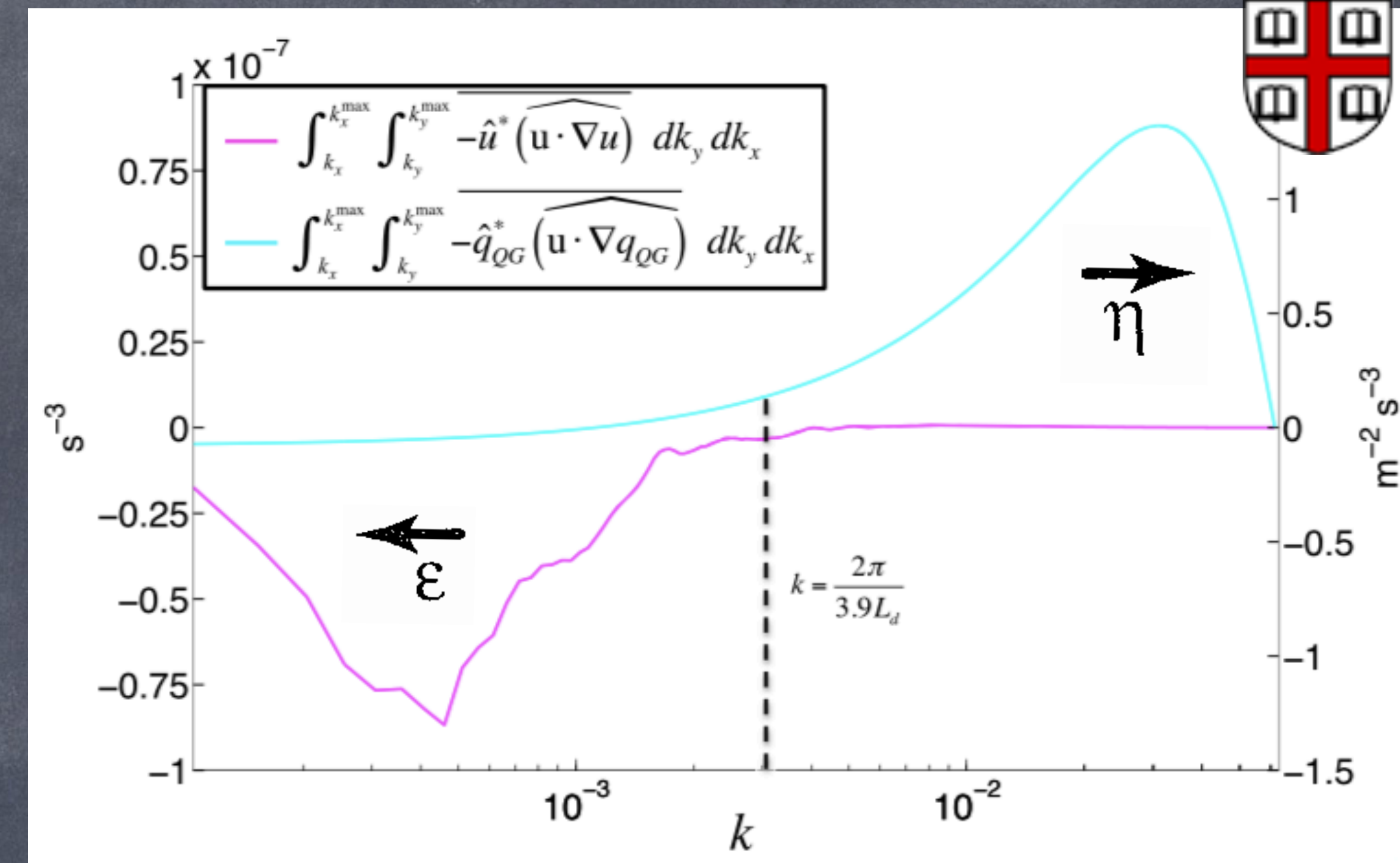
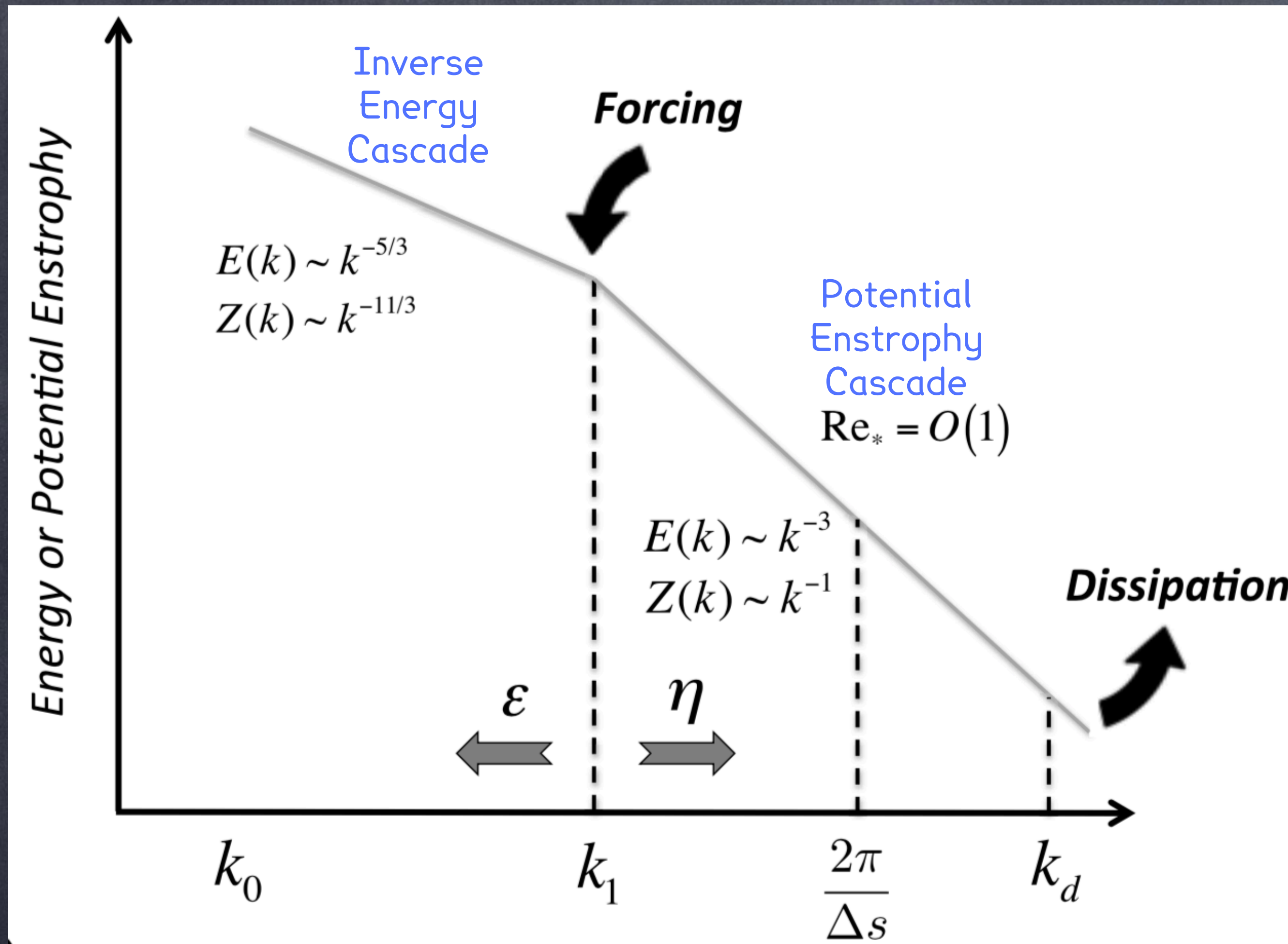
Provides lateral viscosity, diffusivity, and GM coefficient without overdamping resolved eddies

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# QGLeith: Pot'l Enstrophy cascade (efficient versions ready for Hydrostatic Boussinesq at 10km and fine)



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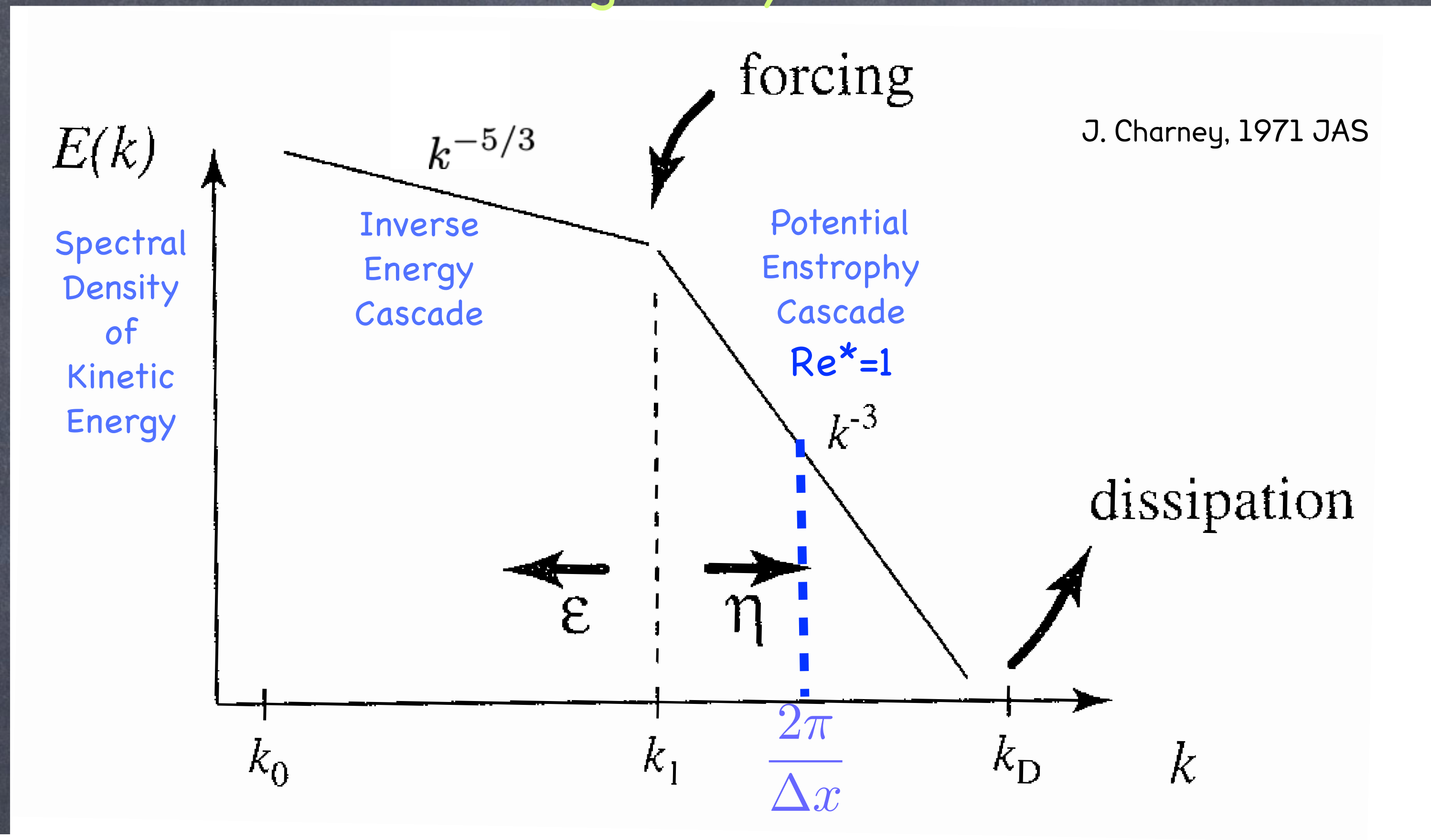


# QG Turbulence: Pot'l Enstrophy cascade

## A Mesoscale Ocean Large Eddy Simulations Closure

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QG Leith Parameterization

$$\nu_{qg} = \left( \frac{\Lambda_{qg} \Delta x}{\pi} \right)^3 |\nabla q_{qg}|$$

But, in QG, PV links buoyancy (diff.) to vorticity (visa.)

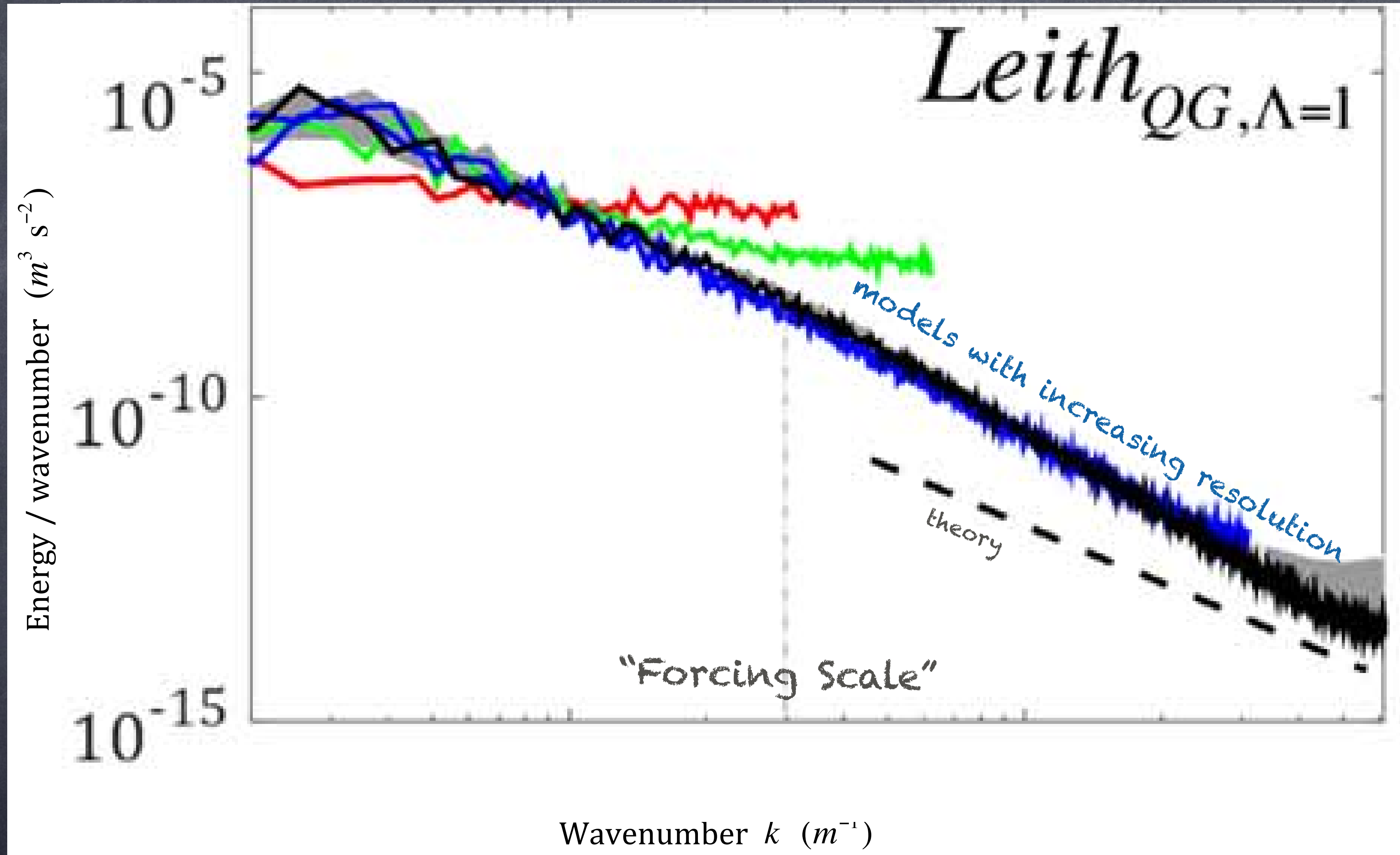
$$q_{2d}^* = f + \hat{k} \cdot \nabla \times u^*$$

$$q_{qg}^* = f + \hat{k} \cdot \nabla \times u^* + \frac{\partial}{\partial z} \frac{f^2}{N^2} b^*$$

$$\nu_{qg} = \kappa_{Redi} = \kappa_{GM} = \left( \frac{\Lambda_{qg} \Delta x}{\pi} \right)^3 |\nabla q_{qg}|.$$

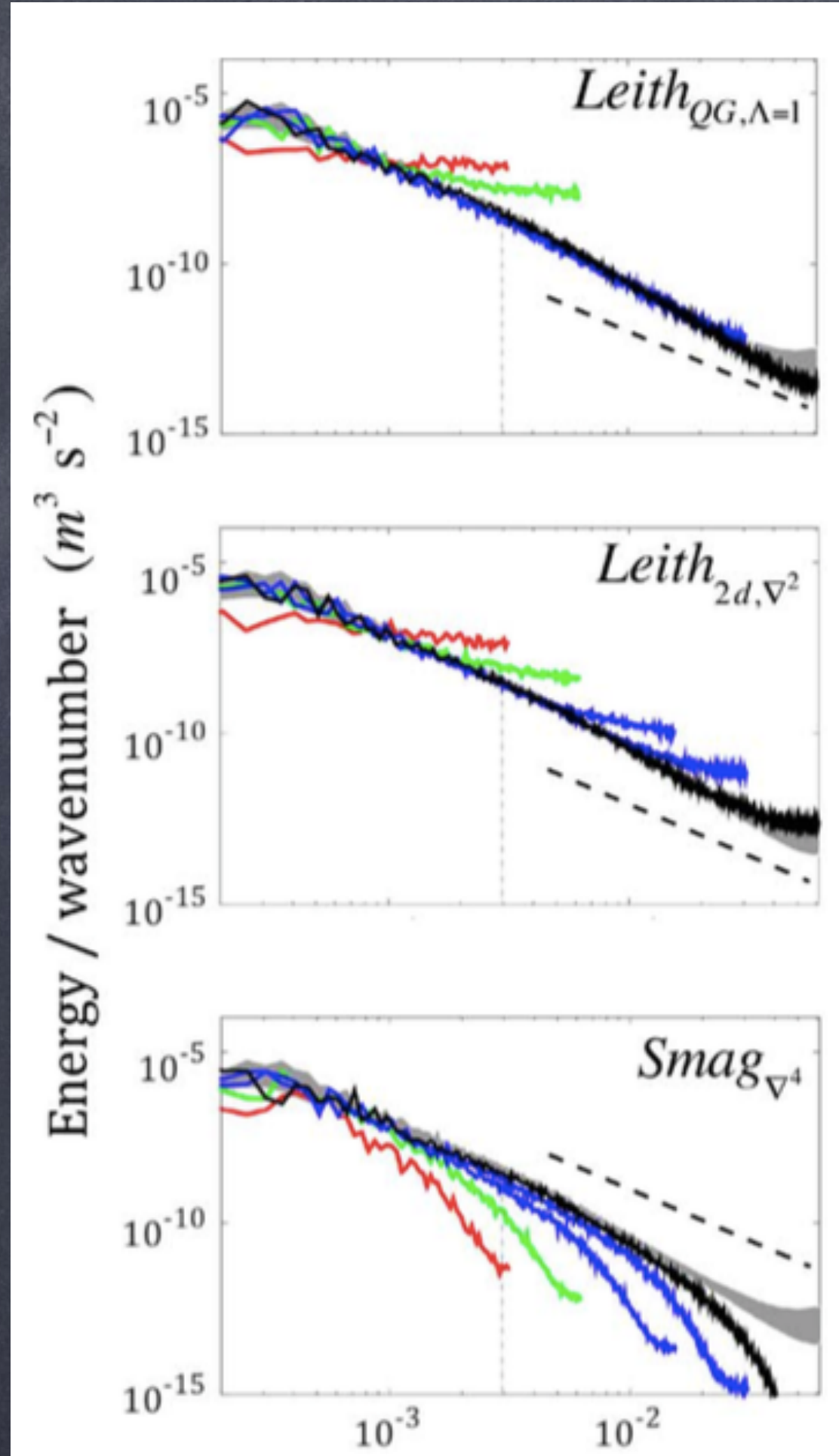
# Where does ocean energy go?

Spectrally speaking



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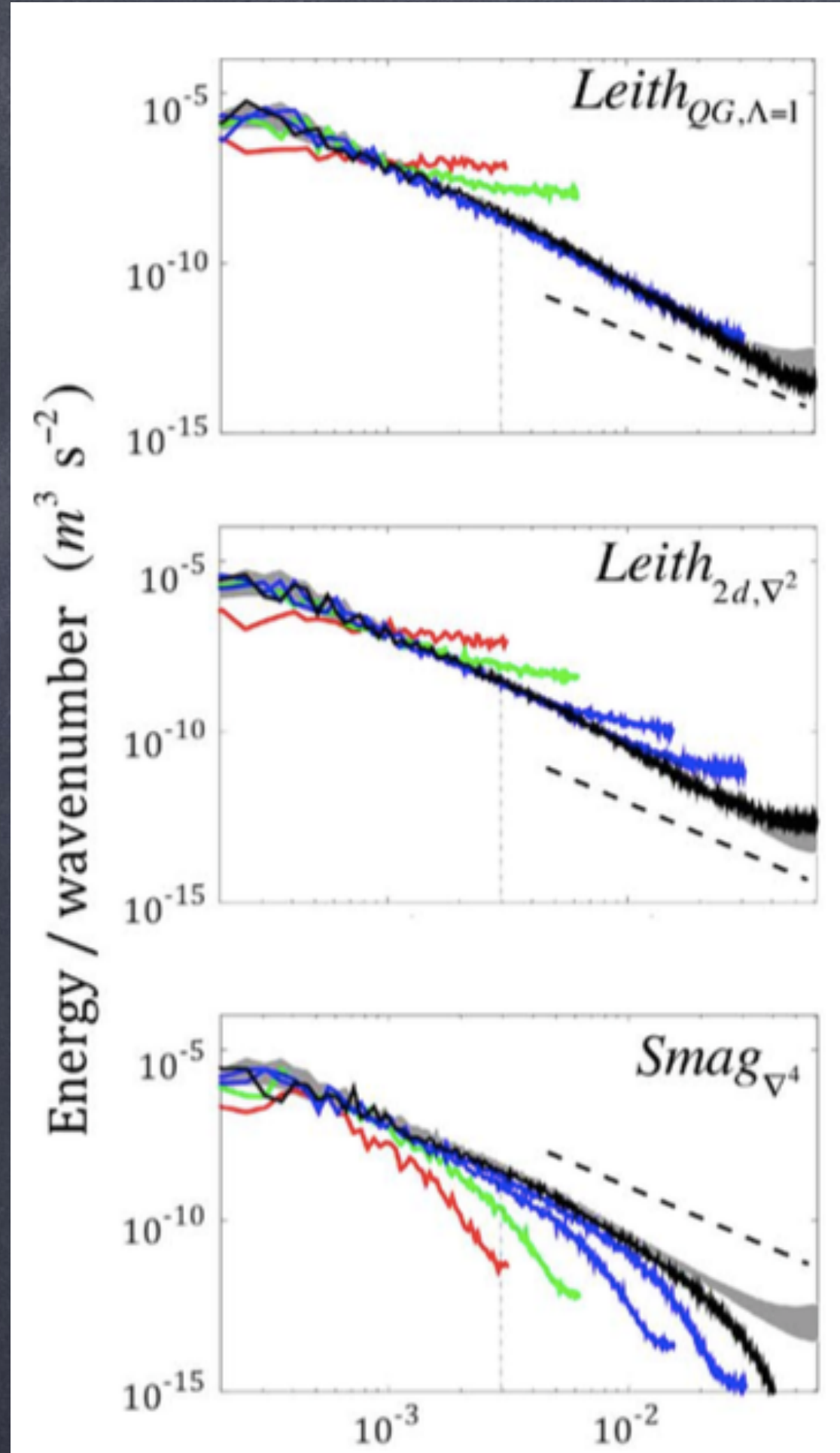
Spectrally speaking



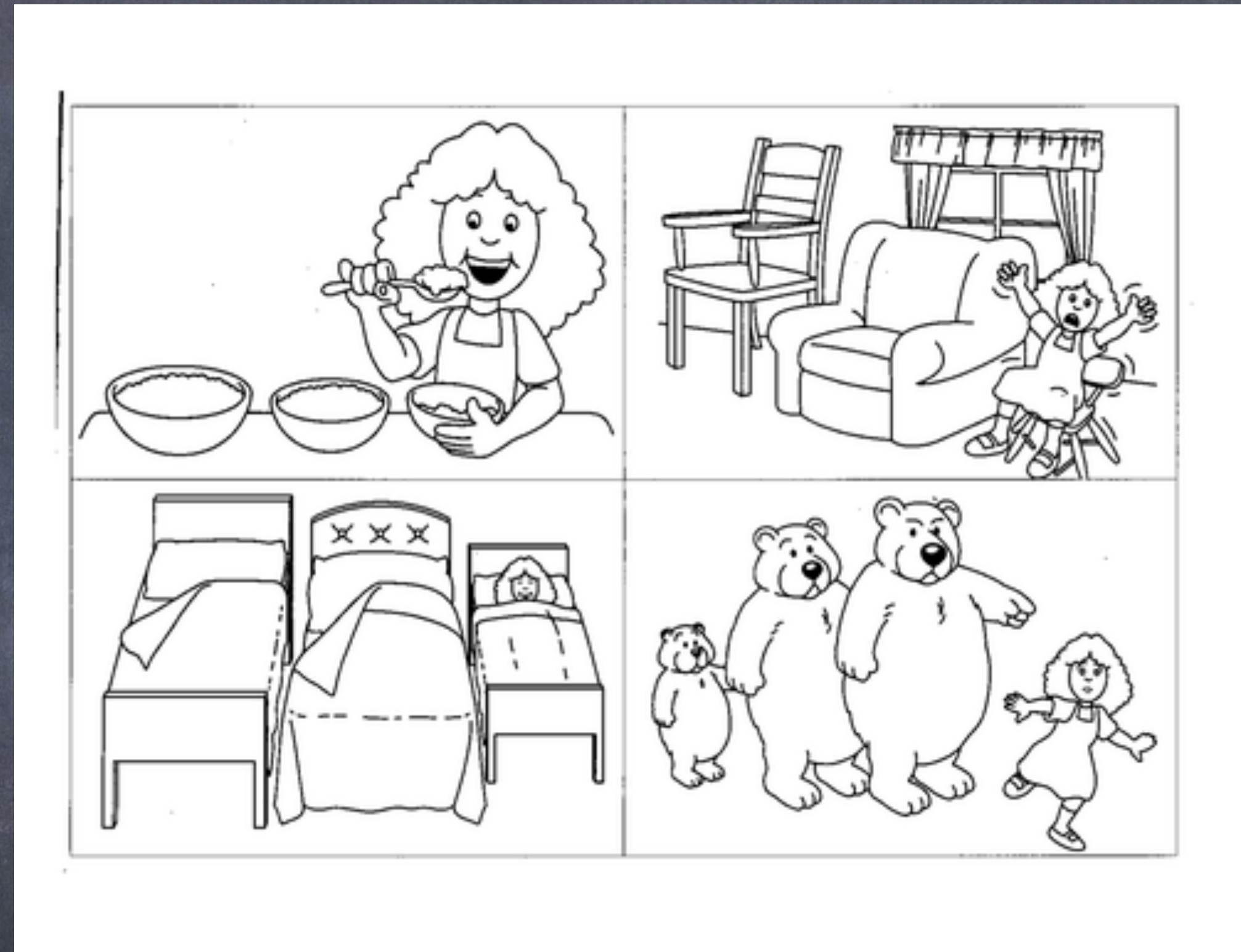
S. D. Bachman, B. Fox-Kemper, and B. Pearson, 2017: A scale-aware subgrid model for quasi-geostrophic turbulence. *Journal of Geophysical Research—Oceans*, 122:1529–1554. URL <http://dx.doi.org/10.1002/2016JC012265>.

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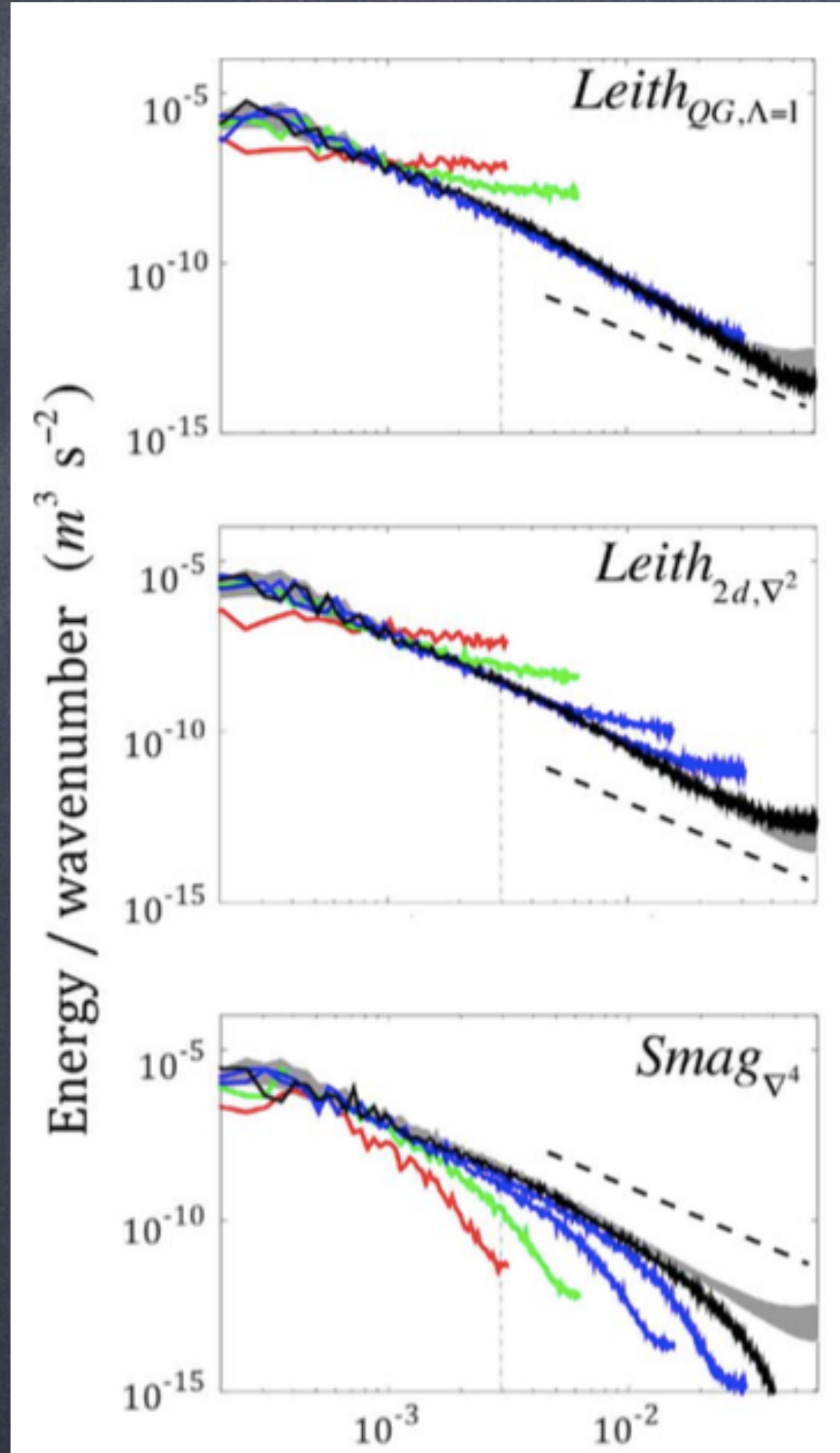
Too Smooth



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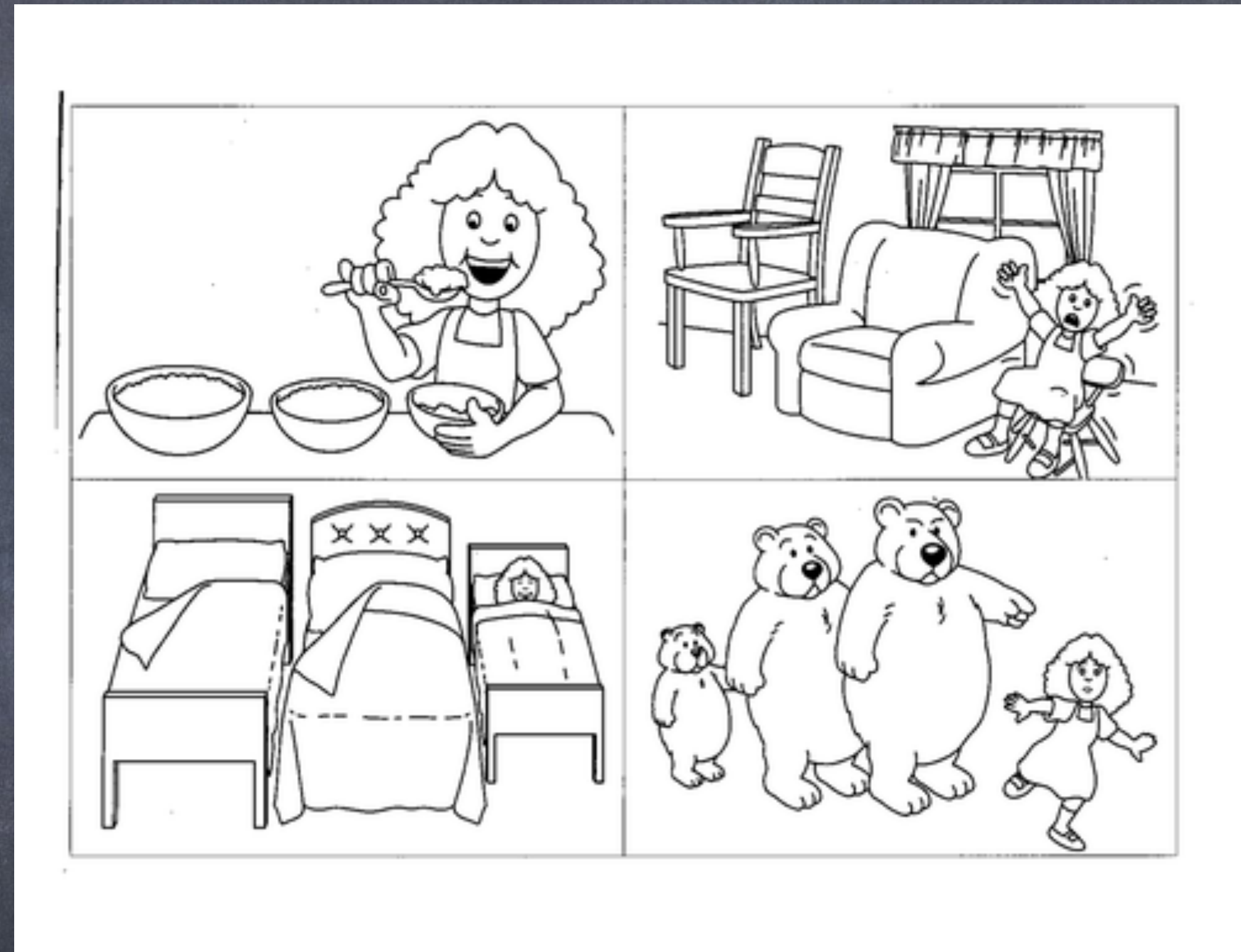
# Where does ocean energy go?

Spectrally speaking



Too Noisy

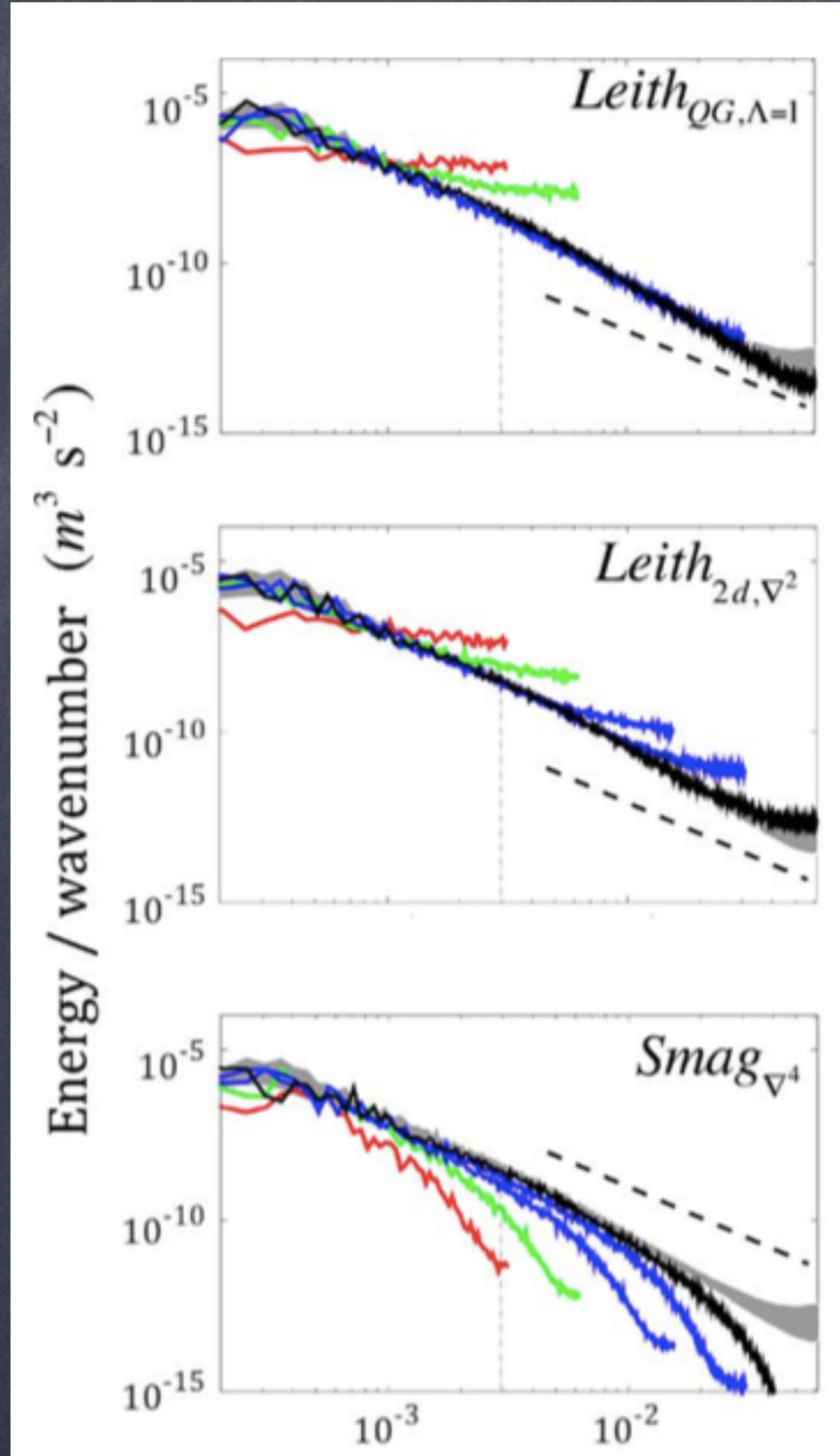
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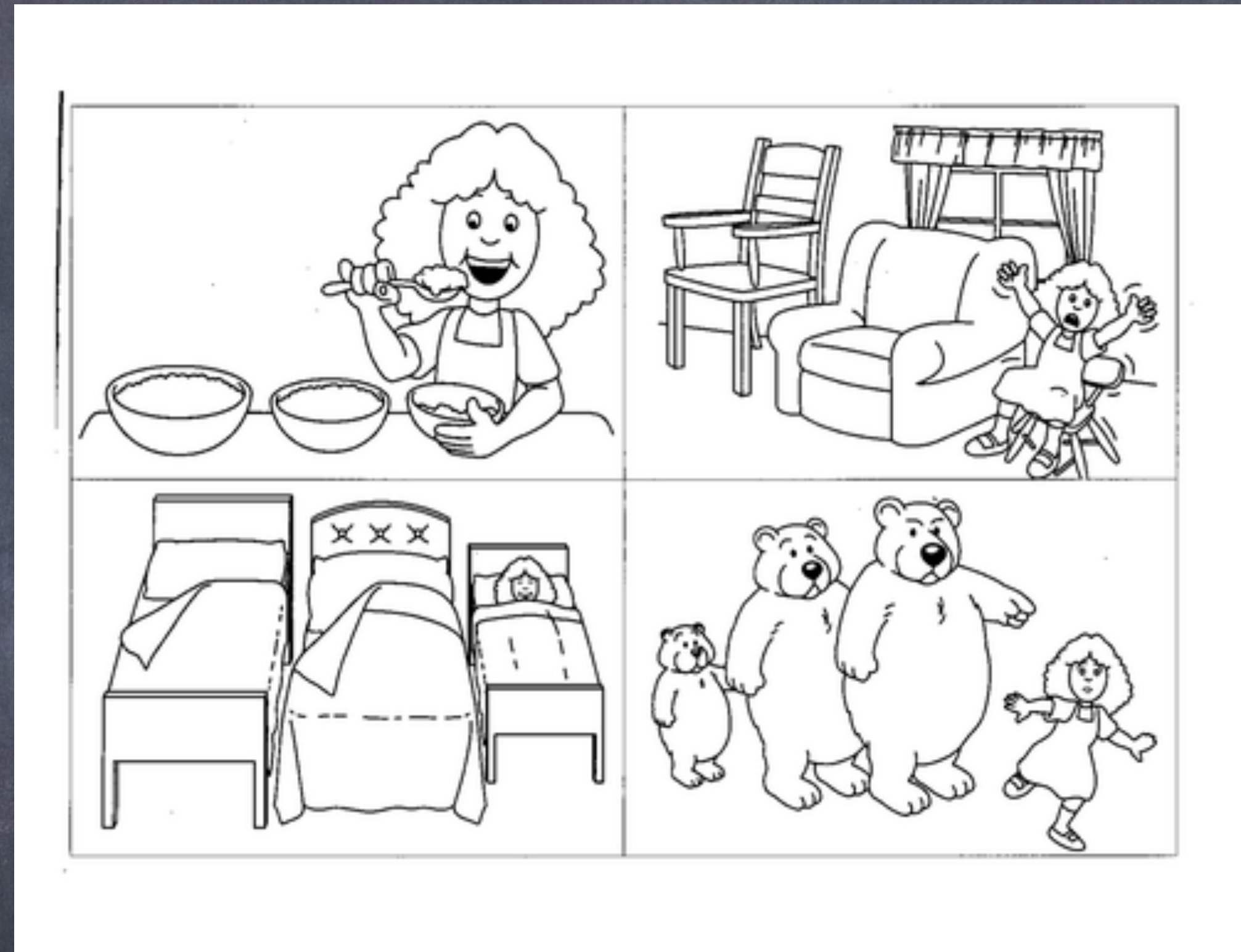
Spectrally speaking



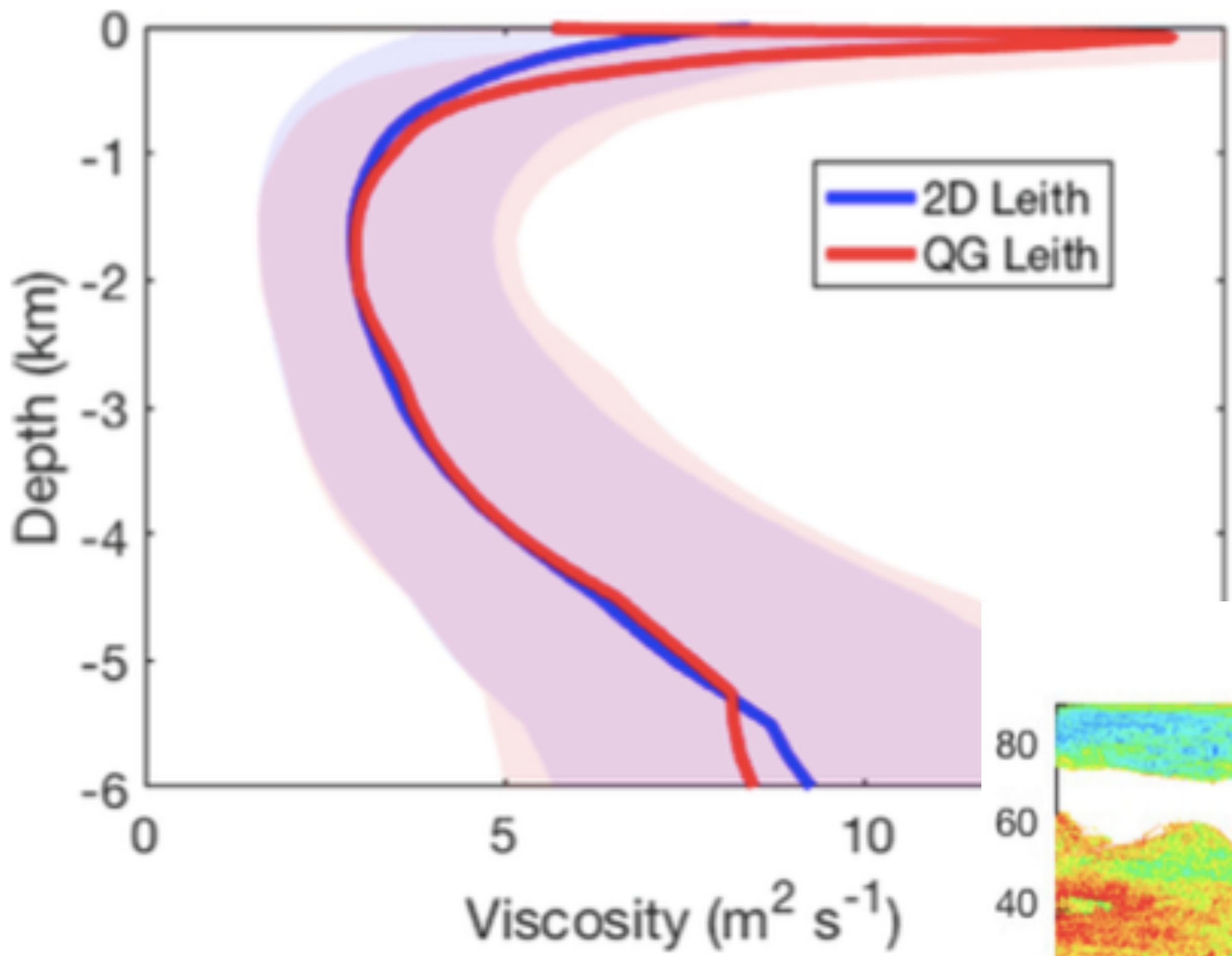
Just Right

Too Noisy

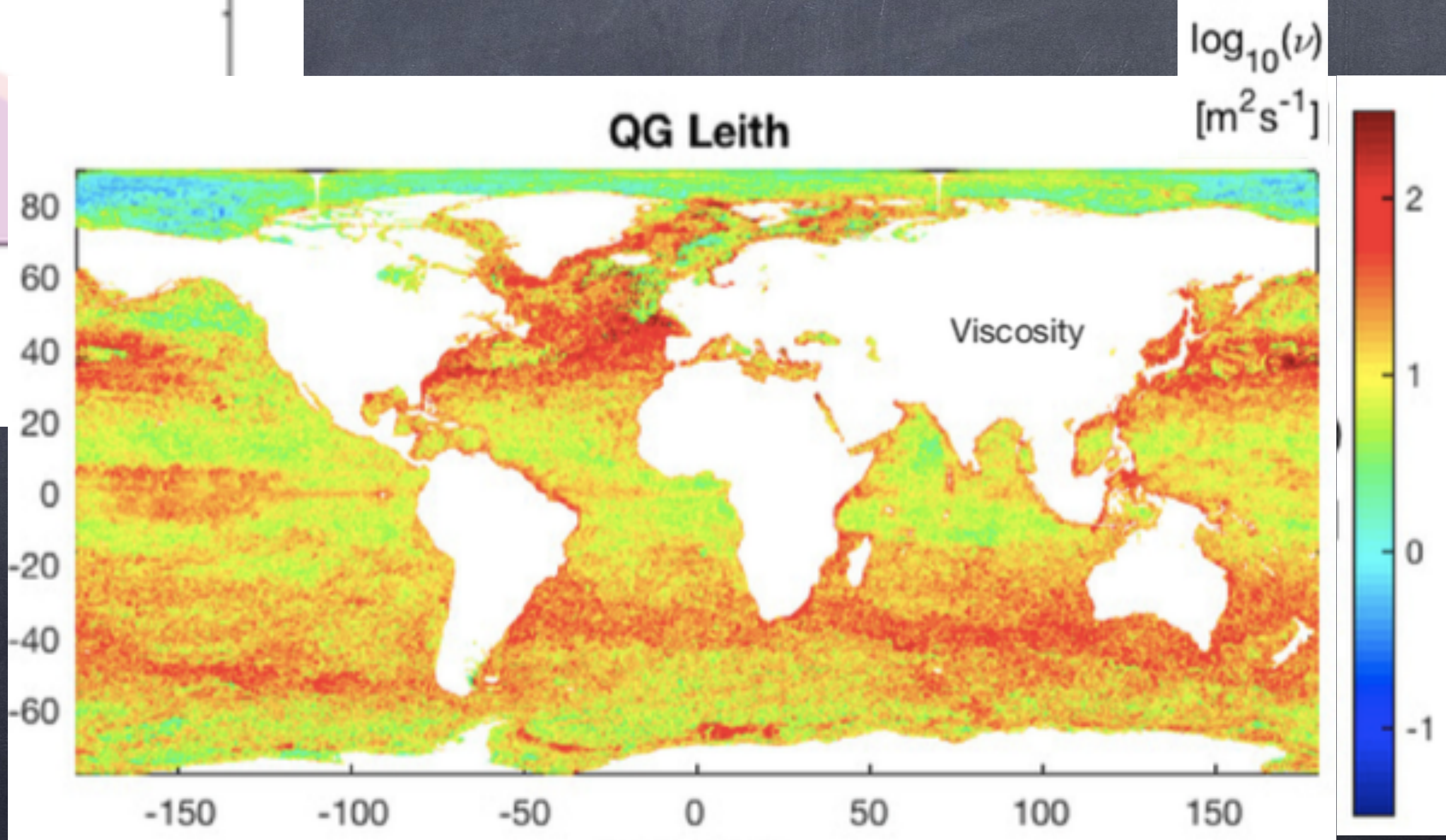
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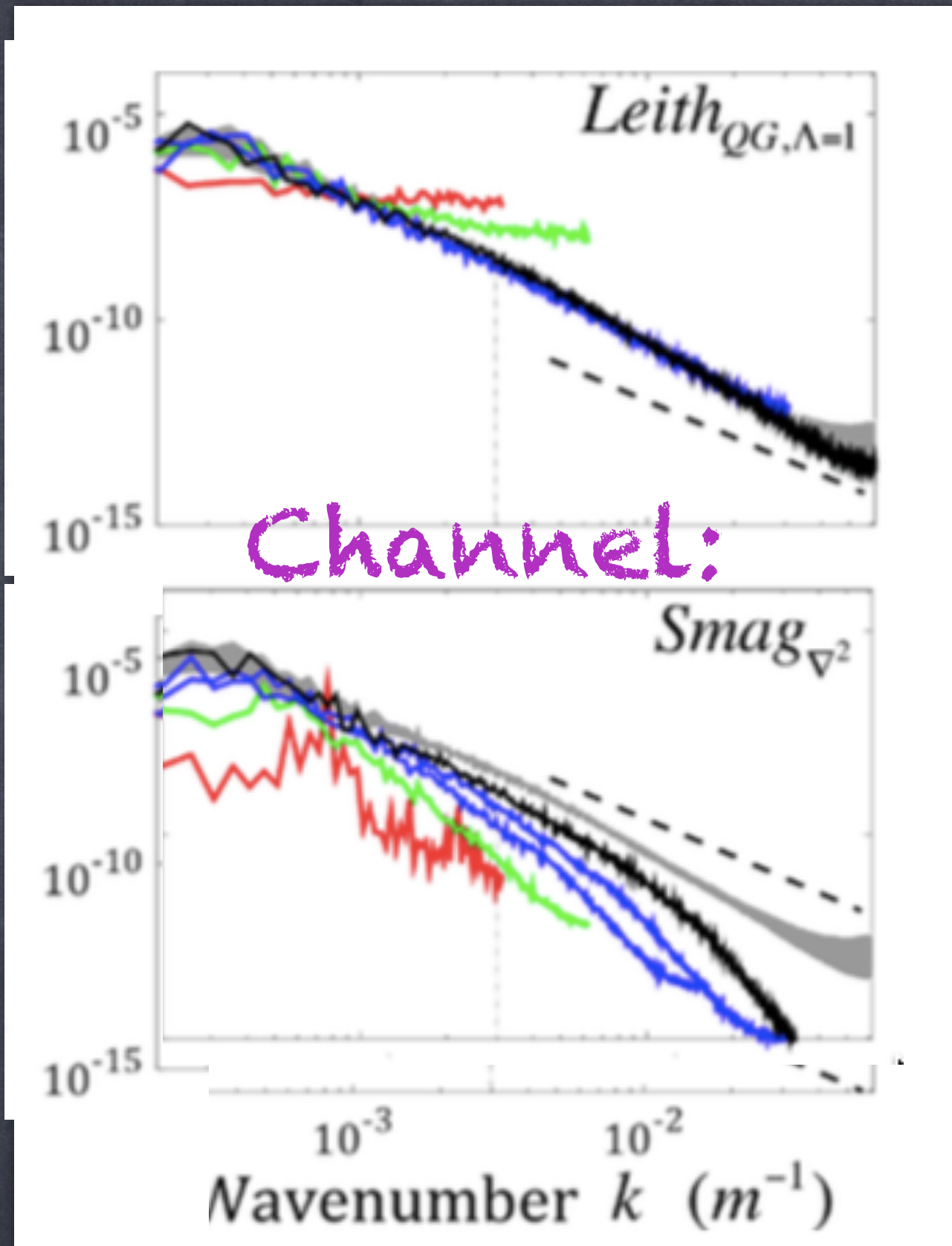
**QG Leith:**  
Works OK in an idealized flow:  
Let's try it in a realistic, global  
model!



B. Pearson, BFK, S. D. Bachman, and F. O. Bryan, 2017: Evaluation of scale-aware subgrid mesoscale eddy models in a global eddy-rich model. *Ocean Modelling*, 115:42–58.



# Mesoscale Ocean LES: QGLEith



LES for Pot'l  
Enstrophy

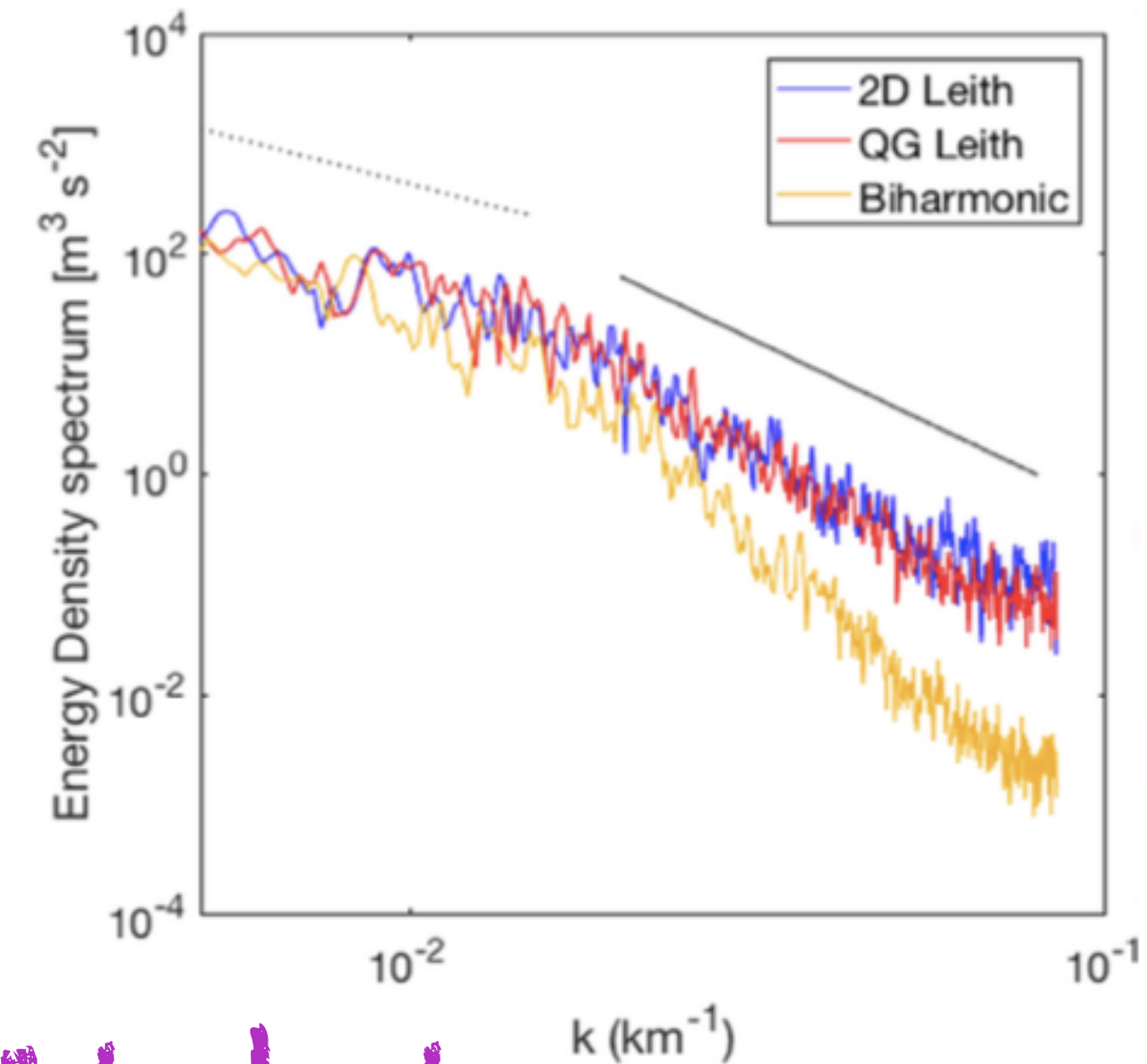
Channel:

LES for EKE

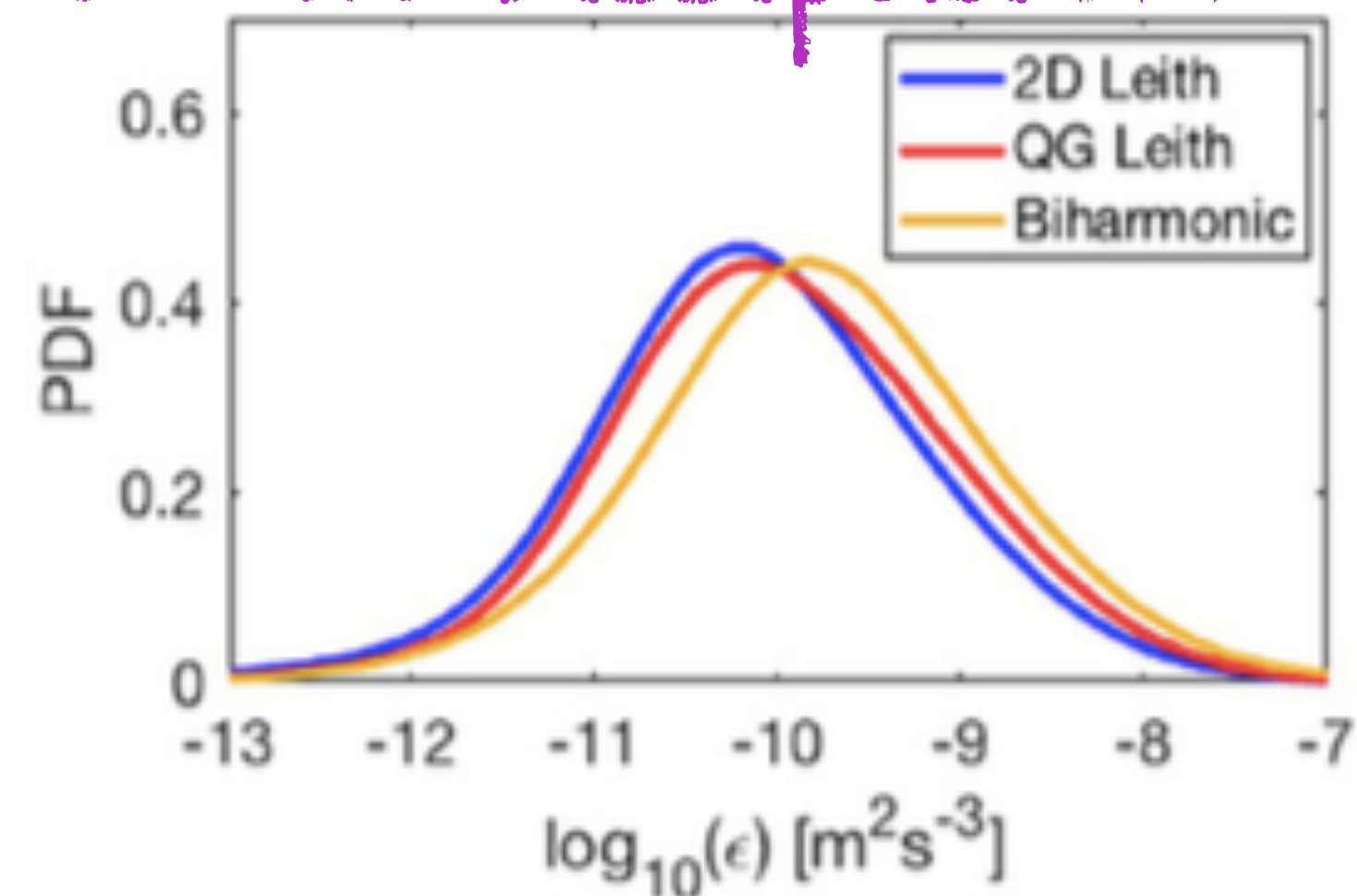
Global, POP, realistic forcing  
10km (nominal) global  
42 vertical levels  
(most in upper 200m)

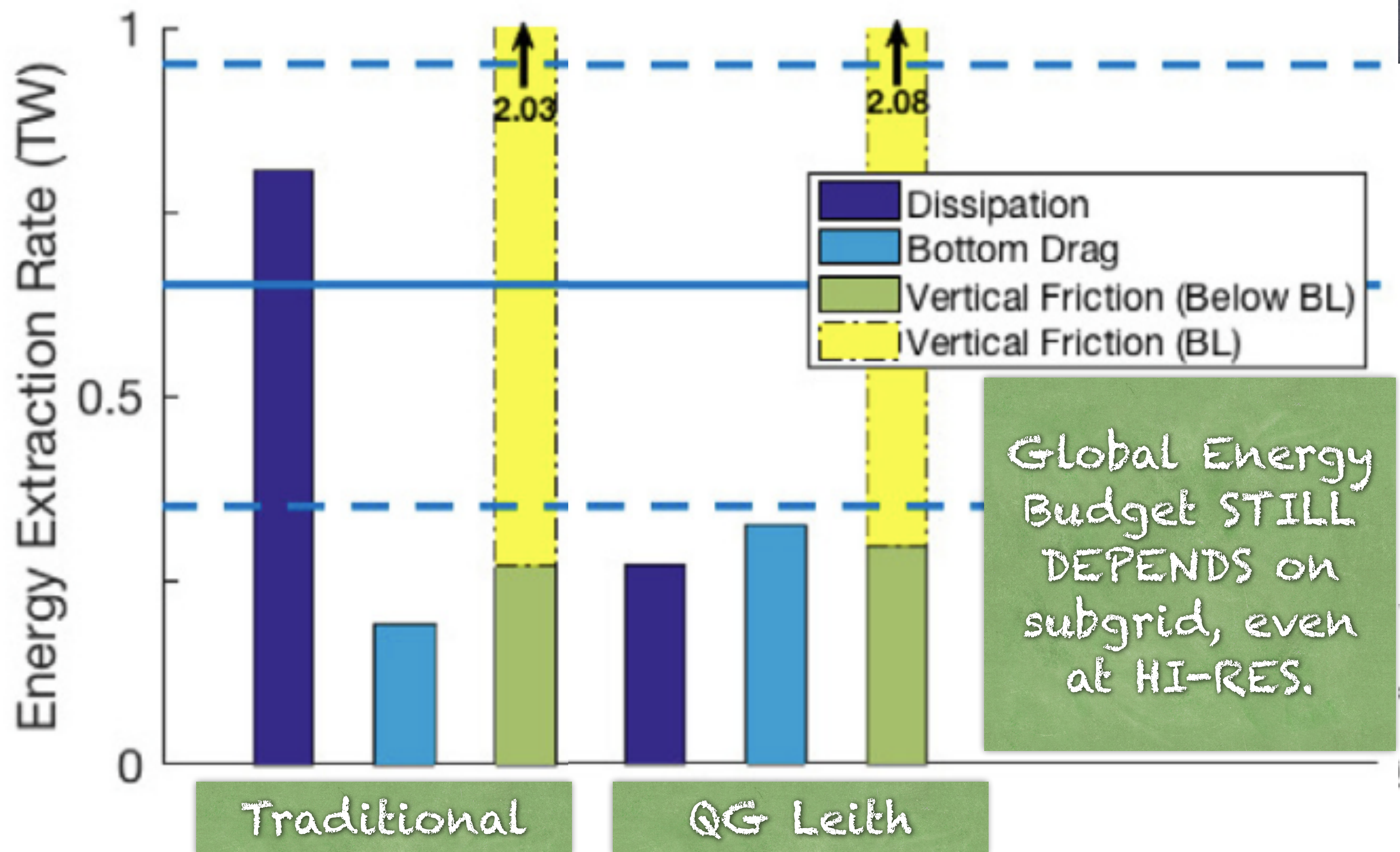
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global eddy-rich model. *Ocean Modelling*, 115:42–58.

ACC in Global!



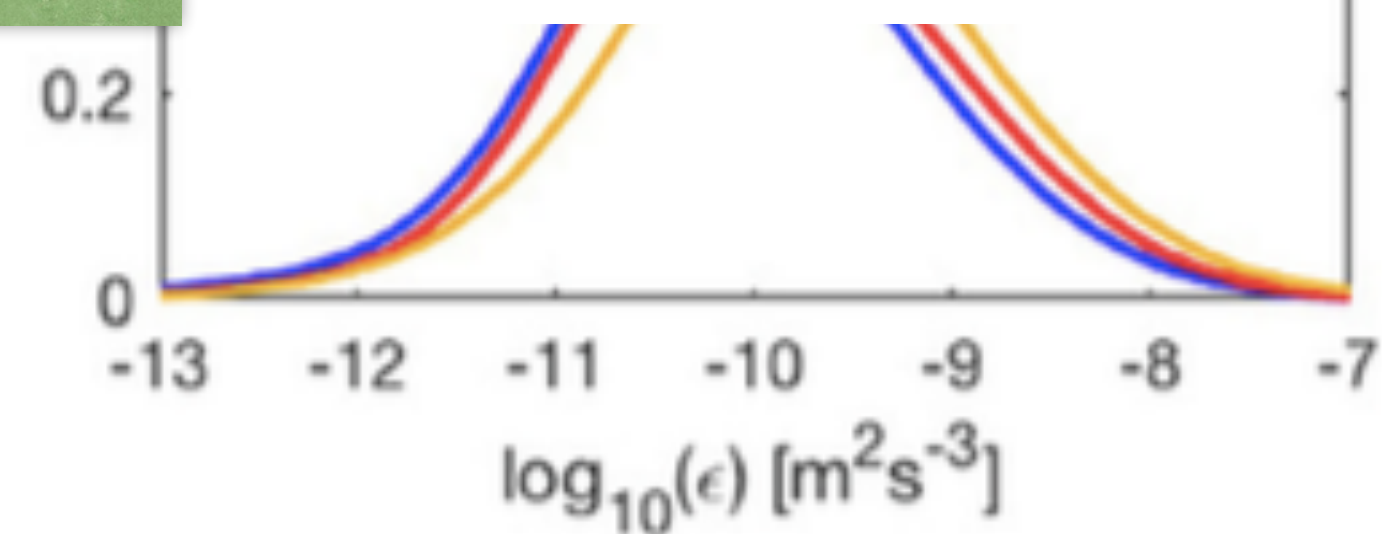
Global:  
100m Dissipation





(most in upper 200m)

B. Pearson, BFK, S. D. Bachman, and F. O. Bryan, 2017: Evaluation of scale-aware subgrid mesoscale eddy models in a global eddy-rich model. *Ocean Modelling*, 115:42–58.





# A fun & meaningful result!

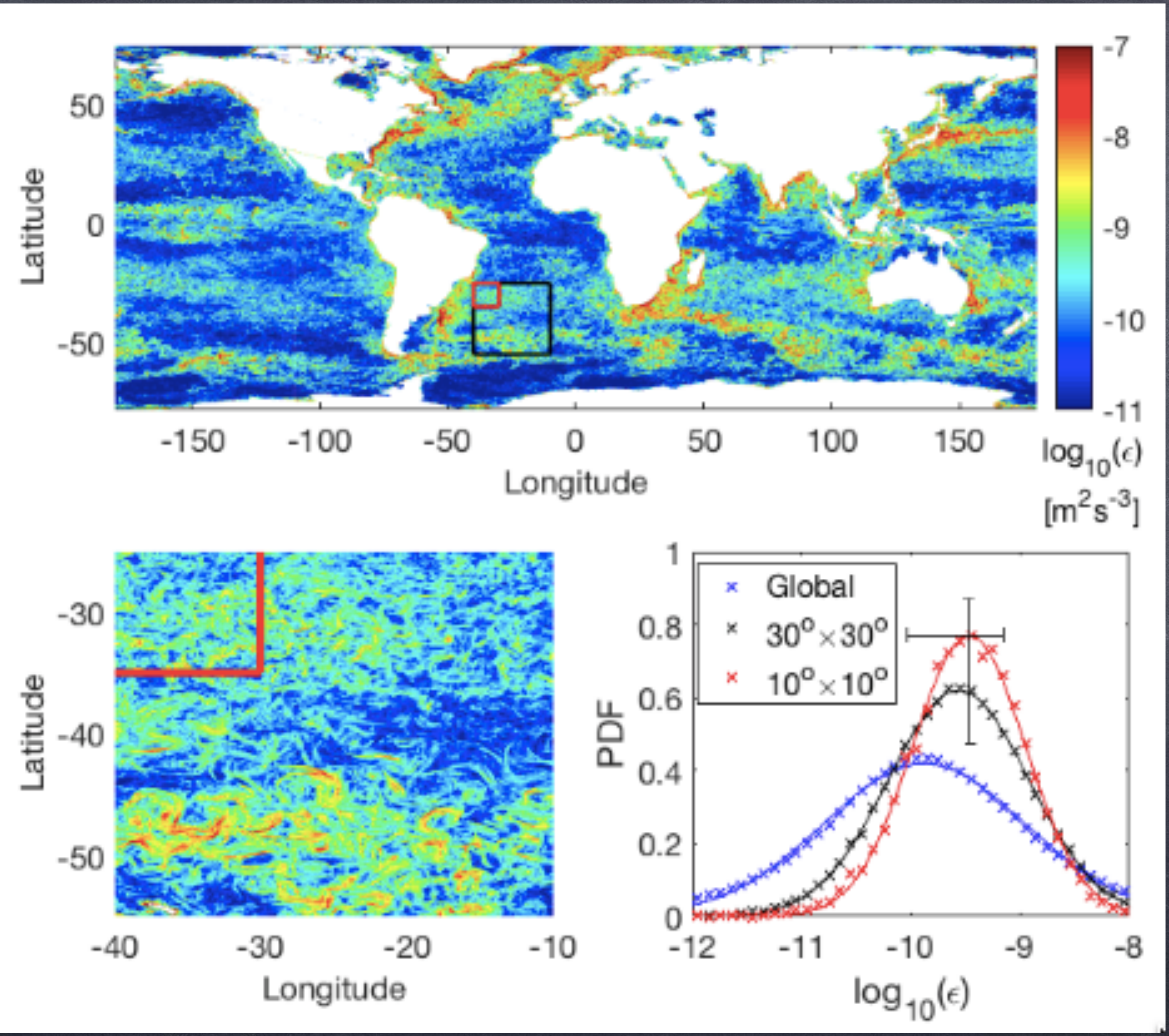
(Credit conversations with Royce in Dresden)

A (weak) dissipation of energy with pot'l enstrophy cascade

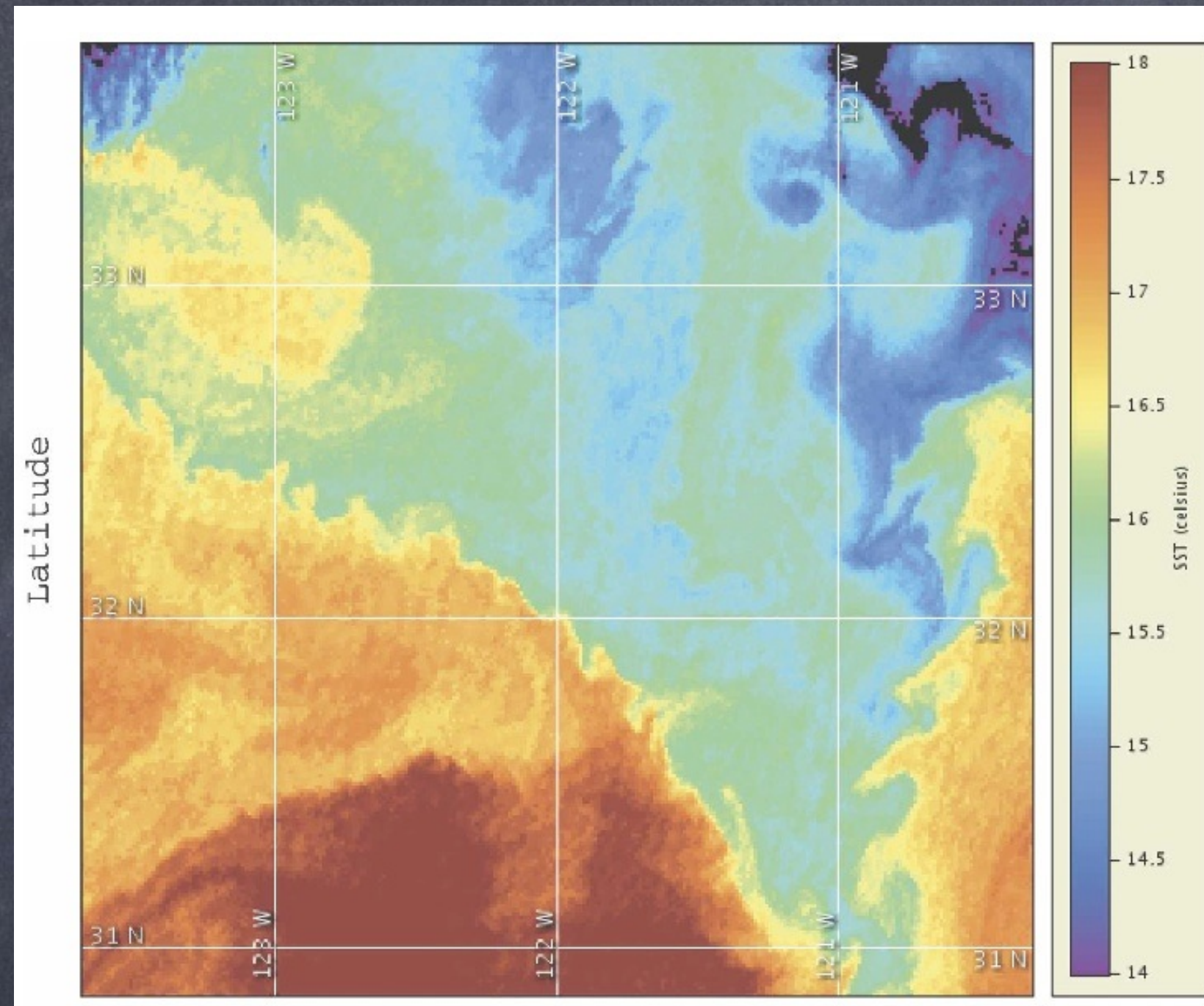
... that's lognormally distributed (super-Yaglom '66)

90% of KE dissipation in 10% of ocean

B. Pearson and BFK. Log-normal turbulence dissipation in global ocean models. Physical Review Letters, 120(9):094501, March 2018.

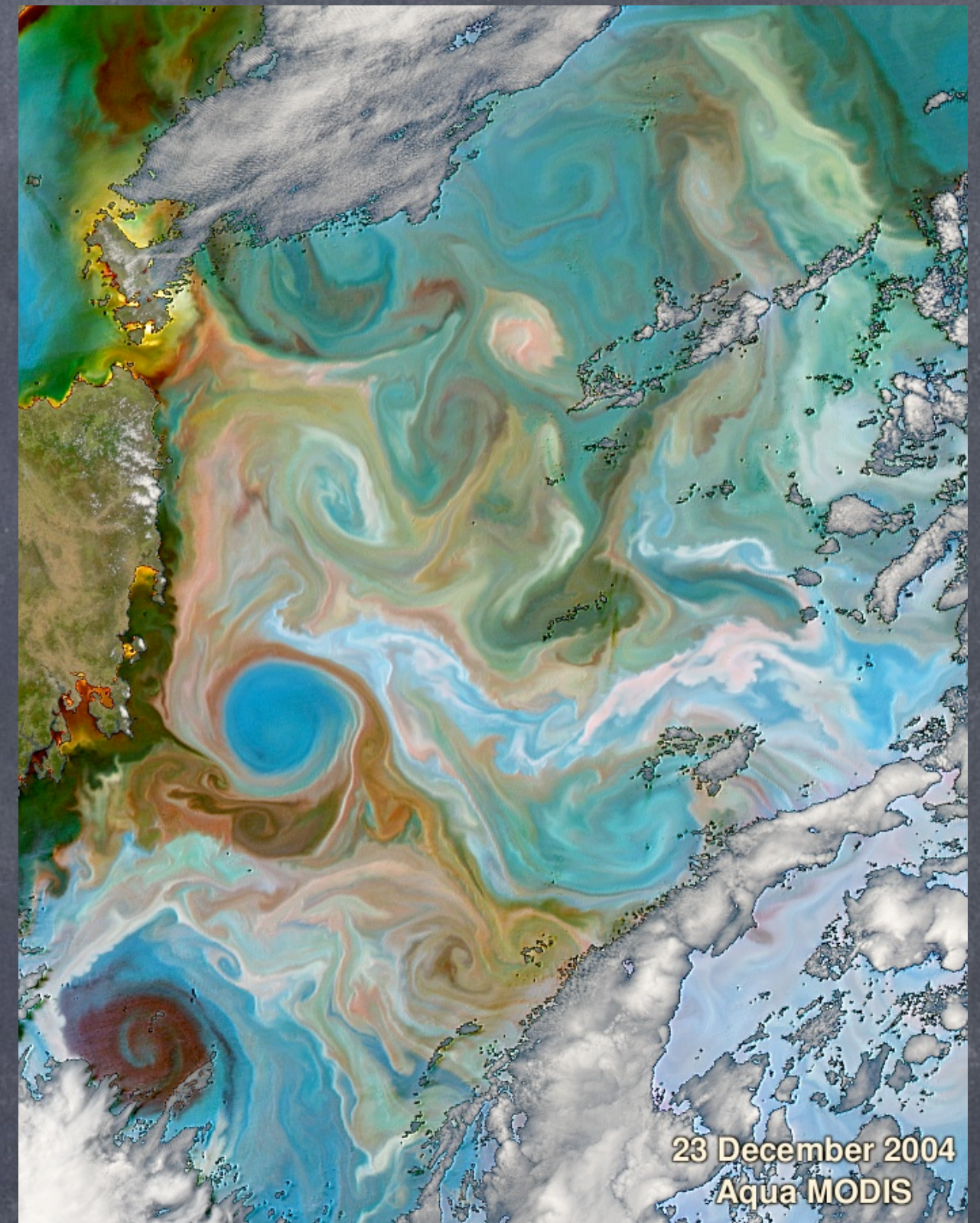


# Character of the Submesoscale

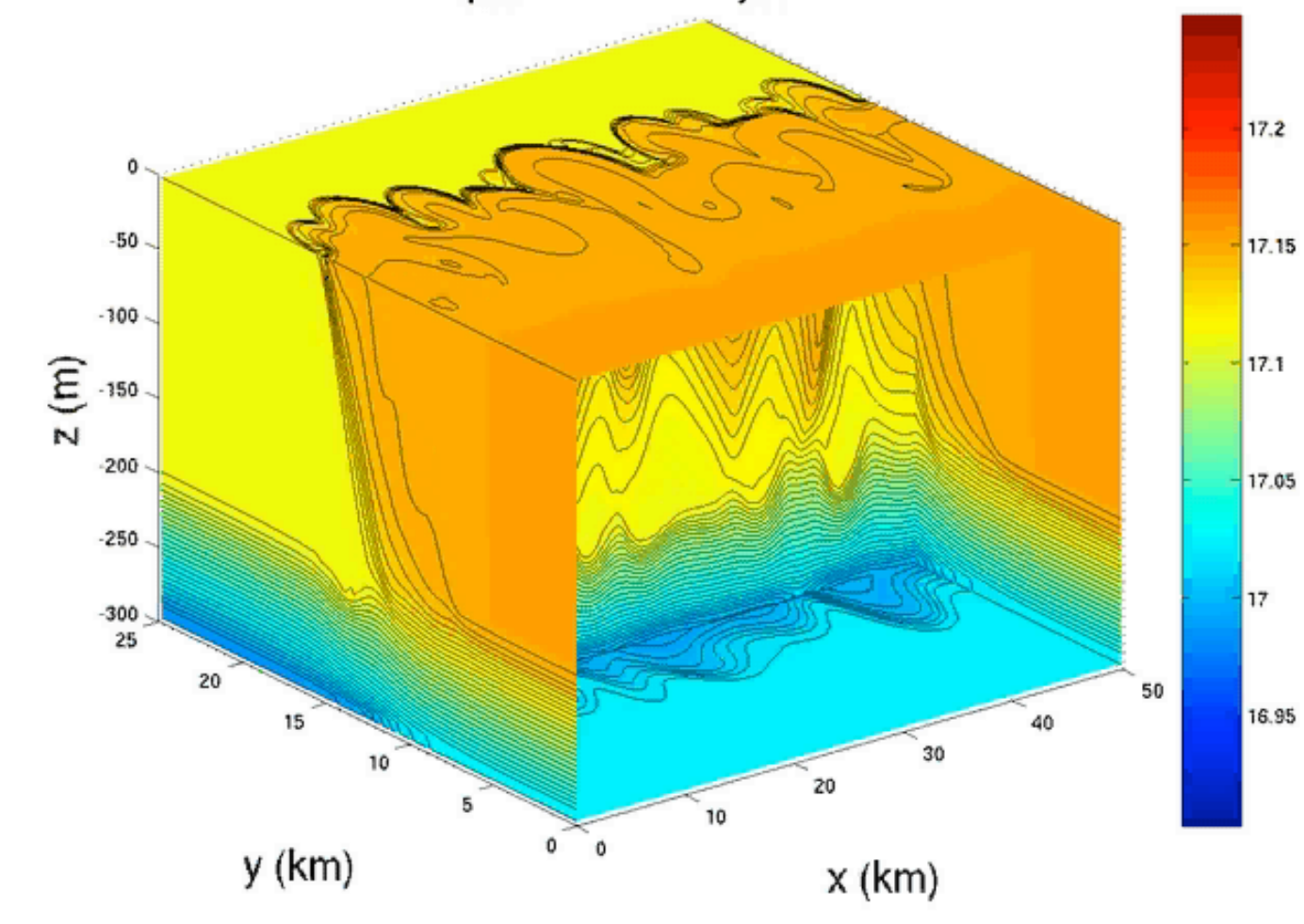


- Fronts
- Eddies
- $Ro=O(1)$
- $Ri=O(1)$
- near-surface ( $H=100m$ )
- 1-10km, days
- $W/H \sim U/L$
- hydrostatic
- Globally resolved in 2070-2100

← 10 km



Temperature on day:17.375

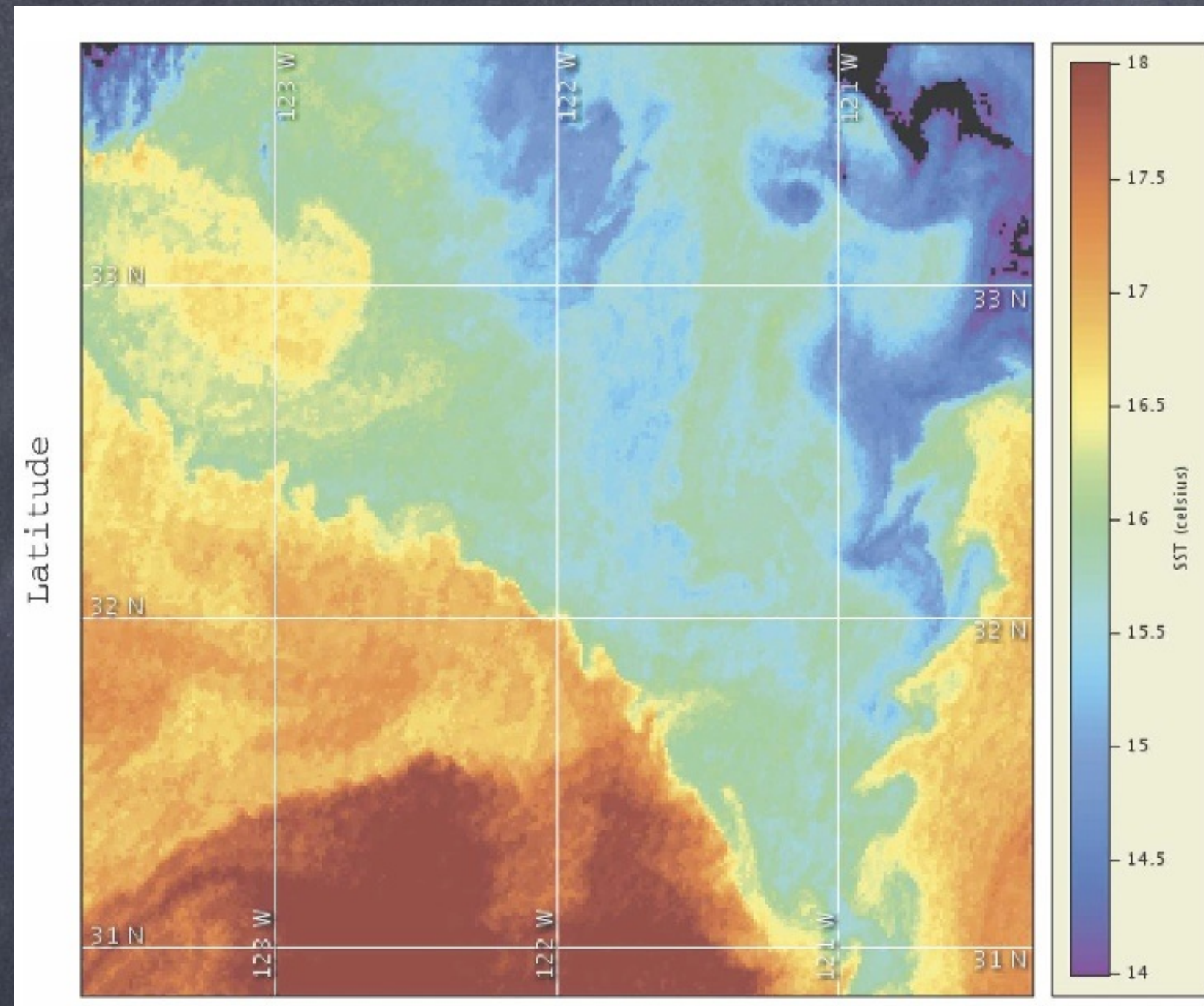


Eddy processes often  
baroclinic instability

BFK, R. Ferrari, and R. W. Hallberg. Parameterization of mixed layer eddies. Part I: Theory and diagnosis. *Journal of Physical Oceanography*, 38(6):1145-1165, 2008

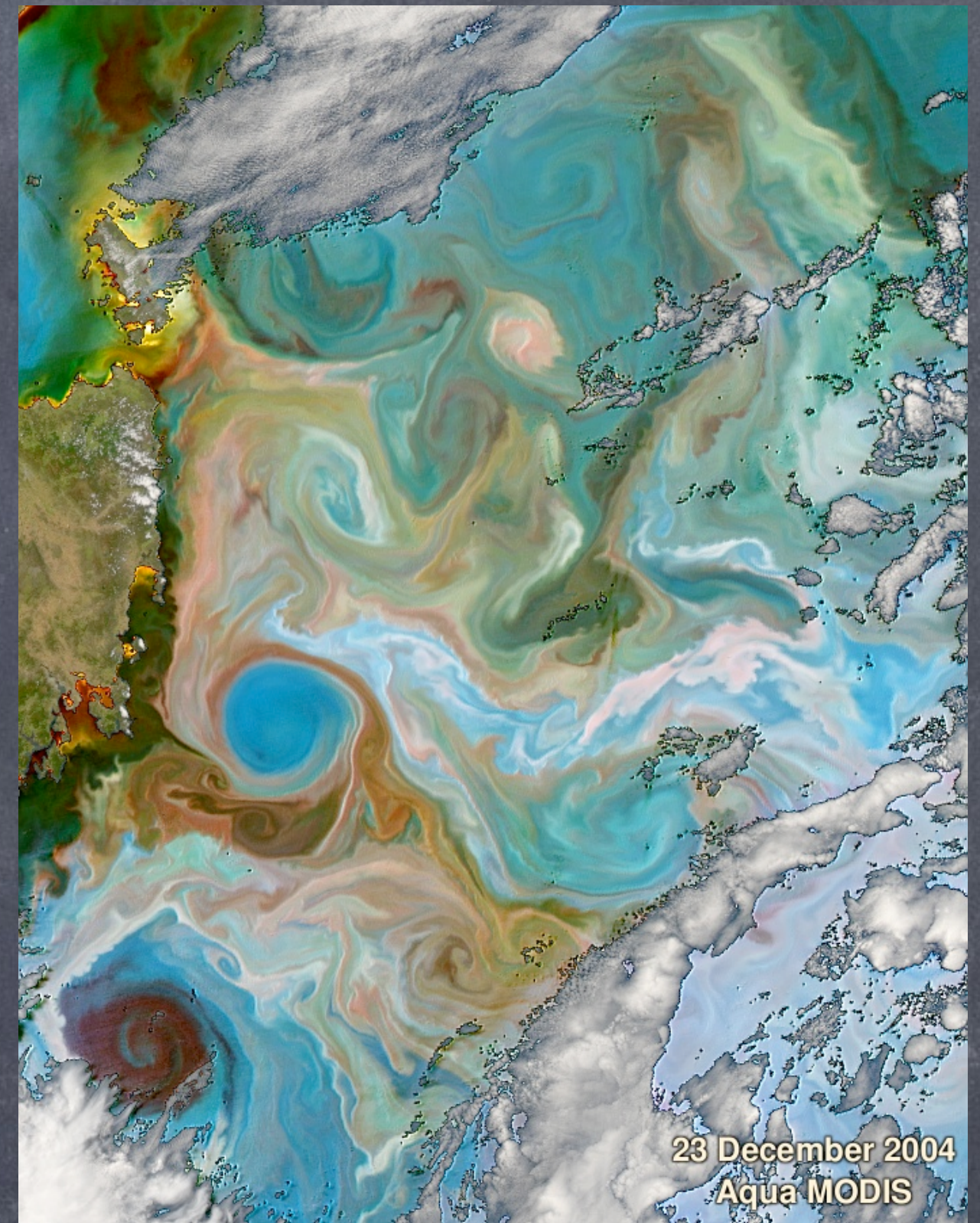
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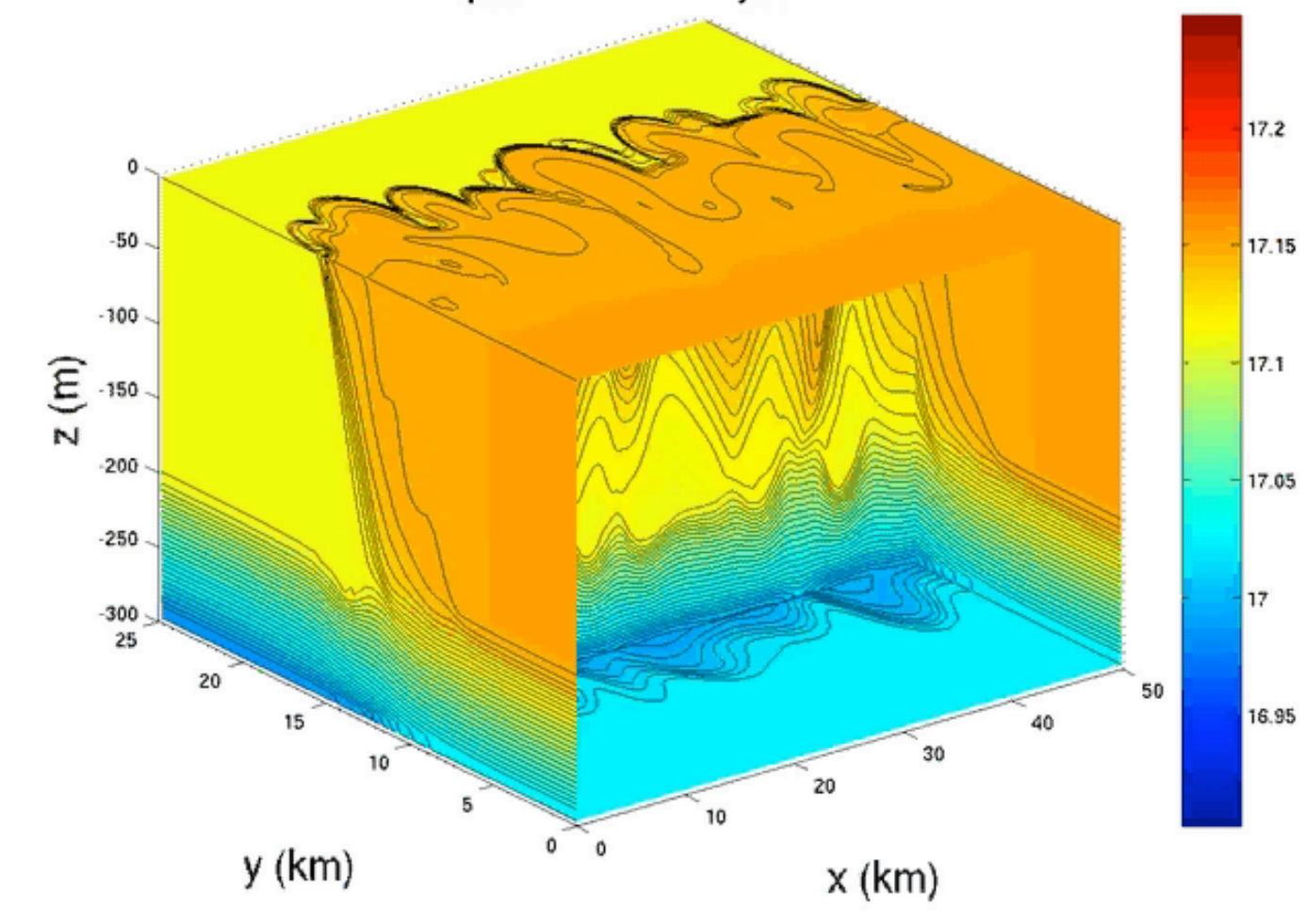


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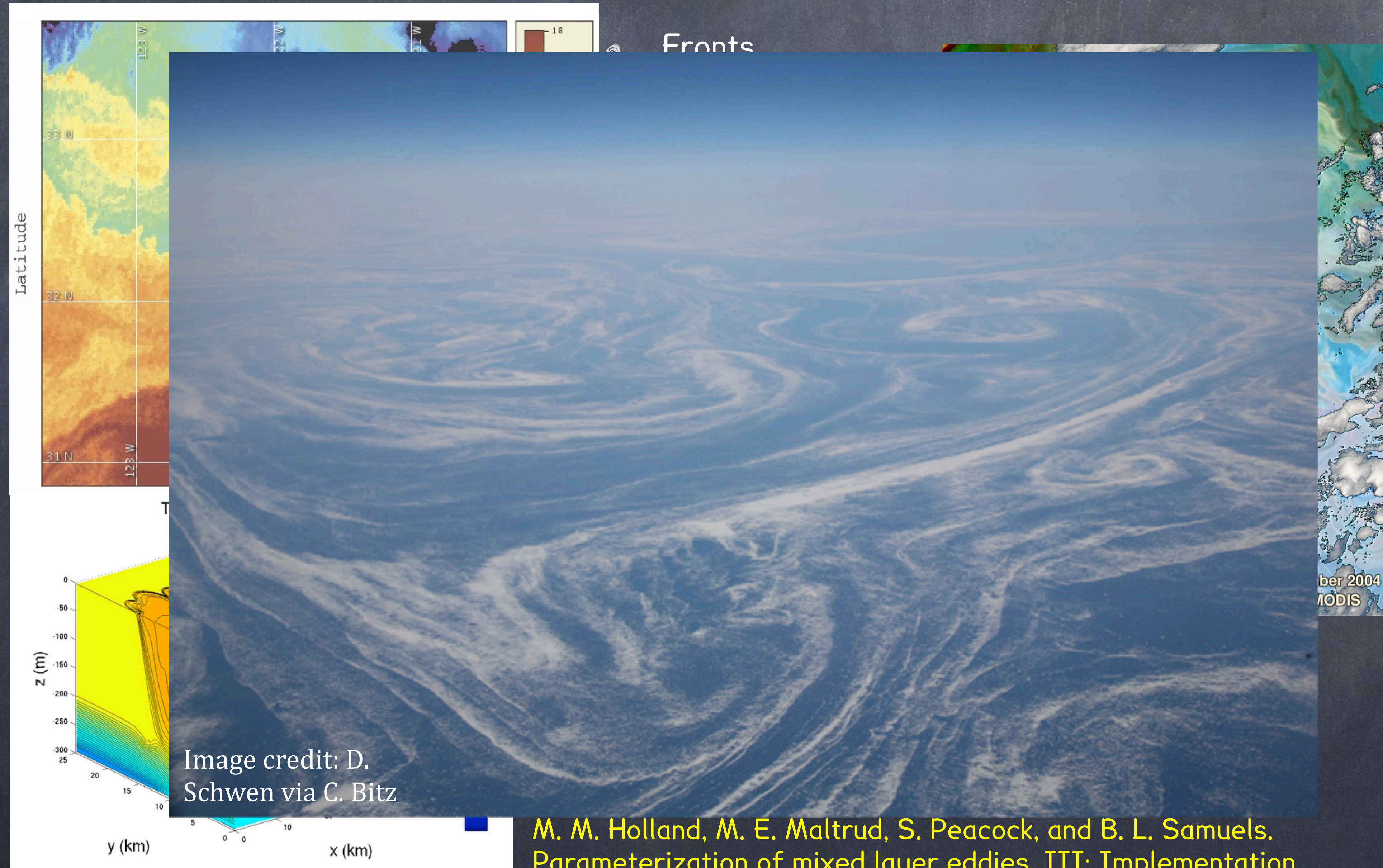


Image credit: D. Schwen via C. Bitz

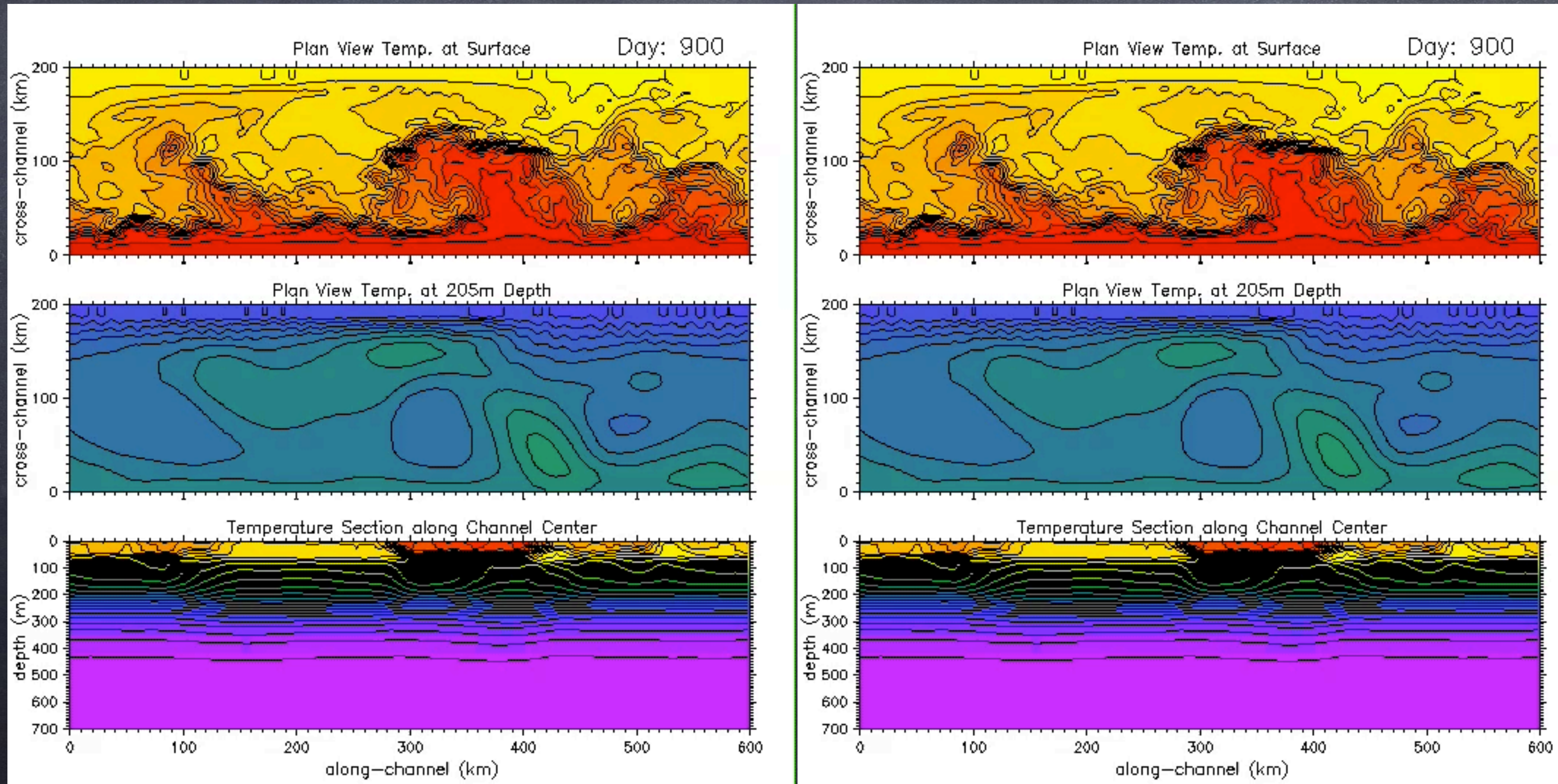
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# Having a Mixed Layer Counts!

The vertical buoyancy flux in the ML ( $\langle w'b' \rangle$ )

without diurnal cycle is **not less** than with cycle (ML)

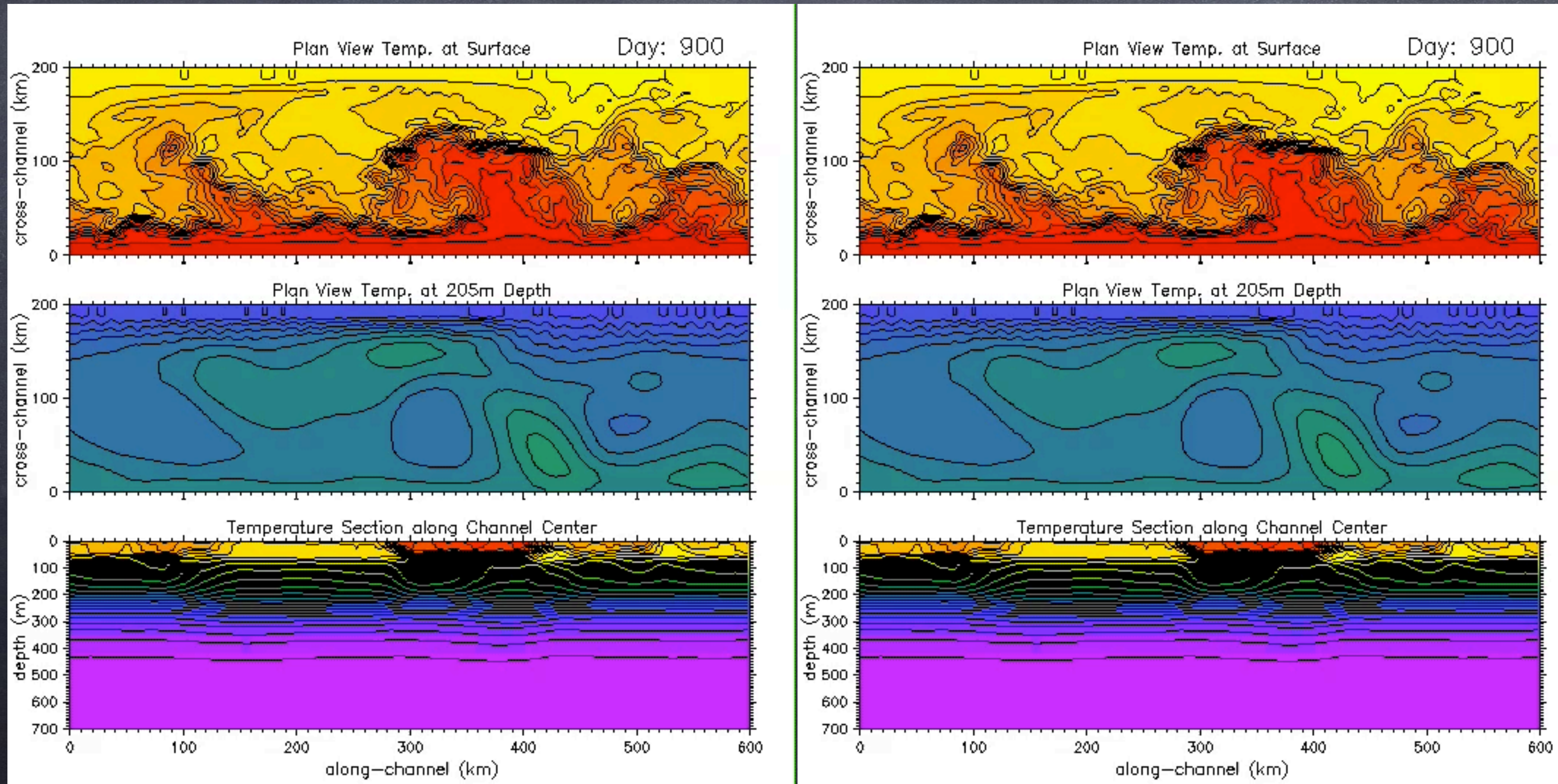


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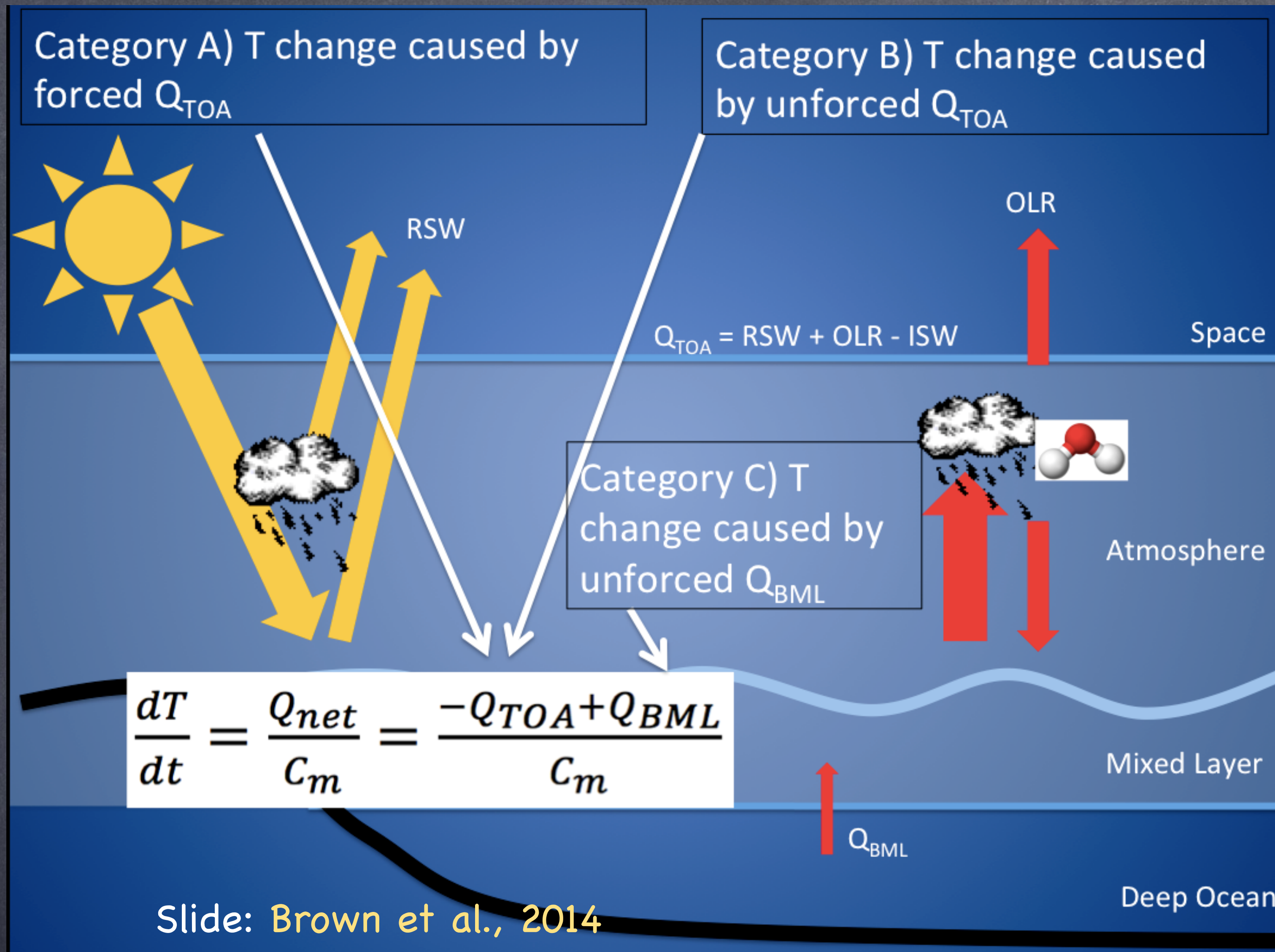
The vertical buoyancy flux in the ML ( $\langle w'b' \rangle$ )

without diurnal cycle is **4x less** than with cycle (ML)





# GMST: Surface Energy Budget=Ocean Heat Content Budget



Top of Atmosphere Imbalance!!

$$341.3 - 101.9 - 238.5 = 0.9$$

This equals net absorbed

Slide: Brown et al., 2014

- 3.4m of ocean has heat capacity of whole atmosphere
- Ocean Mixed Layer is about 100m deep.



# Surface, Mixed Layer, Seasons?

0.7 W/m<sup>2</sup>

=

Atmosphere:

1.9K/yr

=

3.4m Ocean:

1.9K/yr

=

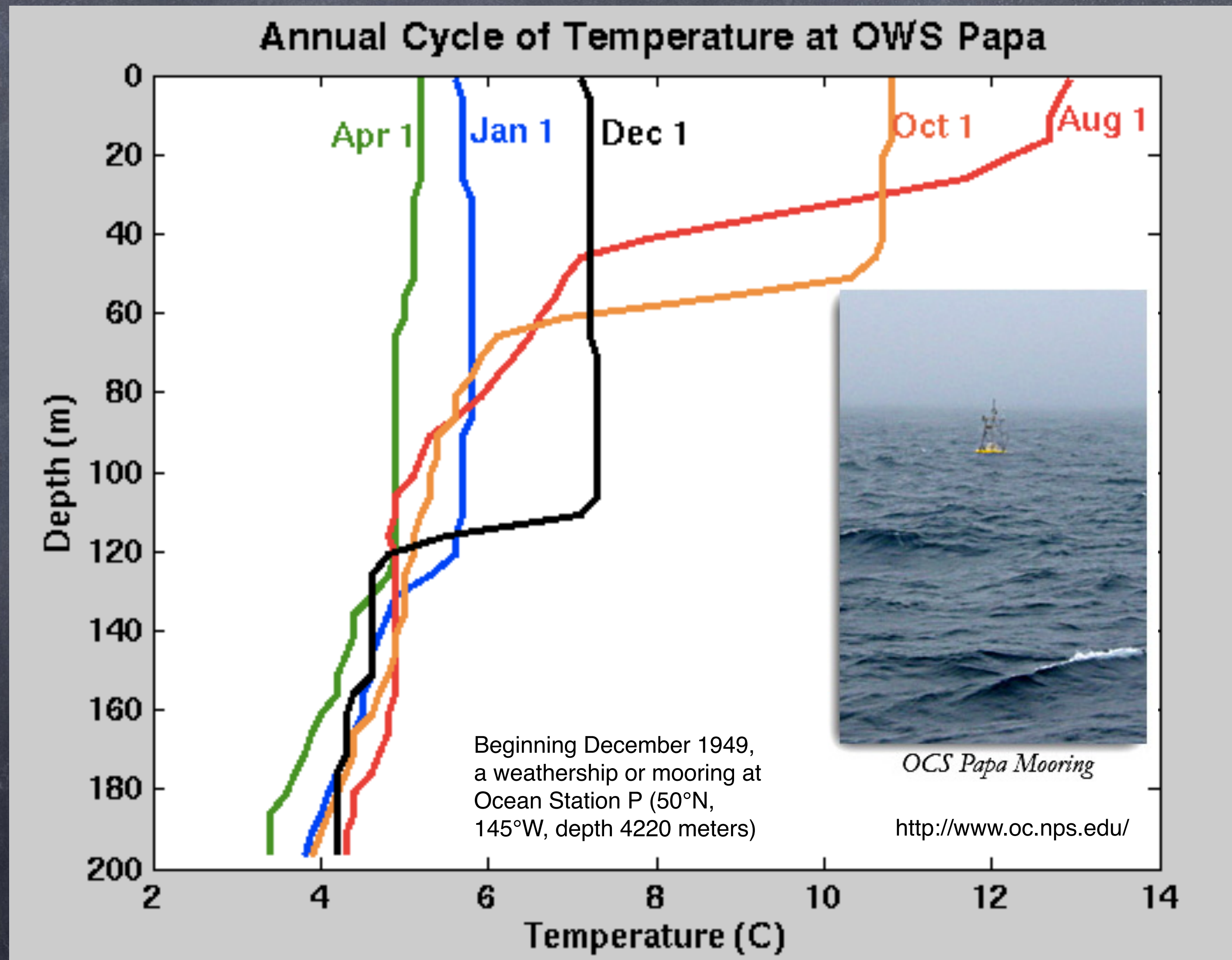
34m Ocean:

0.19K/yr

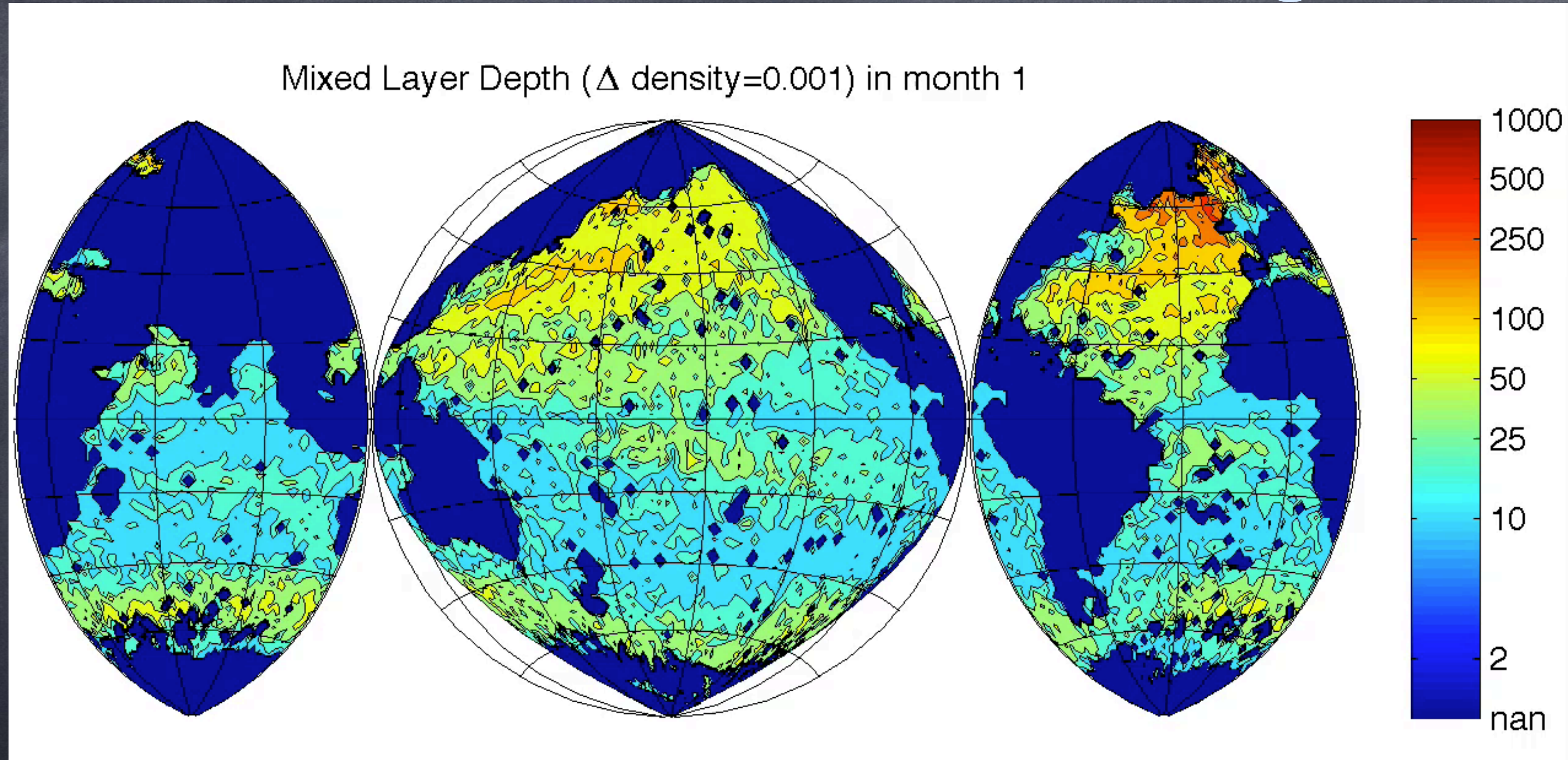
=1% of

mixed layer

seasonality

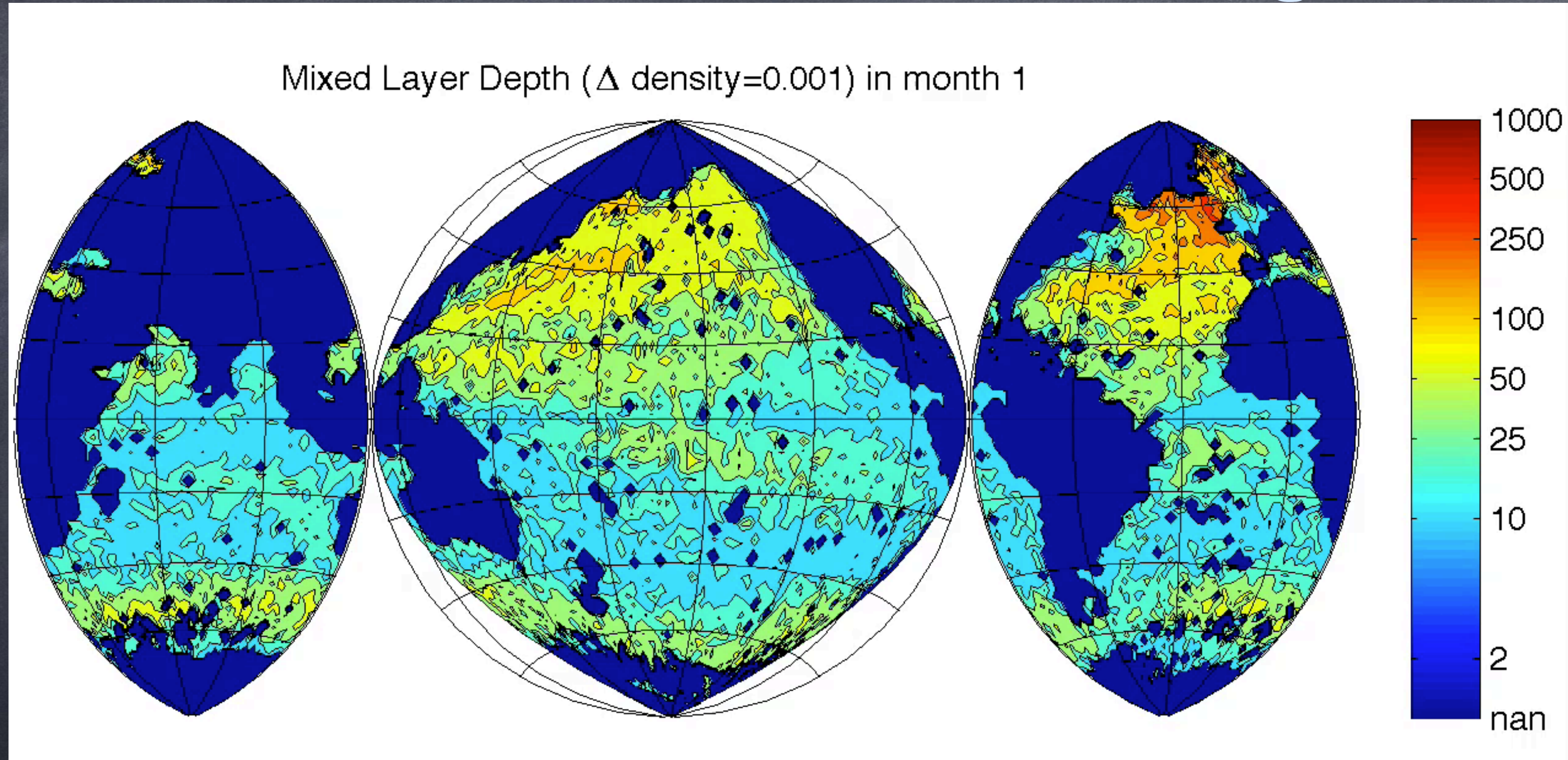


# The Ocean Mixed Layer is home to submesoscales & Langmuir



Mixed Layer Depth climatology  
From Argo float data courtesy C. de Boyer-Montegut

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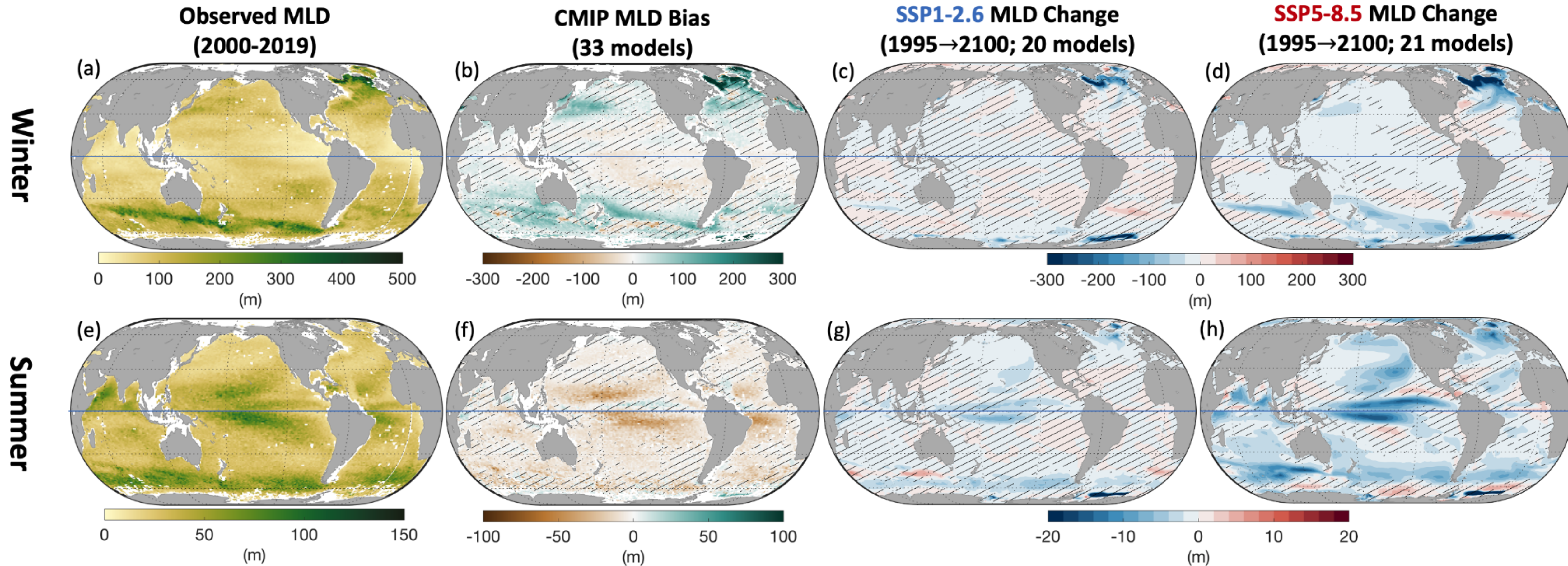


Mixed Layer Depth climatology  
From Argo float data courtesy C. de Boyer-Montegut

# Ocean Mixed Layer Depth (MLD) in Winter and Summer

Observed MLD, model MLD biases, and projected changes in MLD

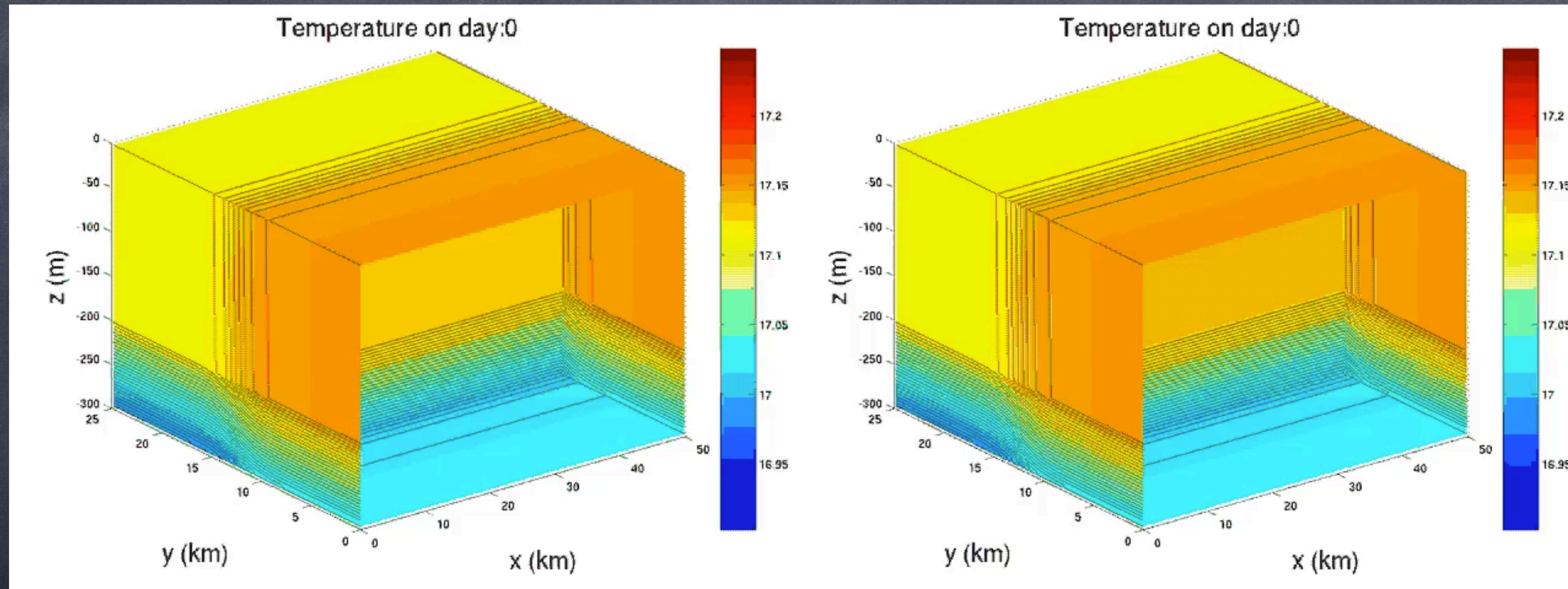
Color High model agreement ( $\geq 80\%$ )  
 Hatched Low model agreement ( $< 80\%$ )



**Mixed layer depth in (a-d) winter and (e-h) summer.** (a, e) Observed climatological mean mixed layer depth (based on density threshold) from the Argo Mixed Layer Depth (Holte et al., 2017) from observations 2000-2019. (b, f) Bias between the observation-based estimate (2000-2019) and the 1995-2014 CMIP6 climatological mean mixed layer depth. (c, d, g, h) Projected MLD change from 1995-2014 to 2081-2100 under (c, g) SSP1-2.6 and (d, h) SSP5-8.5 scenarios.

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# Prototype: Mixed Layer Front Adjustment



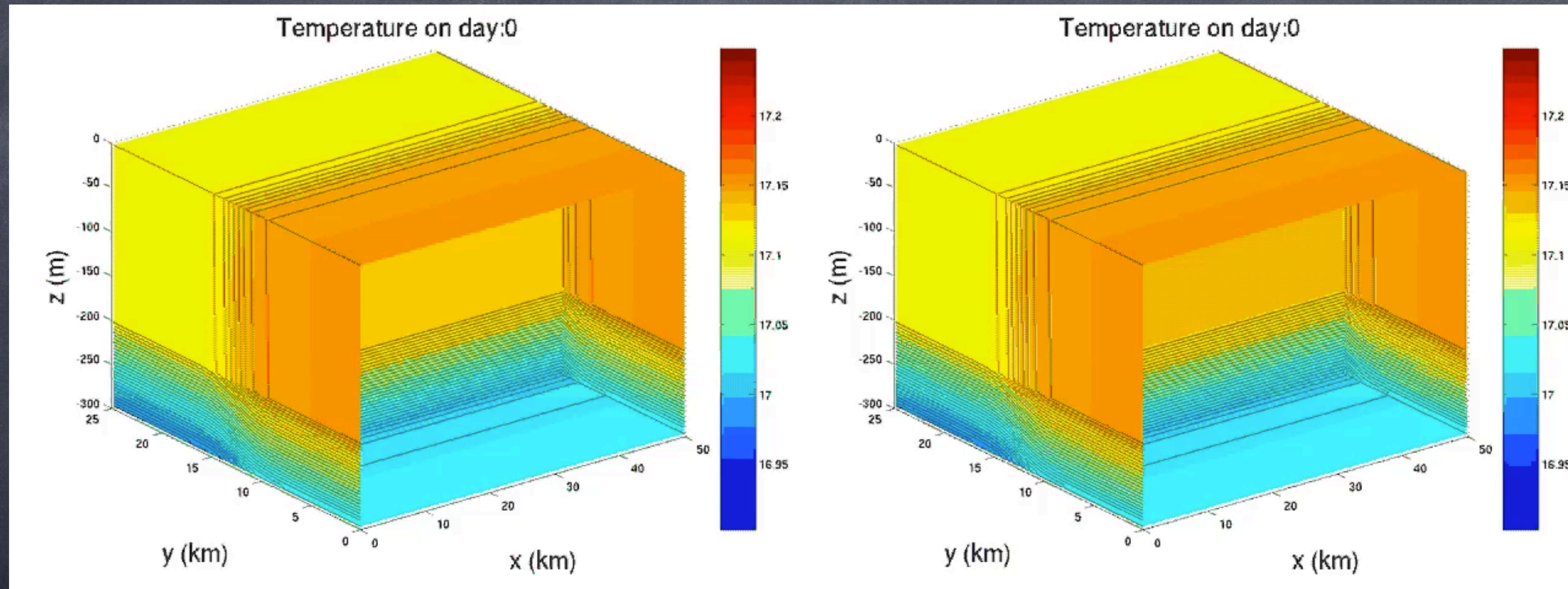
Simple Spindown

Plus, Diurnal Cycle  
and Parameterized Mixing

Note: initial geostrophic adjustment overwhelmed by eddy restratification

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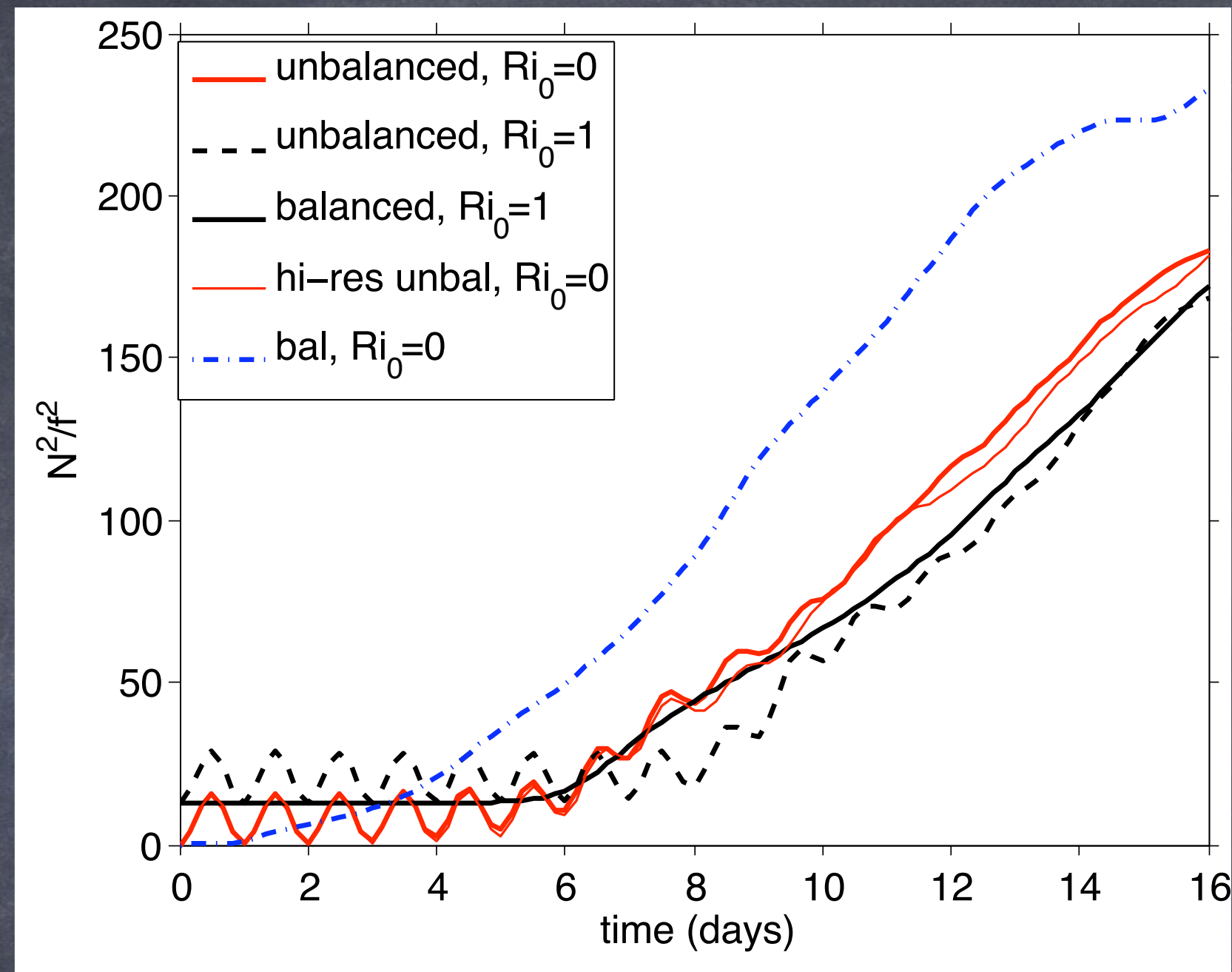


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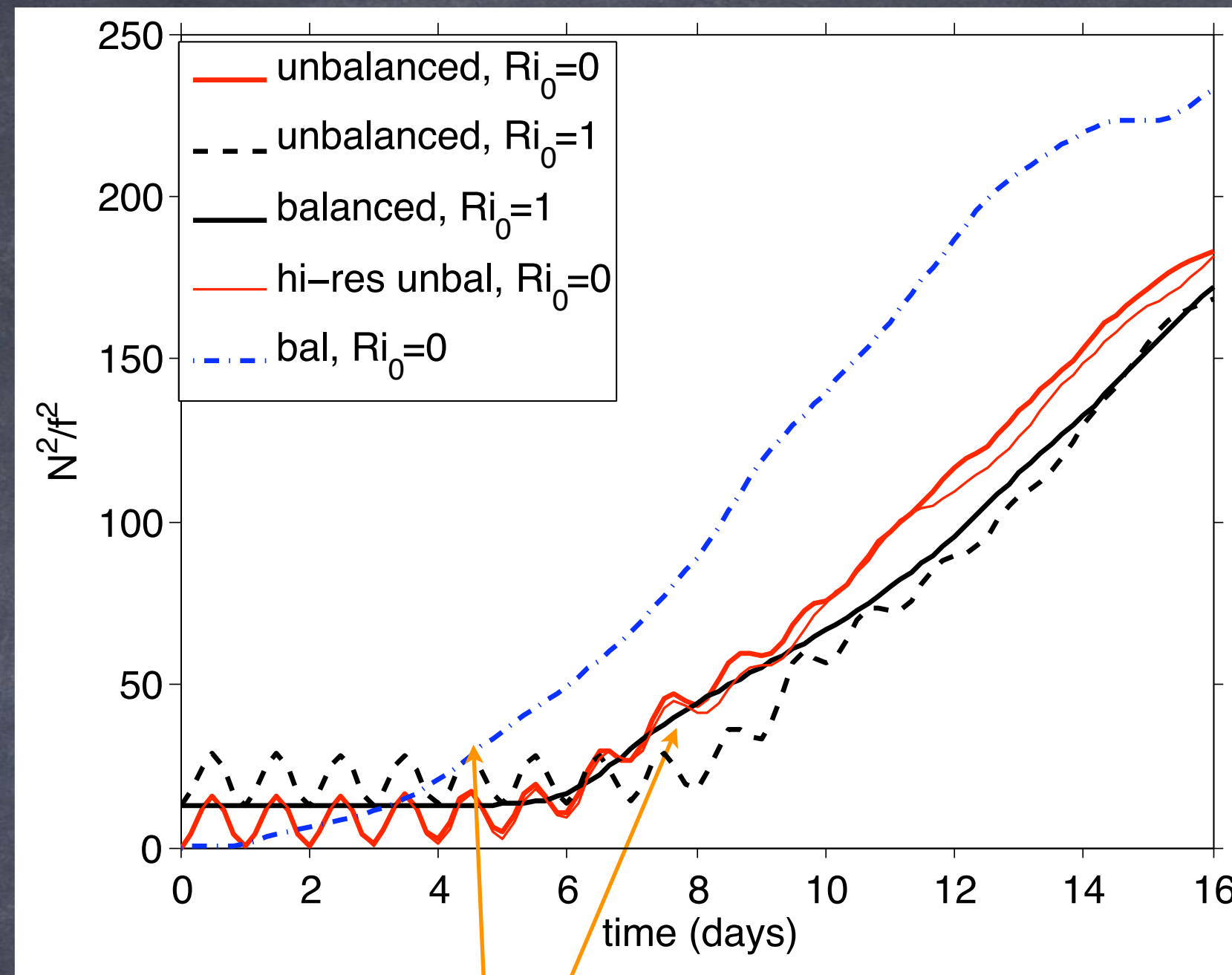
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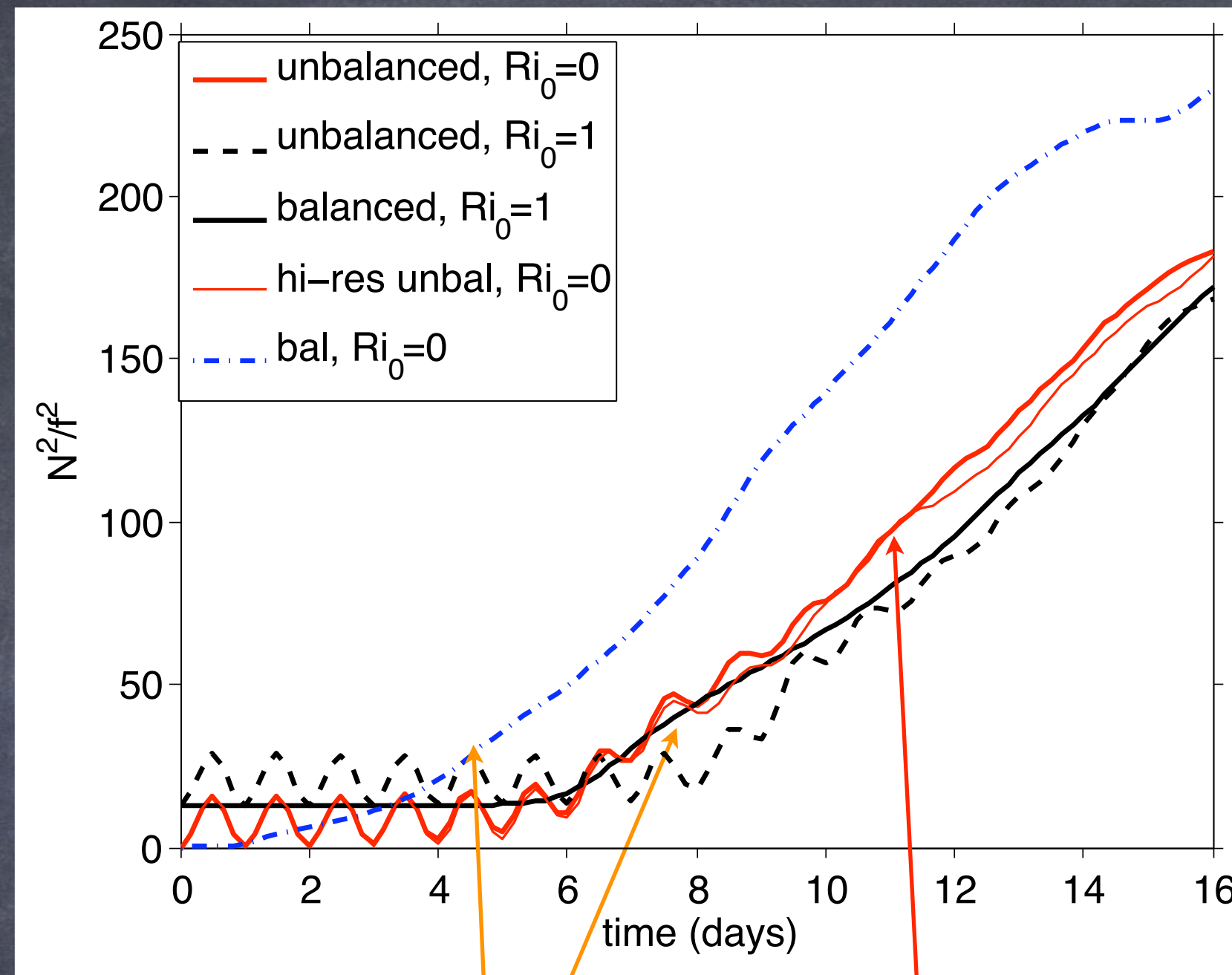
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Eddies at Finite  
Amplitude

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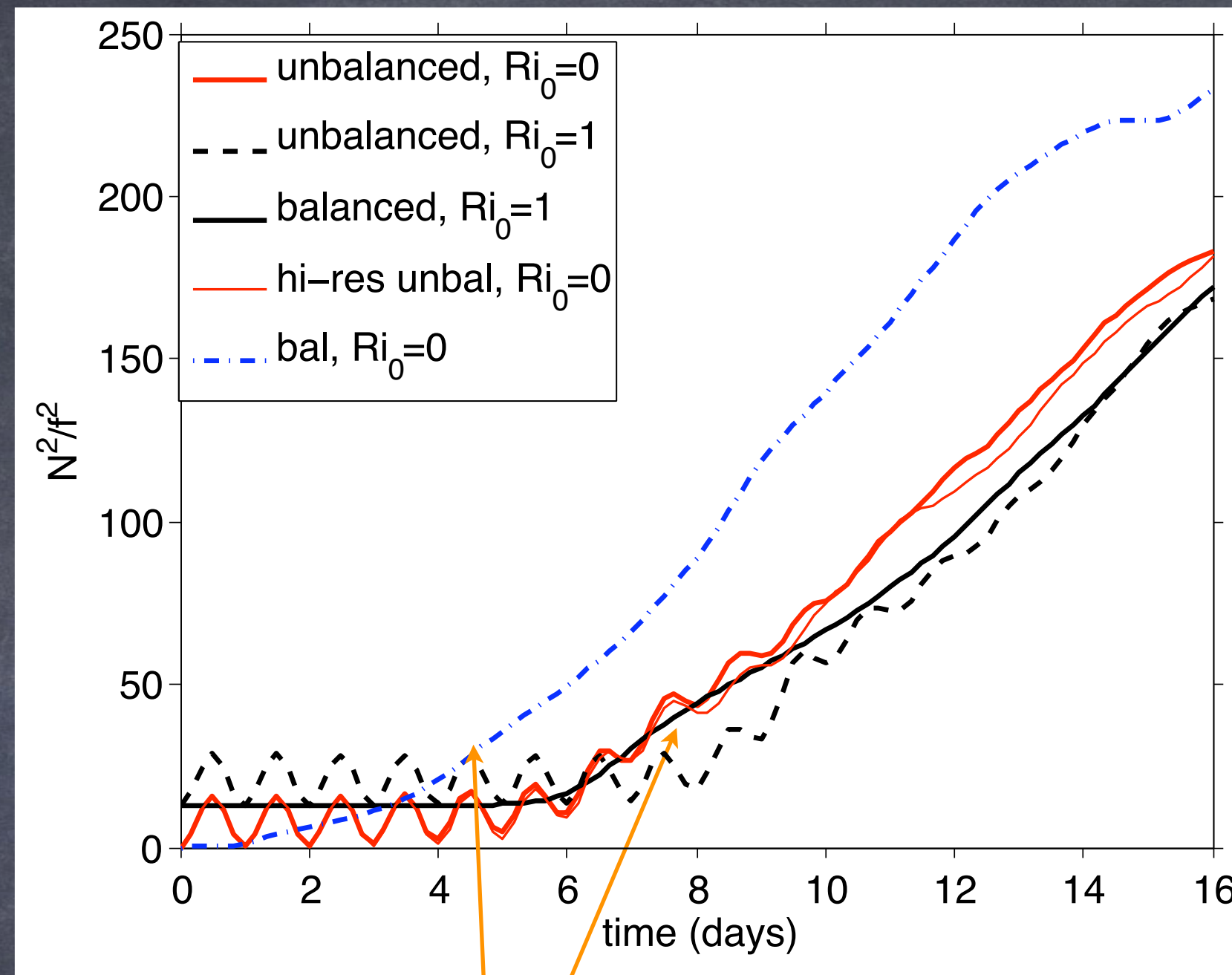


Eddies at Finite  
Amplitude

Resolution  
Convergence

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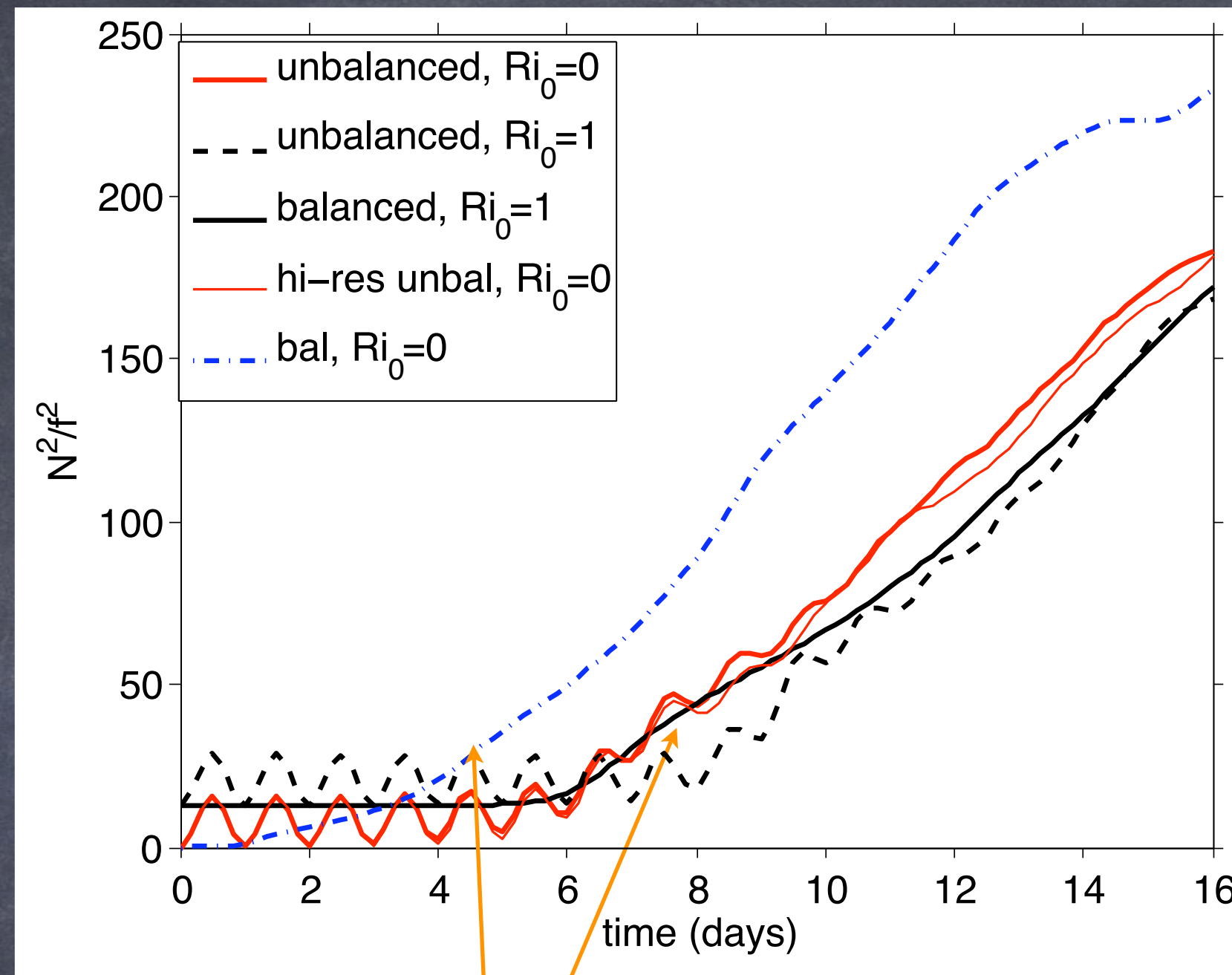
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## Power Spectrum of KE

Eddies at Finite  
Amplitude

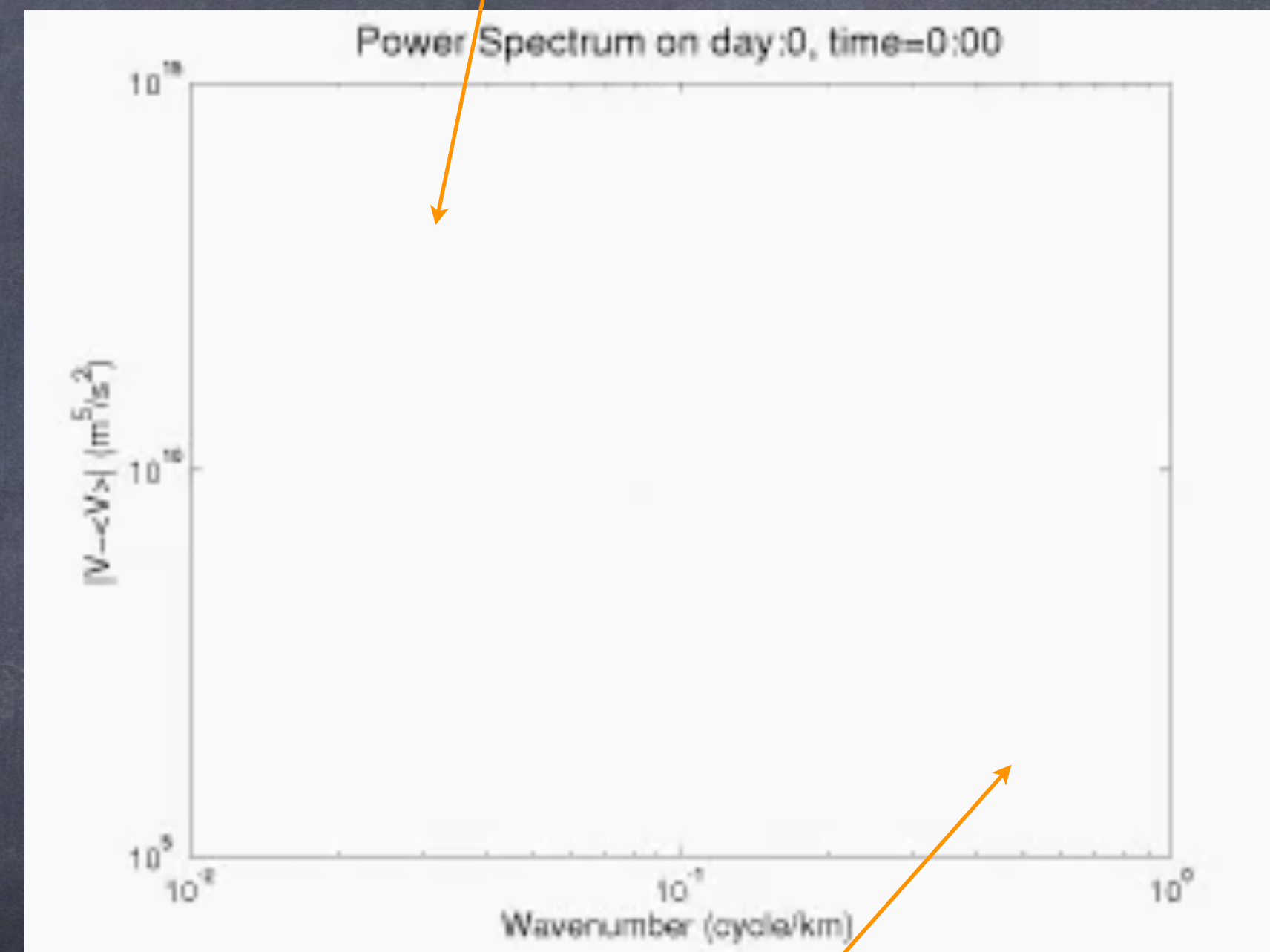
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Eddies at Finite Amplitude

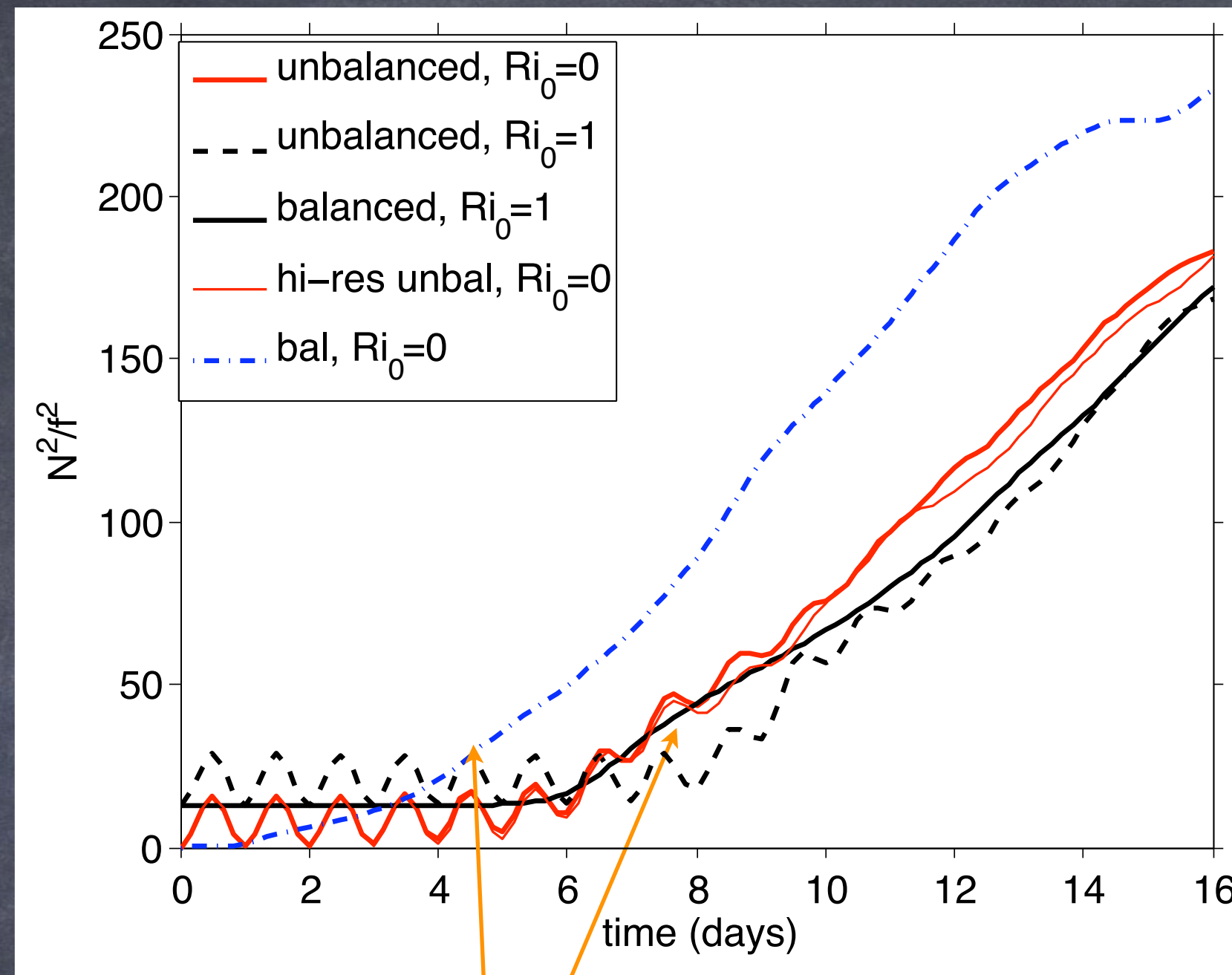
## Power Spectrum of KE

At Finite Amplitude  
Horizontal Scale Unclear



Initially, Linear Prediction of  
Lengthscale good

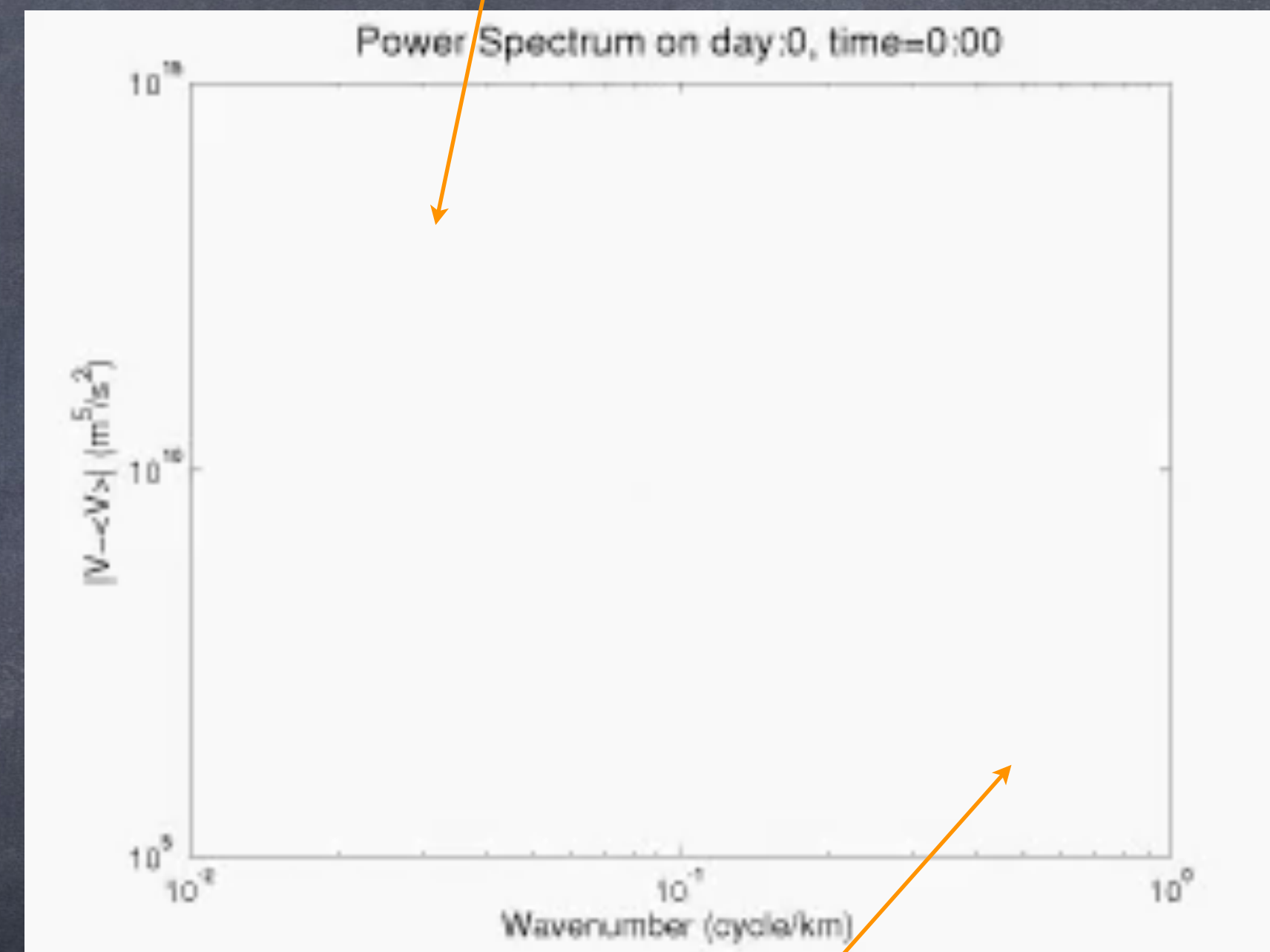
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Eddies at Finite Amplitude

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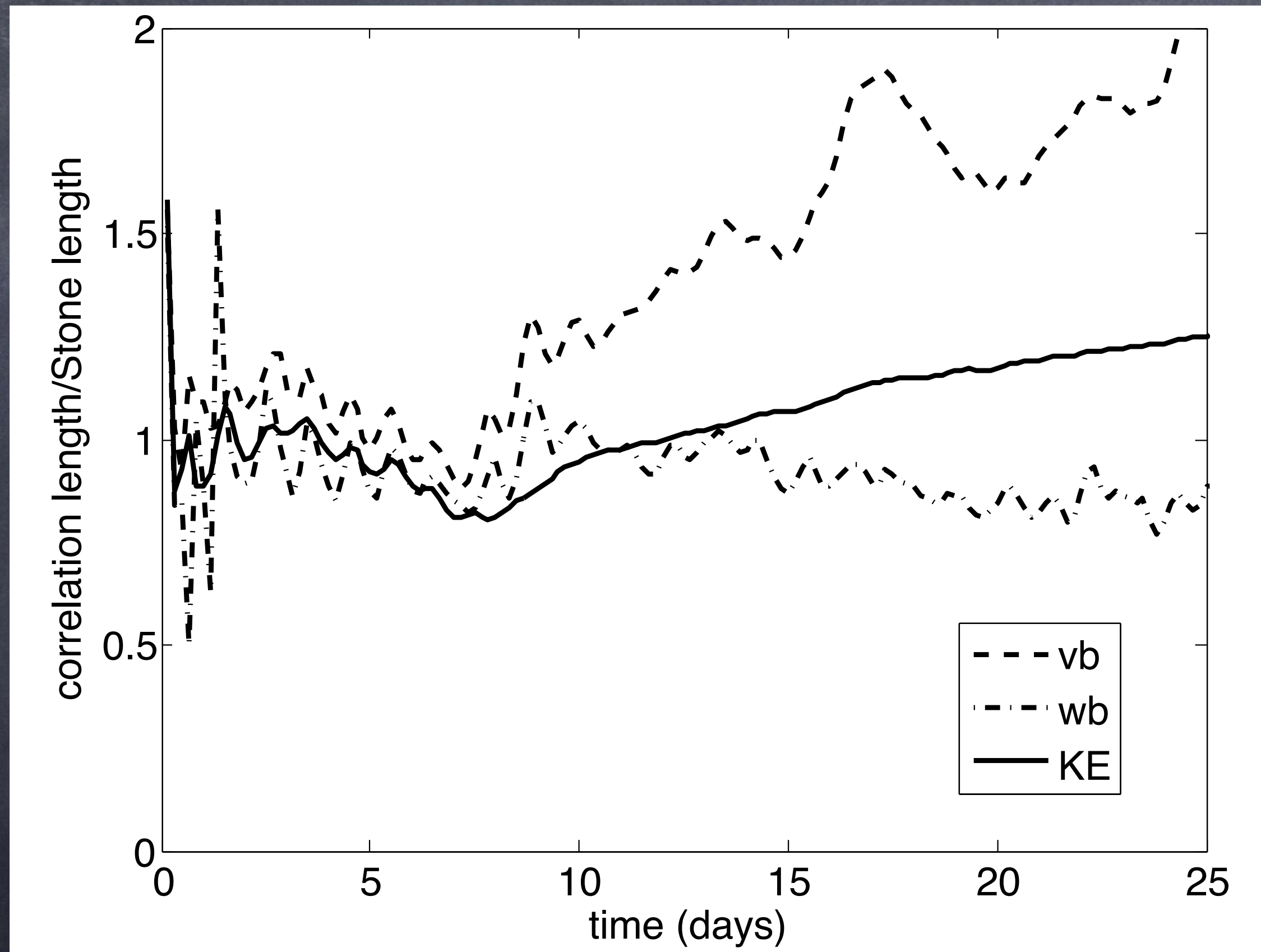
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Inverse Cascade => No Results from Linear Instability Ingredients

# What lengthscale dominates $\langle w'b' \rangle$ ?

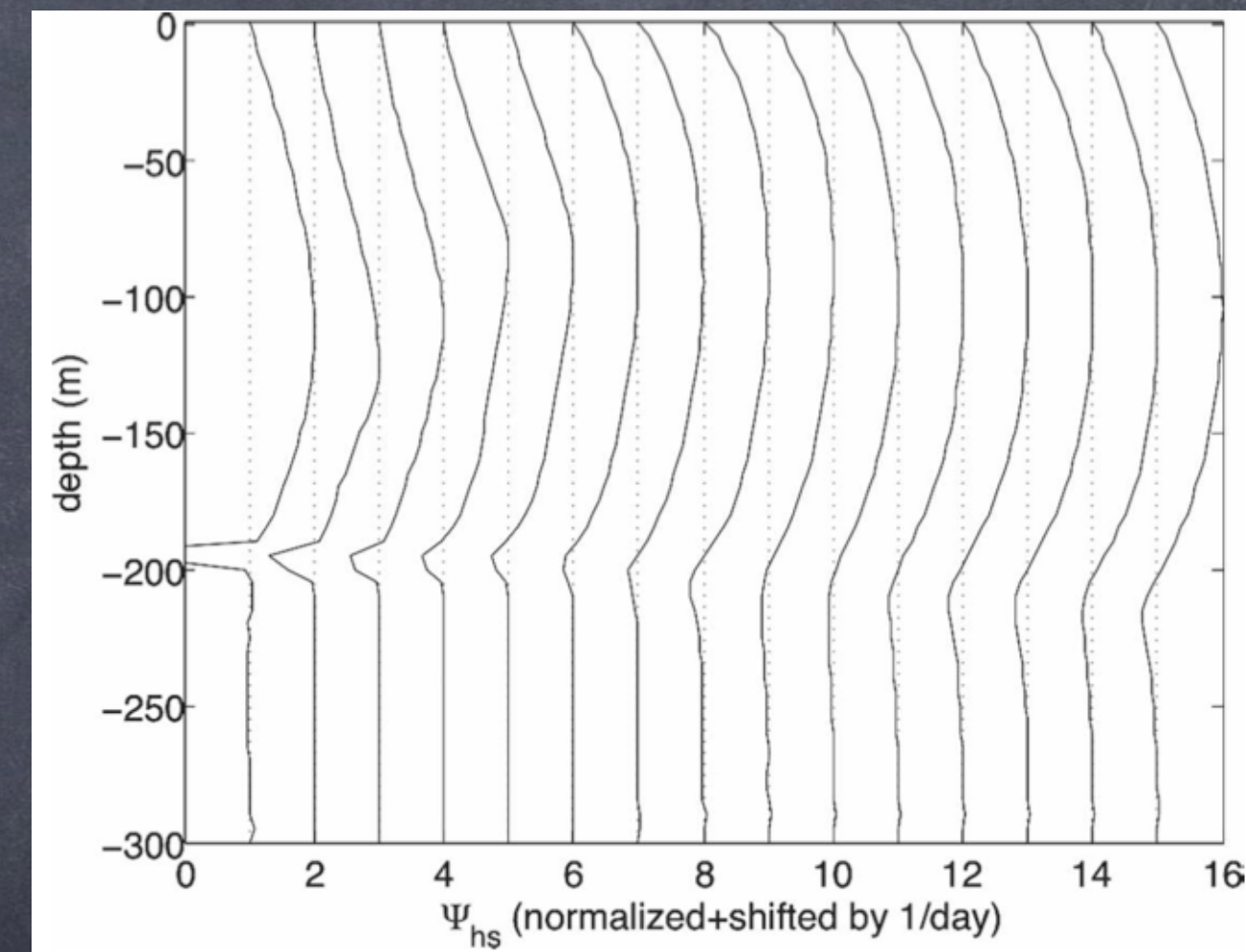
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vb

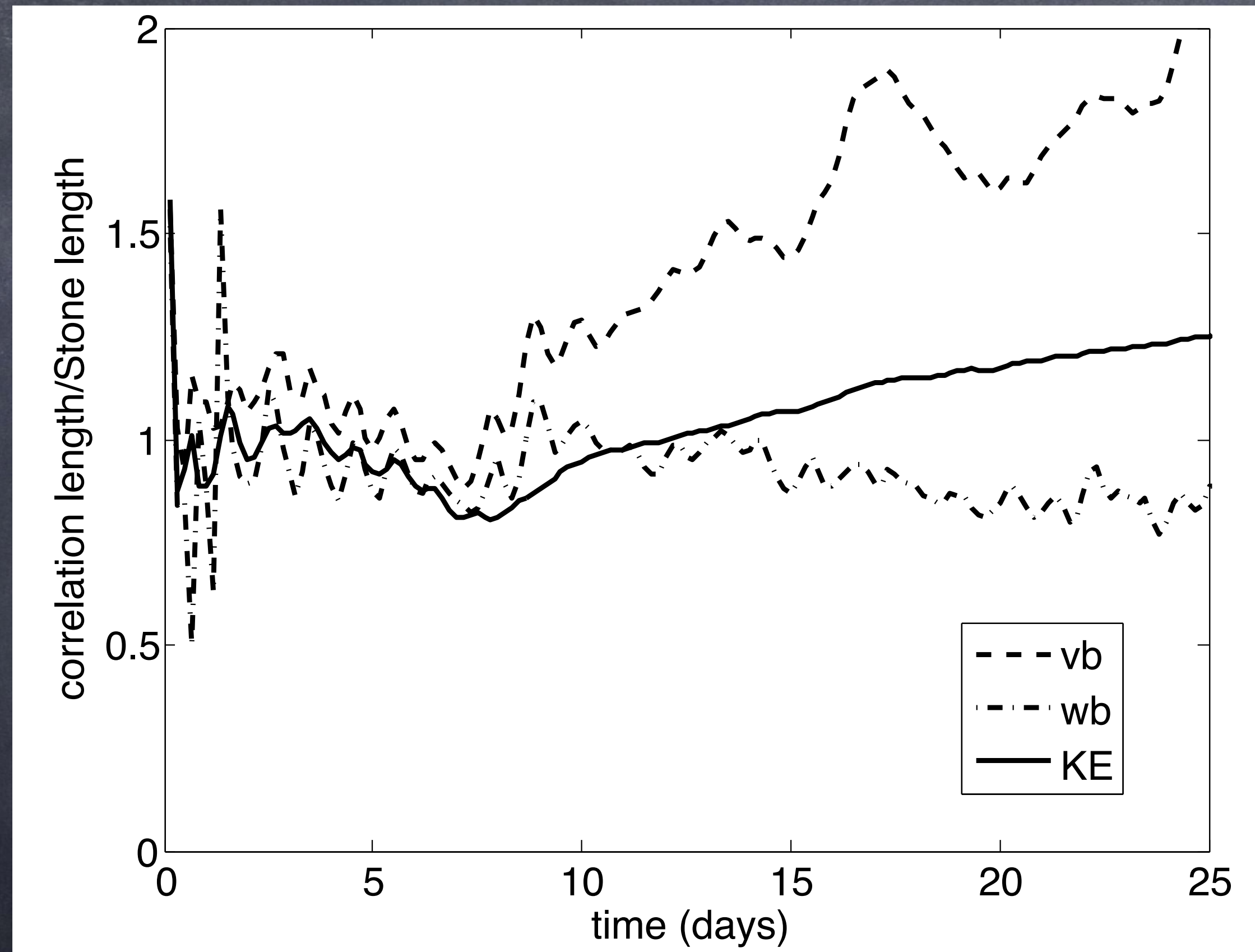
KE

wb



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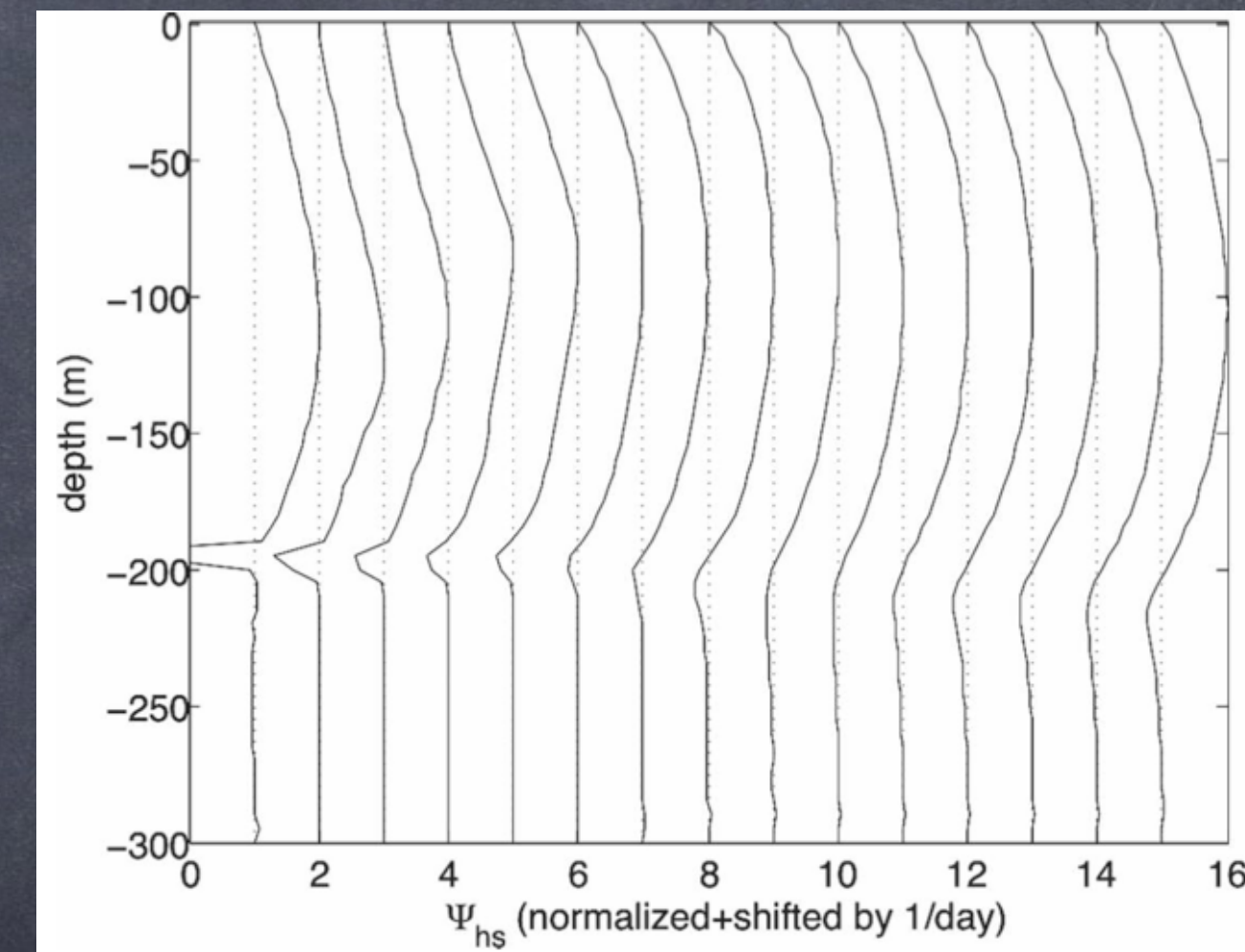
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vb

KE

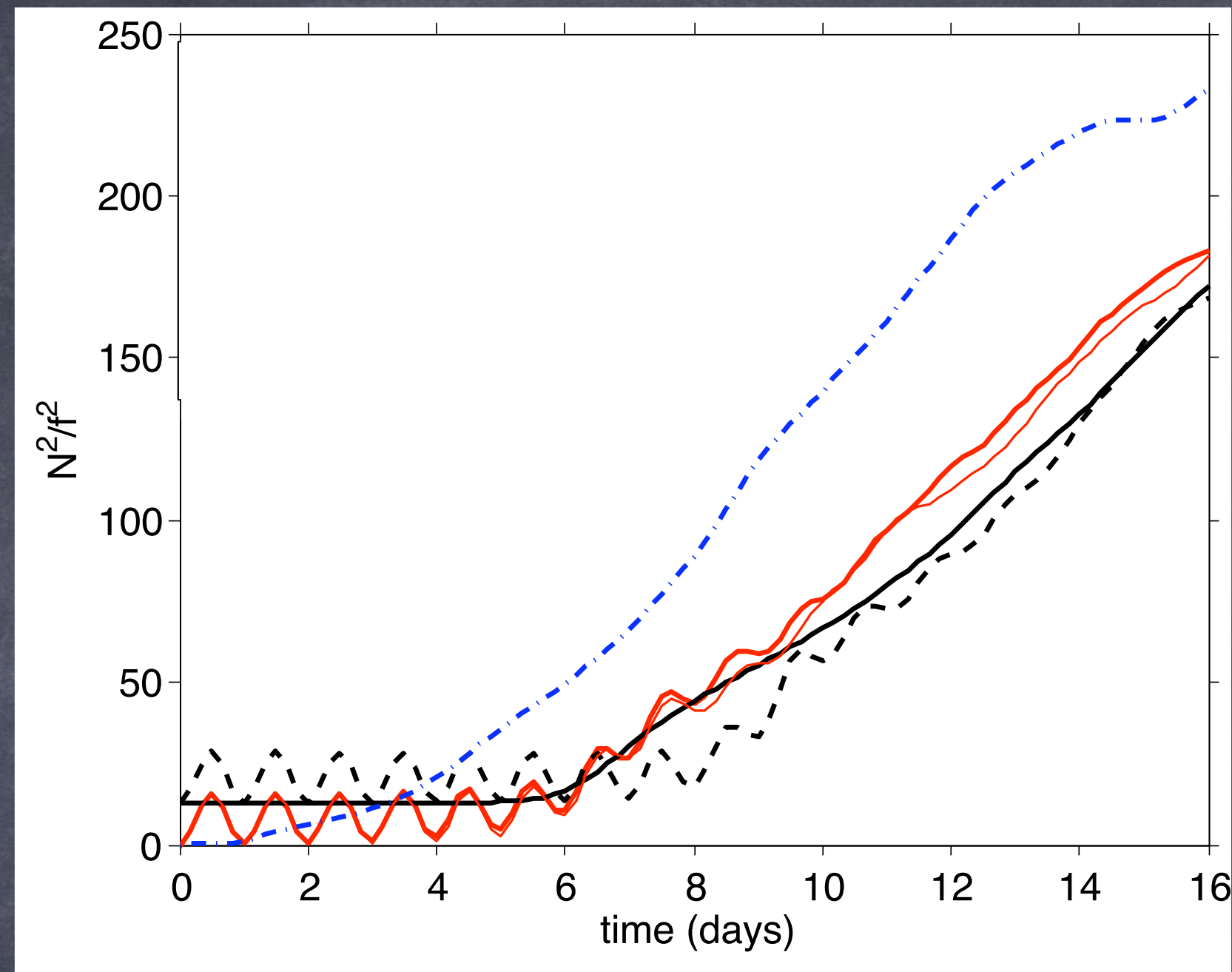
wb



Stone fastest-mode Soln OK!

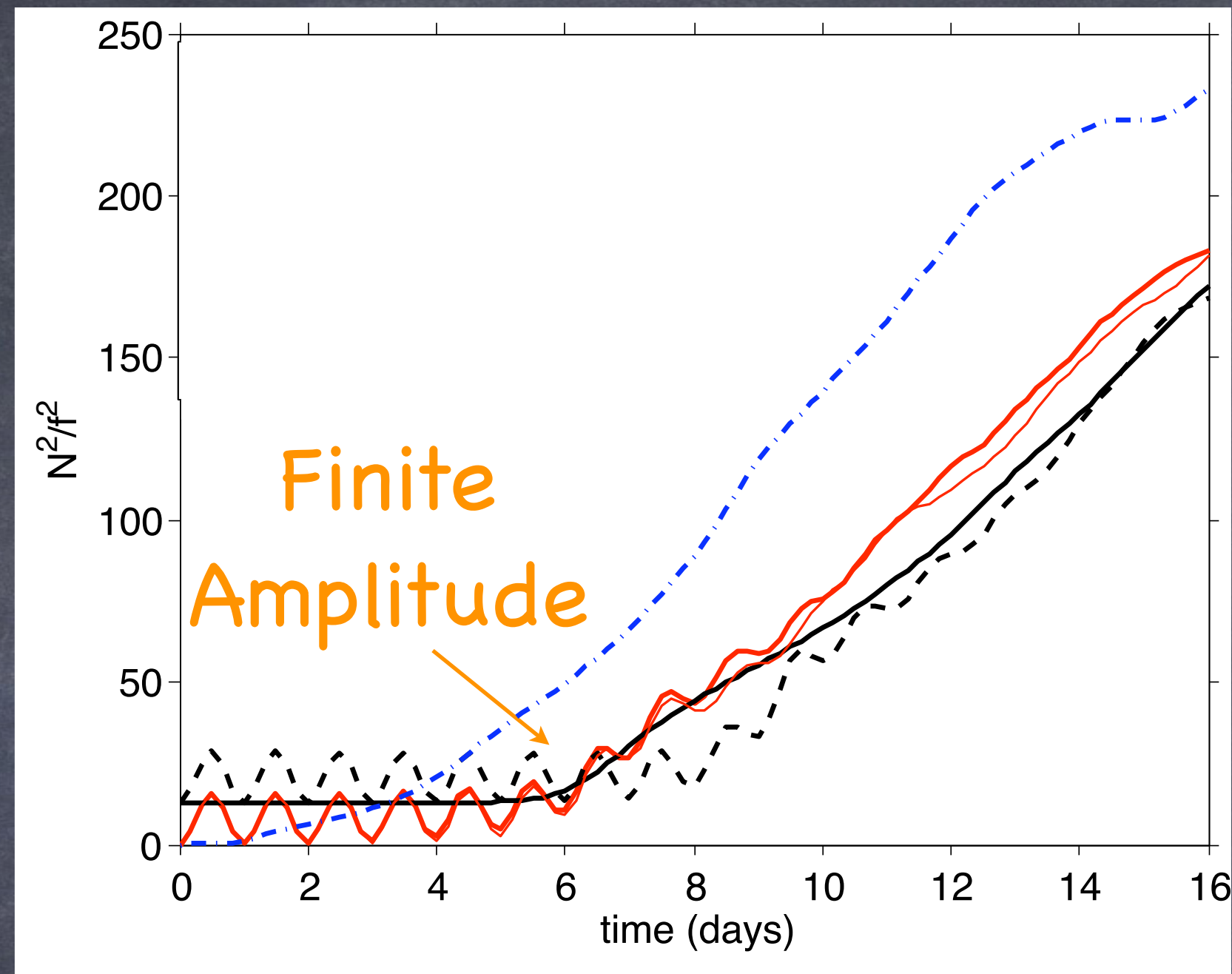
$$\mu(z) = \left[ 1 - \left( \frac{2z}{H} + 1 \right)^2 \right] \left[ 1 + \frac{5}{21} \left( \frac{2z}{H} + 1 \right)^2 \right]$$

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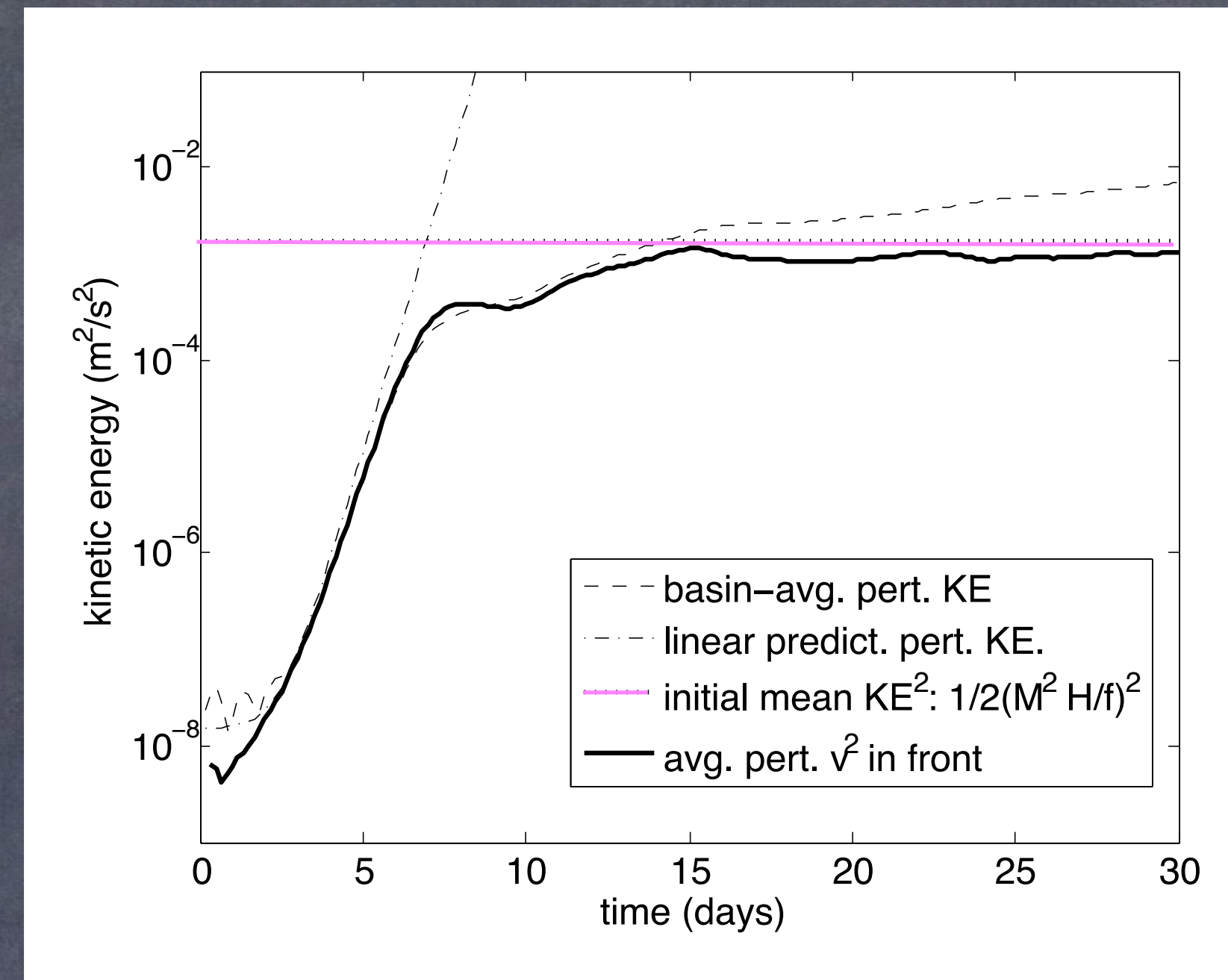
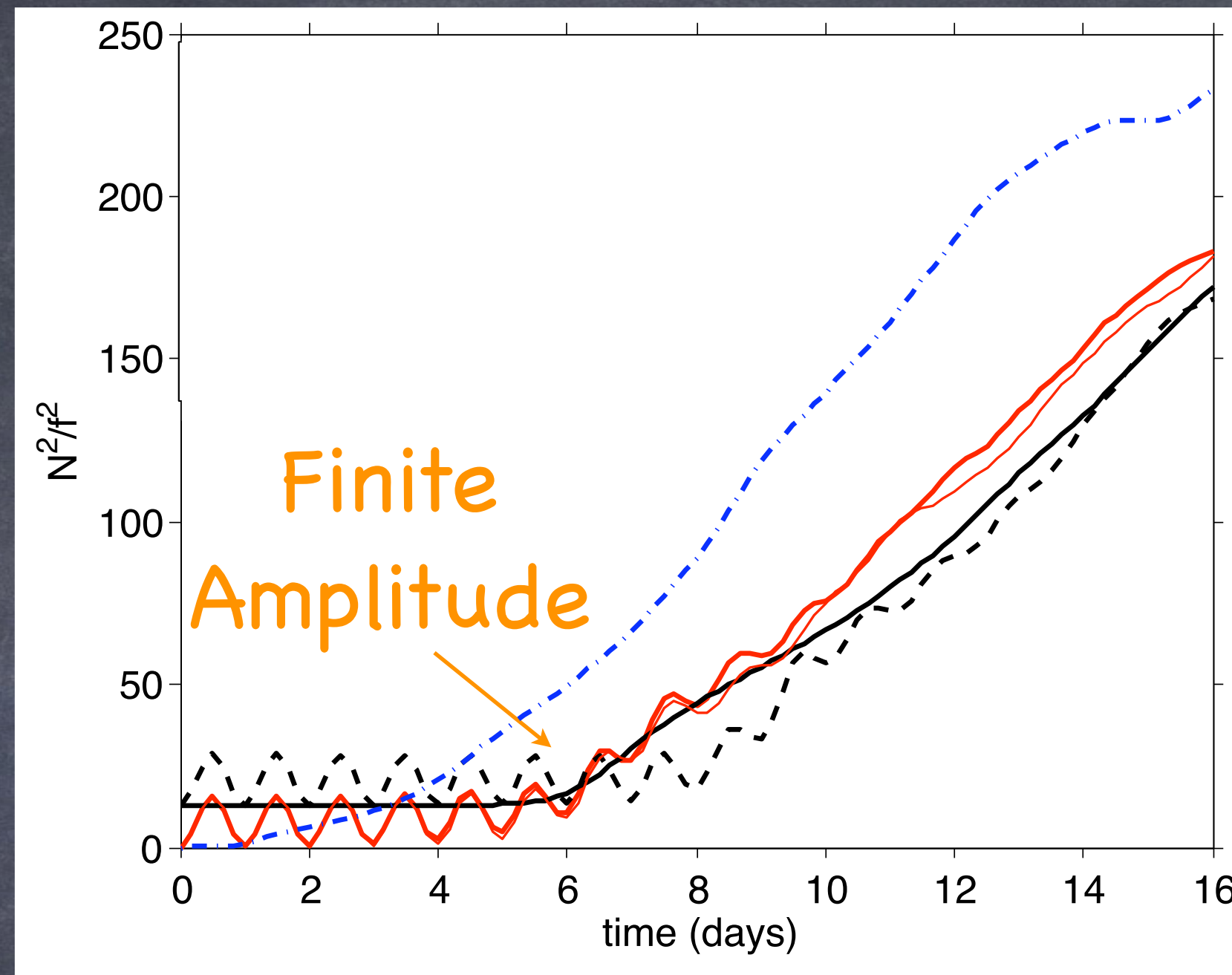




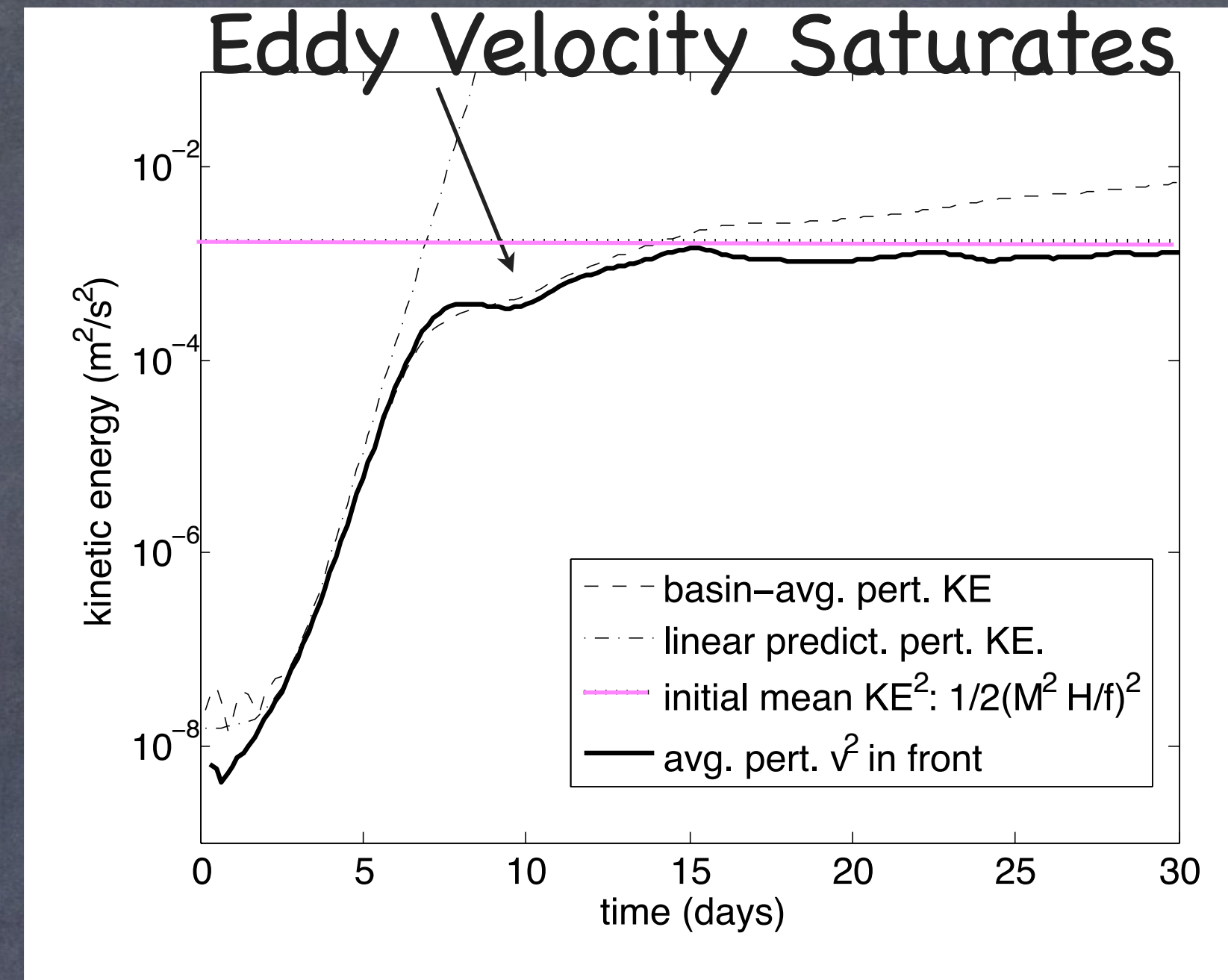
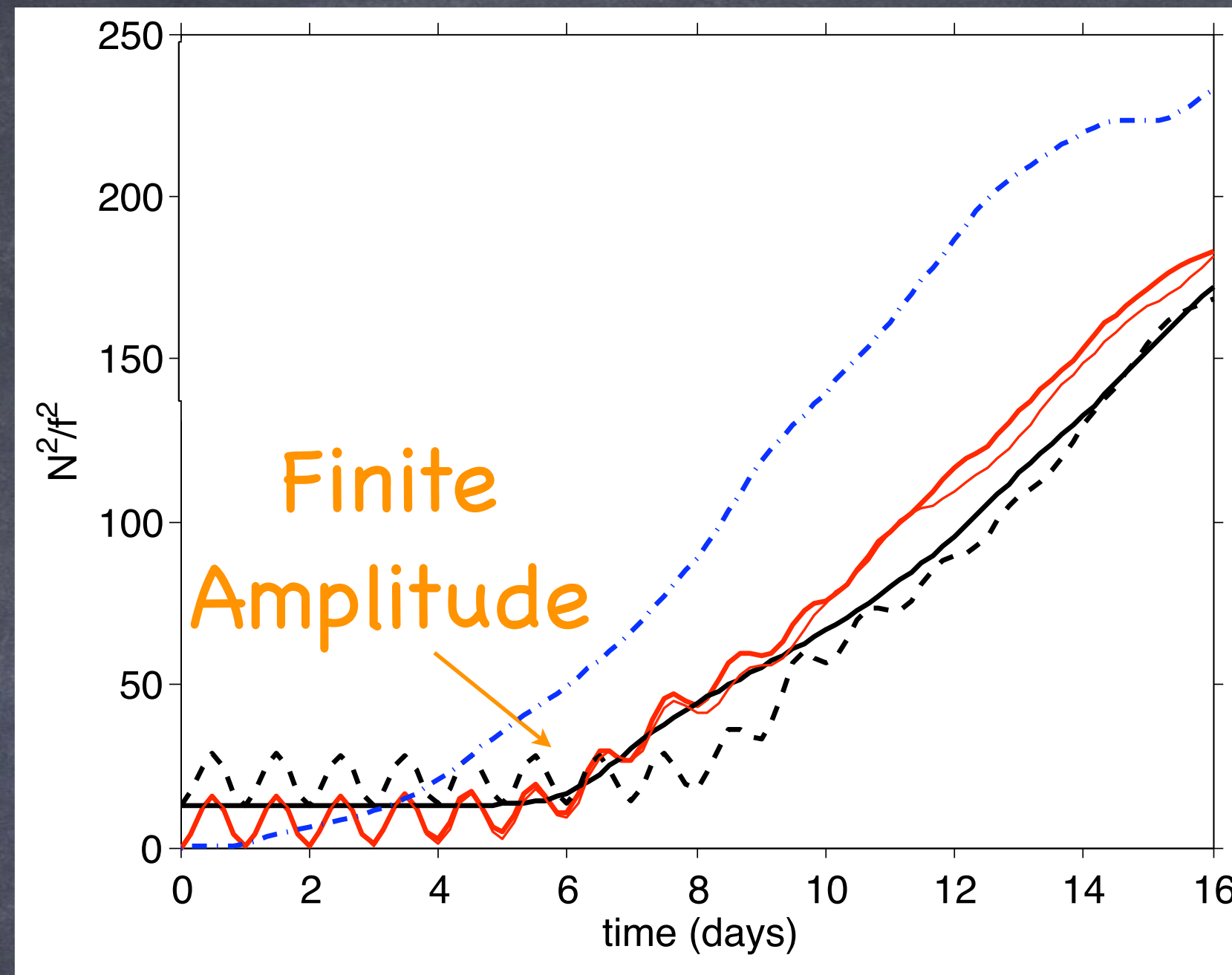
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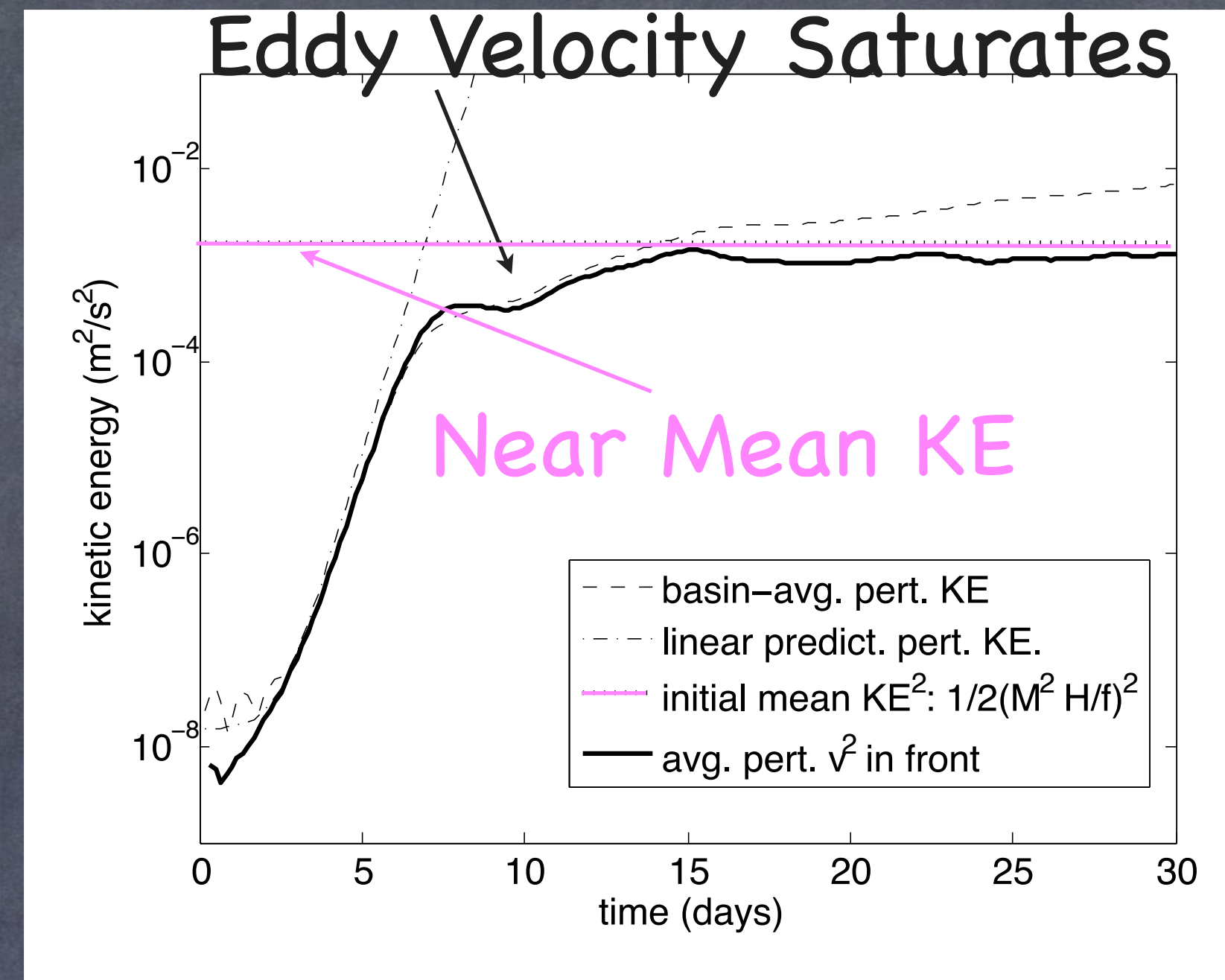
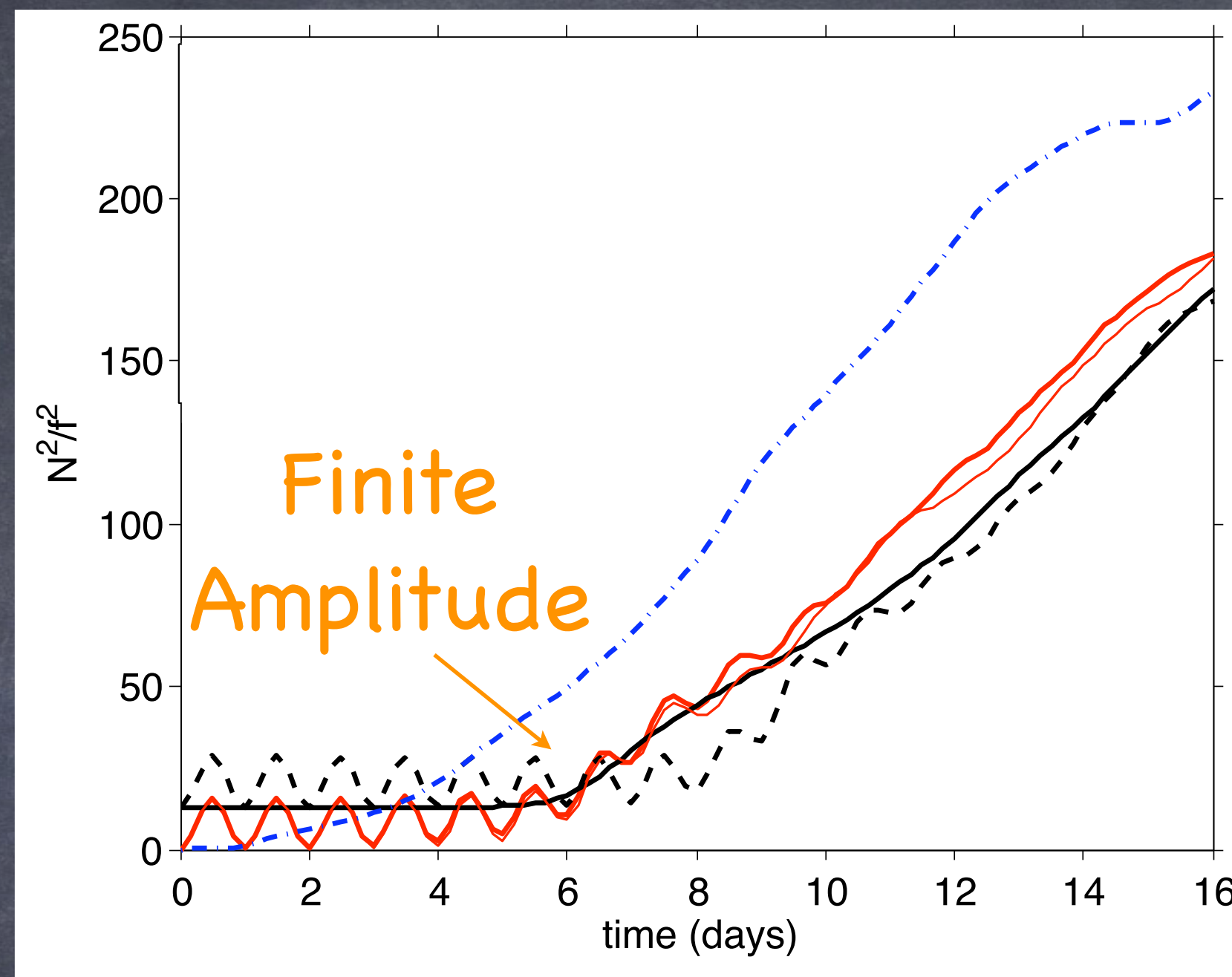
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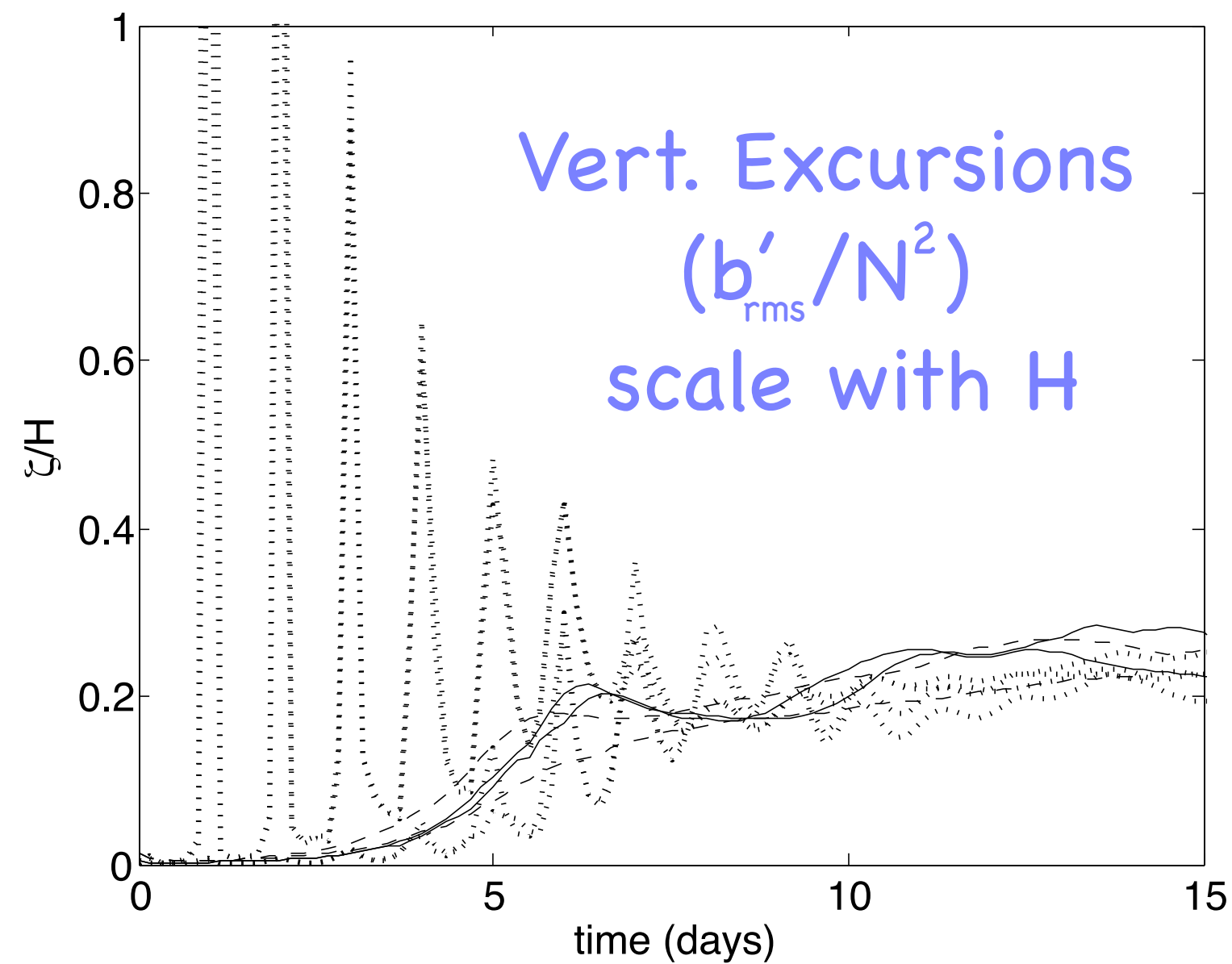
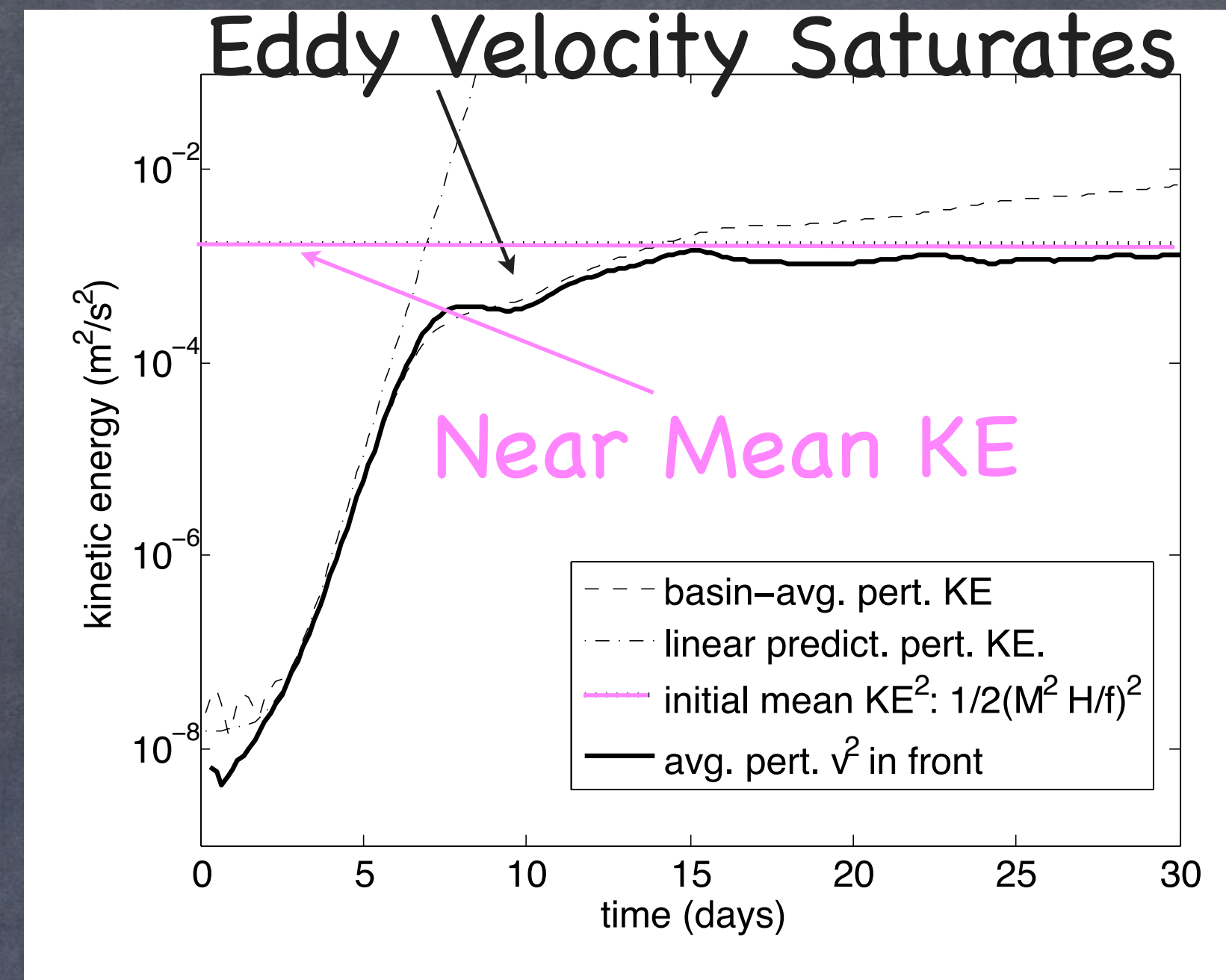
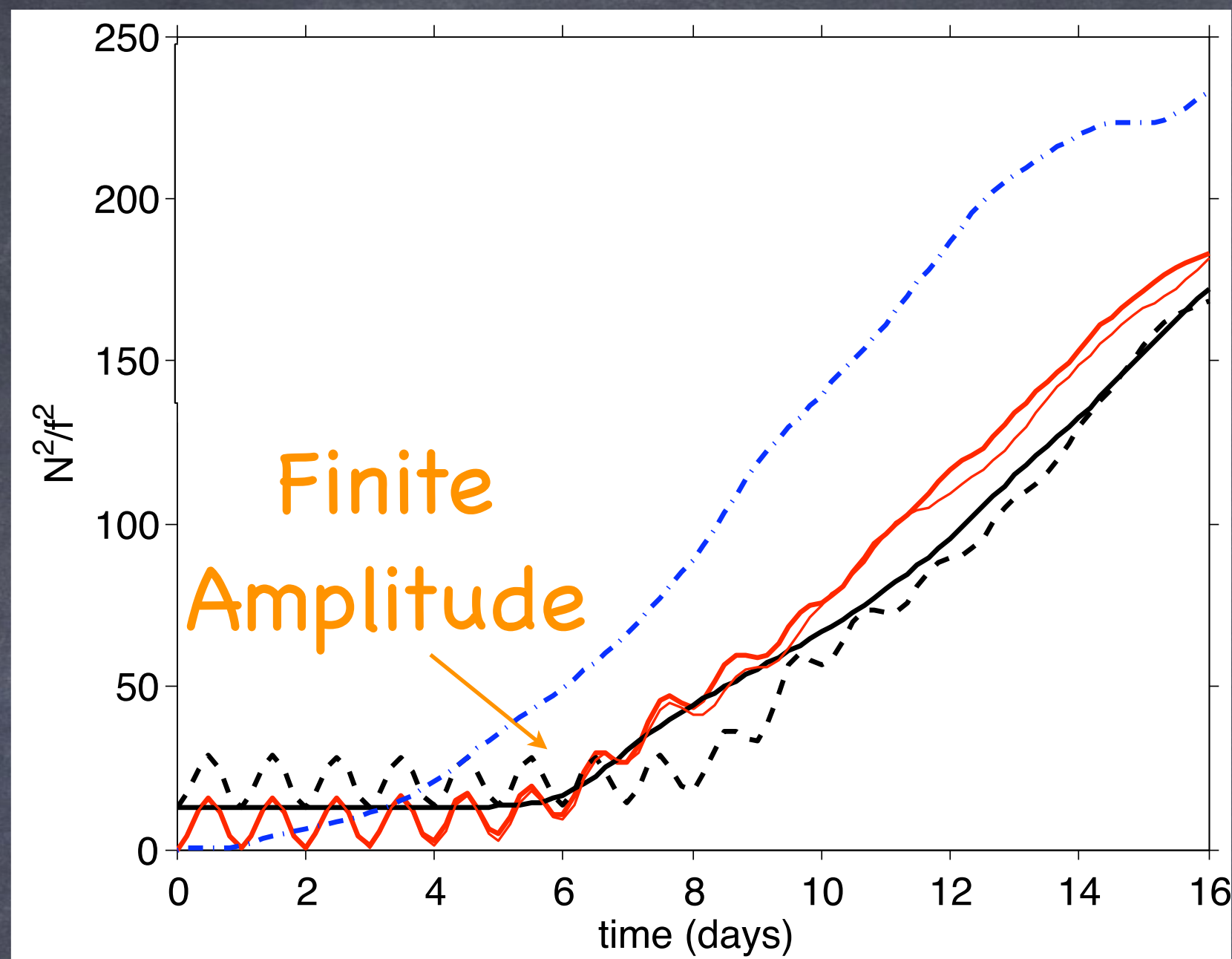
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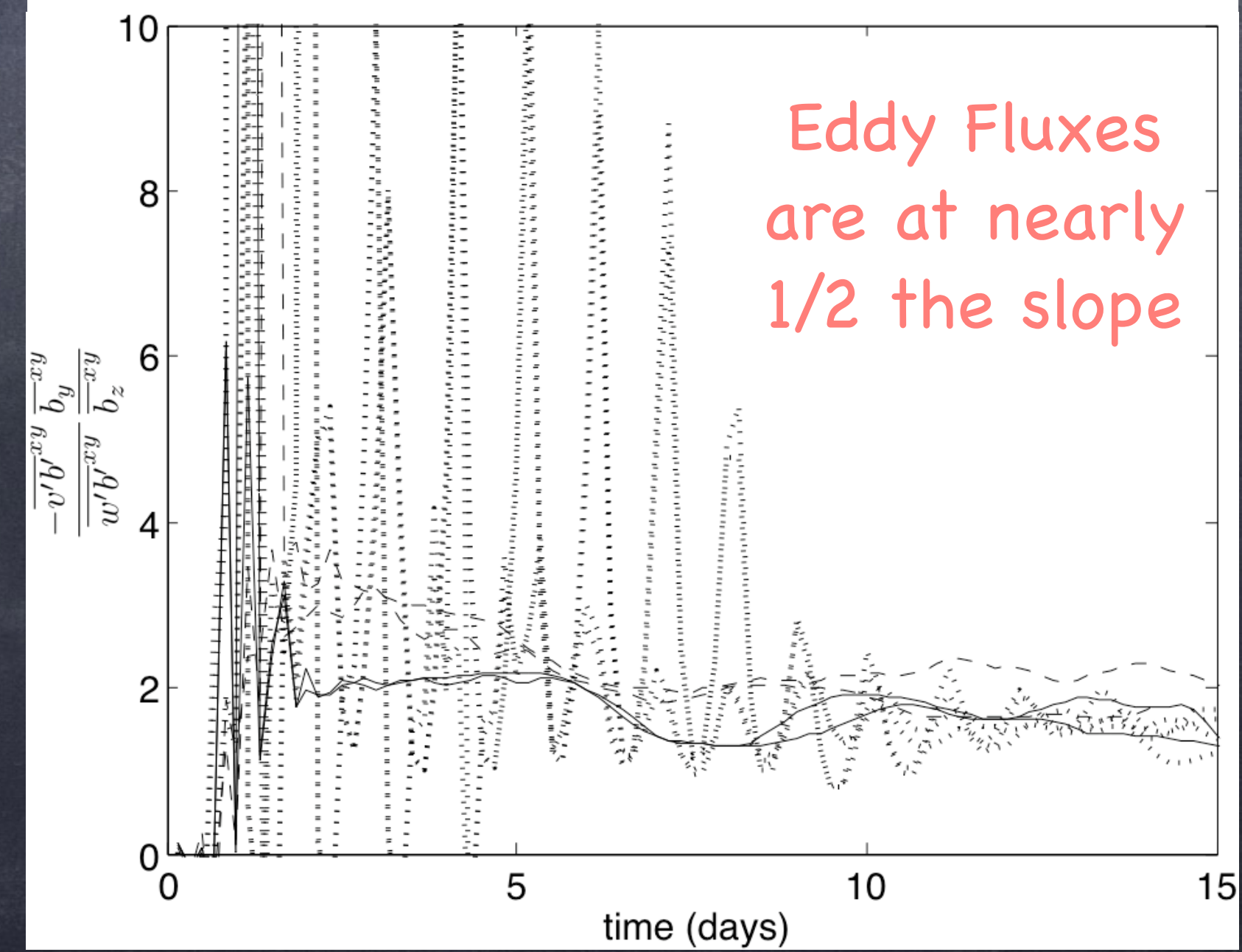
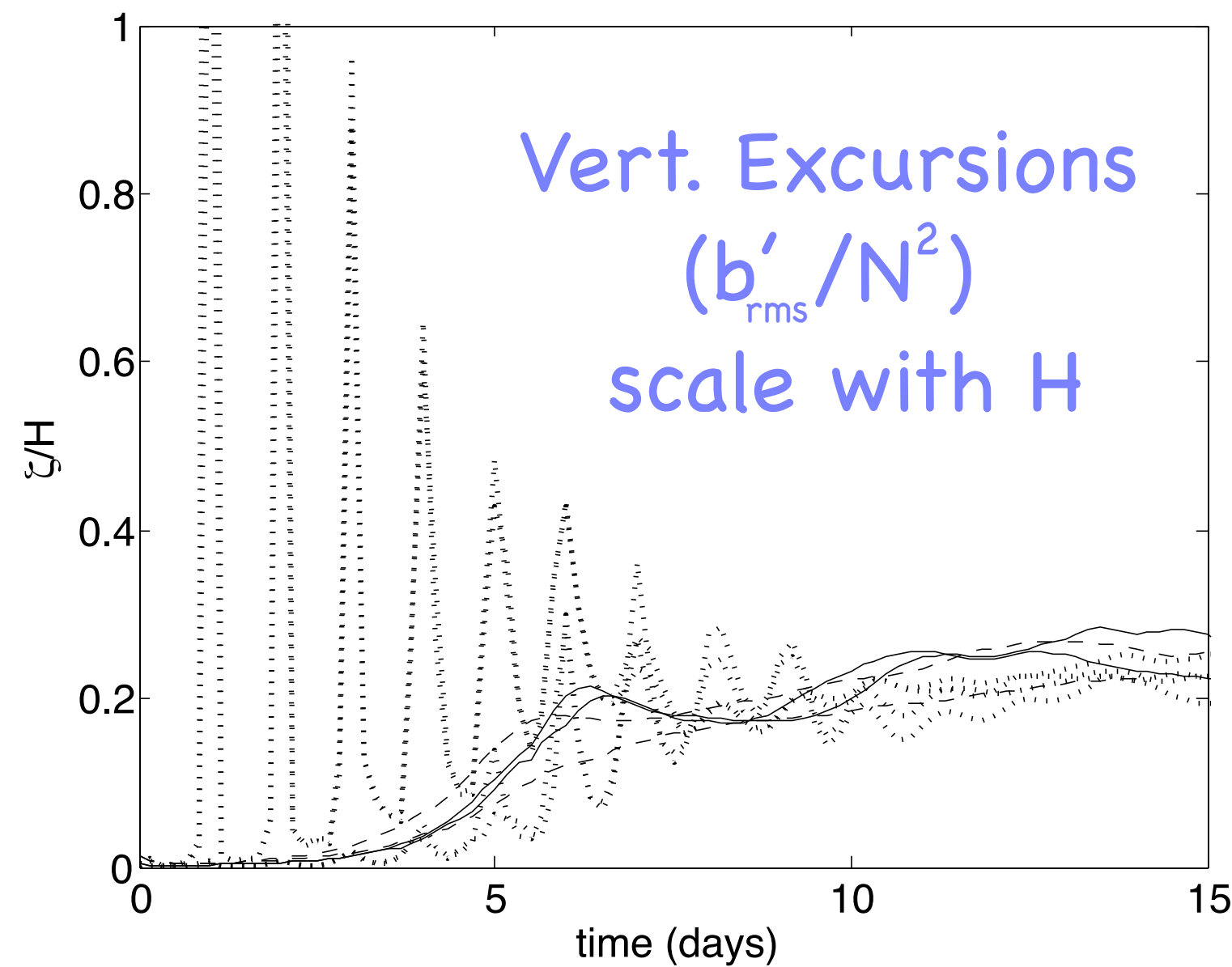
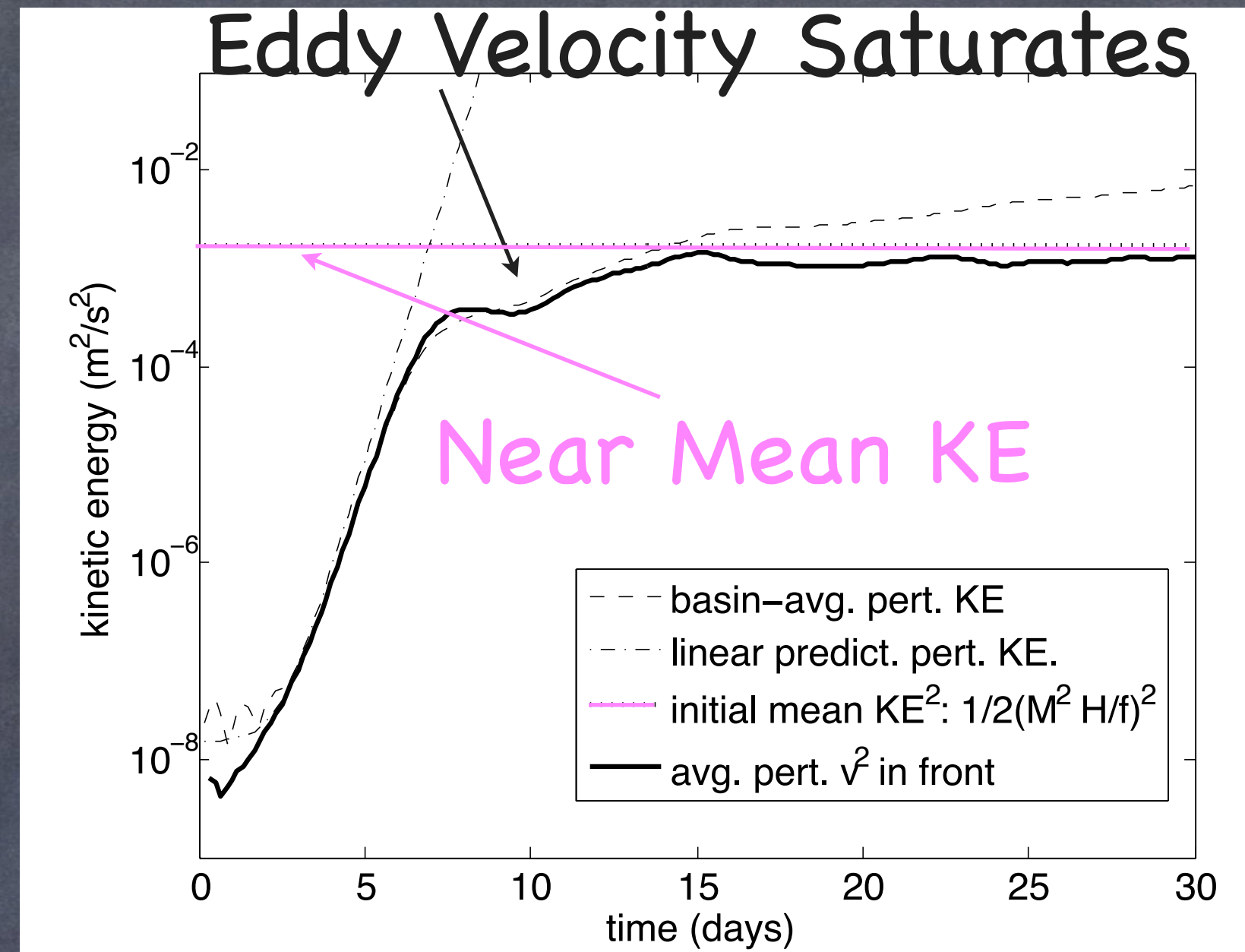
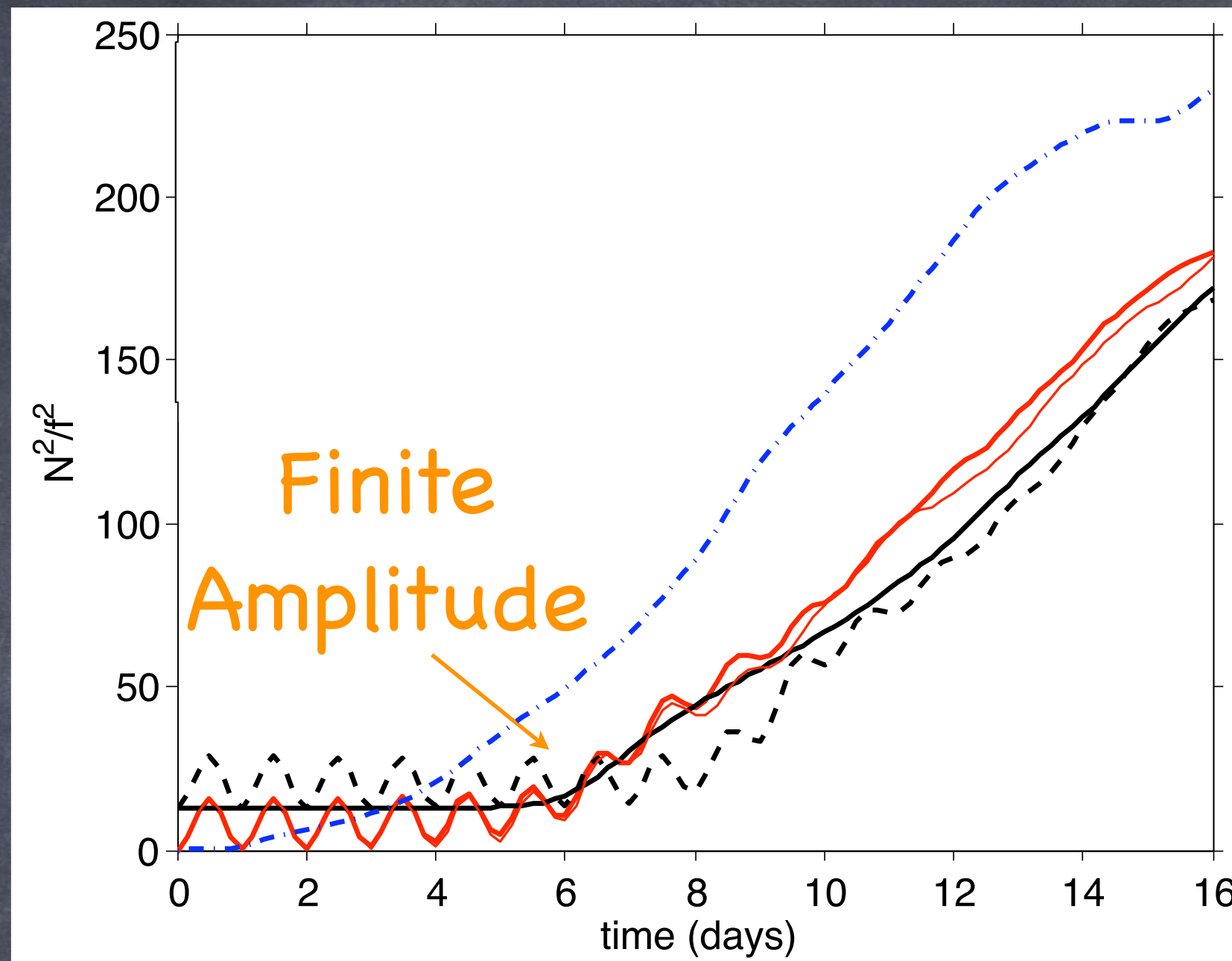
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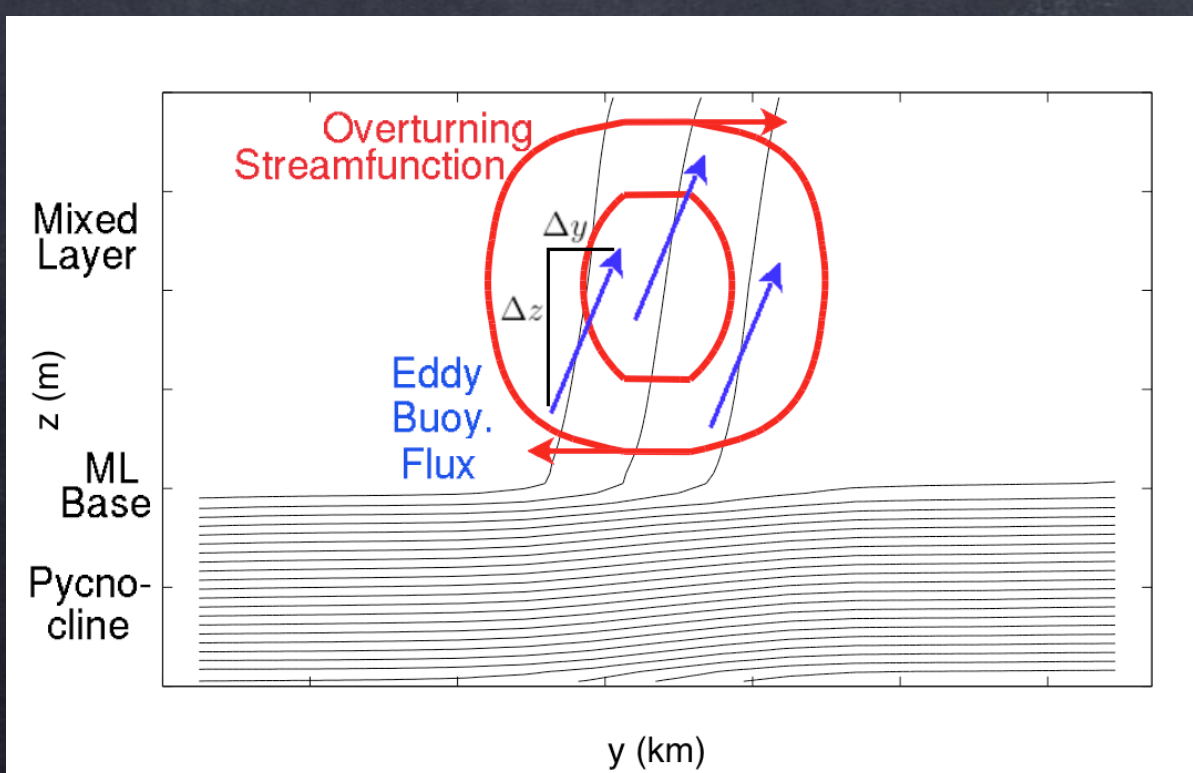


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# Magnitude Analysis: Vert. Fluxes

Extraction of potential energy by submesoscale eddies:

$$-\langle wb \rangle = \frac{\partial \langle PE \rangle}{\partial t} \approx \frac{\Delta PE}{\Delta t} \propto \frac{\Delta z \Delta b}{\Delta t}$$



$$\langle wb \rangle \propto \frac{-\Delta z \Delta b}{\Delta t}$$

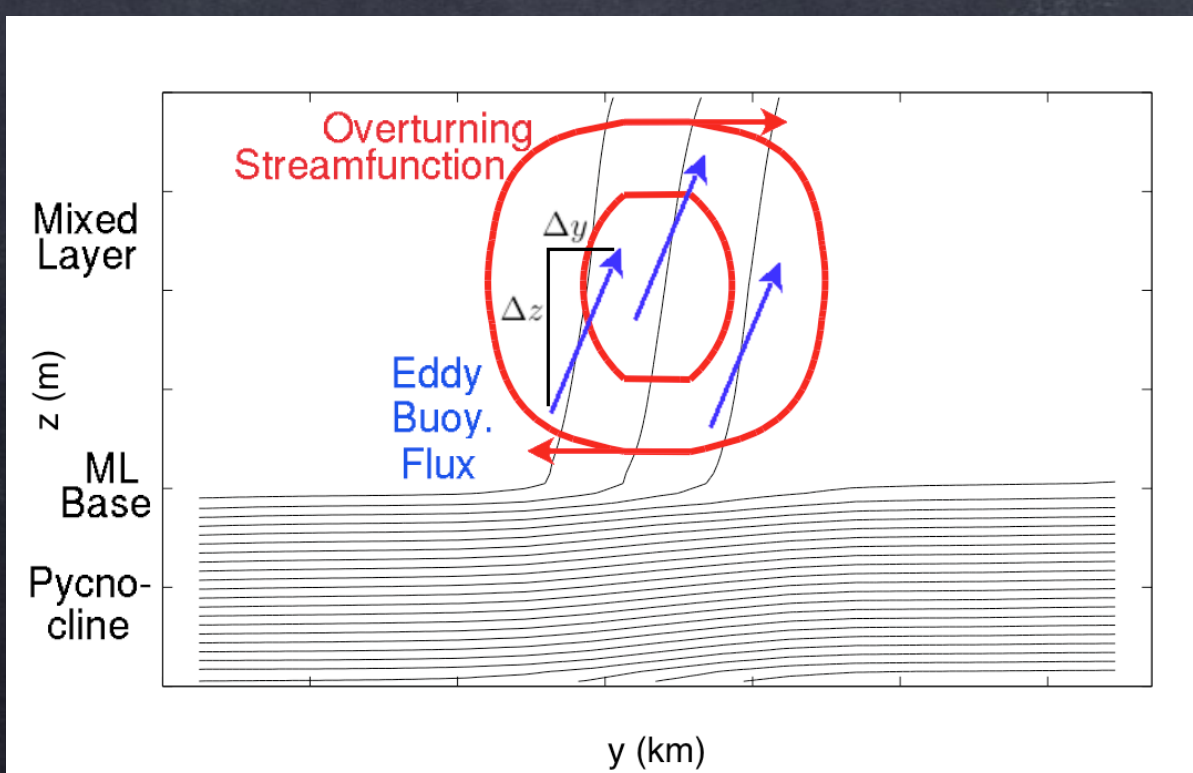
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Buoy. diff just parcel exchange of large-scale buoy.



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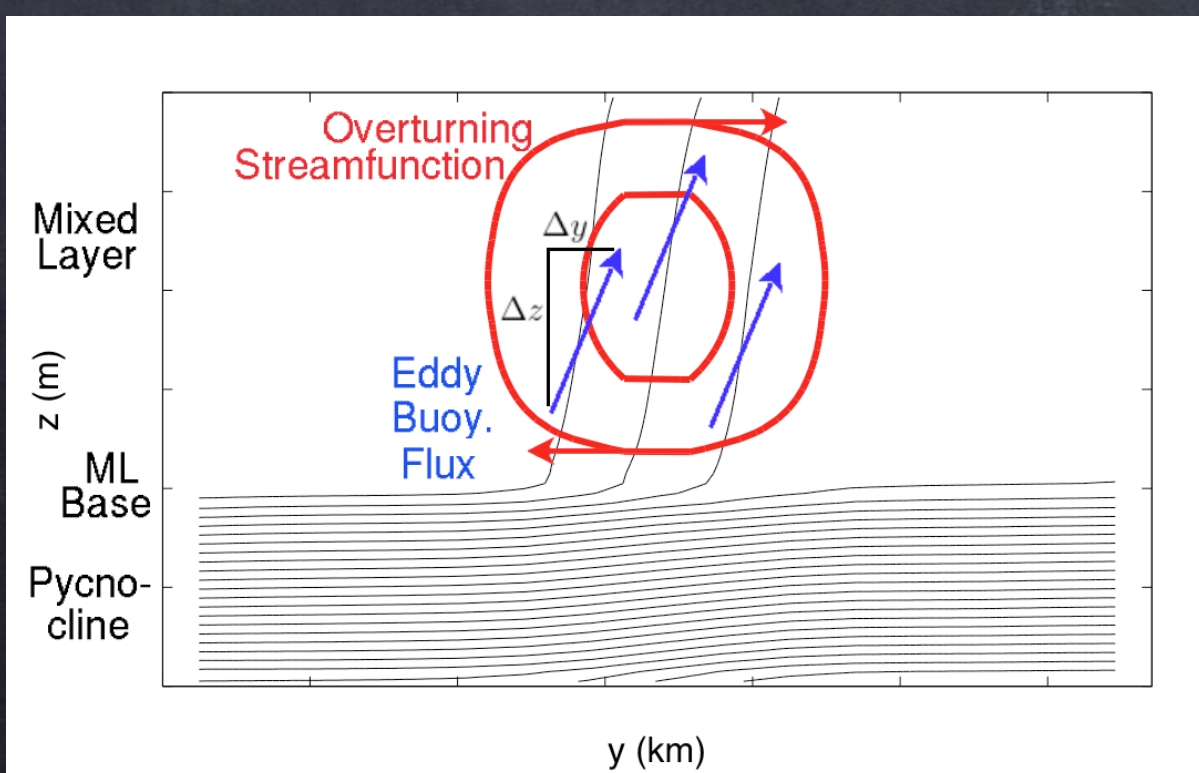
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$$\langle wb \rangle \propto \frac{-\Delta z \left( \Delta y \frac{\partial \bar{b}}{\partial y} + \Delta z \frac{\partial \bar{b}}{\partial z} \right)}{\Delta t}$$

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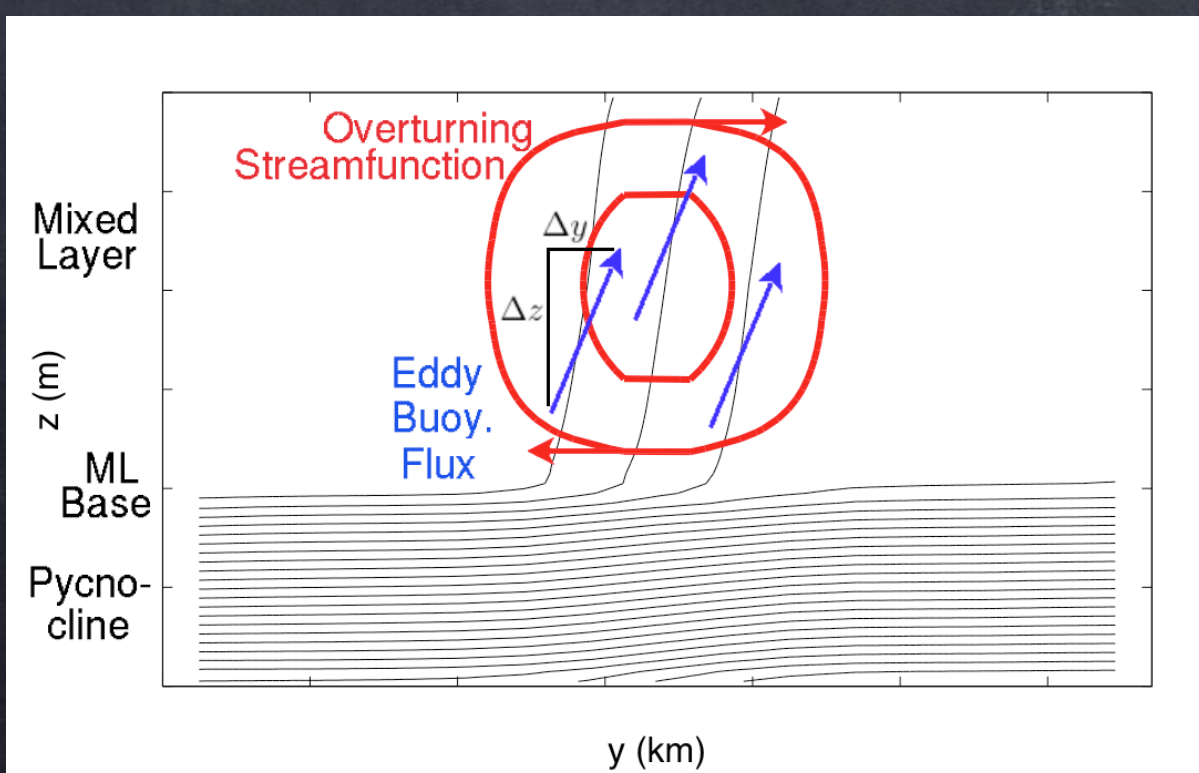
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Flux slope scales with the buoy. slope:  $\frac{\Delta y}{\Delta z} \propto \frac{-\frac{\partial \bar{b}}{\partial z}}{\frac{\partial \bar{b}}{\partial y}}$



$$\langle wb \rangle \propto \frac{-\Delta z \left( \Delta y \frac{\partial \bar{b}}{\partial y} + \Delta z \frac{\partial \bar{b}}{\partial z} \right)}{\Delta t}$$

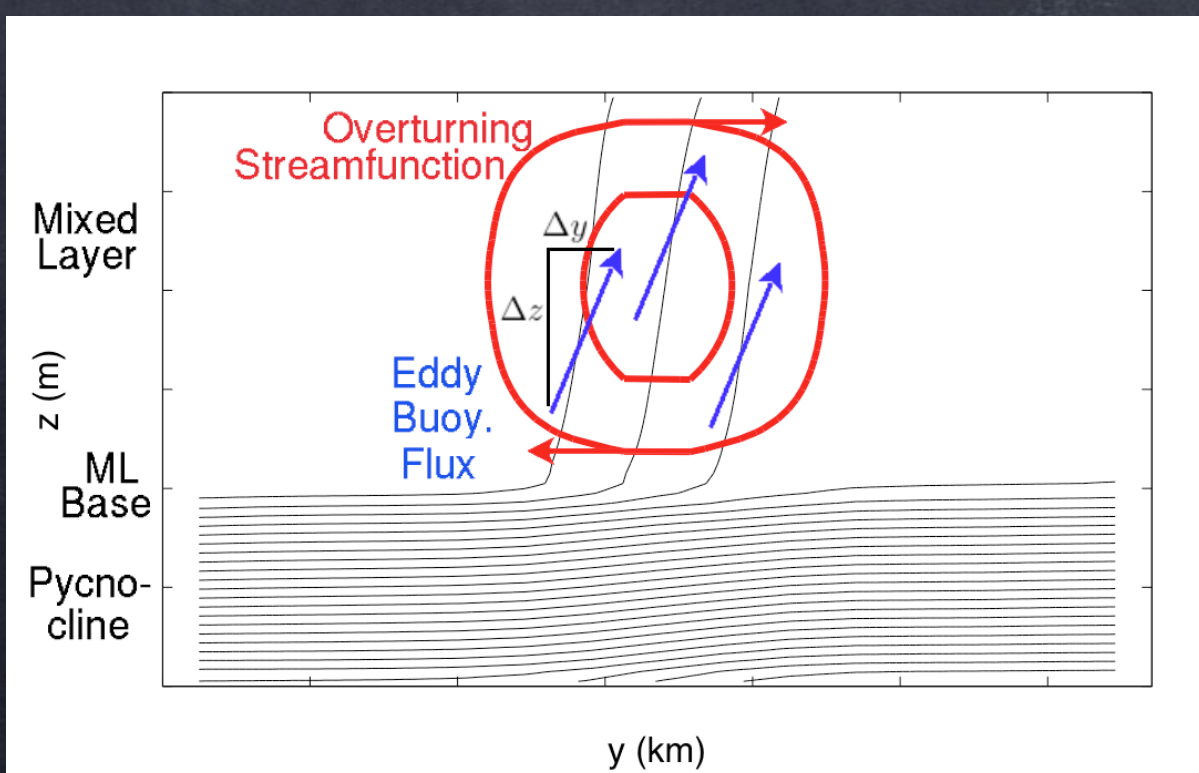
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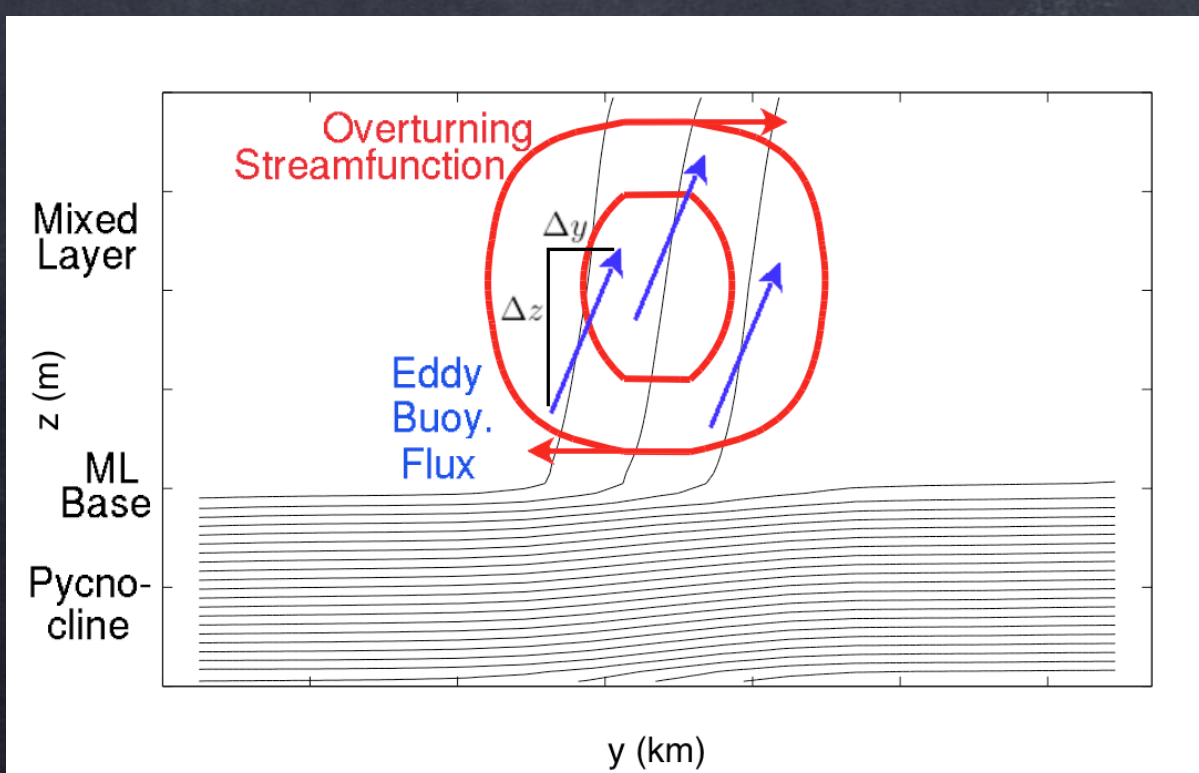
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Flux slope scales with the buoy. slope:  $\frac{\Delta y}{\Delta z} \propto \frac{-\frac{\partial \bar{b}}{\partial z}}{\frac{\partial \bar{b}}{\partial y}}$

Time scale is turnover time



$$\langle wb \rangle \propto \frac{\Delta z \Delta y \frac{\partial \bar{b}}{\partial y}}{\Delta t}$$

# Magnitude Analysis: Vert. Fluxes

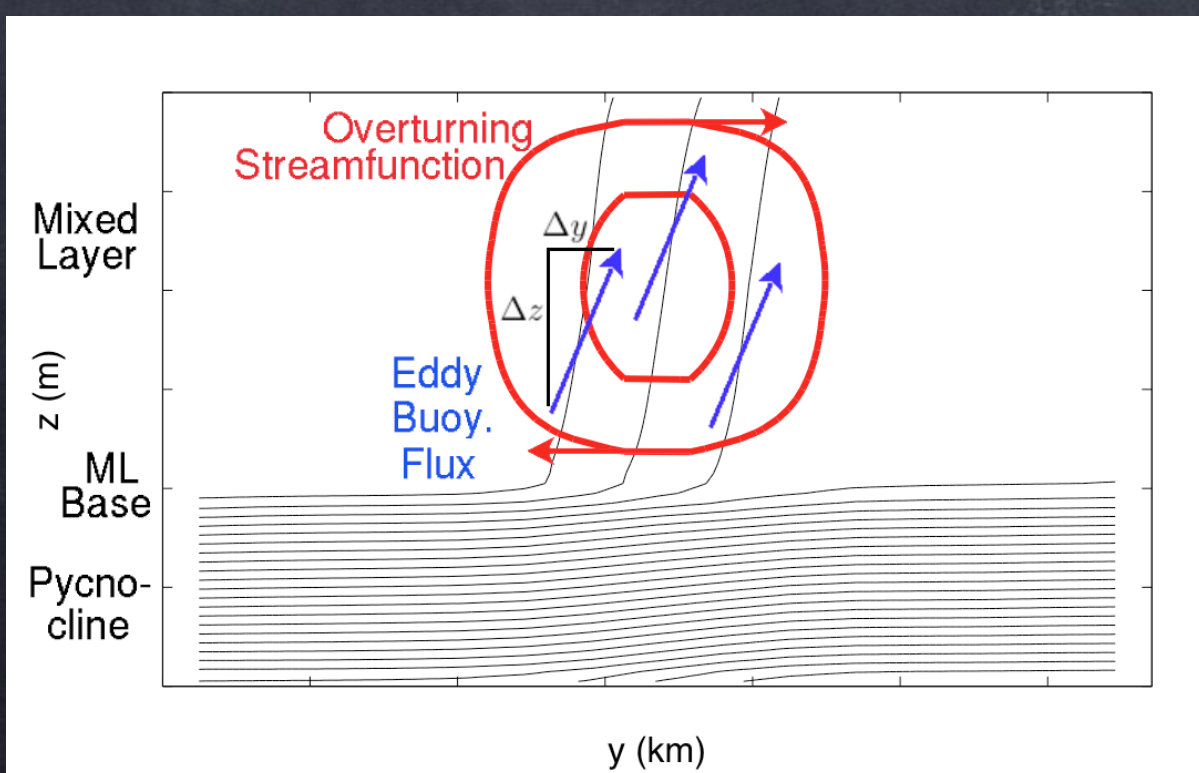
Extraction of potential energy by submesoscale eddies:

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$$\langle wb \rangle \propto \frac{\Delta z \Delta y \frac{\partial \bar{b}}{\partial y}}{\Delta y / V}$$

# Magnitude Analysis: Vert. Fluxes

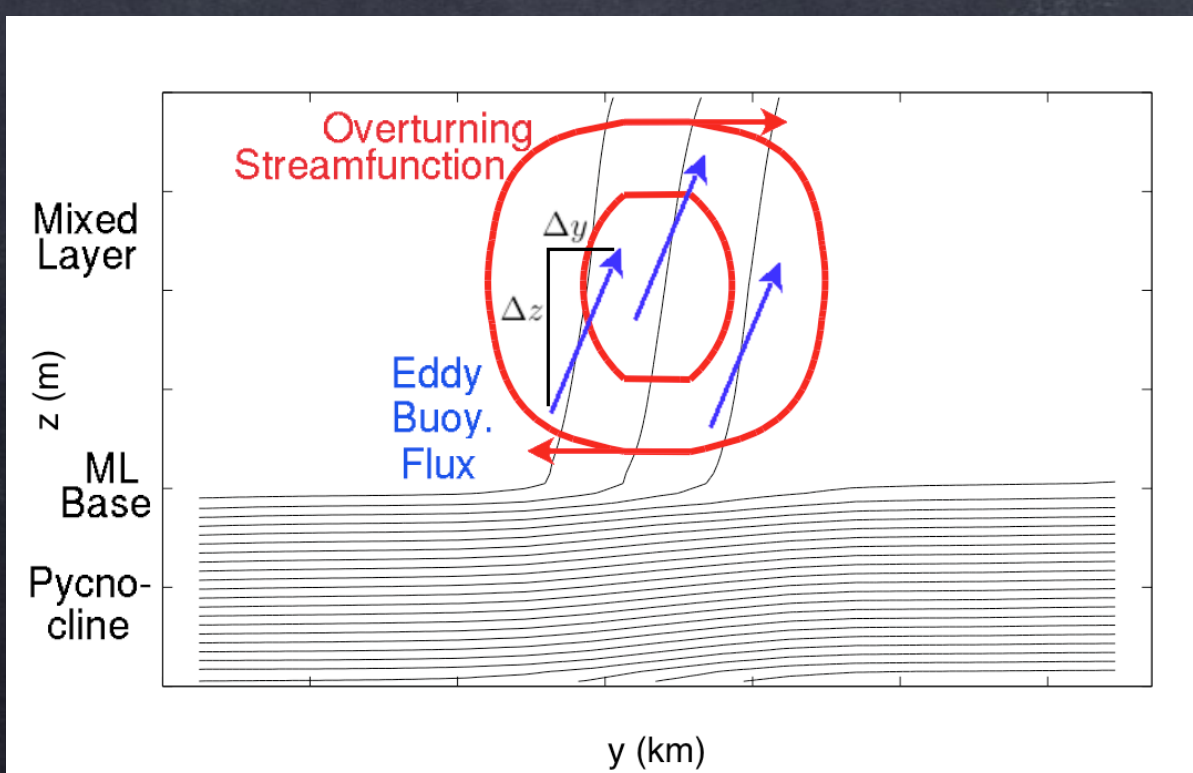
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$$\langle wb \rangle \propto \frac{\Delta z H}{|f|} \left[ \frac{\partial \bar{b}}{\partial y} \right]^2$$

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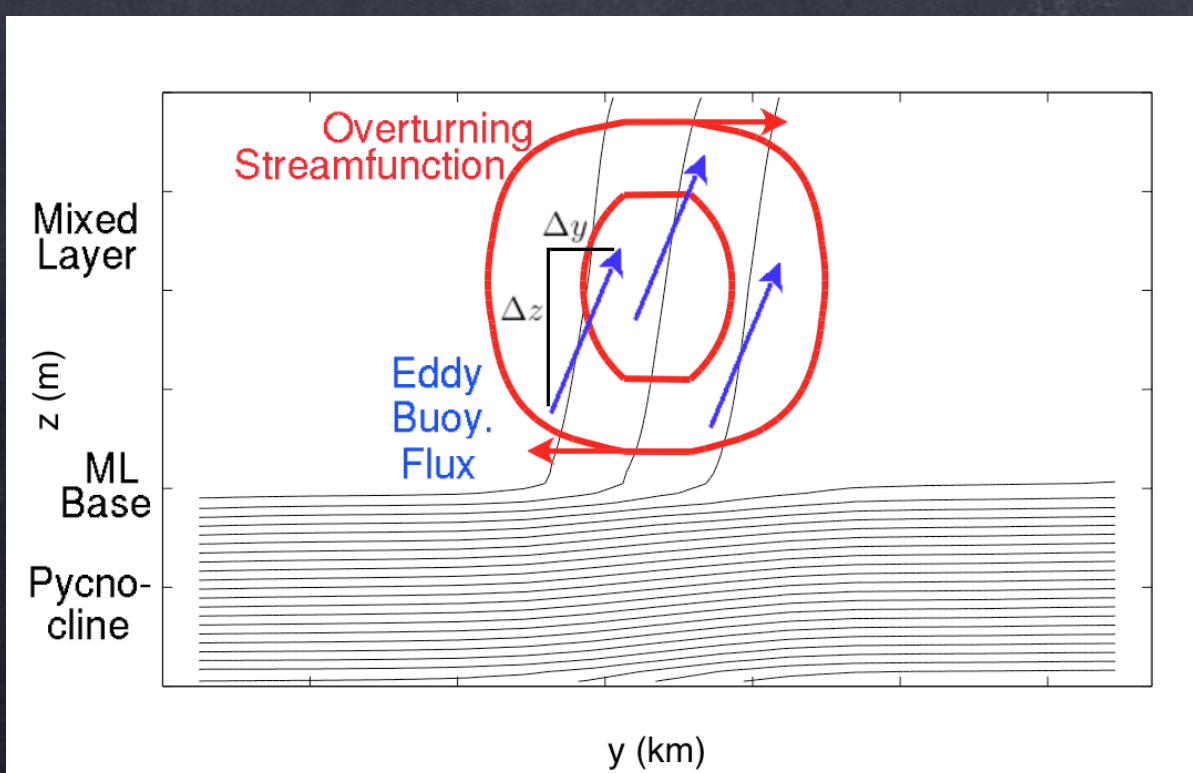
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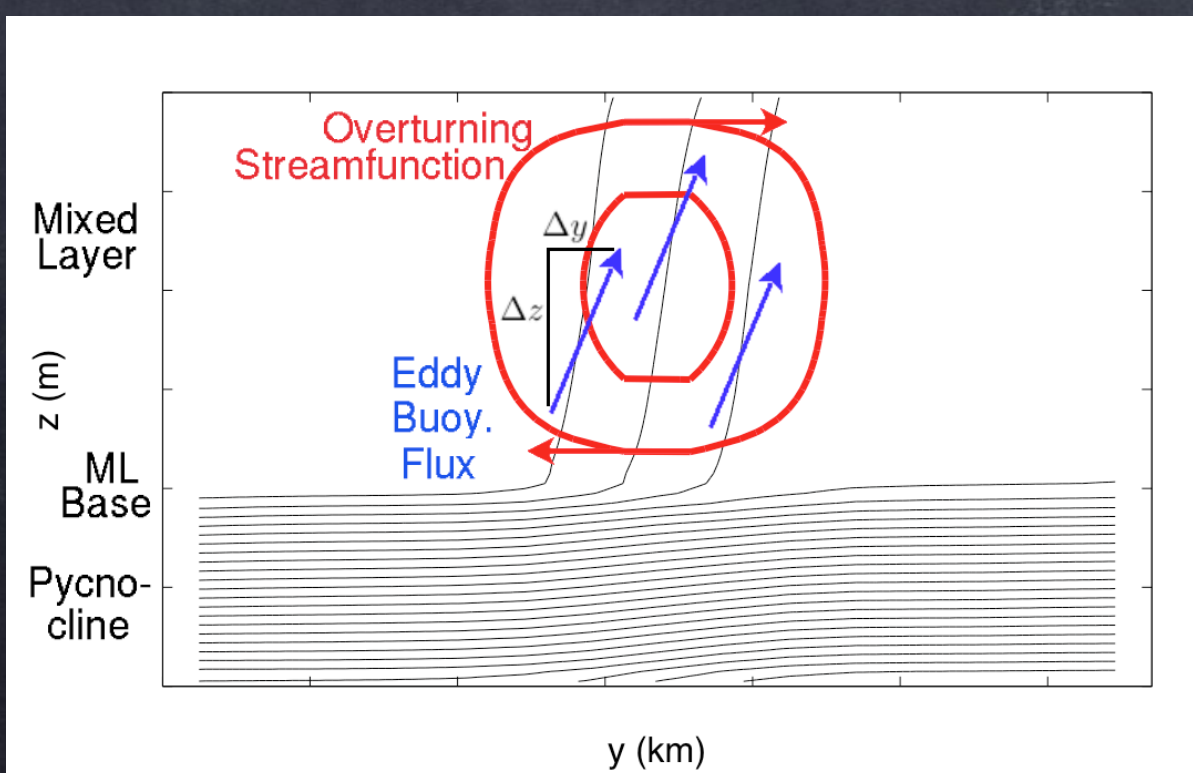
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$$\langle wb \rangle \propto \frac{H^2}{|f|} \left[ \frac{\partial \bar{b}}{\partial y} \right]^2$$



Eddies effect a largely adiabatic transfer:  
thus representable by a **streamfunction**

$$\Psi \propto \frac{H^2 \nabla \bar{b} \times \hat{\mathbf{z}}}{|f|} \longrightarrow \overline{\mathbf{u}'b'} \equiv \Psi \times \nabla \bar{b}$$

$$\overline{w'b'} \propto \frac{H^2}{|f|} |\nabla_H \bar{b}|^2$$

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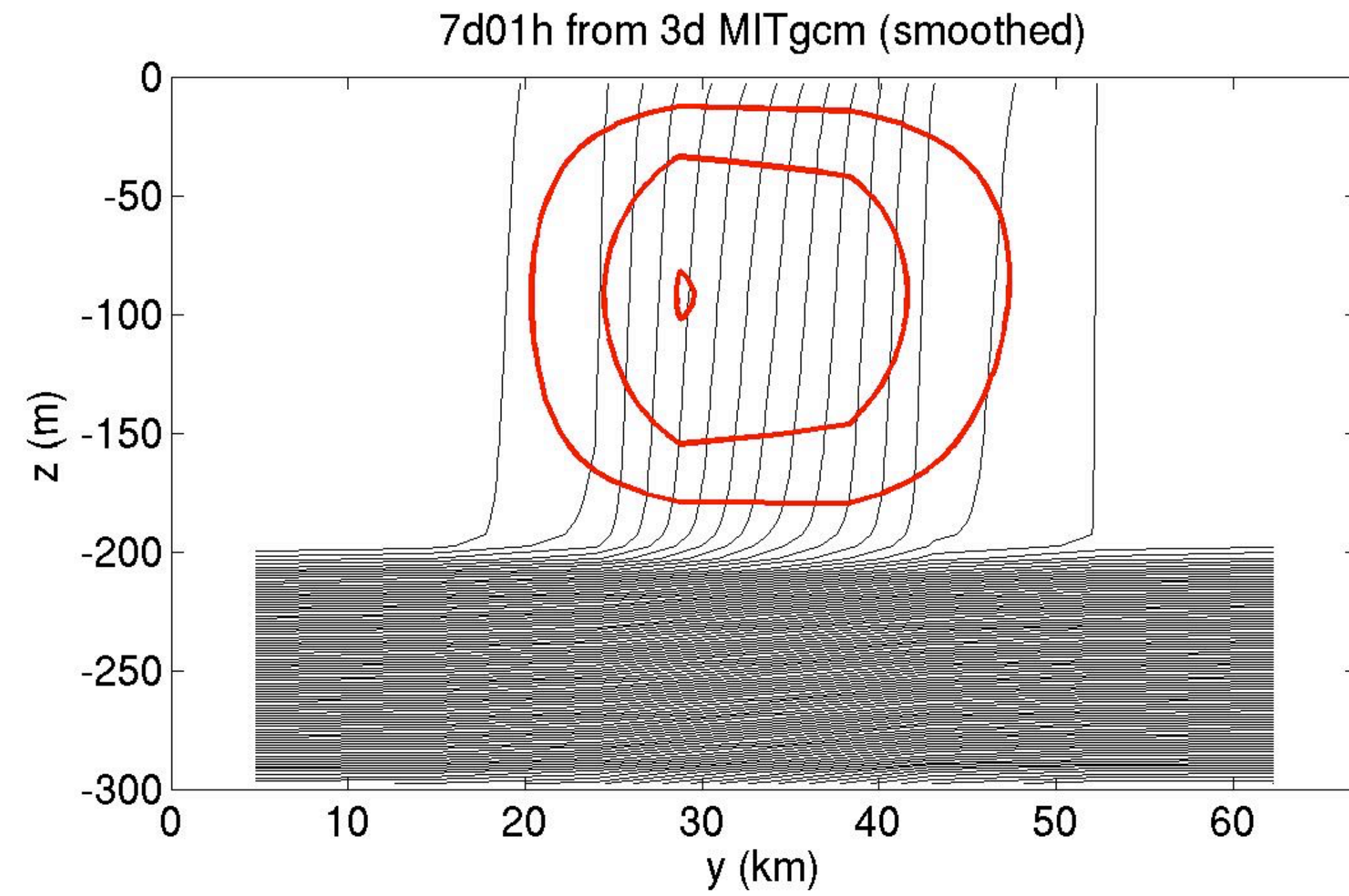
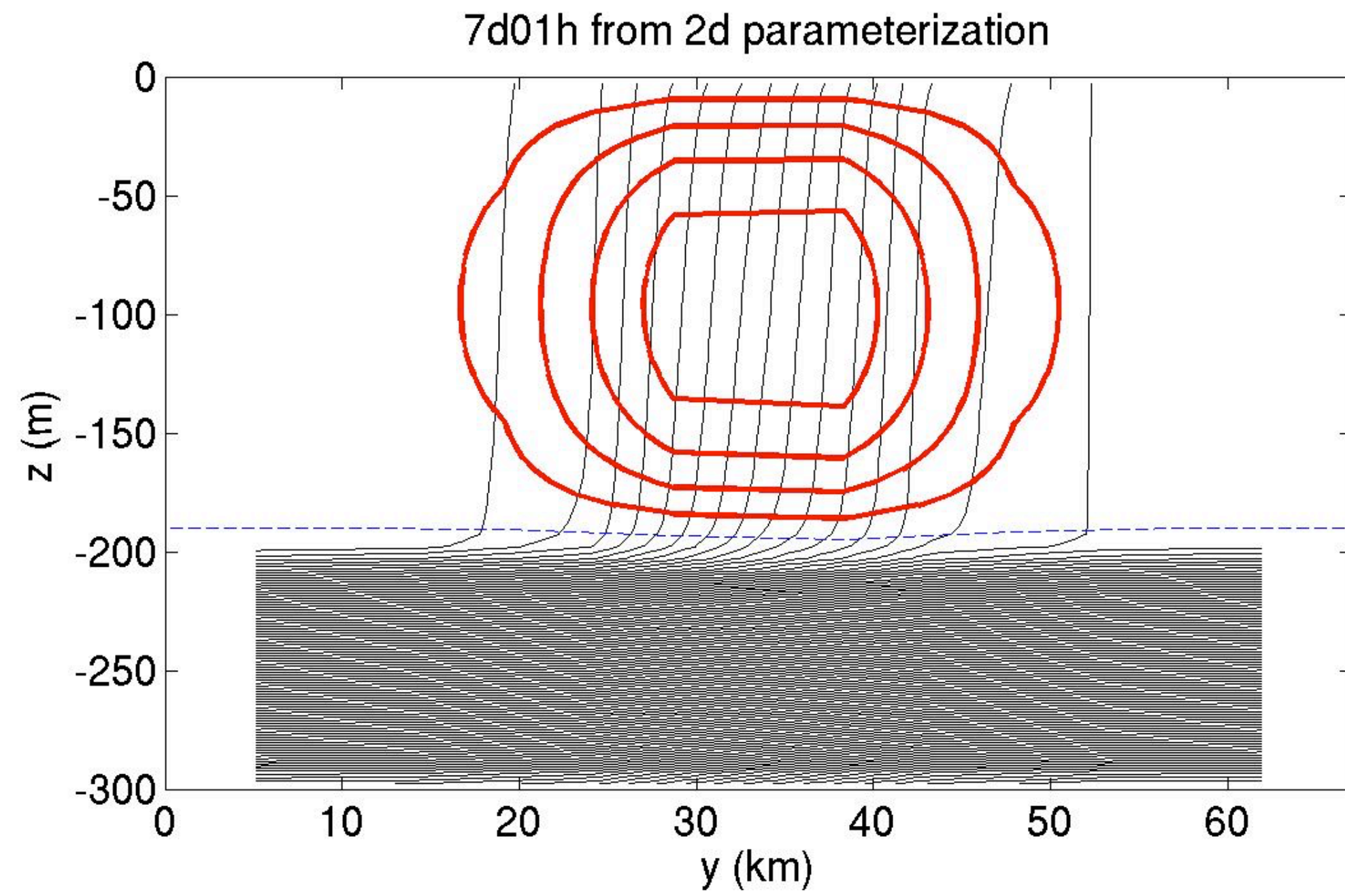
For a consistently upward,

$$\overline{w'b'} \propto \frac{H^2}{|f|} |\nabla_H \bar{b}|^2$$

And horizontally downgradient flux.

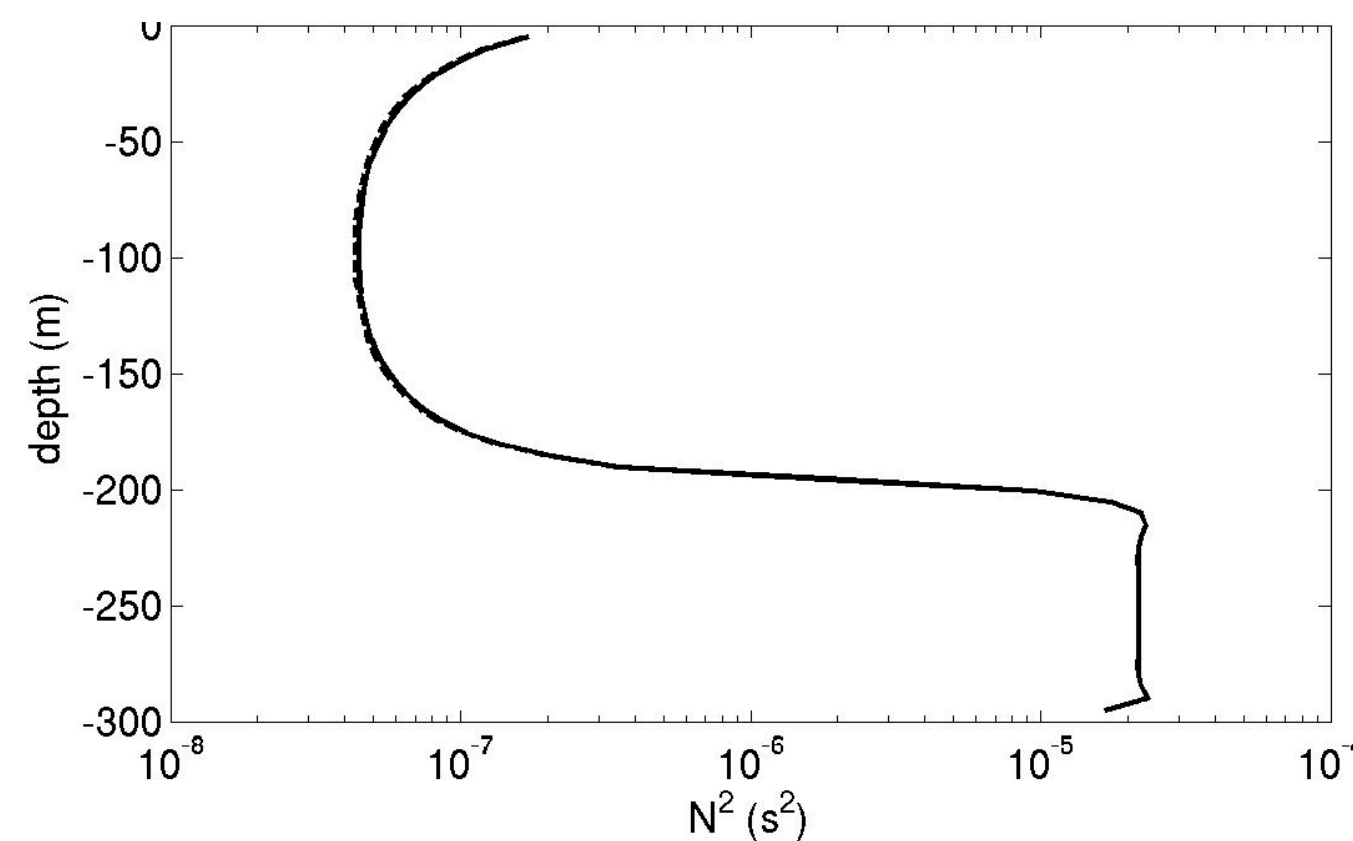
$$\overline{\mathbf{u}'_H b'} \propto \frac{-H^2 \frac{\partial \bar{b}}{\partial z}}{|f|} \nabla_H \bar{b}$$

# What does it look like?

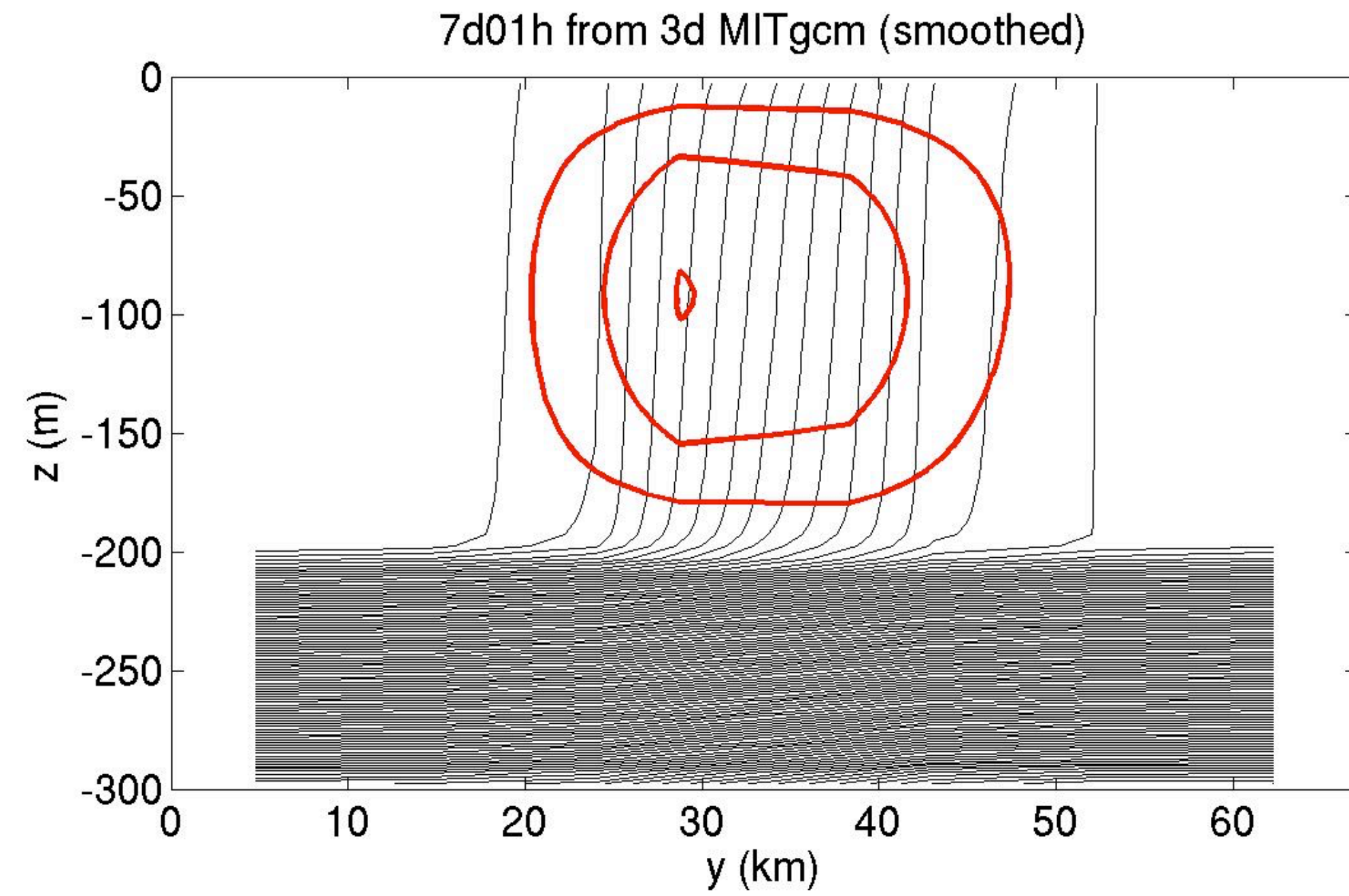
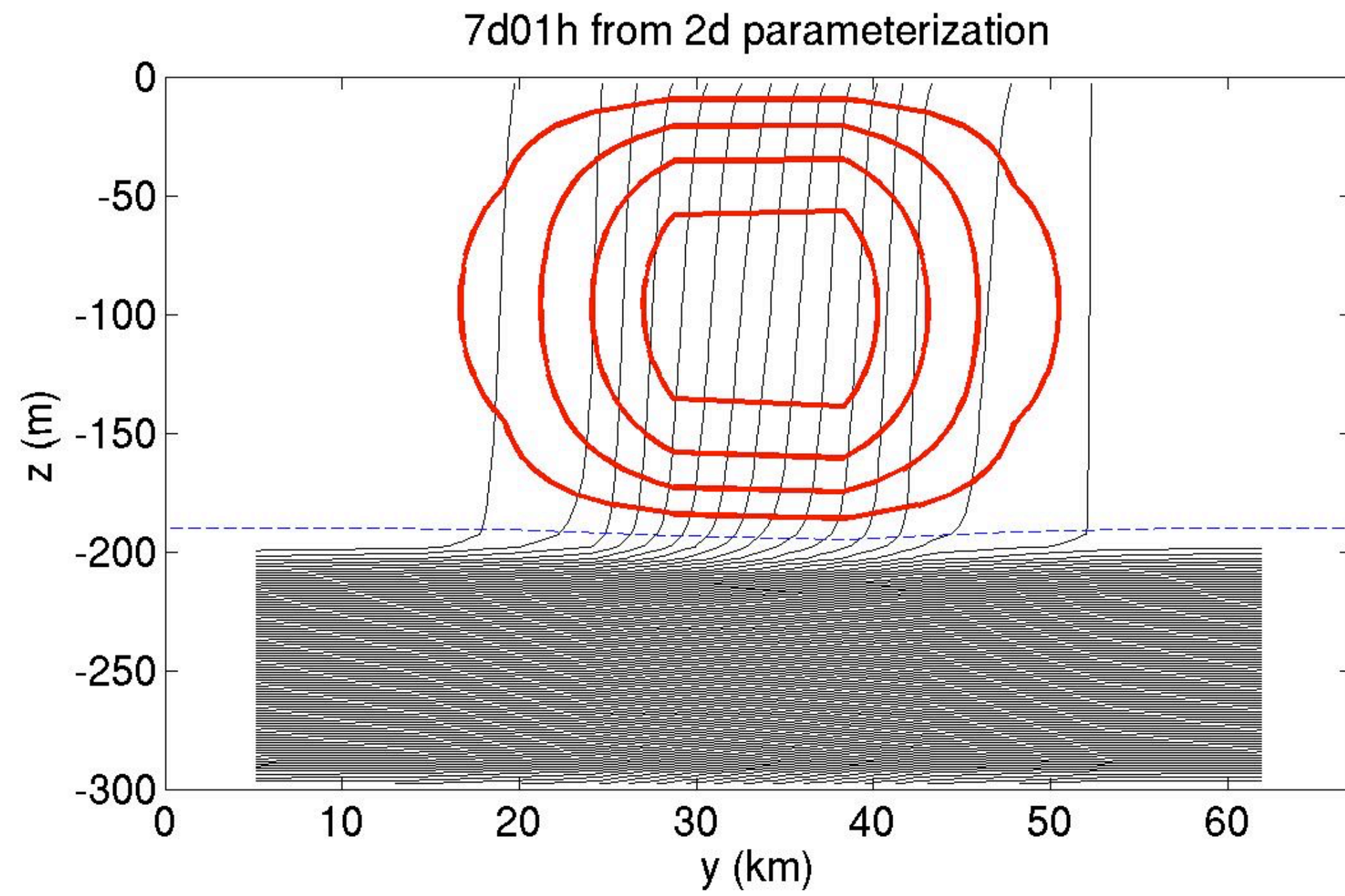


$N^2$

B. Fox-Kemper and R. Ferrari. Parameterization of mixed layer eddies. Part II: Prognosis and impact. *Journal of Physical Oceanography*, 38(6):1166-1179, 2008.

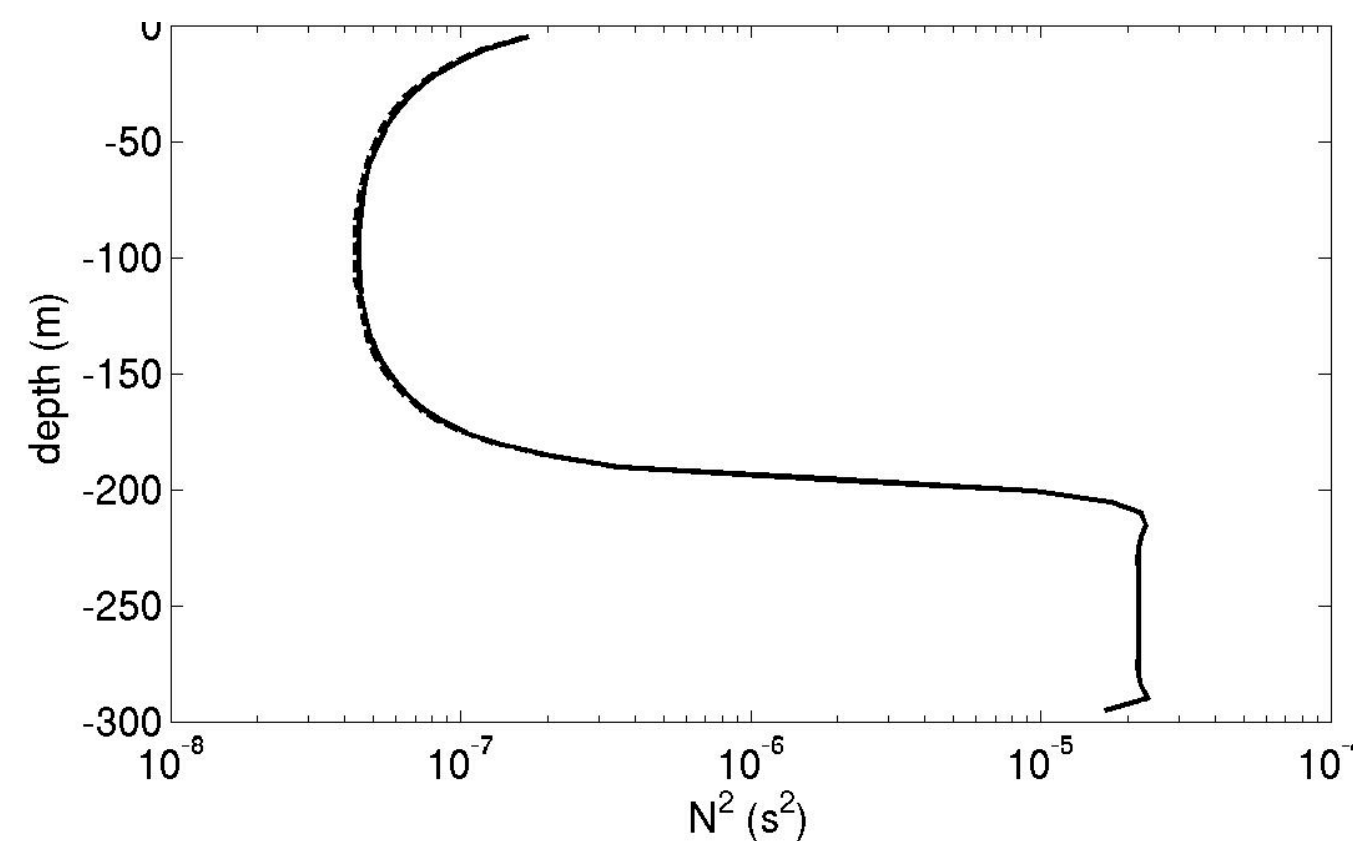


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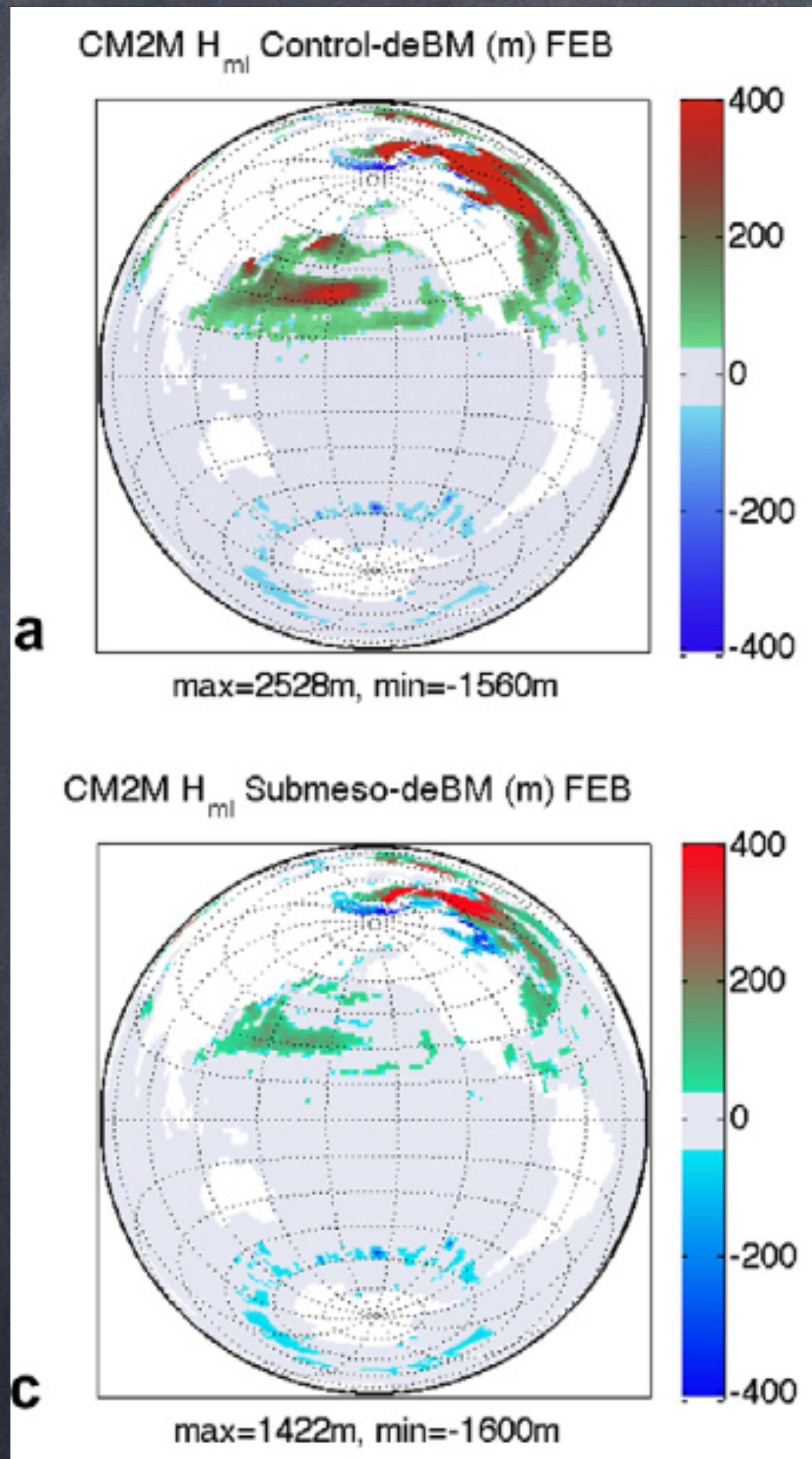


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Global Ocean Climate is SENSITIVE to these Mixed Layer Eddies! At least as parameterized Implemented in IPCC AR5 & 6: NCAR, GFDL, Hadley, NEMO,...

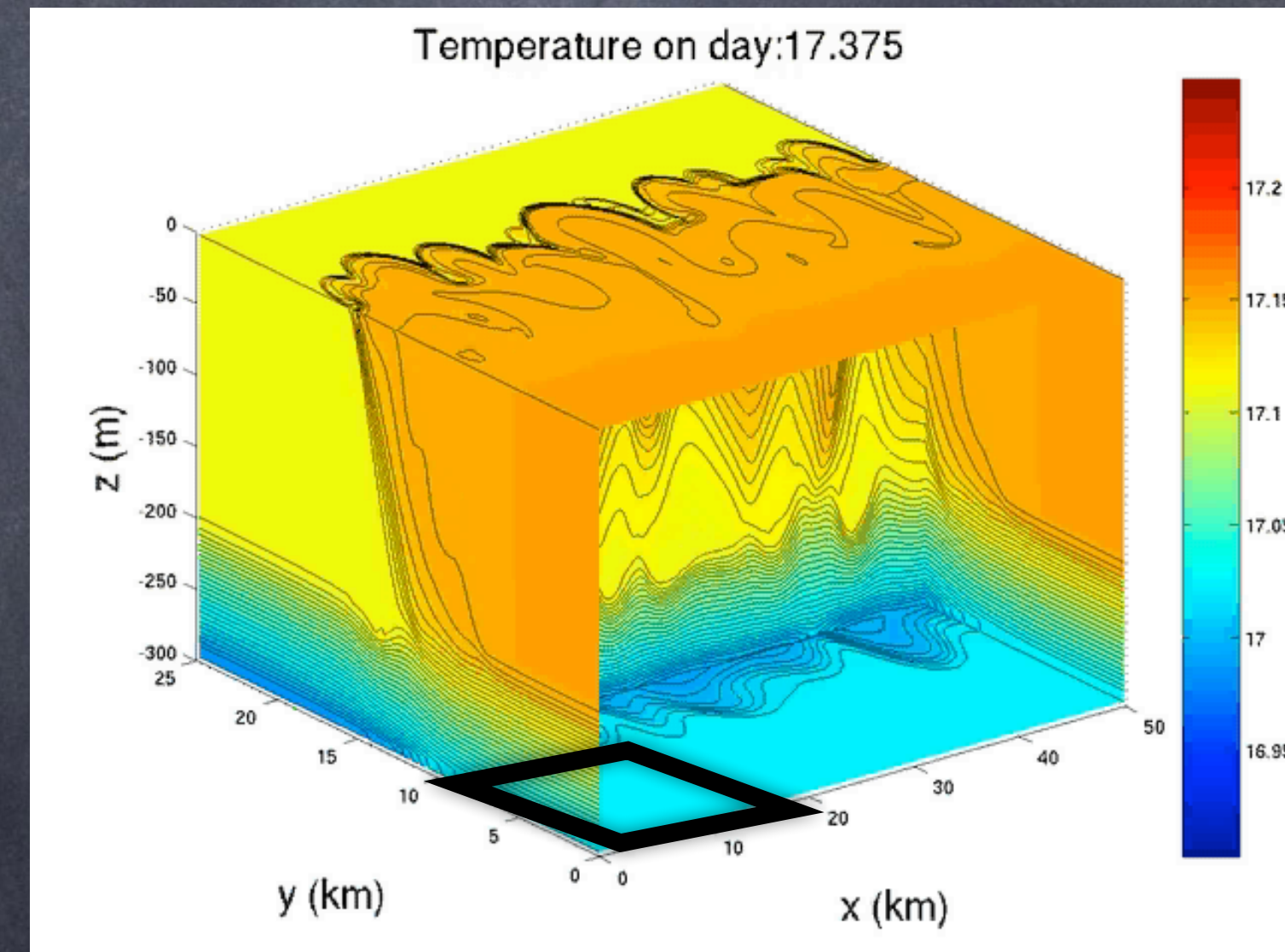


February Mixed layer depth Bias w/o MLE

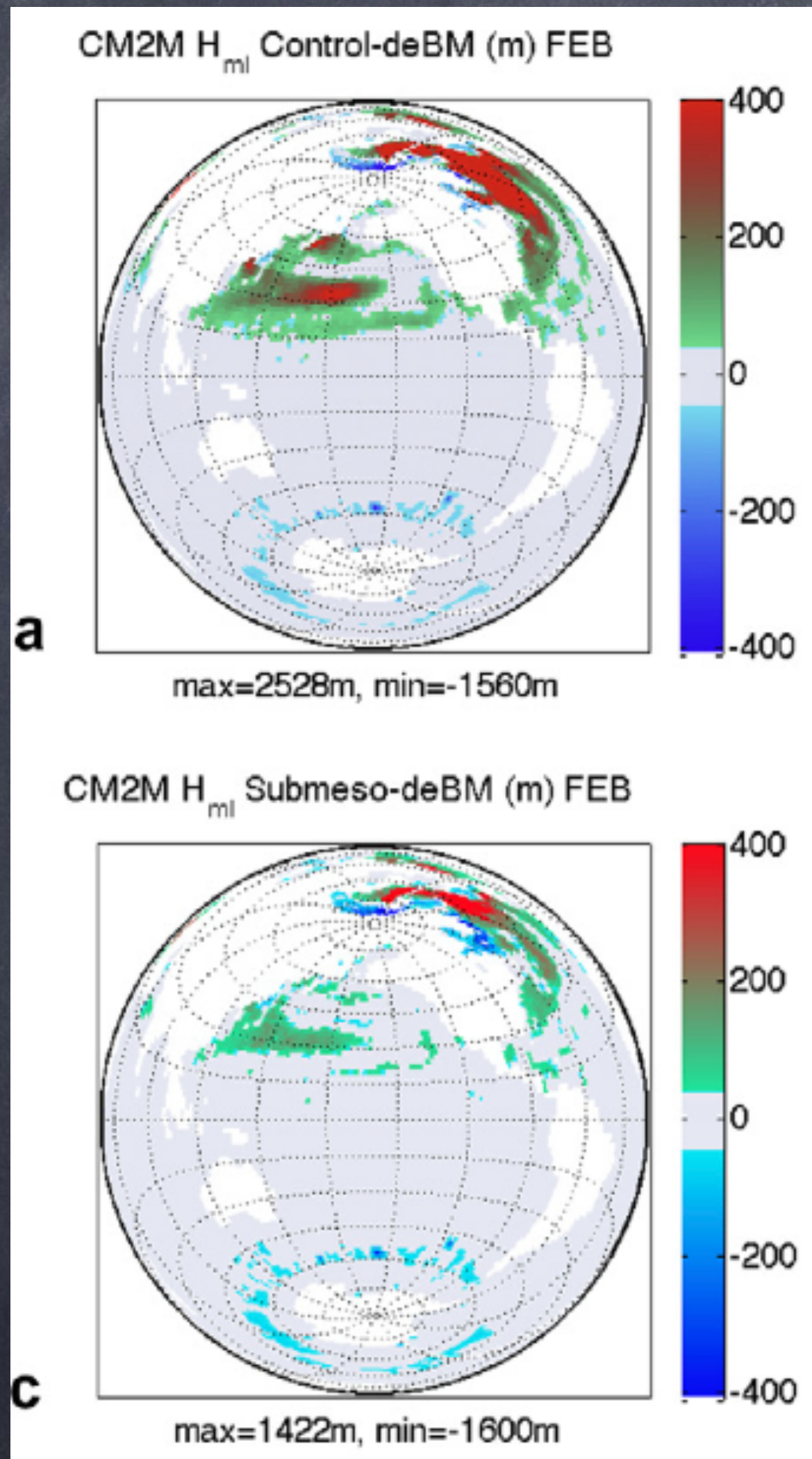
Deep Mixed Layer Bias reduced

February MLD Bias With MLE Parameterization

$O(0.1 \text{ W/m}^2)$  change to global mean net fluxes, Regional: 5 to 50  $\text{W/m}^2$



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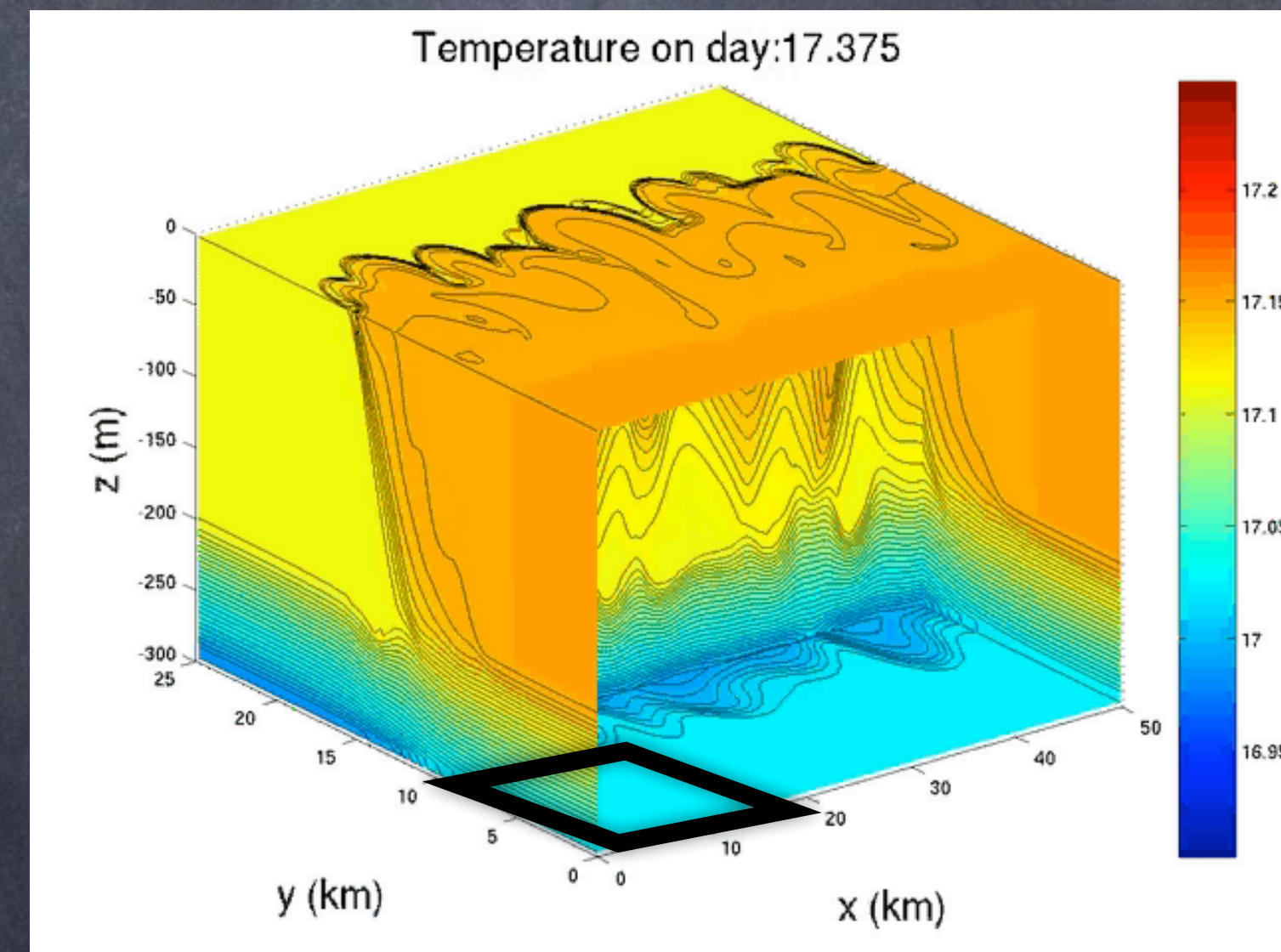


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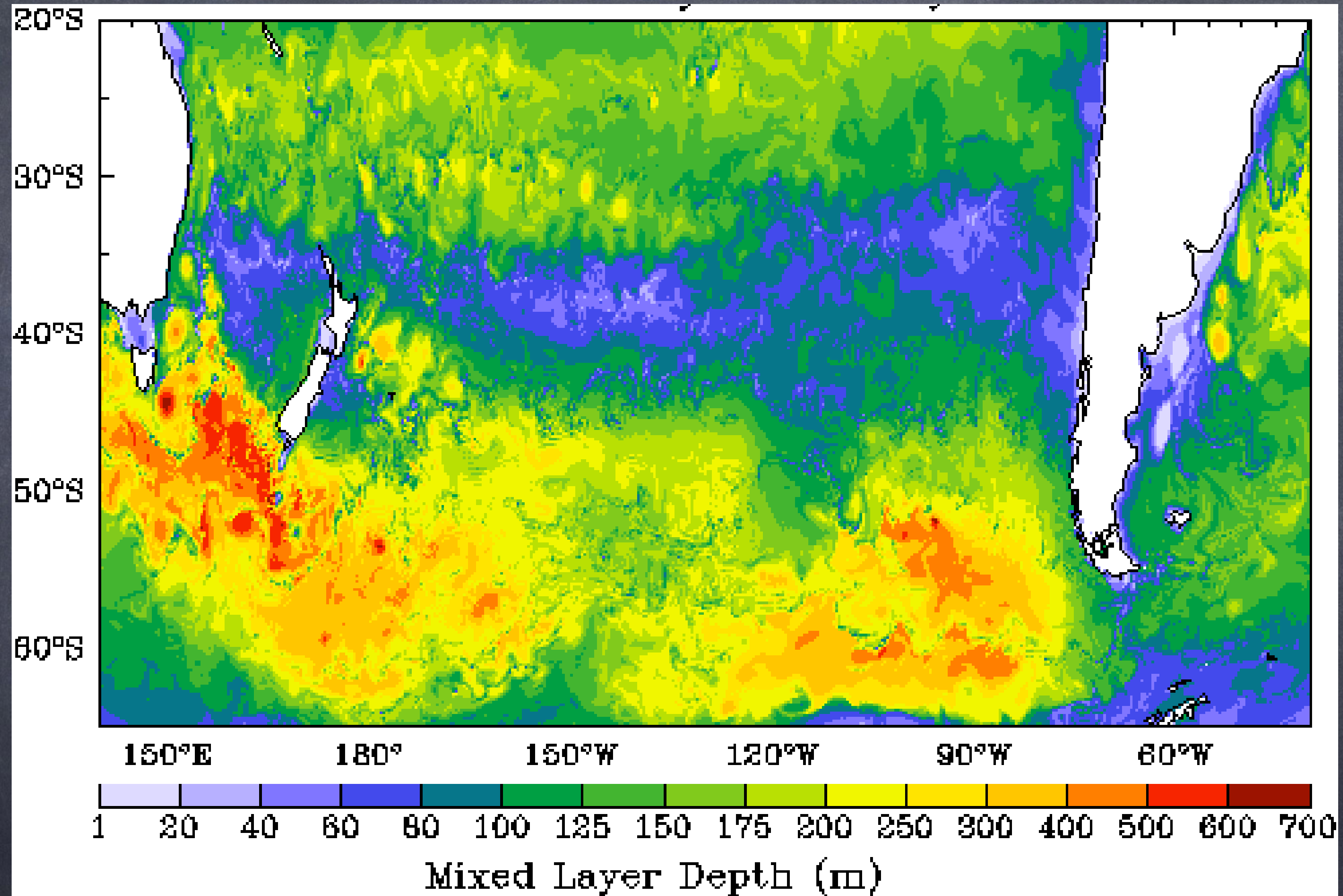
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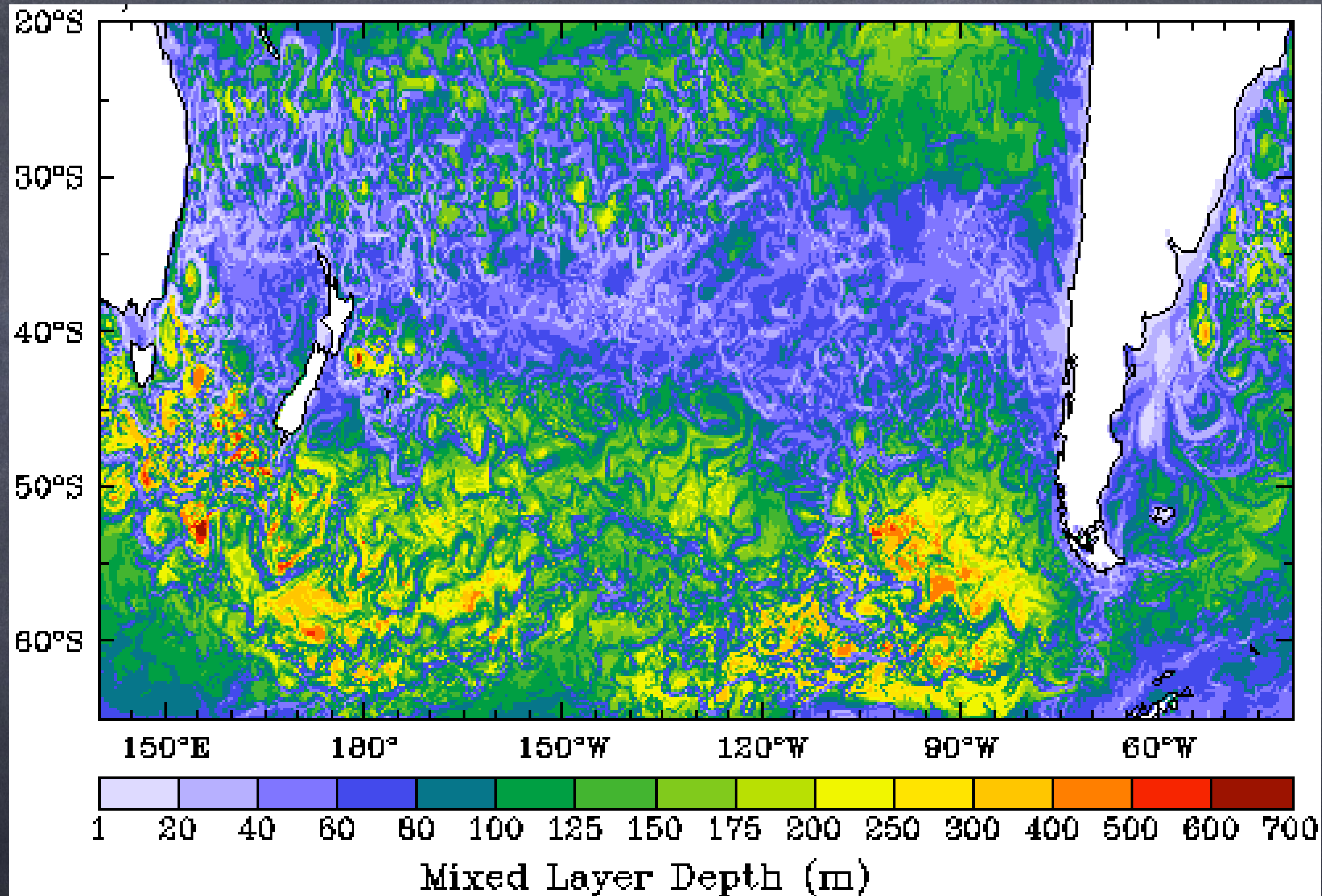
# MLEs on Mixing Layer Depth in Eddy-Resolving Southern Ocean Model



BFK, G. Danabasoglu, R. Ferrari, S. M. Griffies, R. W. Hallberg, M. M. Holland, M. E. Maltrud, S. Peacock, and B. L. Samuels. Parameterization of mixed layer eddies. III: Implementation and impact in global ocean climate simulations. *Ocean Modelling*, 39:61-78, 2011.

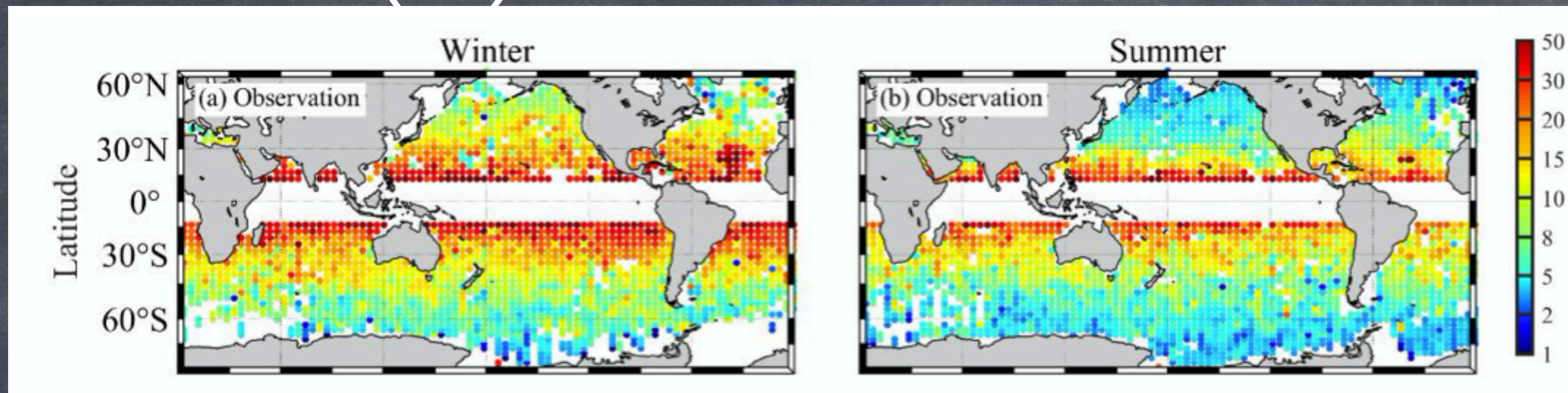


# MLEs on Mixing Layer Depth in Eddy-Resolving Southern Ocean Model



SI are roughly 10x smaller than MLI

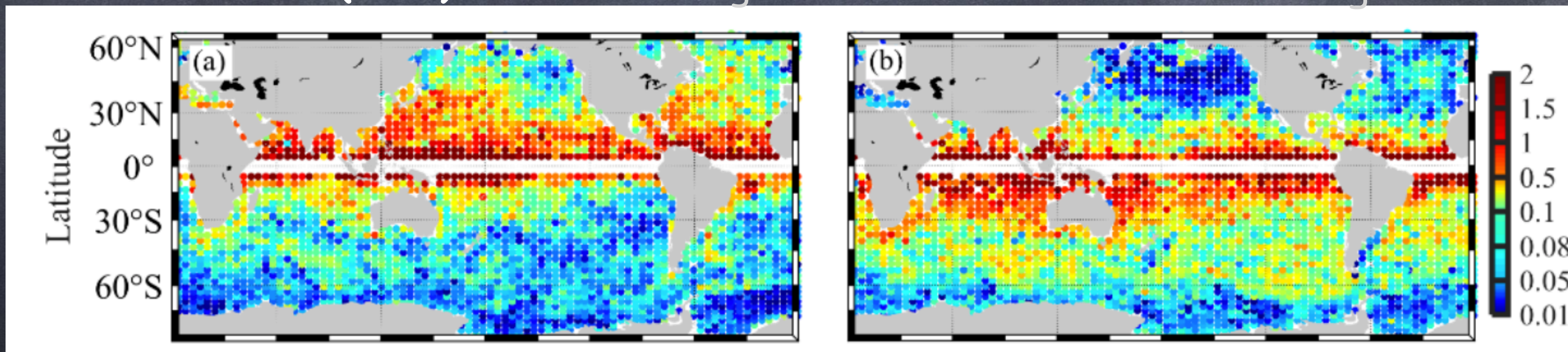
MLI Scale (km)



SI Scale (km)

February

August



J. Dong, BFK, H. Zhang, and C. Dong. The scale of submesoscale baroclinic instability globally. *JPO*, 50(9):2649-2667, 2020. [dx.doi.org/10.1175/JPO-D-20-0043.1](https://doi.org/10.1175/JPO-D-20-0043.1)

J. Dong, BFK, H. Zhang, and C. Dong. The Scale and Activity of Symmetric Instability Estimated from a Global Submesoscale-Permitting Ocean Model.

*JPO*, 2021. [dx.doi.org/10.1175/JPO-D-20-0159.1](https://doi.org/10.1175/JPO-D-20-0159.1)

# Sensitivity of Climate to Submeso: AMOC & Cryosphere Impacts

May Stabilize AMOC

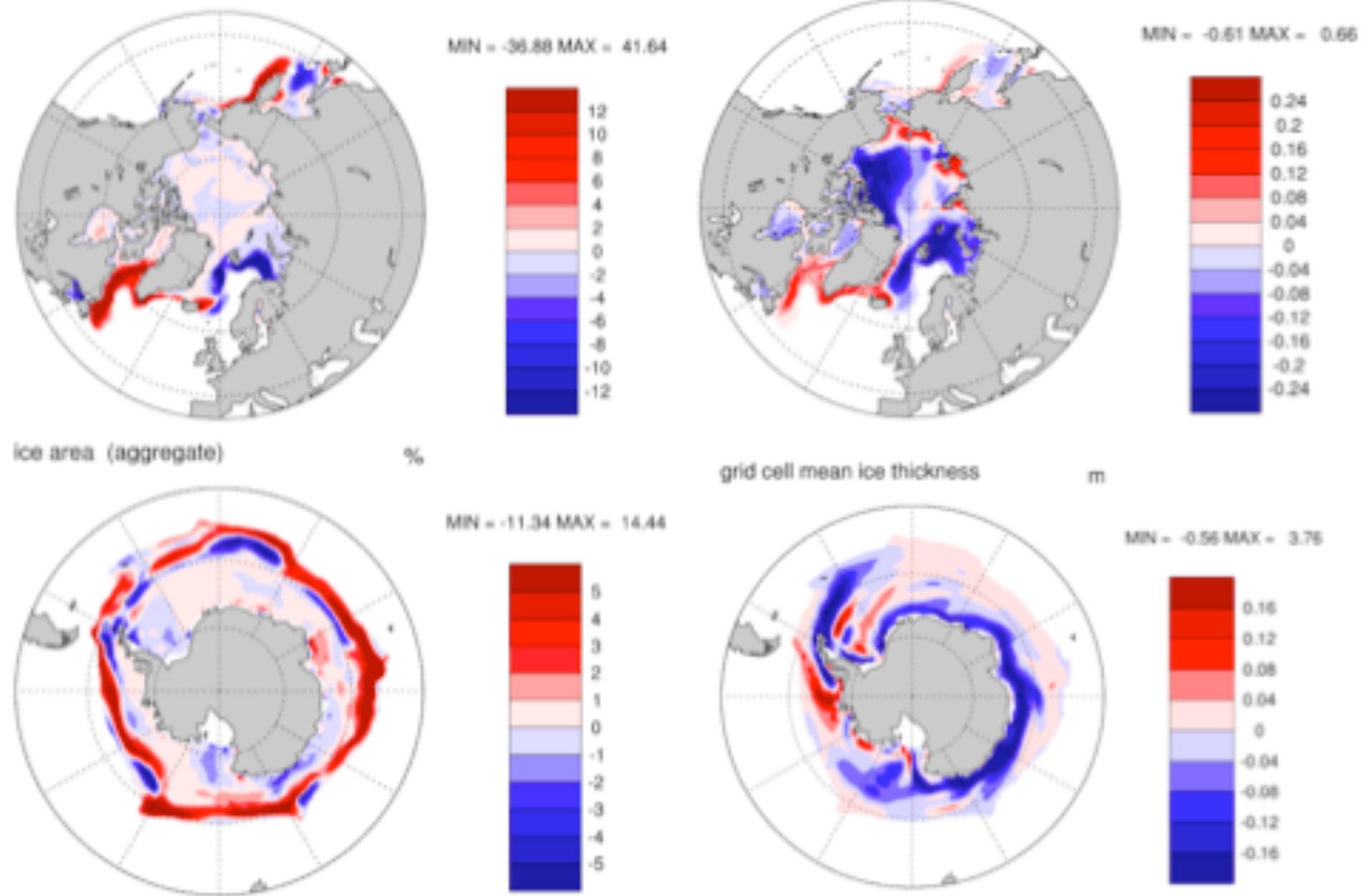
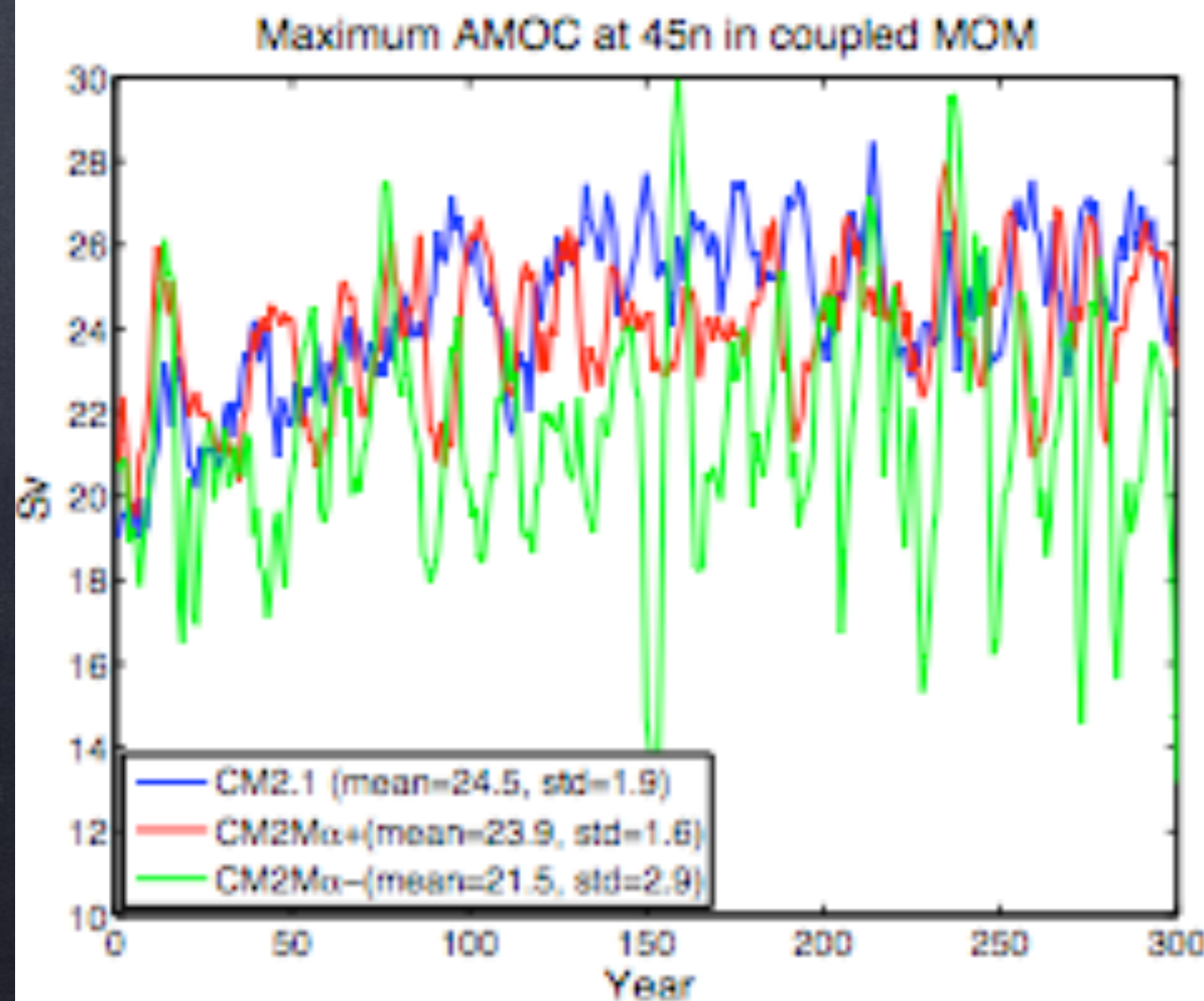


Figure 10: Wintertime sea ice sensitivity to introduction of MLE parameterization (CCSM<sup>+</sup> minus CCSM<sup>-</sup>): January to March Northern Hemisphere a) ice area and b) thickness and July to September Southern Hemisphere c) ice area and d) thickness.

Affects sea ice

**NO RETUNING  
NEEDED!!!**

B. Fox-Kemper, G. Danabasoglu, R. Ferrari, S. M. Griffies, R. W. Hallberg, M. M. Holland, M. E. Maltrud, S. Peacock, and B. L. Samuels. Parameterization of mixed layer eddies. III: Implementation and impact in global ocean climate simulations. *Ocean Modelling*, 39:61-78, 2011.

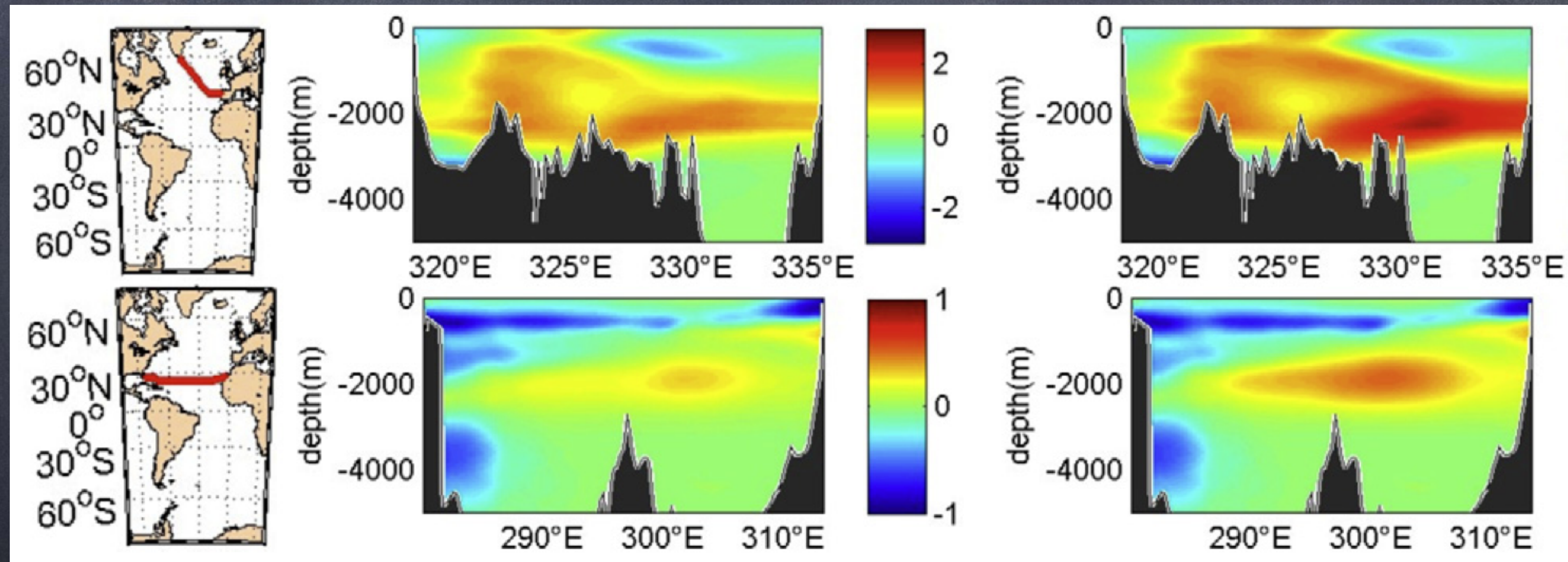
These are impacts:  
bias change unknown

# Physical Sensitivity of Ocean Climate to Submesoscale Mixed Layer Eddy Restratification:

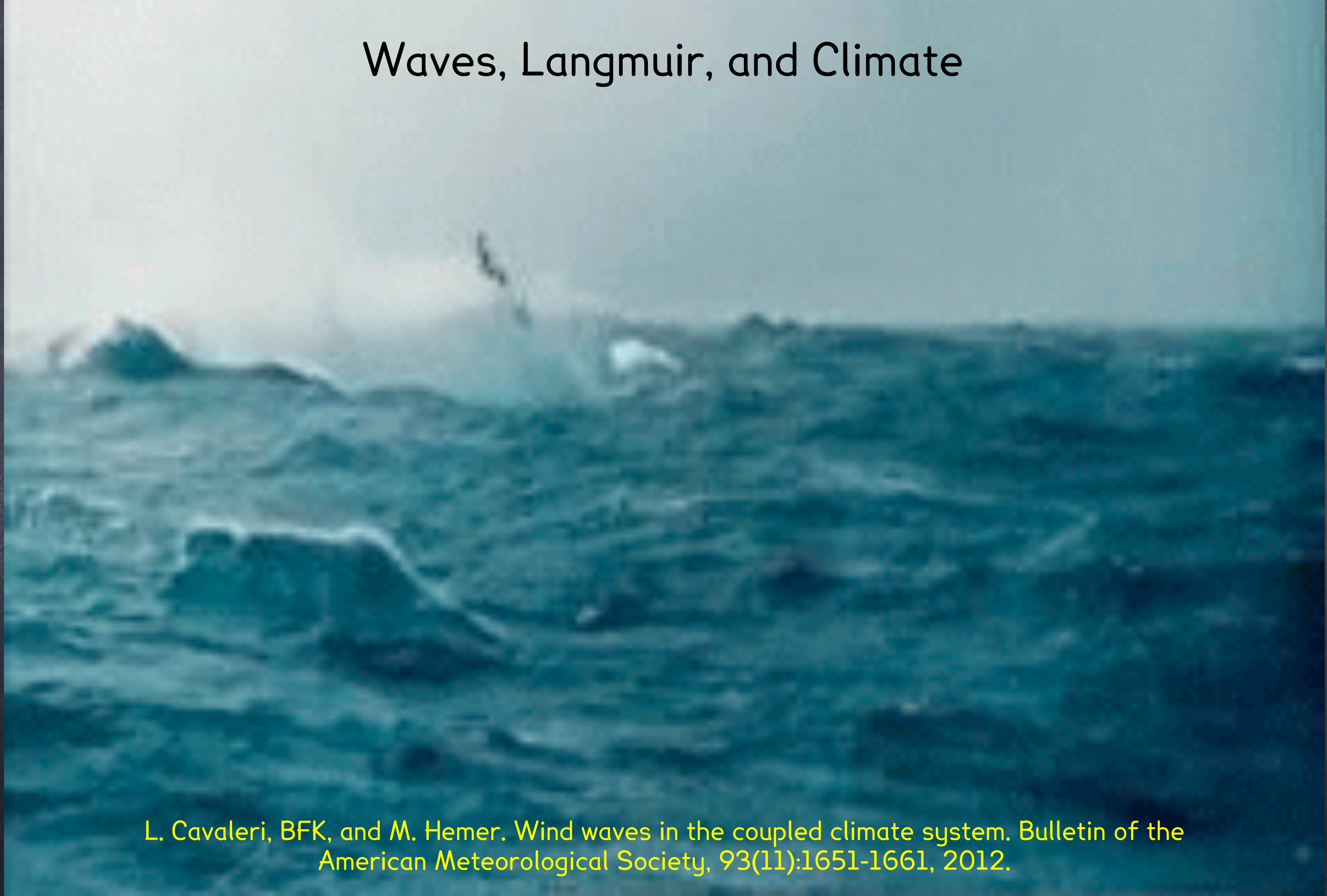
Improves CFC uptake (water masses)  
So, affects ocean heat & carbon uptake, too!

With MLE  
Parameterization

Bias w/o MLE

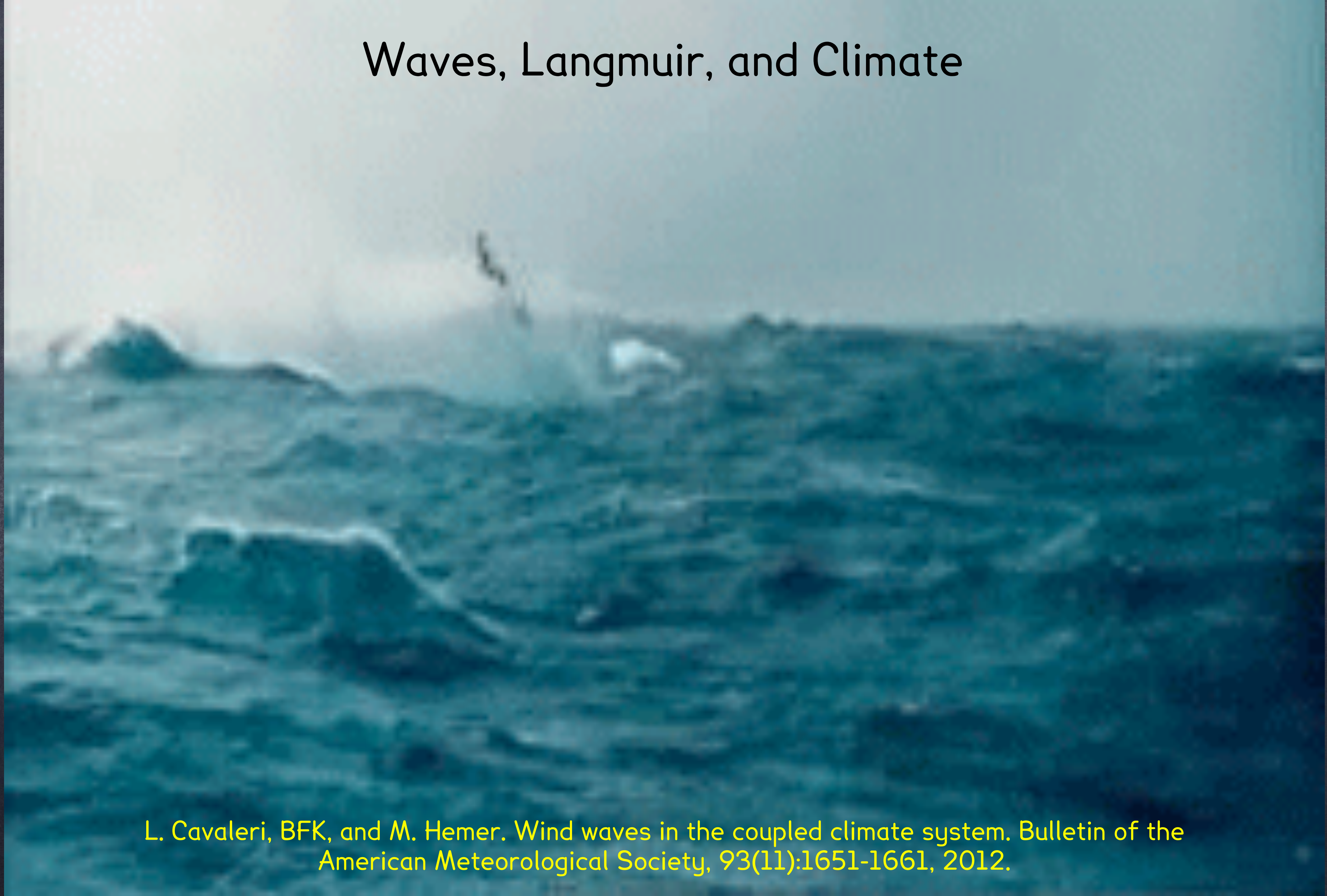


# Waves, Langmuir, and Climate



L. Cavaleri, BFK, and M. Hemer. Wind waves in the coupled climate system. *Bulletin of the American Meteorological Society*, 93(11):1651-1661, 2012.

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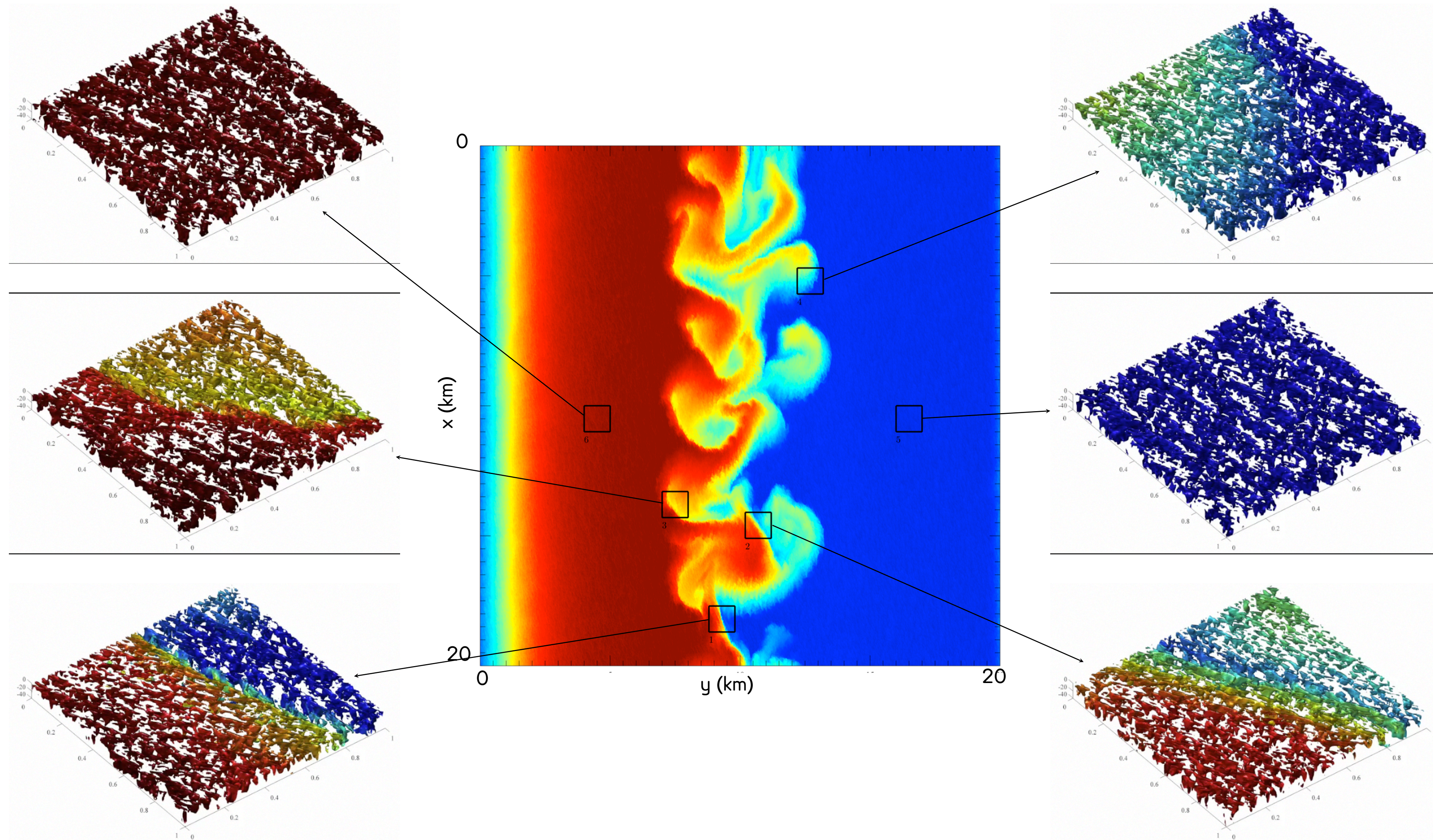
L. Cavaleri, BFK, and M. Hemer. Wind waves in the coupled climate system. *Bulletin of the American Meteorological Society*, 93(11):1651-1661, 2012.

# The Character of Langmuir Turbulence



- Near-surface
- Langmuir Cells & Langmuir Turb.
- $Ro \gg 1$
- $Ri < 1$
- 1-100m ( $H=L$ )
- 10s to 1hr
- $w \sim u = O(10\text{cm/s})$
- Nonhydrostatic
- Stokes drift
- Eqtns: Wave-Averaged aka Craik-Leibovich

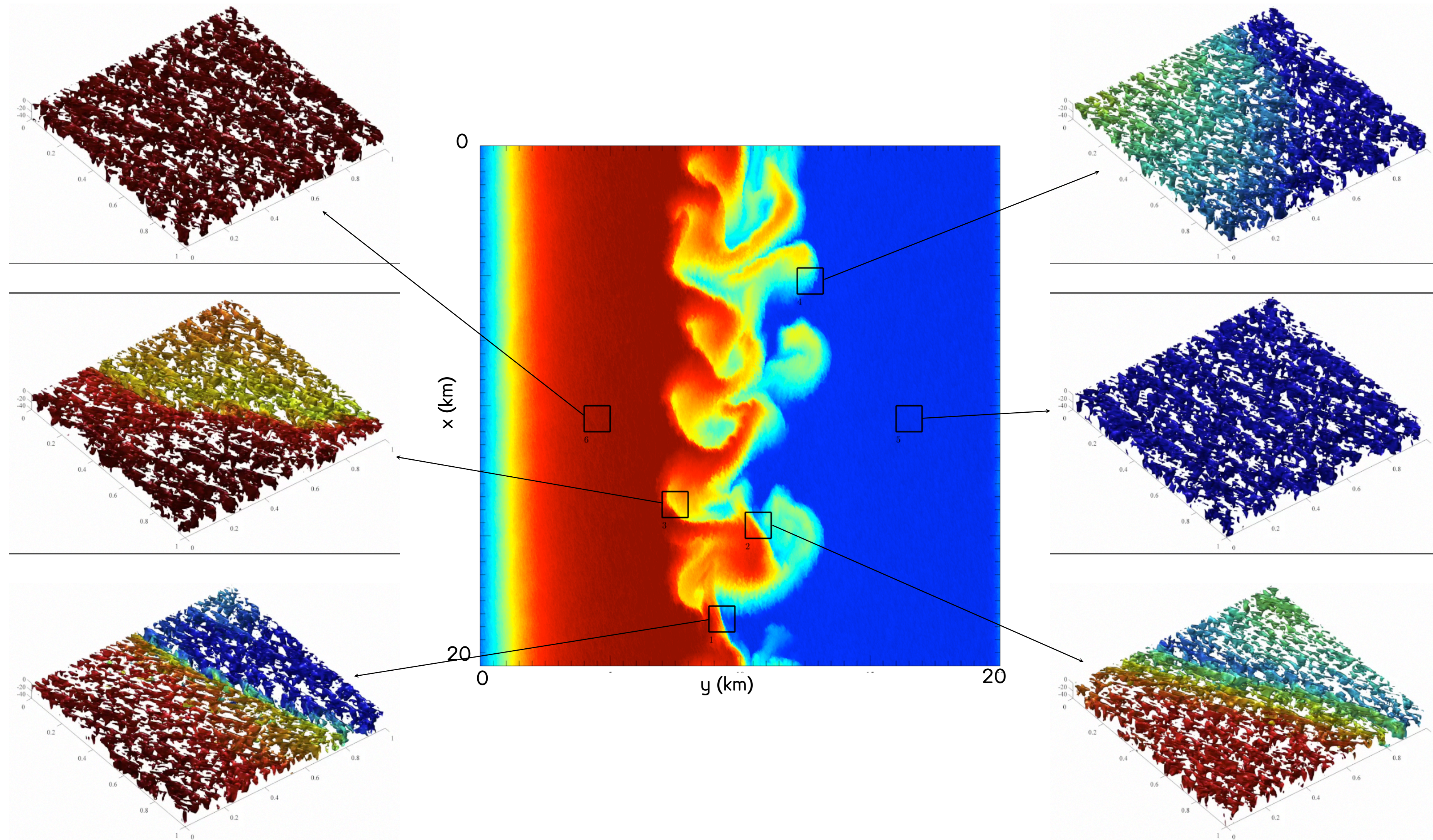
# Diverse types of interaction: Stronger Langmuir (small) Turbulence, Fronts vary from place to place, one orientation is stronger?



P. E. Hamlington, L. P. Van Roekel, BFK, K. Julien, and G. P. Chini. Langmuir-submesoscale interactions: Descriptive analysis of multiscale frontal spin-down simulations. *Journal of Physical Oceanography*, 44(9):2249-2272, September 2014.



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A photo of Langmuir windrows impinging a front in Narragansett Bay, Rhode Island (courtesy P. Cornillon)

# What Pays the Bills:

## Parameterizing Turbulence to Improve Climate and Weather Models

- These turbulent phenomena aren't only pretty—they accumulate into sizable effects onto global properties, especially boundary layer depth.
- The boundary layer in turn filters the exchange of energy, carbon, momentum, etc., between the changing atmosphere and the ocean reservoir.
- Langmuir Turbulence is an excellent recent example, which energizes the boundary layer turbulence, so entrainment and mixing are faster.

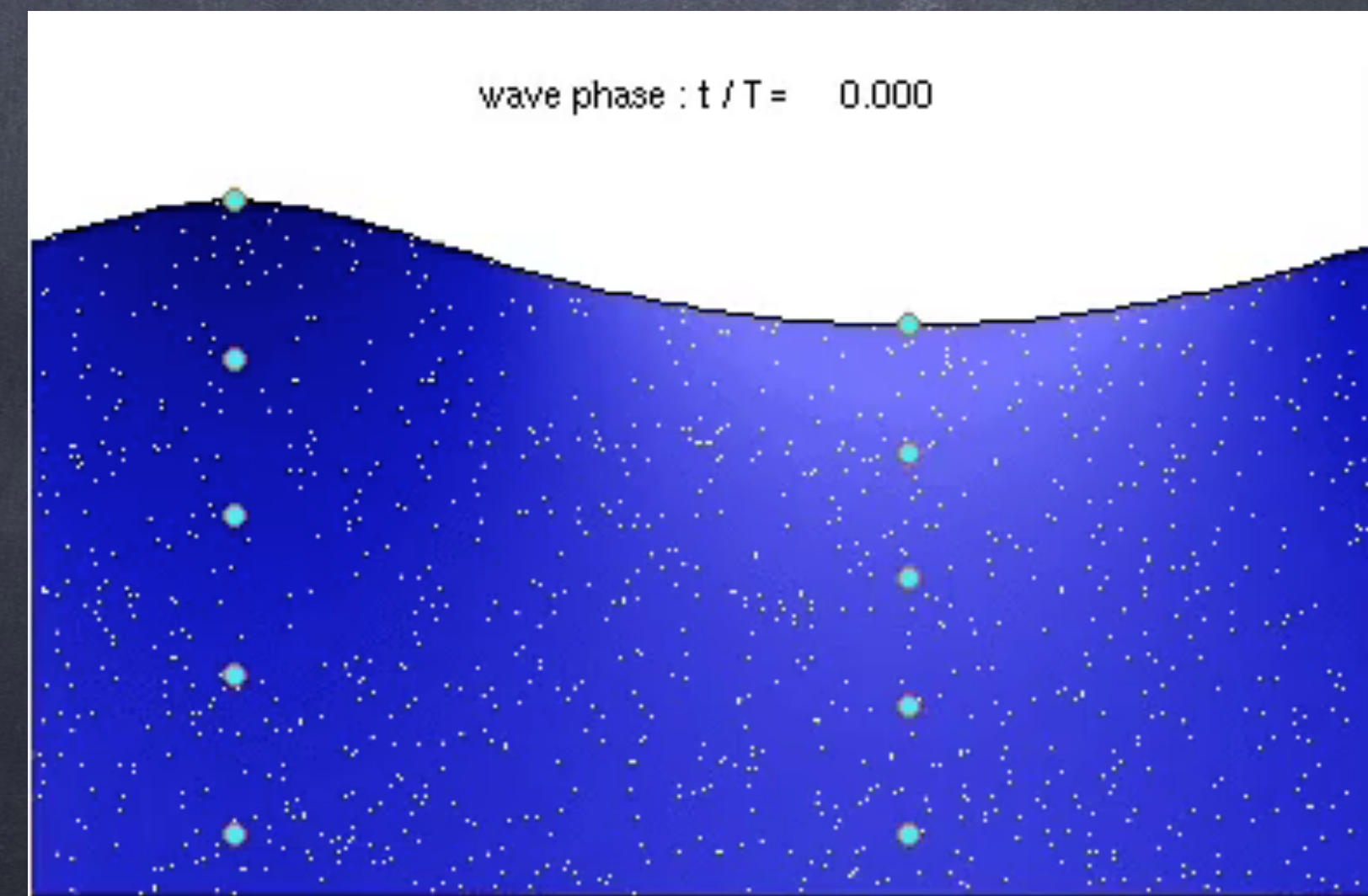
3 Effects Dominate open ocean  
“Wave-Averaged Equations”:  
(Craik, Leibovich, McWilliams et al. 1997)  
All rely only on Stokes drift of waves

1: Stokes Advection: parcels, tracers, momentum move with Lagrangian, not Eulerian flow

2: Stokes Coriolis: water parcels experience Coriolis force during this motion

3: Stokes Shear Force

N. Suzuki and BFK. Understanding Stokes forces in the wave-averaged equations. *Journal of Geophysical Research-Oceans*, 121:1-18, 2016.



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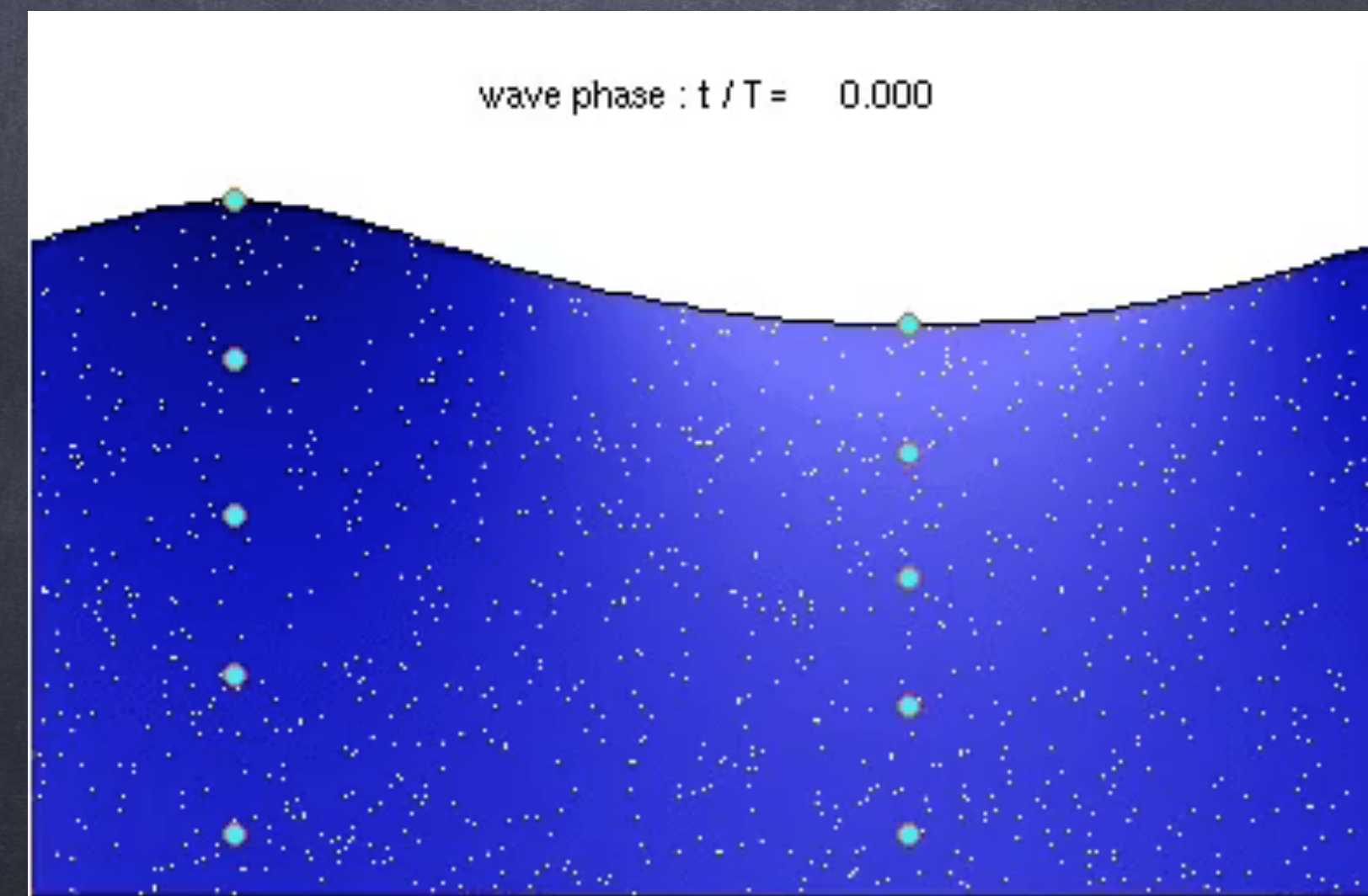
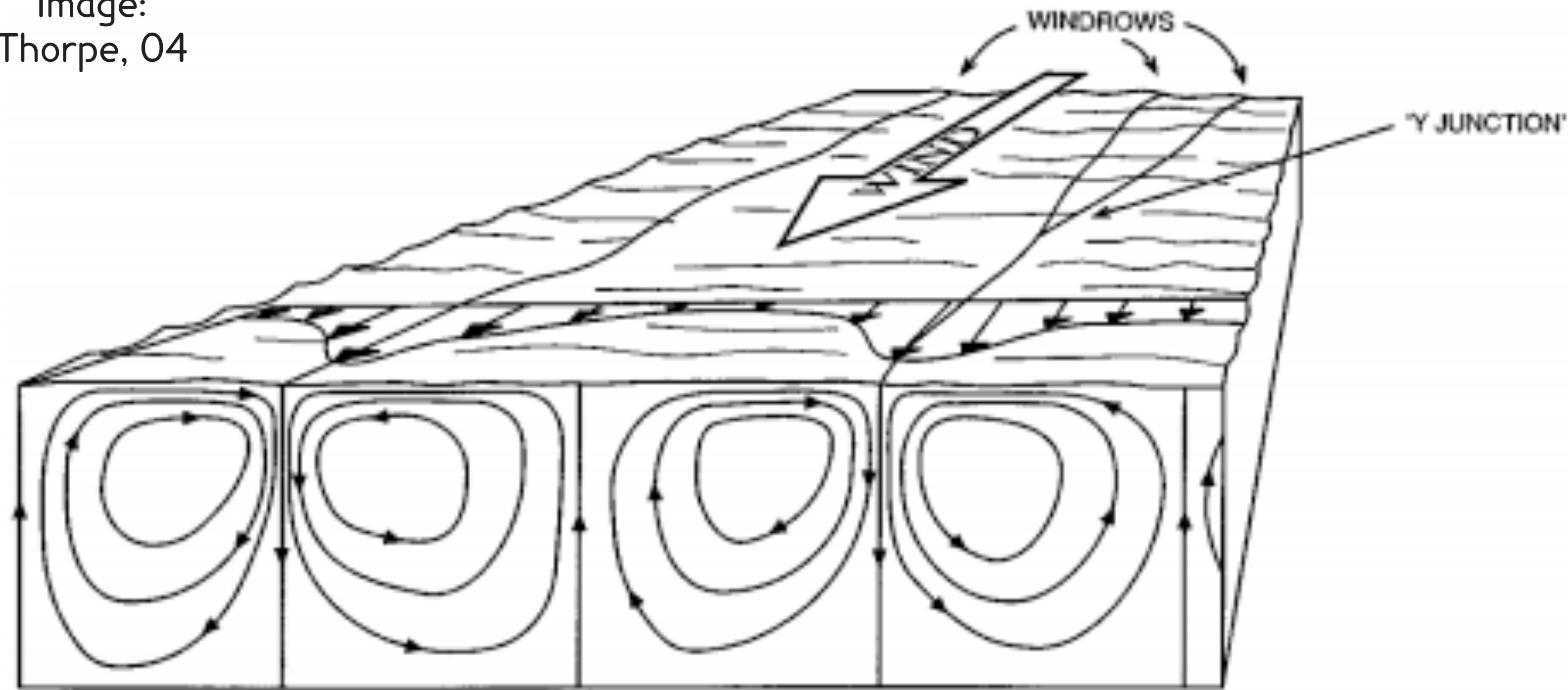


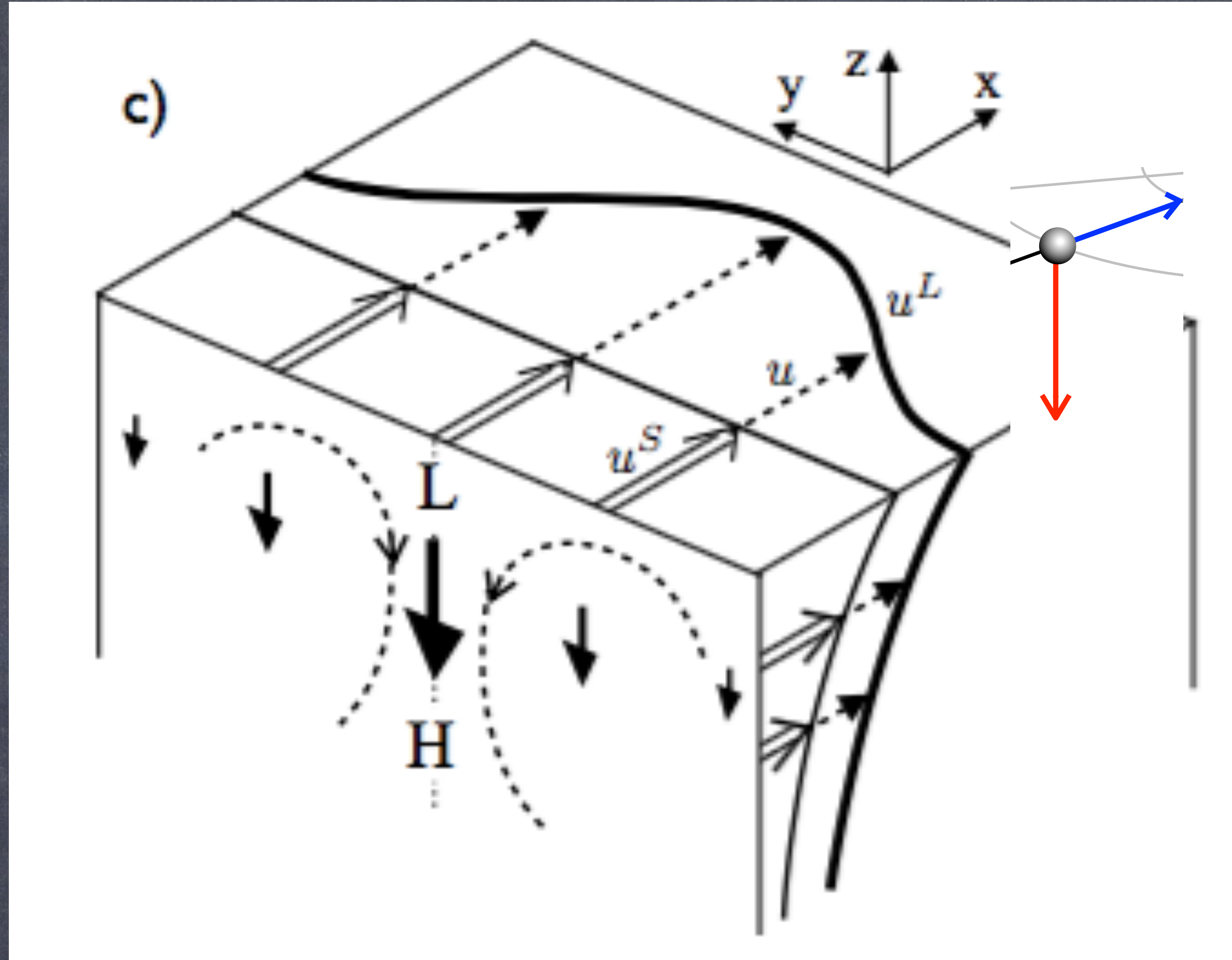
image:  
Thorpe, 04



**Figure 1** Sketch showing the pattern of mean flow in idealized Langmuir circulation. The windrows may be 2 m to 300 m apart, and the cell form is roughly square (as shown). In practice the flow is turbulent, especially near the water surface, and the windrows (Figure 2) amalgamate and meander in space and time. Bands of bubbles or buoyant algae may form within the downward-going (or downwelling) flow (see Figure 3).

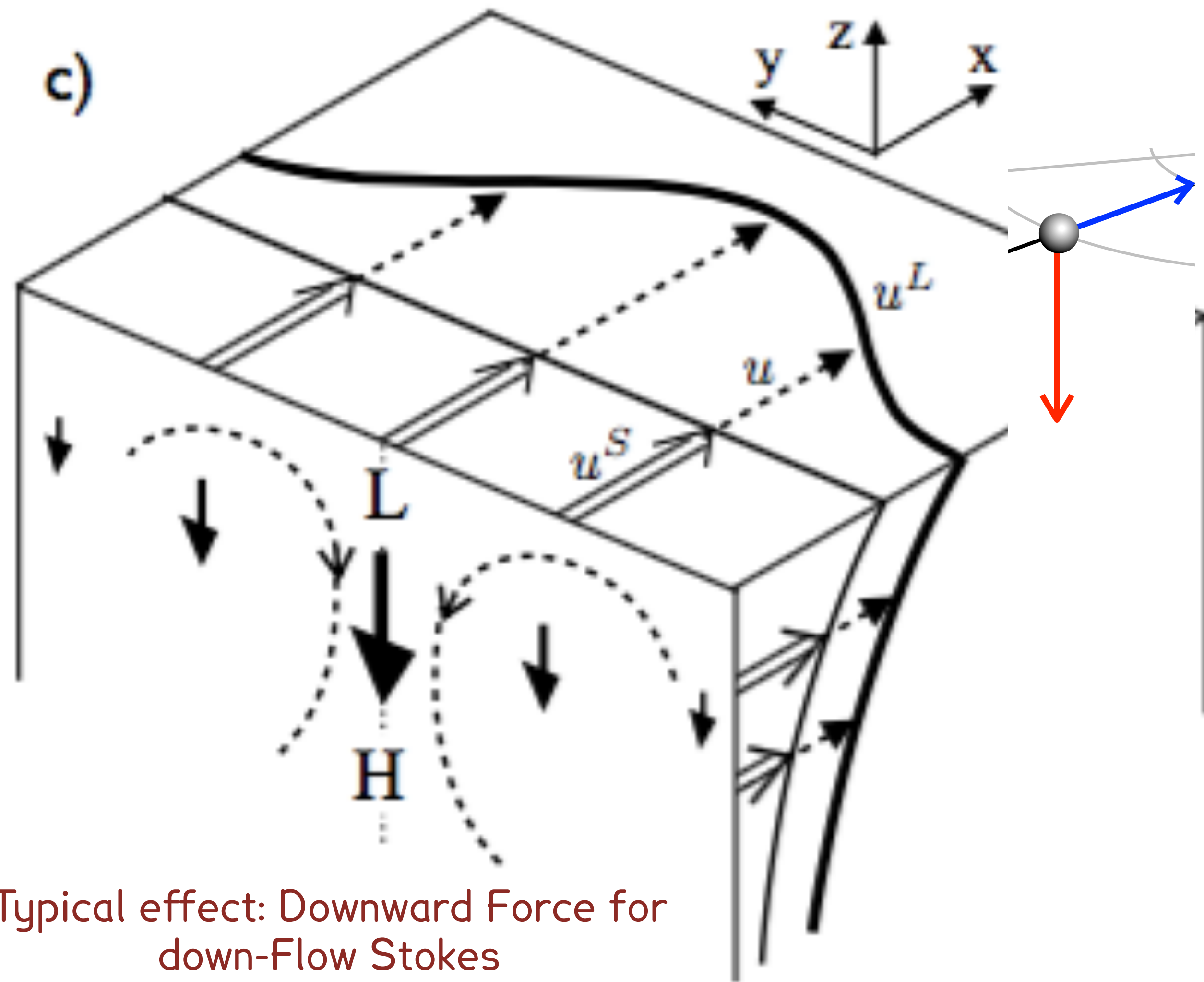
$$\frac{\alpha^2}{Re} \left[ w_{,t} + v_j^L w_{,j} + \frac{M_{Ro}}{RoRe} w w_{,z} \right] = -\pi_{,z} + b - \epsilon v_j^L v_{j,z}^s + \frac{\alpha^2}{ReRe} w_{,jj}$$

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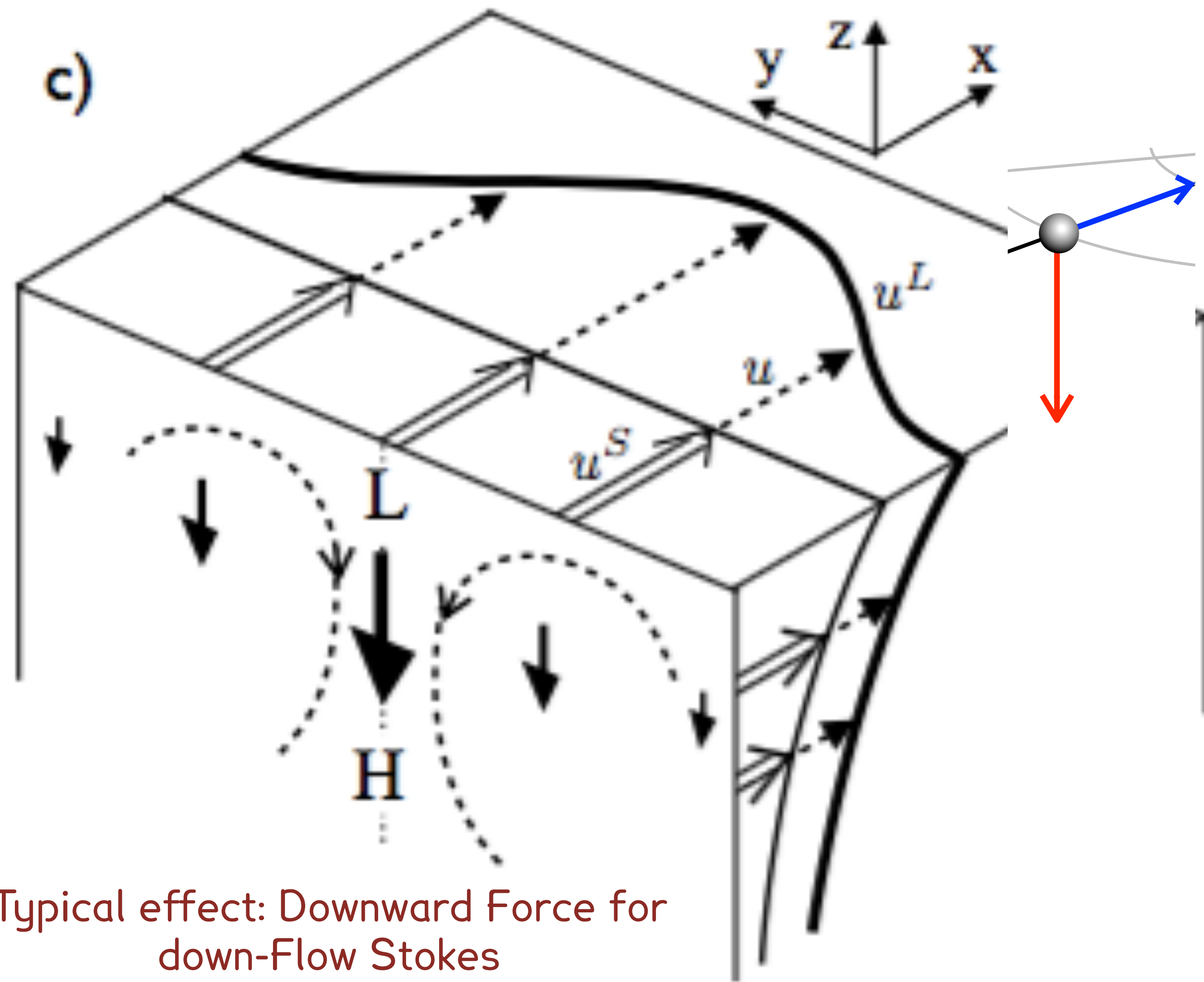
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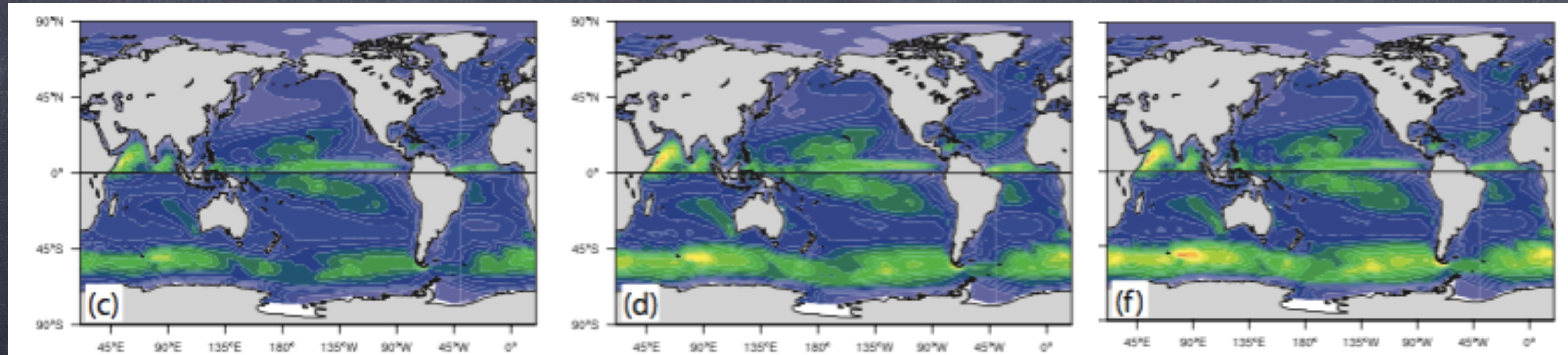
“wavy hydrostatic” if  $\epsilon \gg 1$

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# Langmuir Mixing in Climate: Boundary Layer Depth Improved (RMS error tabulated)

Case	Summer			Winter			
	Global	South of 30°S	30°S-30°N	Global	South of 30°S	30°S-30°N	
Control	CTRL	10.28 ± 0.29	16.00 ± 0.48	6.57 ± 0.23	50.24 ± 1.42	52.52 ± 0.54	15.89 ± 0.33
Old Scheme Mixing	VR12-MA	9.31 ± 0.28	10.64 ± 0.49	9.60 ± 0.33	47.65 ± 1.15	48.47 ± 0.49	22.98 ± 0.42
Bad Entrain.	VR12-EN	11.65 ± 0.29	11.91 ± 0.83	12.79 ± 0.39	56.85 ± 0.93	61.30 ± 1.21	33.60 ± 0.55
New Entrain.	LF17	8.48 ± 0.24	8.92 ± 0.39	9.15 ± 0.30	47.78 ± 1.08	49.98 ± 0.77	22.43 ± 0.43



CTRL (No Lang.)

Mixing w/o Entrain Eval.

Mixing & Refined Entrainment

L. P. Van Roekel, BFK, P. P. Sullivan, P. E. Hamlington, and S. R. Haney. The form and orientation of Langmuir cells for misaligned winds and waves. *Journal of Geophysical Research-Oceans*, 117:C05001, 22pp, May 2012.

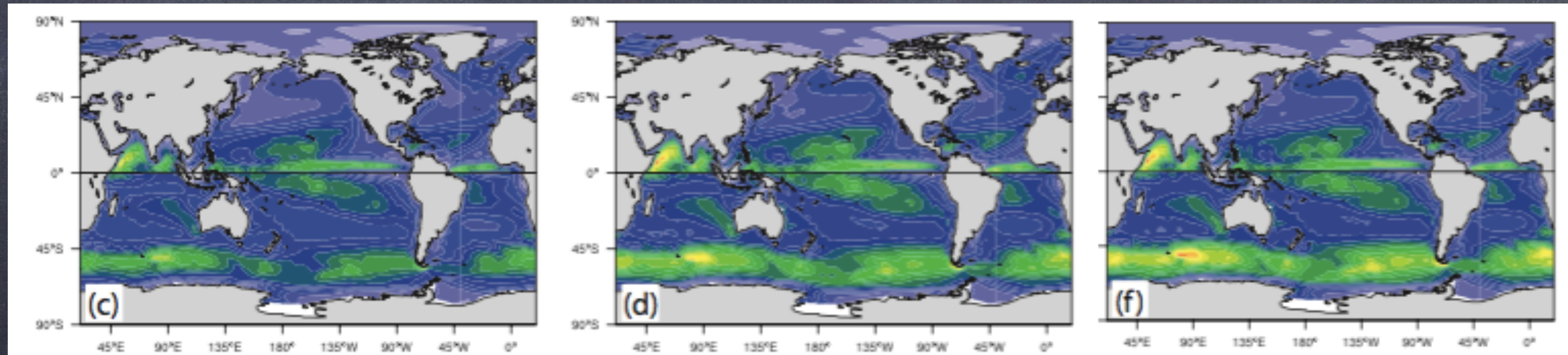
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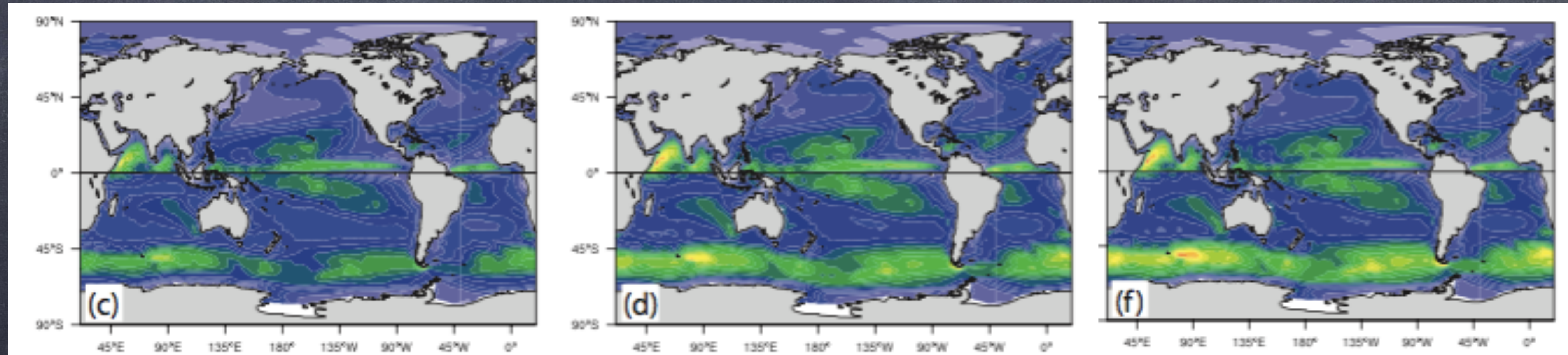
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Old Scheme Mixing	9.31 ± 0.28	10.64 ± 0.49	9.60 ± 0.33	47.65 ± 1.15	48.47 ± 0.49	22.98 ± 0.42
Bad Entrain.	11.65 ± 0.29	11.91 ± 0.83	12.79 ± 0.39	56.85 ± 0.93	61.30 ± 1.21	33.60 ± 0.55
New Entrain.	8.48 ± 0.24	8.92 ± 0.39	9.15 ± 0.30	47.78 ± 1.08	49.98 ± 0.77	22.43 ± 0.43



CTRL (No Lang.)

Mixing w/o Entrain Eval.

Mixing & Refined Entrainment

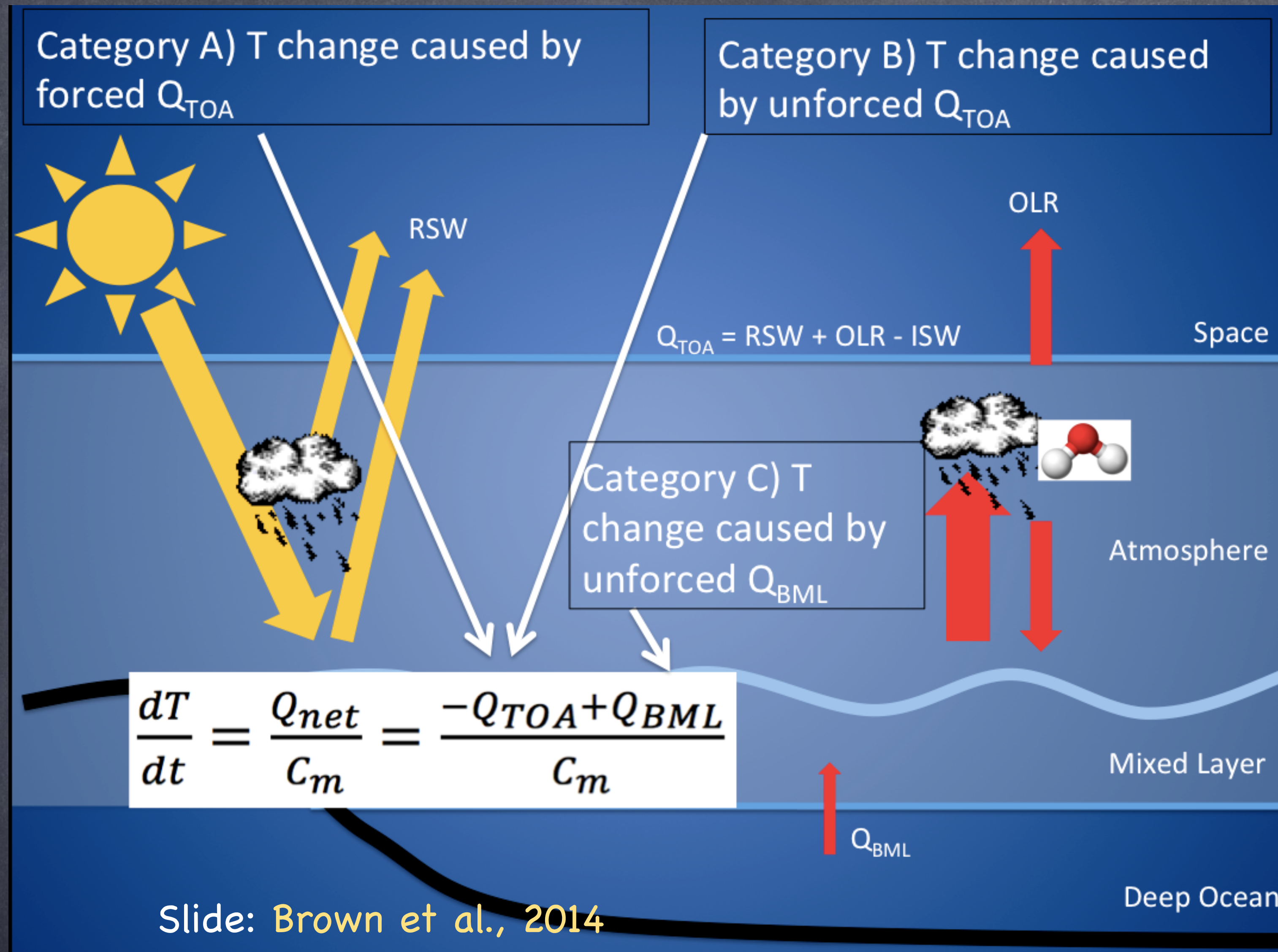
L. P. Van Roekel, BFK, P. P. Sullivan, P. E. Hamlington, and S. R. Haney. The form and orientation of Langmuir cells for misaligned winds and waves. *Journal of Geophysical Research-Oceans*, 117:C05001, 22pp, May 2012.

Q. Li, A. Webb, BFK, A. Craig, G. Danabasoglu, W. G. Large, and M. Vertenstein. Langmuir mixing effects on global climate: WAVEWATCH III in CESM. *Ocean Modelling*, 103:145-160, July 2016.

Q. Li & BFK. Assessing the effects of Langmuir turbulence on the entrainment buoyancy flux in the ocean surface boundary layer. *Journal of Physical Oceanography*, 47:2863-2886, December 2017.

Q. Li and B. Fox-Kemper. Anisotropy of Langmuir turbulence and the Langmuir-enhanced mixed layer entrainment. *Physical Review Fluids*, 5:013803, January 2020.

# GMST: Surface Energy Budget=Ocean Heat Content Budget



So, we've improved the mixed layer, shouldn't it improve the climate projections somehow?

- 3.4m of ocean has heat capacity of whole atmosphere
- Ocean Mixed Layer is about 100m deep.



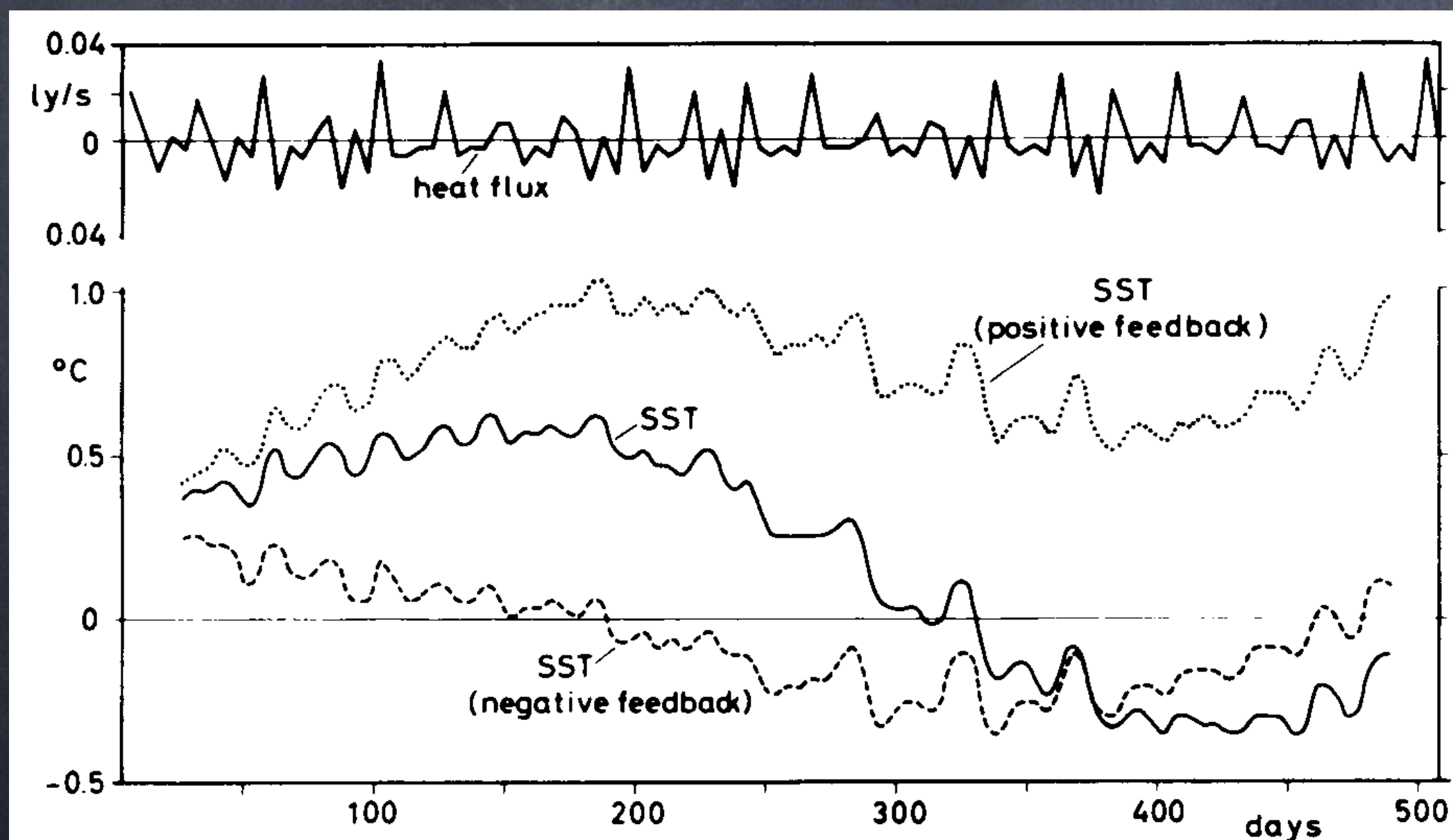
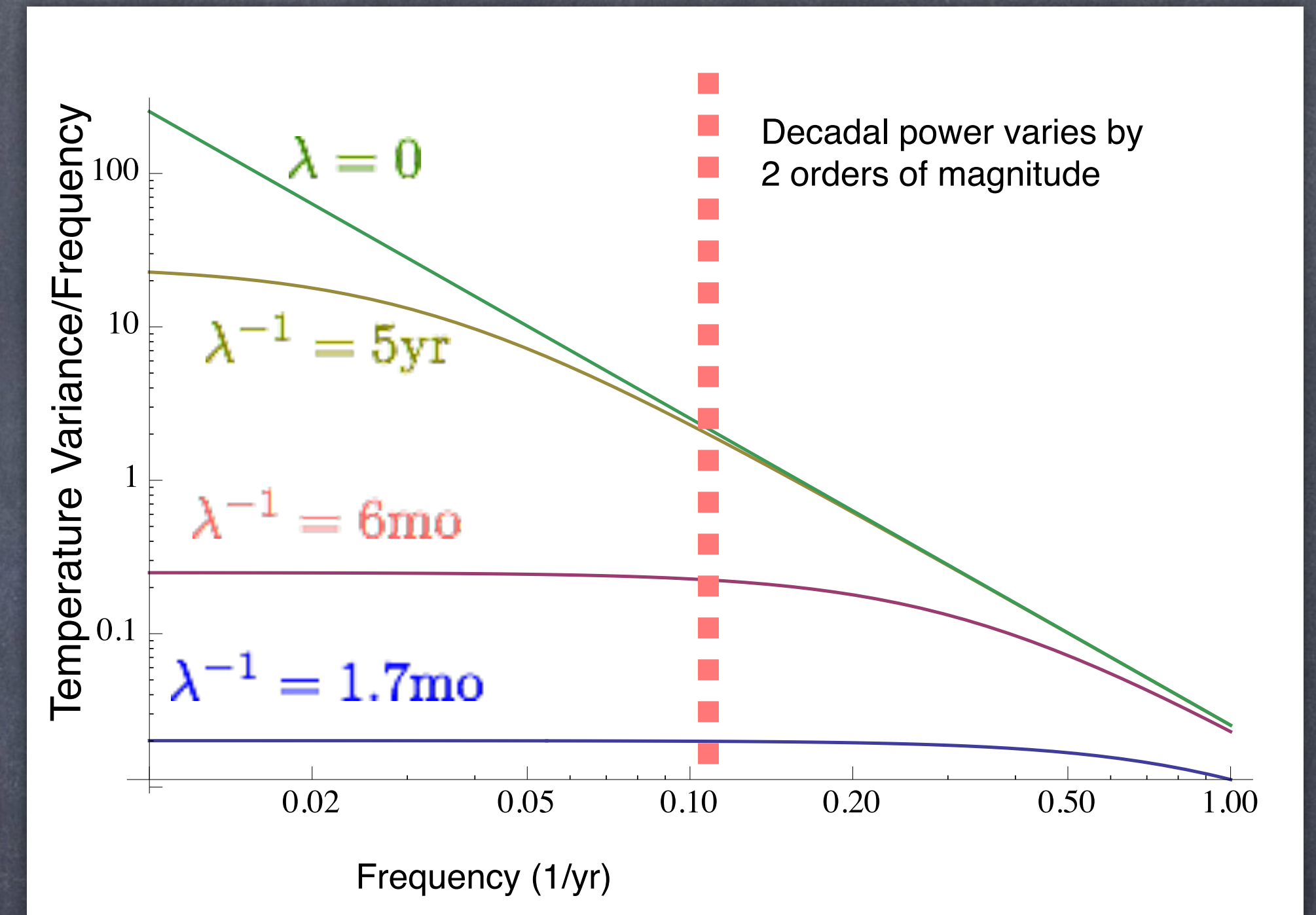
# Modeling of variability



A stochastic, predictable persistence model:  
Frankignoul & Hasselmann (77)

$$\frac{dT}{dt} = \frac{f_1}{h} - \lambda T$$

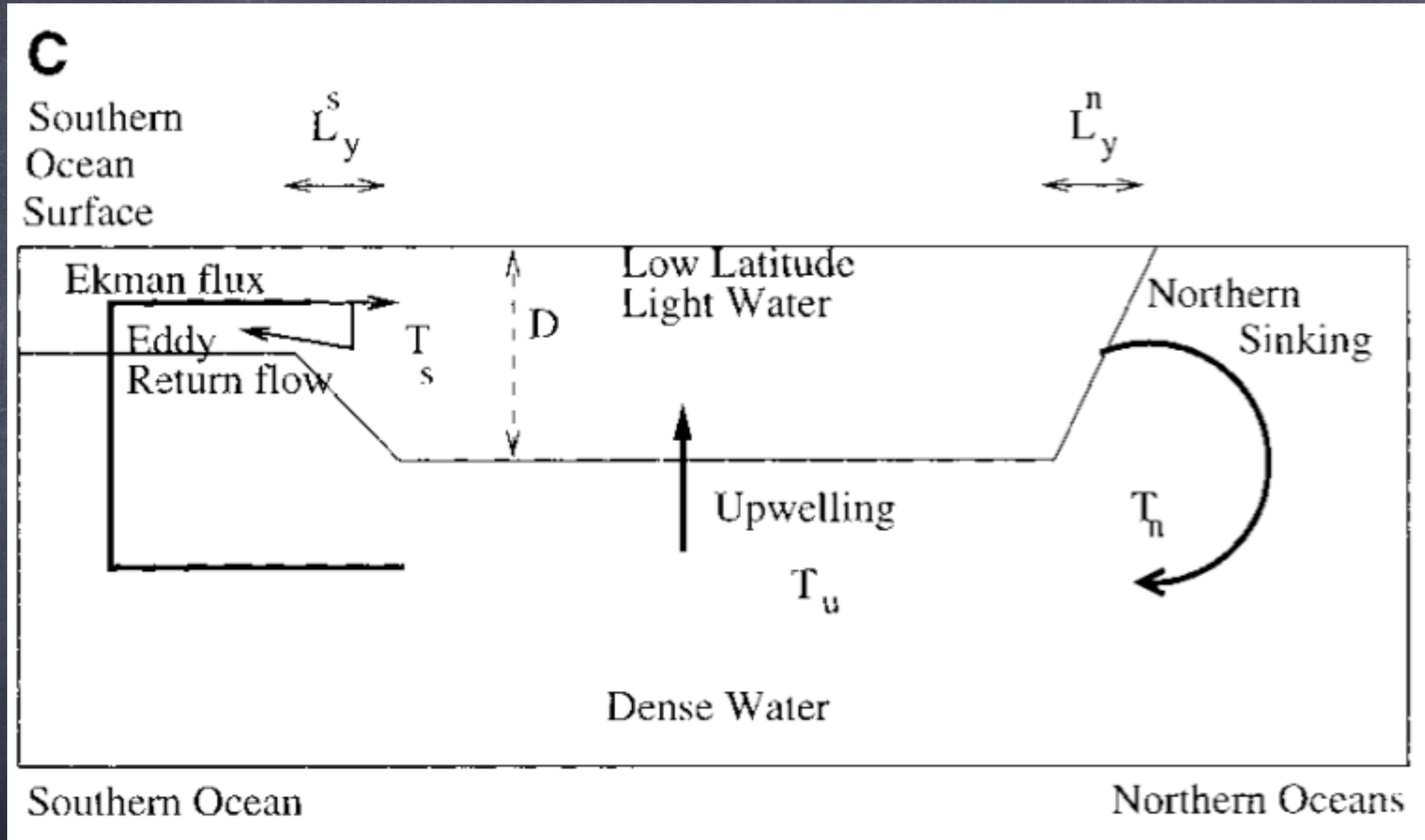
Temp Change      Random Atmosphere      Restoring  
*h* Mixed Layer



$$\lambda = \rho^a C_p^a (\rho^w C_p^w)^{-1} C_H (1 + B) \langle |U| \rangle h^{-1}$$

$$= (1.7 \text{ month})^{-1}$$

# World Ocean



Gnanadesikan

SCIENCE VOL 283 26 MARCH 1999

Even Simpler: 2-Layer Homogeneous Energy Balance Model (Gregory, 2000)

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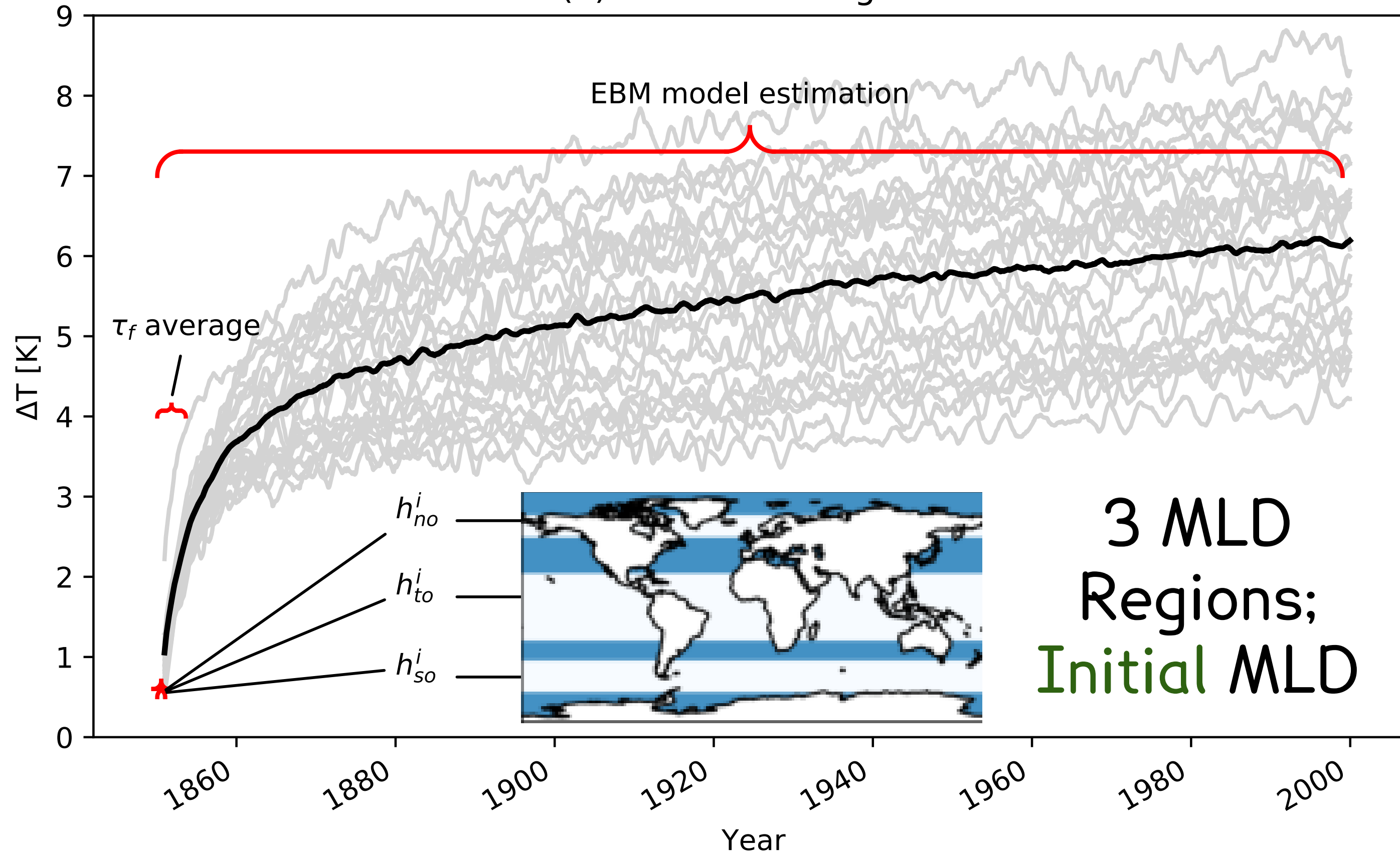
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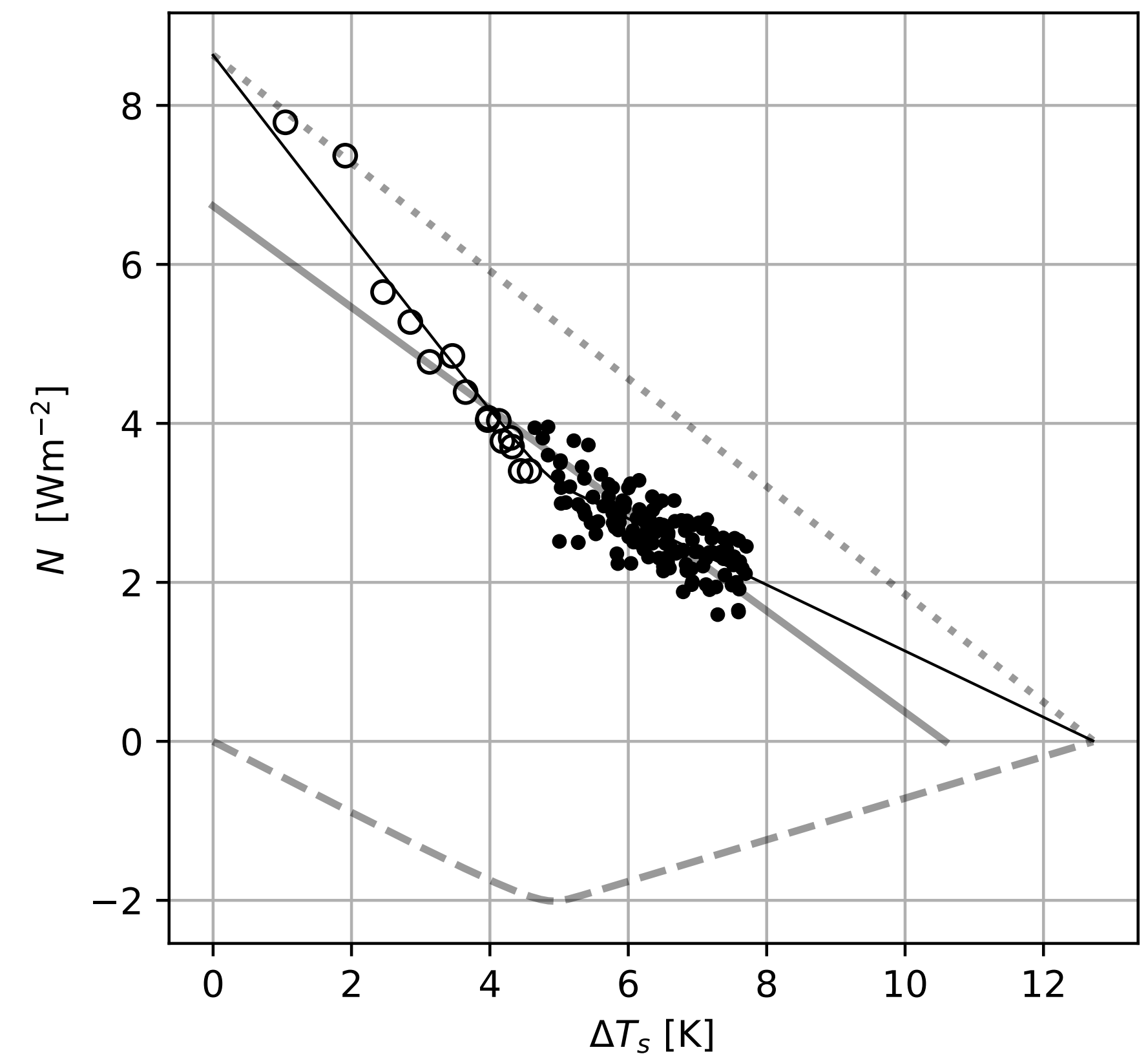
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(a) Timescale diagram



(b) Example EBM- $\varepsilon$  fit: CESM2





# We can use the mixed layer depth as An EMERGENT CONSTRAINT

That is, an observable that correlates or constrains the other properties useful for projections.

## CMIP6 GCMs give MLD & 2-Layer Model Properties

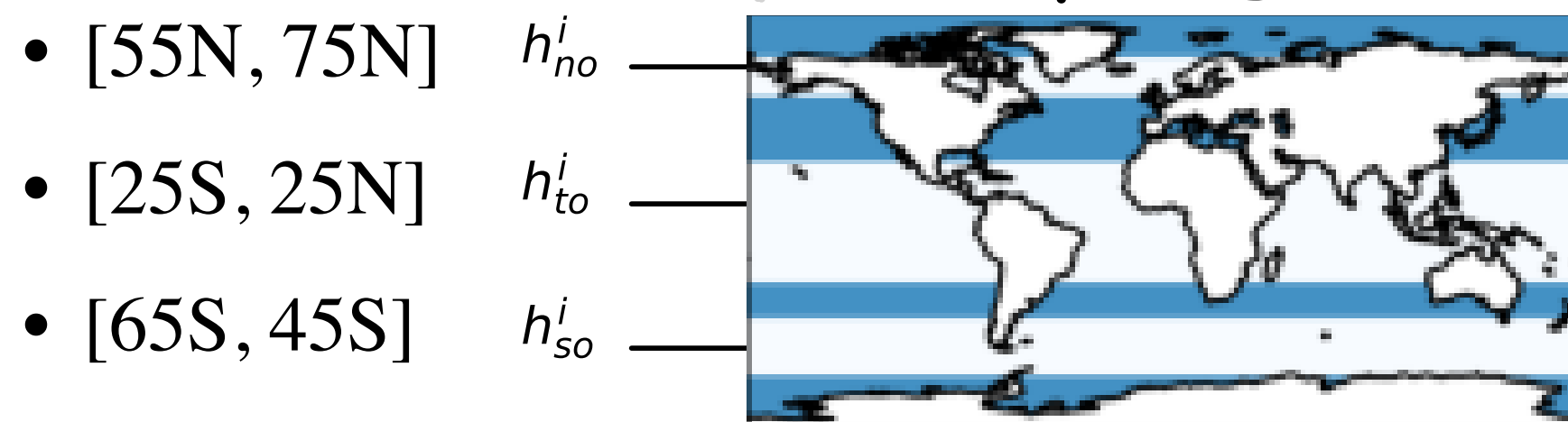
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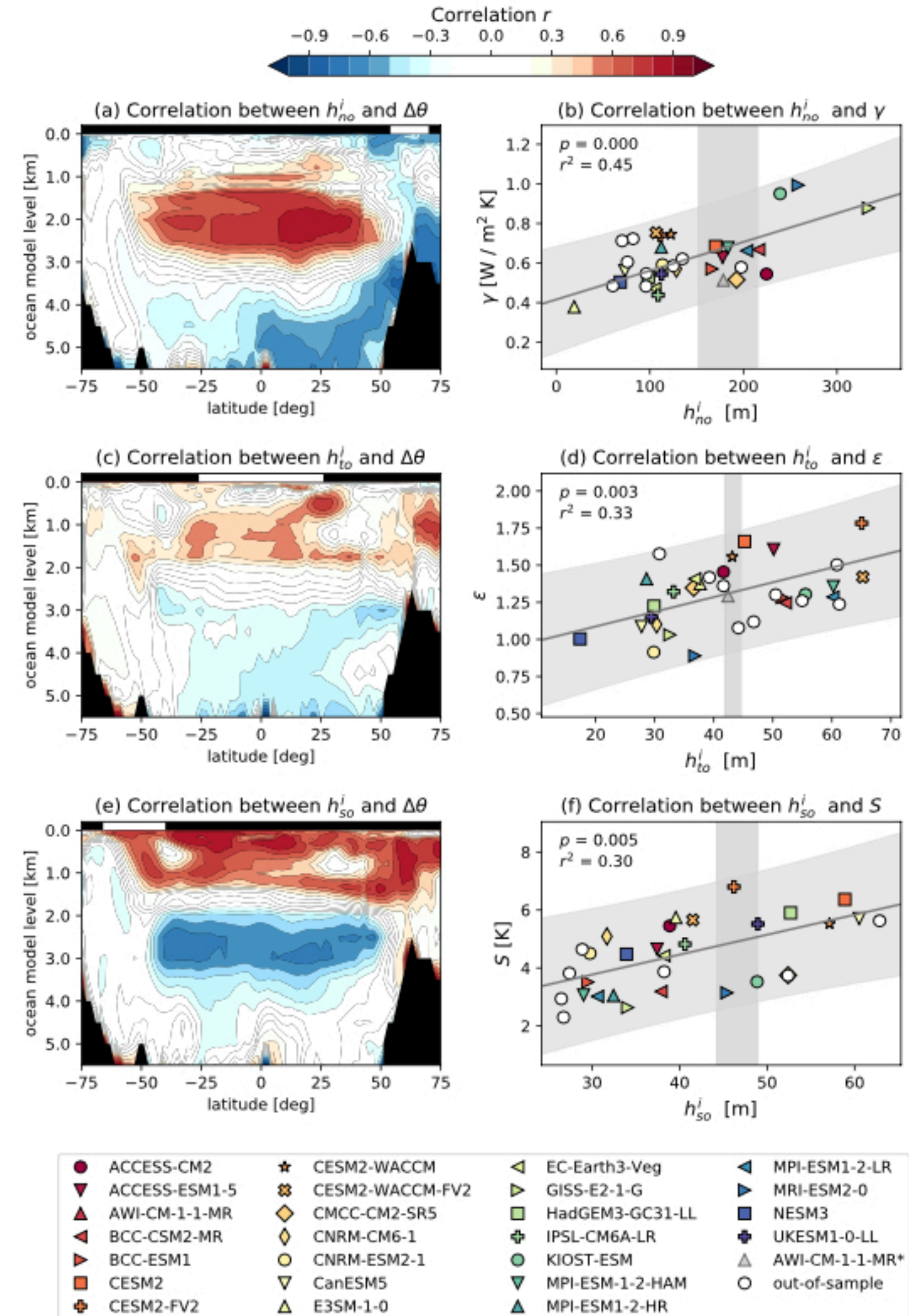
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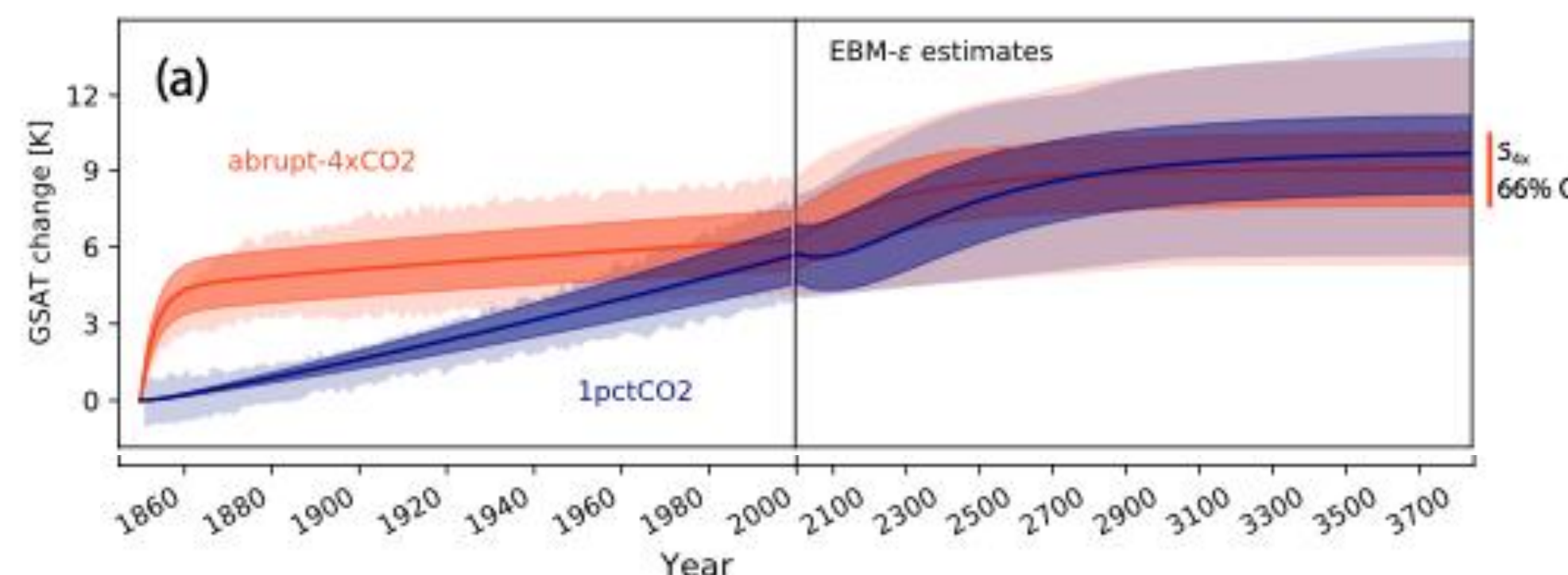
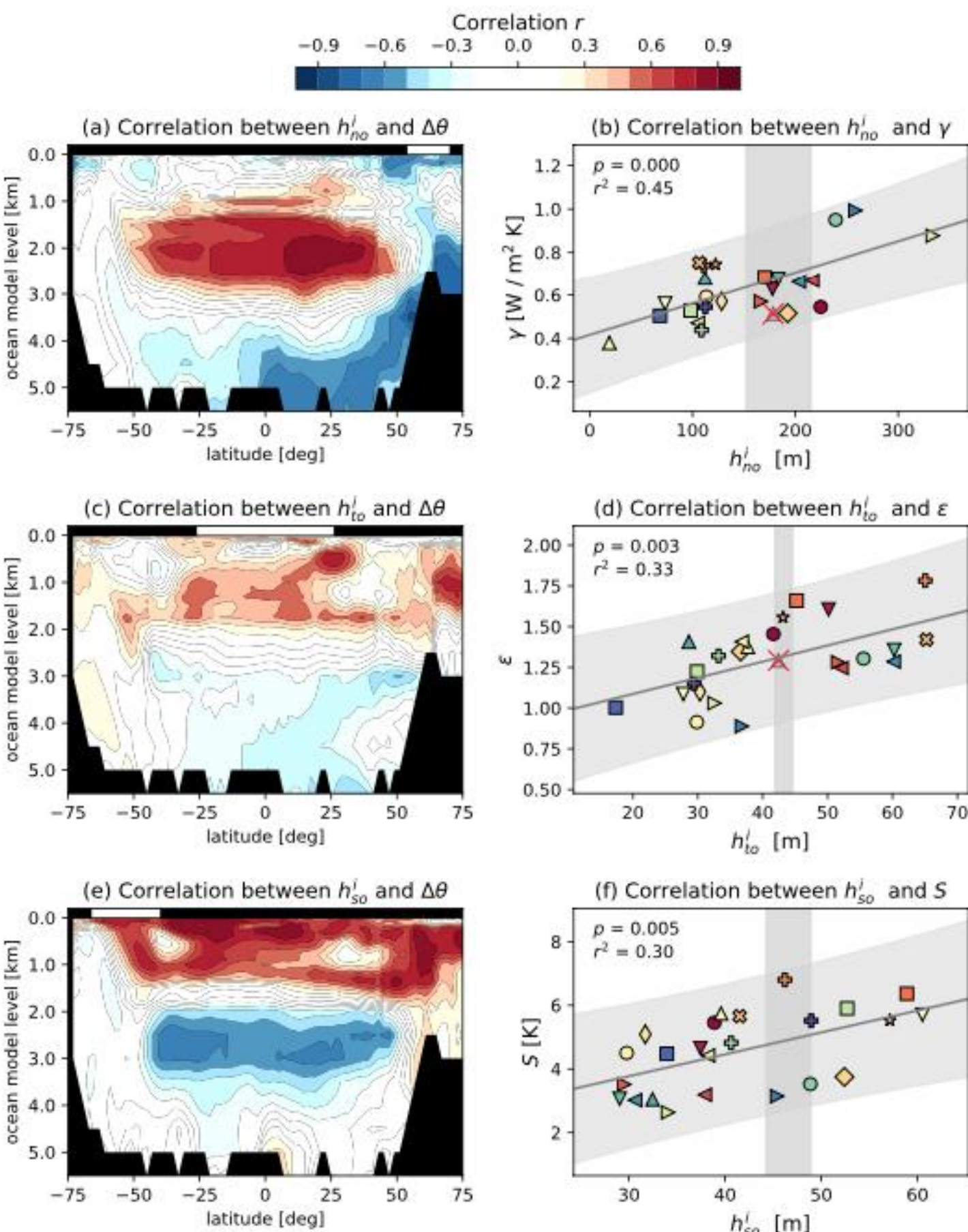


3 MLD Regions; Initial MLD



# Galen Hall: Physics Honors Thesis, May 2020

Regional mixed layer depth as a climate diagnostic and emergent constraint



$$C_S \frac{d\Delta T}{dt} = F - \lambda \Delta T - \epsilon \gamma (\Delta T - \Delta T_D),$$

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“Using these correlations and observations from the Argo float network, we revise the ensemble mean and narrow the 66% range of equilibrium climate sensitivity (ECS) for the particular CMIP6 model collection from 4.51 (3.13–5.71) °C, to 4.66 (3.88–5.43) °C, amounting to a 40% reduction in the span of the uncertainty range.”

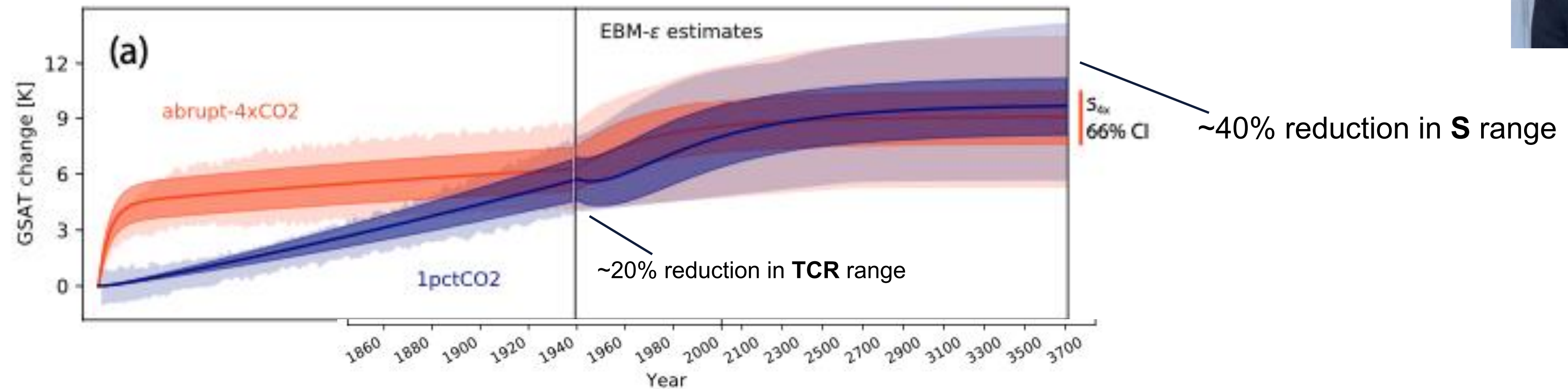
- ACCESS-CM2
- ACCESS-ESM1-5
- BCC-CSM2-MR
- BCC-ESM1
- CESM2
- CESM2-FV2
- CESM2-WACCM
- CESM2-WACCM-FV2
- CMCC-CM2-SR5
- CNRM-CM6-1
- CNRM-ESM2-1
- CanESM5
- E3SM-1-0
- EC-Earth3-Veg
- GISS-E2-1-G
- HadGEM3-GC31-LL
- IPSL-CM6A-LR
- KIOST-ESM
- MPI-ESM1-2-HAM
- MPI-ESM1-2-HR
- MPI-ESM1-2-LR
- MRI-ESM2-0
- NESM3
- UKESM1-0-LL
- AWI-CM-1-1-MR\*

Hall & Fox-Kemper, Submitted to *GRL*, 2021

**PANGEO**  
 Emulate CMIP6 model oceans via pangeo.io  
 Understand emulator parameter space

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Regional mixed layer depth as a climate diagnostic and emergent constraint

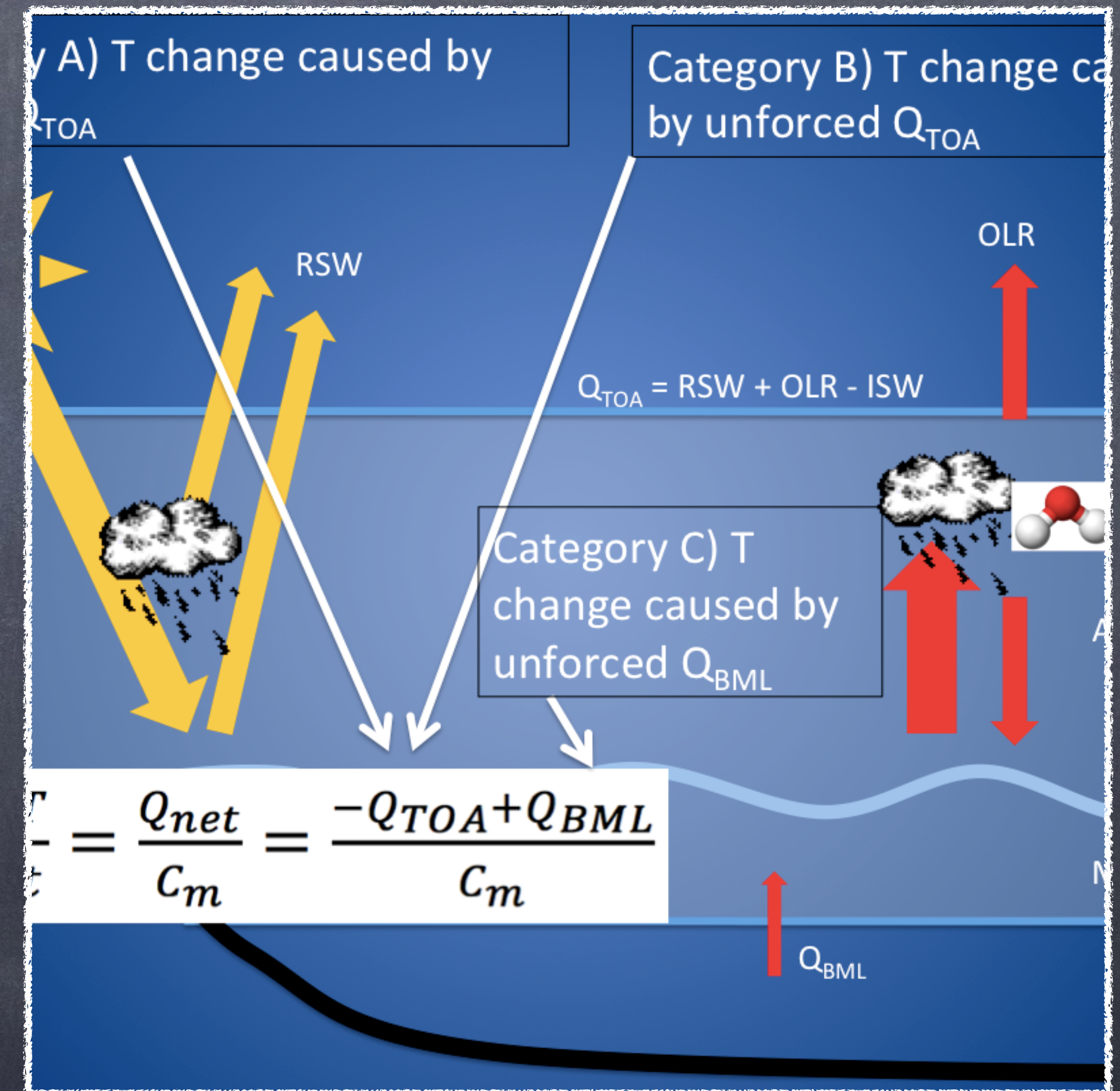


Approximately halving the uncertainty range for [transient climate response] has a net present value of about \$10.3 trillion (year 2005 US\$) if accomplished in time for emissions to be adjusted in 2020, falling to \$9.7 trillion if accomplished by 2030.

-C. Hope, 2015, *Phil. Trans. A.*, <https://doi.org/10.1098/rsta.2014.0429>

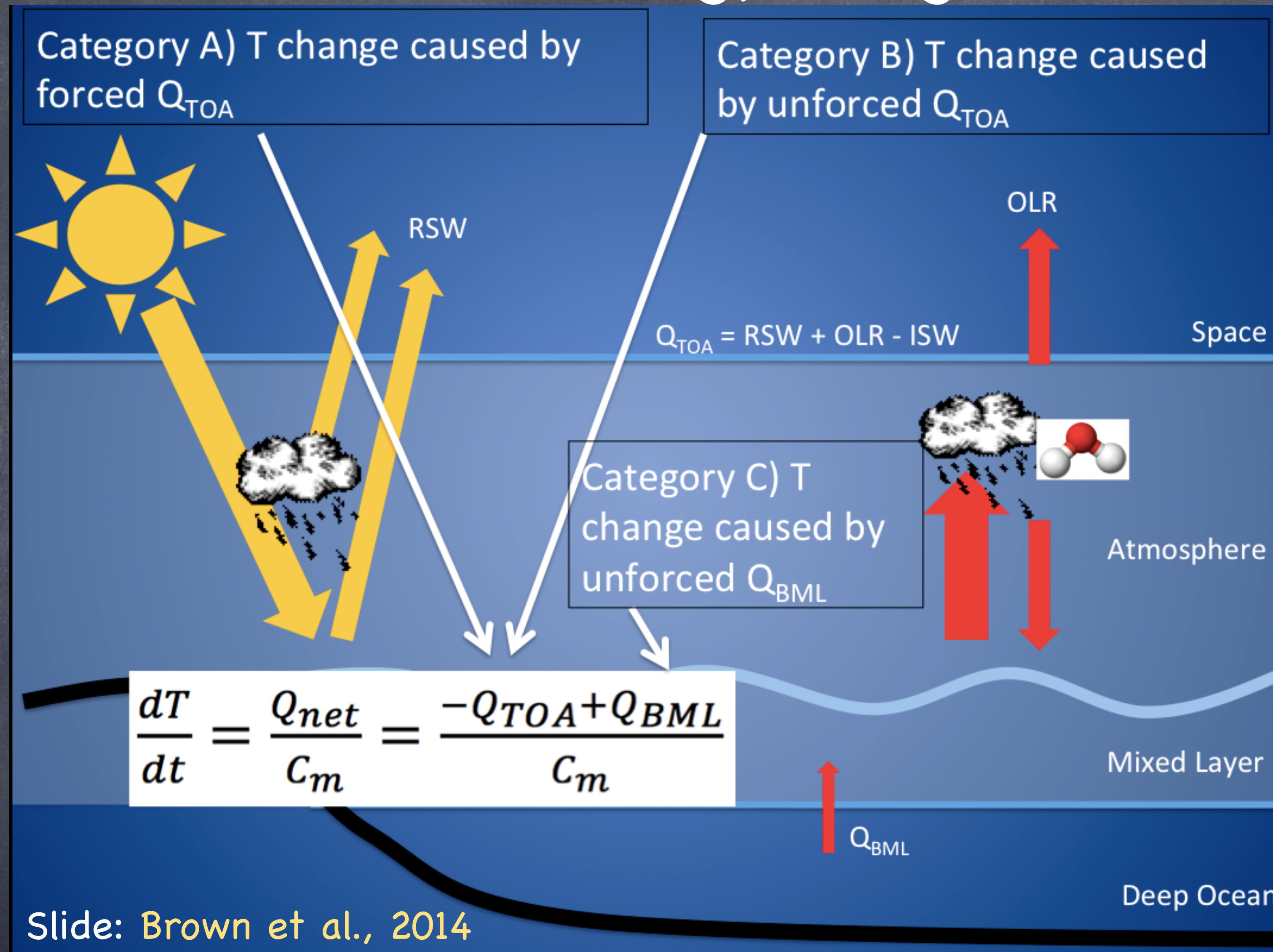
# For Today: Small-scales affect Climate

- Key Ocean Climate Questions
- Large Eddy Simulation Closures
  - Smagorinsky, Leith, QG Leith
- Effects on Global Kinetic Energy
- Submesoscale affects Mixed Layer
- Wave-Driven Turbulence affects Mixed Layer
- Regional Mixed Layer Depth affects Climate Sensitivity



Extras

# Surface Energy Budget



Slide: Brown et al., 2014

- $O(2W/m^2)$  change to  $Q_{BML}$  as important as GHG
- Slight oversimplification—sensitivity + budget



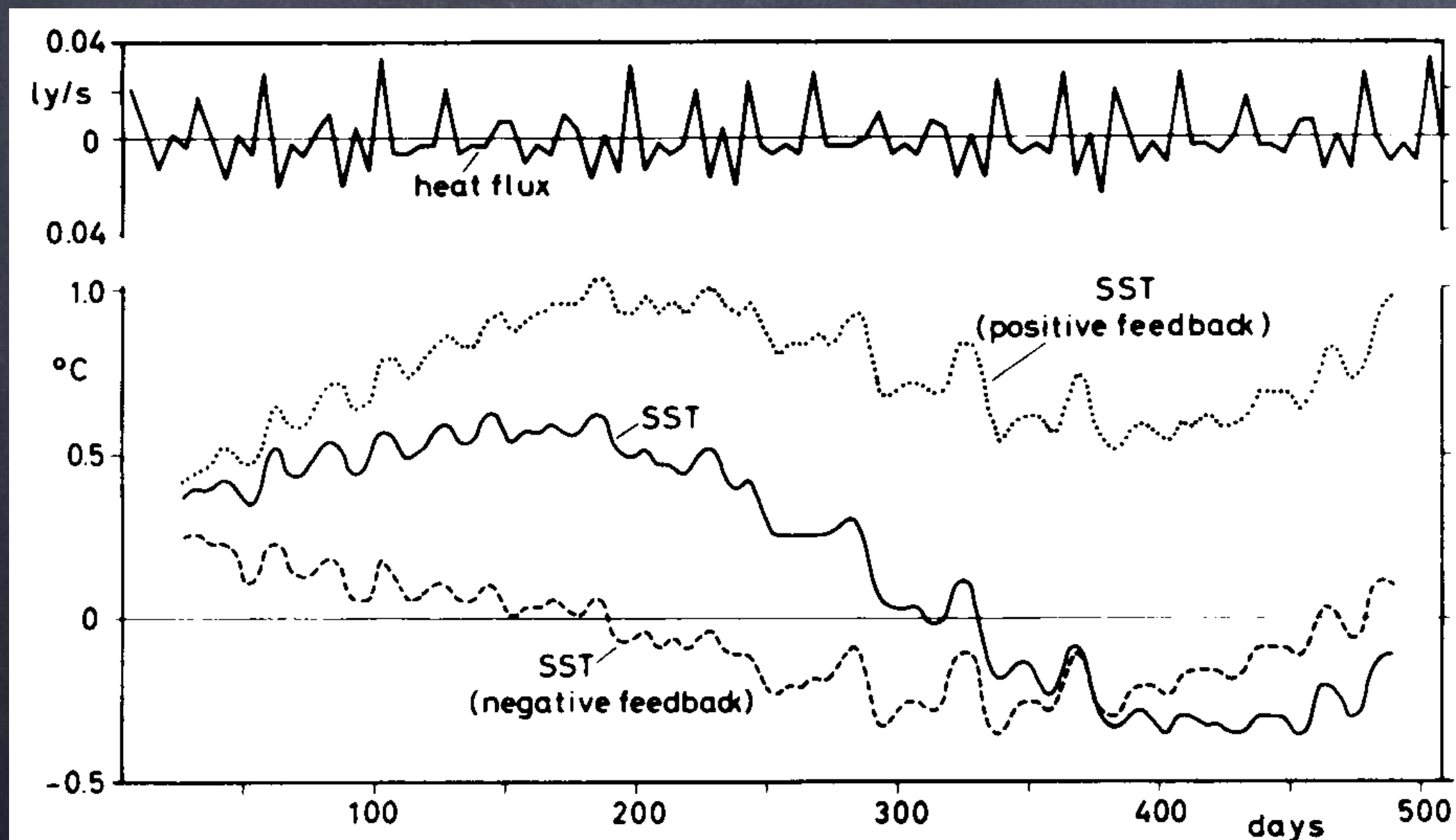
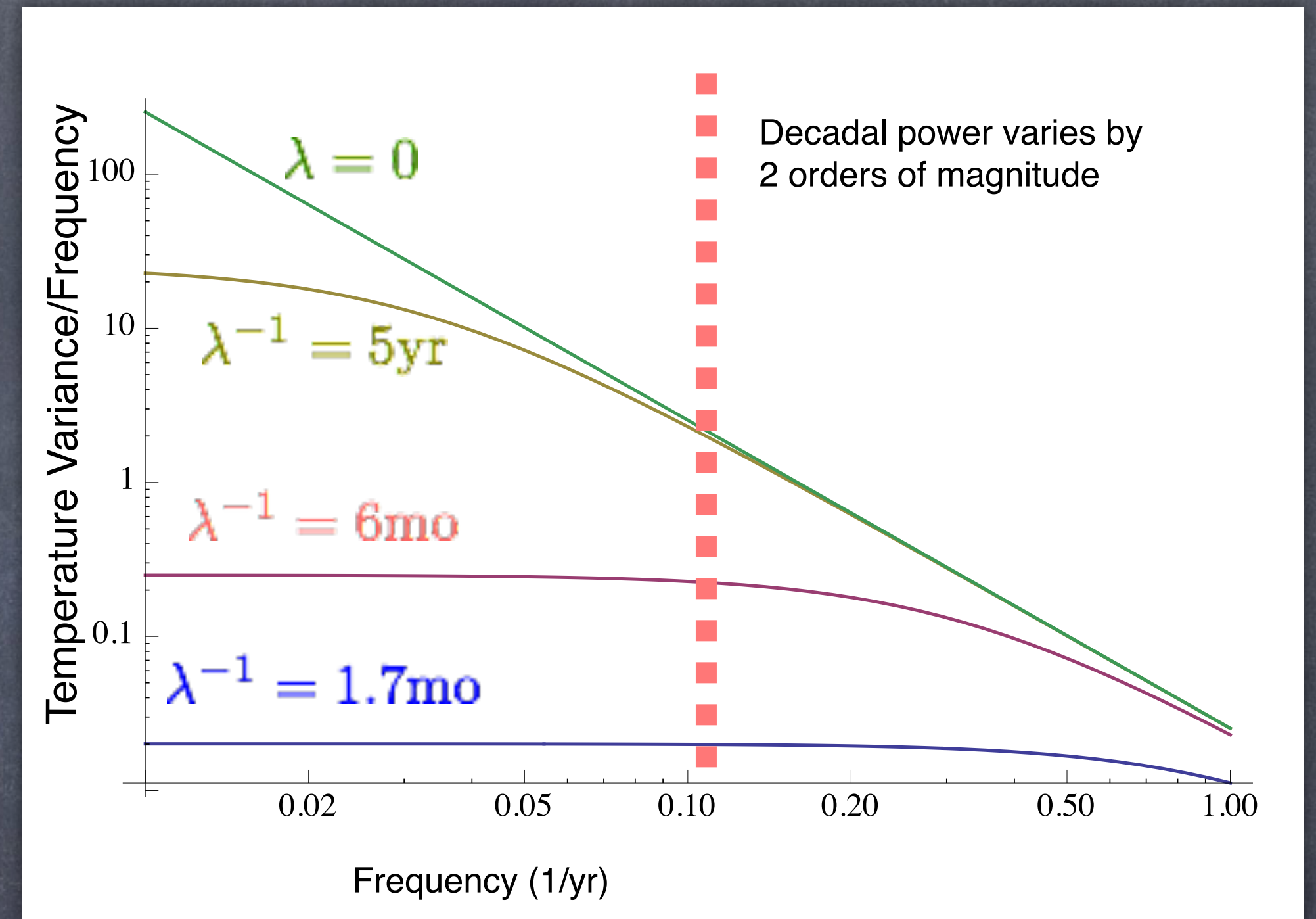
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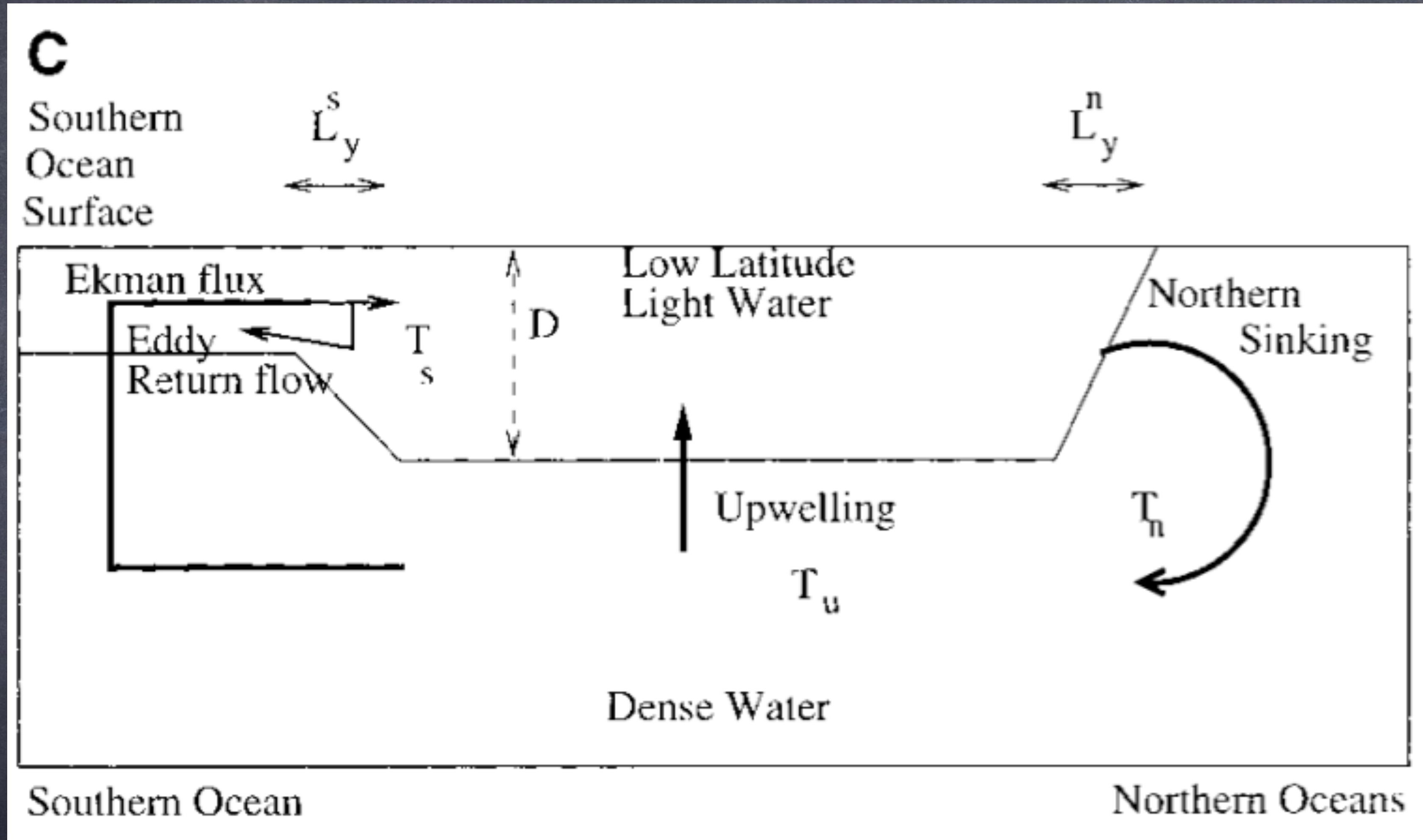
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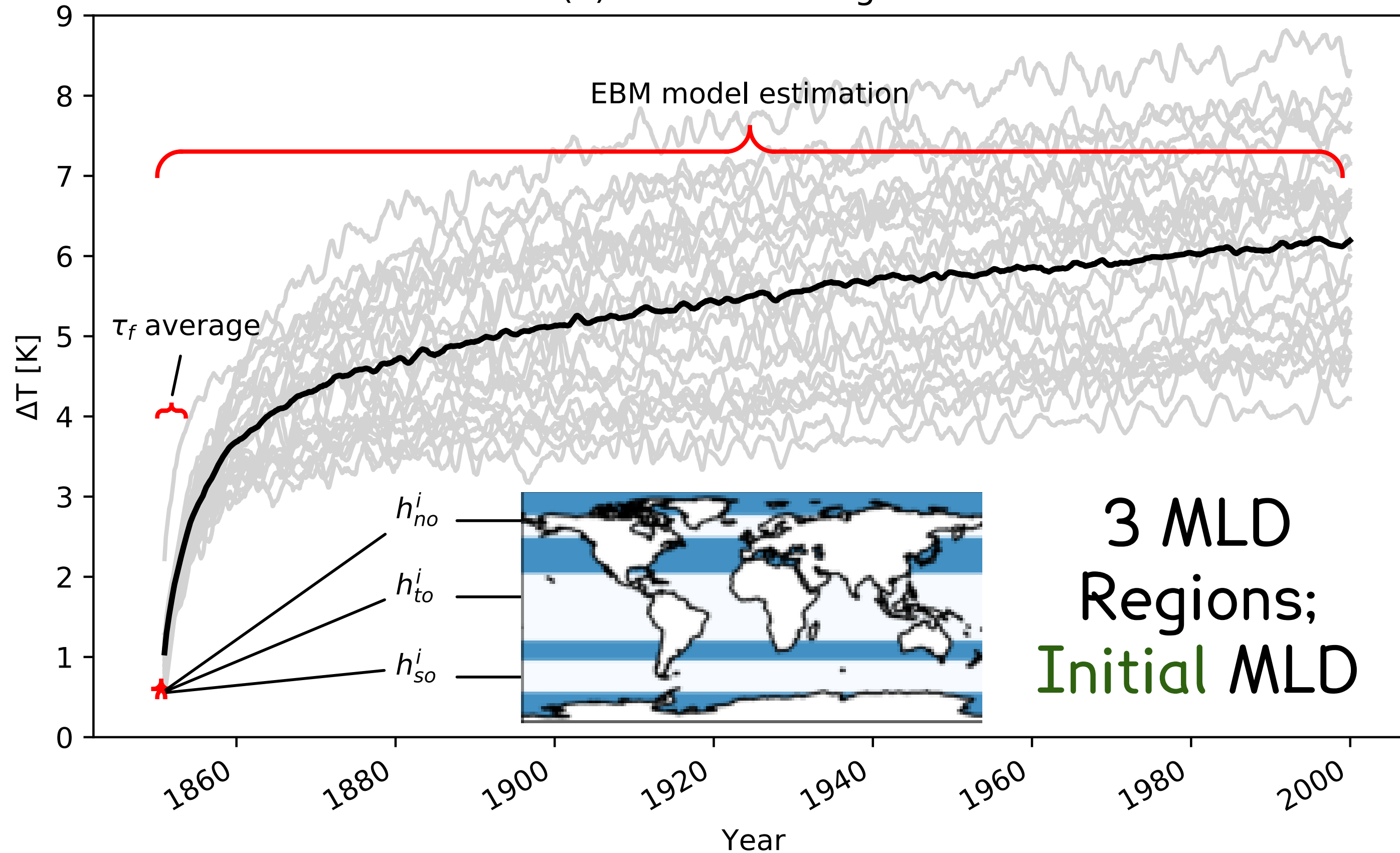
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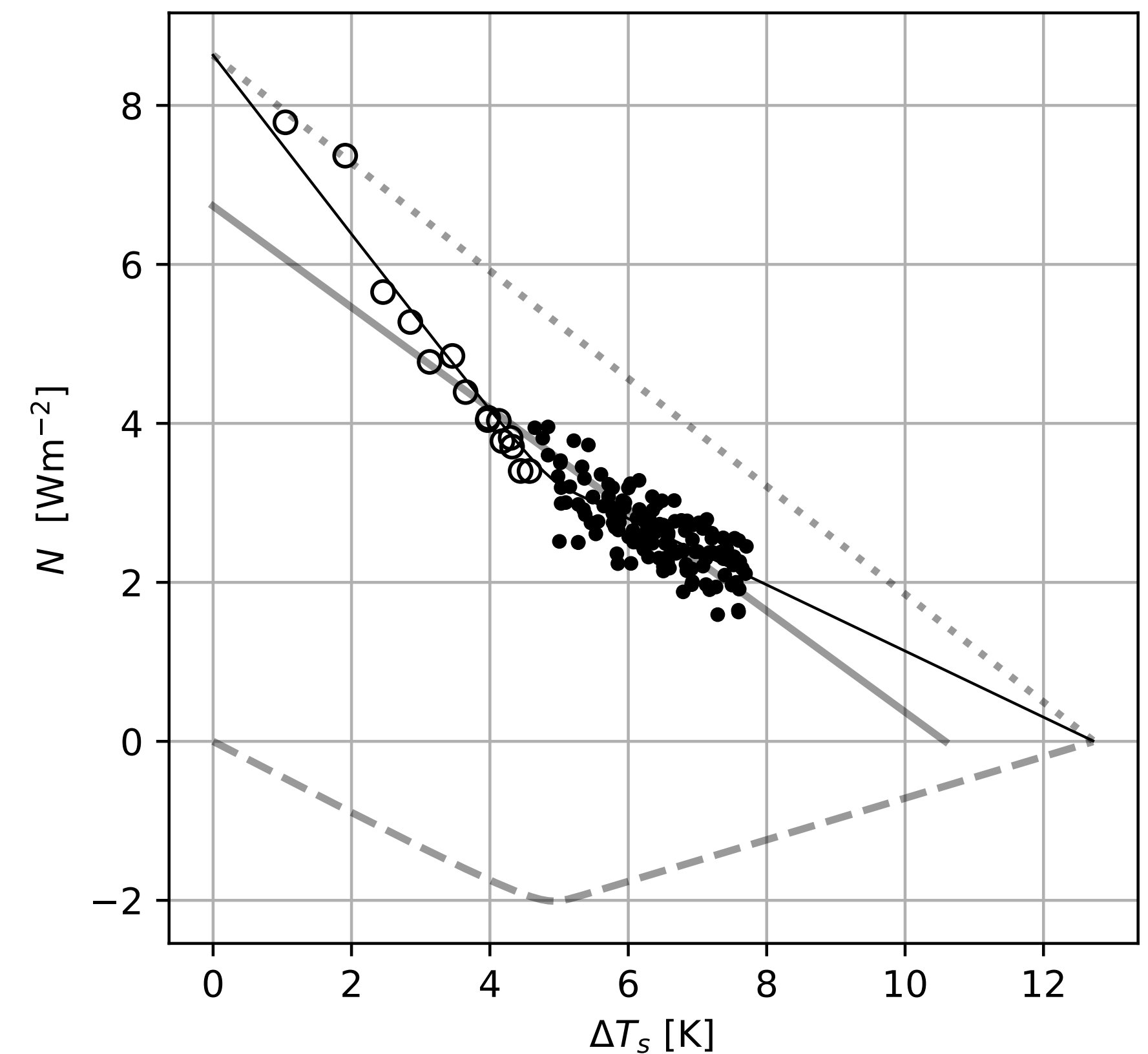
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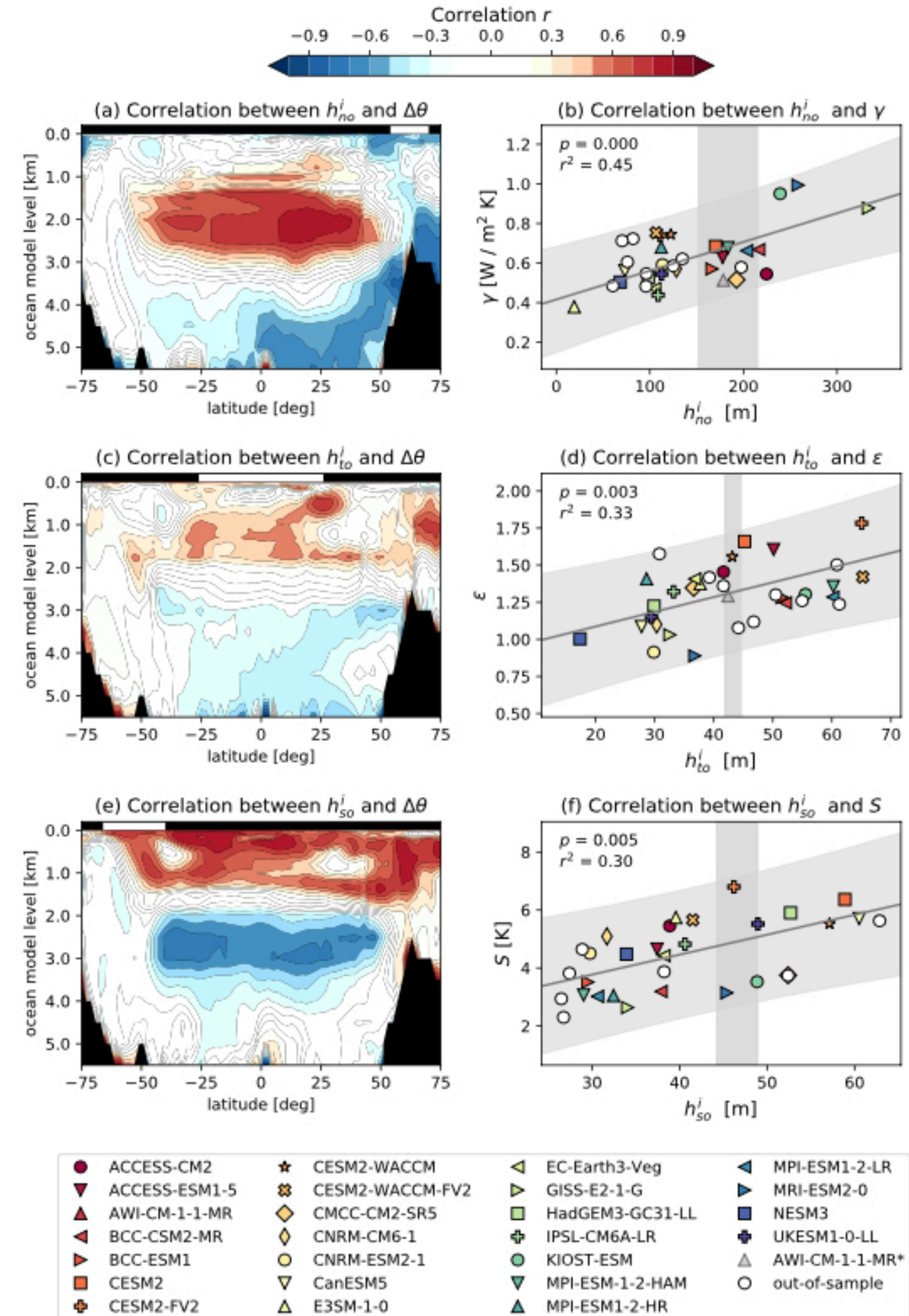
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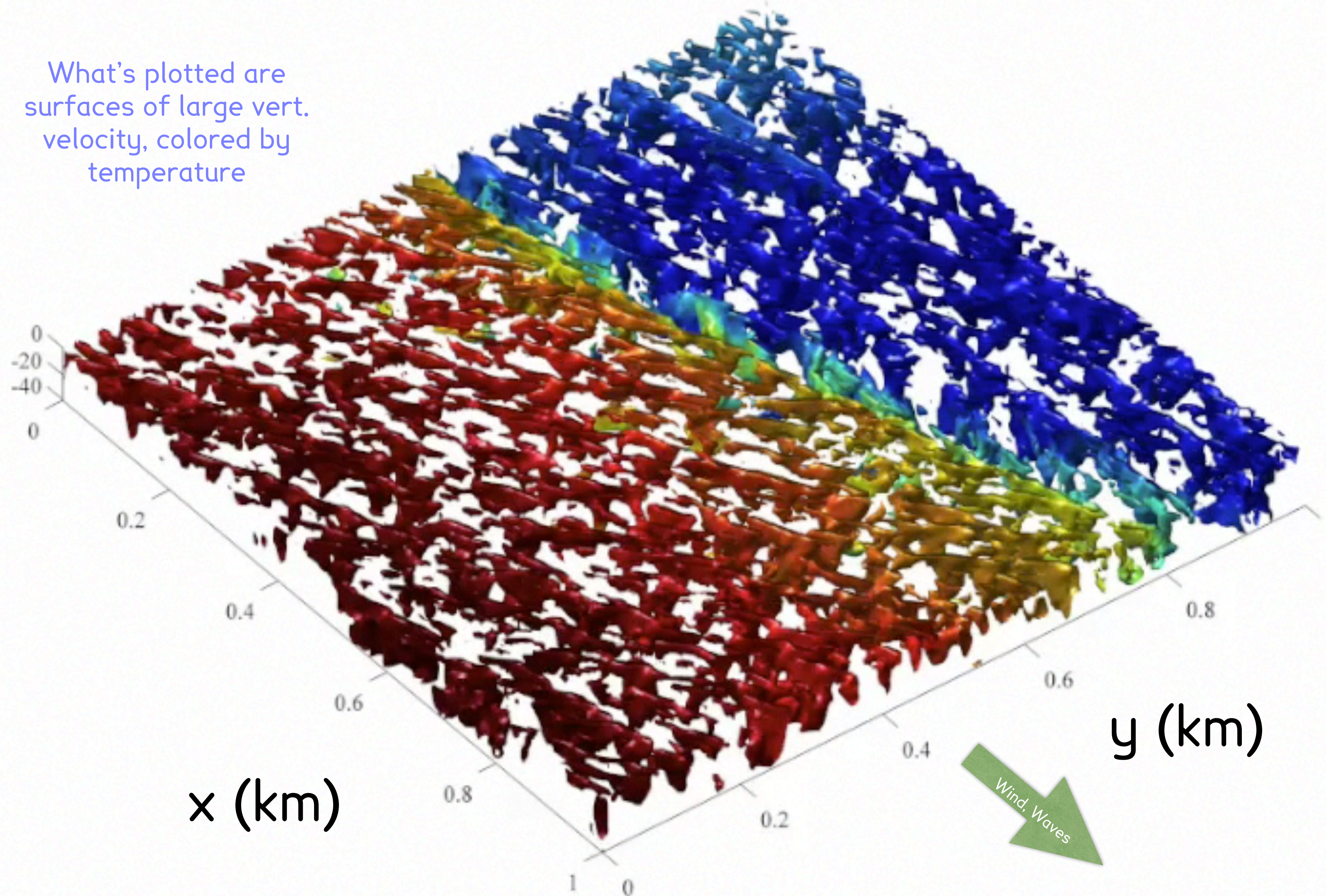
- [55N, 75N]  $h_{no}^i$
  - [25S, 25N]  $h_{to}^i$
  - [65S, 45S]  $h_{so}^i$
- 

3 MLD Regions;  
Initial MLD



# Zoom: Submeso-Langmuir Interaction!

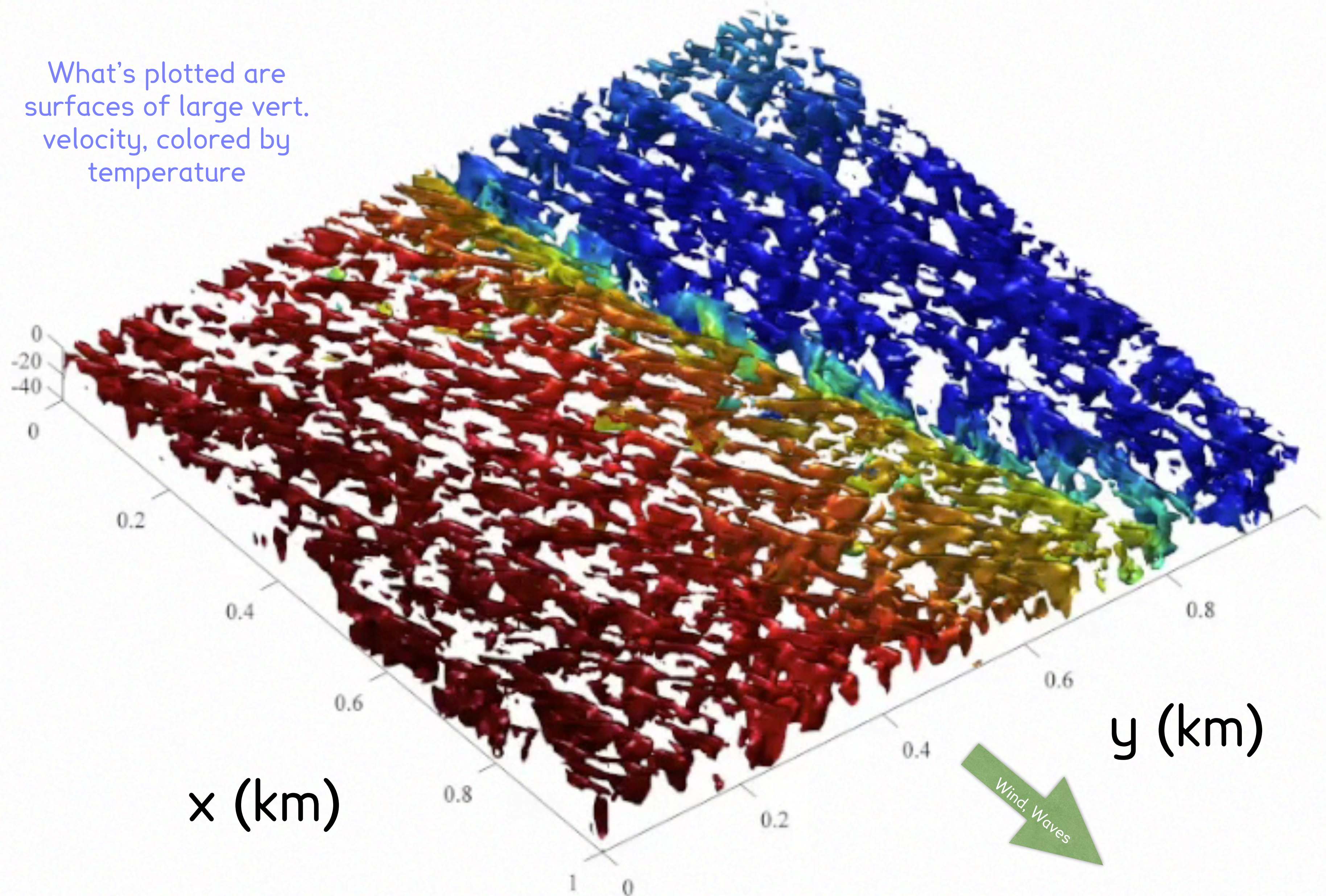
What's plotted are surfaces of large vert. velocity, colored by temperature



P. E. Hamlington, L. P. Van Roekel, BFK, K. Julien, and G. P. Chini. Langmuir-submesoscale interactions: Descriptive analysis of multiscale frontal spin-down simulations. *JPO*, 44(9):2249-2272, 2014.

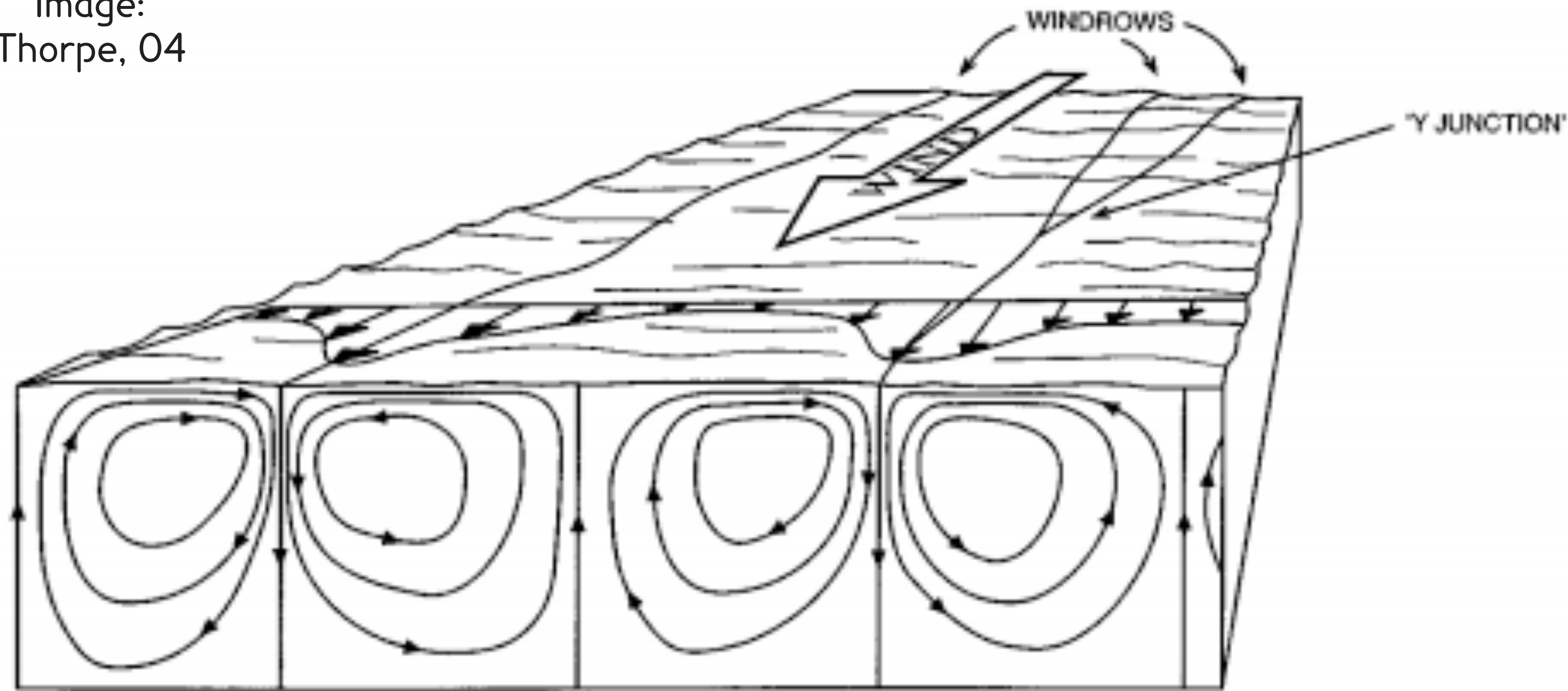
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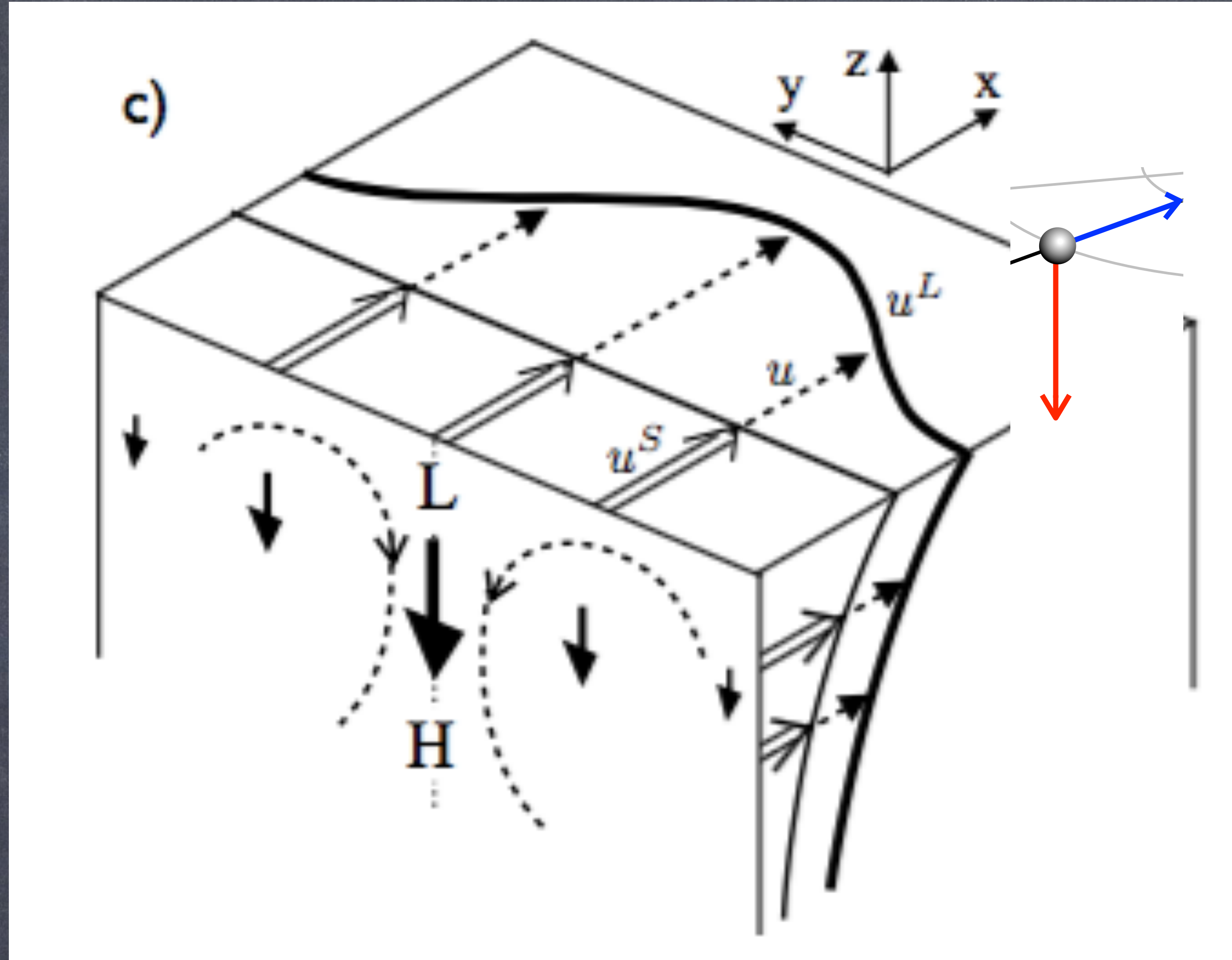
image:  
Thorpe, 04



**Figure 1** Sketch showing the pattern of mean flow in idealized Langmuir circulation. The windrows may be 2 m to 300 m apart, and the cell form is roughly square (as shown). In practice the flow is turbulent, especially near the water surface, and the windrows (Figure 2) amalgamate and meander in space and time. Bands of bubbles or buoyant algae may form within the downward-going (or downwelling) flow (see Figure 3).

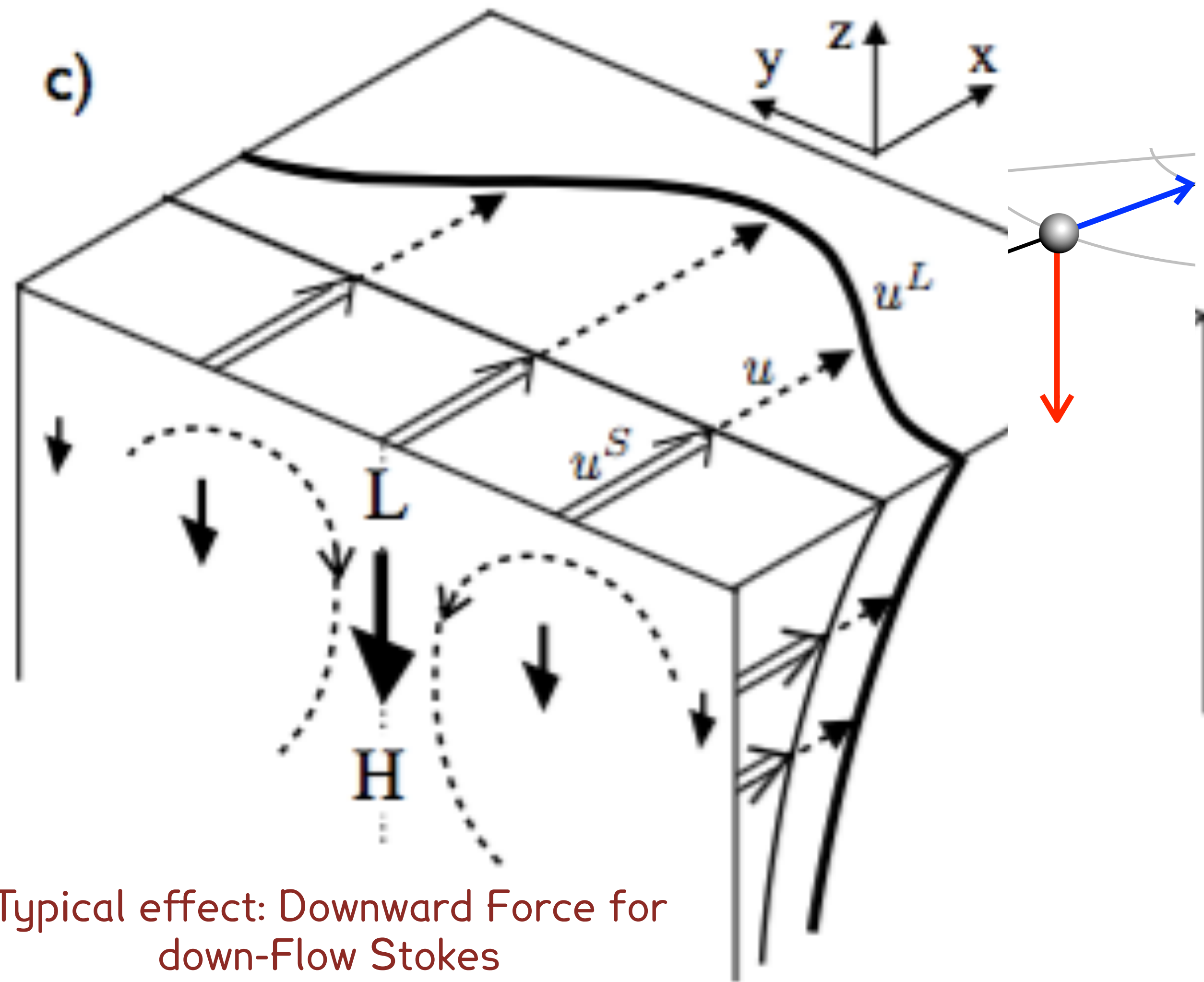
$$\frac{\alpha^2}{Re} \left[ w_{,t} + v_j^L w_{,j} + \frac{M_{Ro}}{RoRe} w w_{,z} \right] = -\pi_{,z} + b - \epsilon v_j^L v_{j,z}^s + \frac{\alpha^2}{ReRe} w_{,jj}$$

N. Suzuki and BFK. Understanding Stokes forces in the wave-averaged equations. *Journal of Geophysical Research-Oceans*, 121:1-18, 2016.



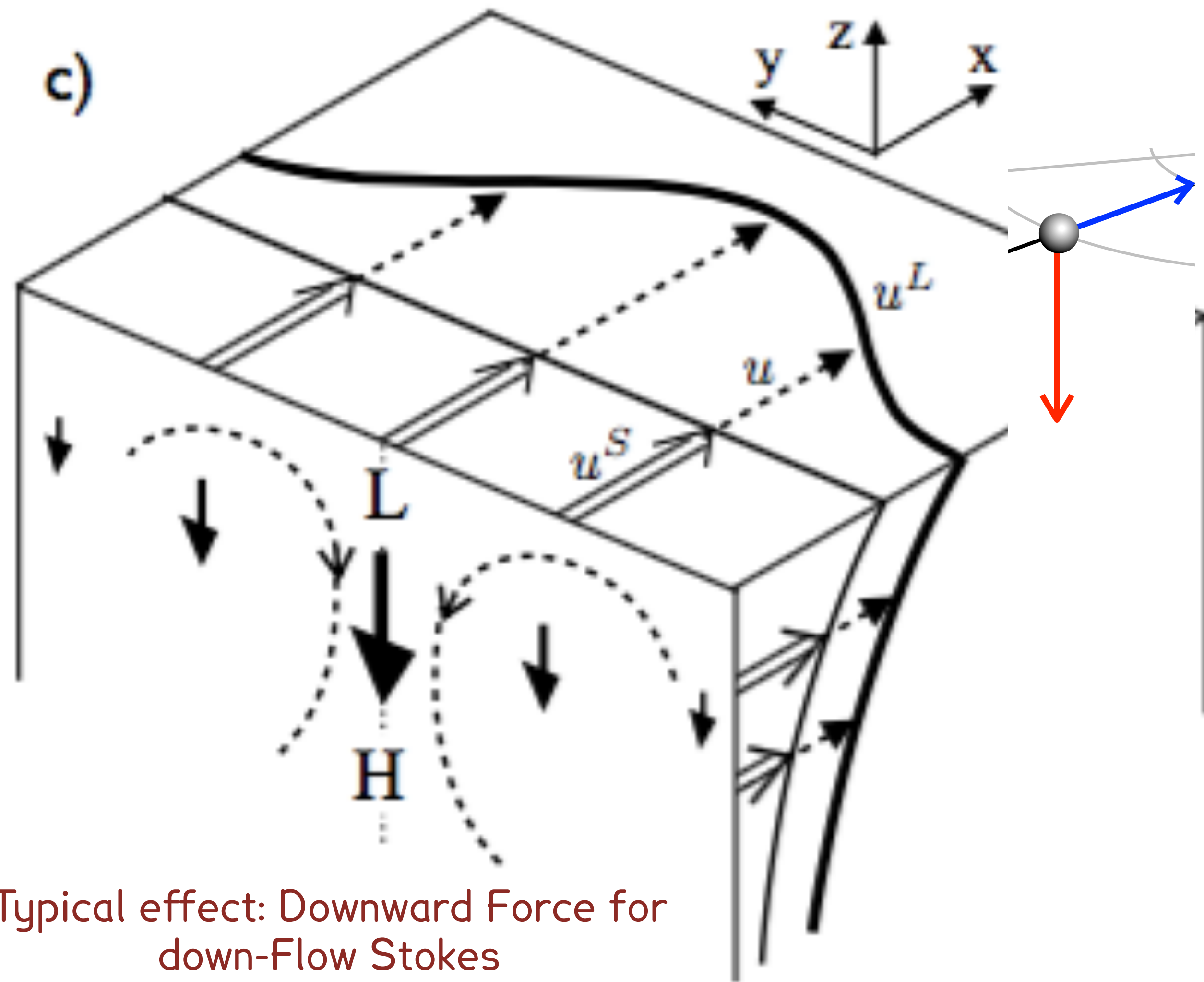
$$\frac{\alpha^2}{Ri} \left[ w_{,t} + v_j^L w_{,j} + \frac{M_{Ro}}{Ro Ri} w w_{,z} \right] = -\pi_{,z} + b - \varepsilon v_j^L v_{j,z}^s + \frac{\alpha^2}{Re Ri} w_{,jj}$$

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Typical effect: Downward Force for down-Flow Stokes

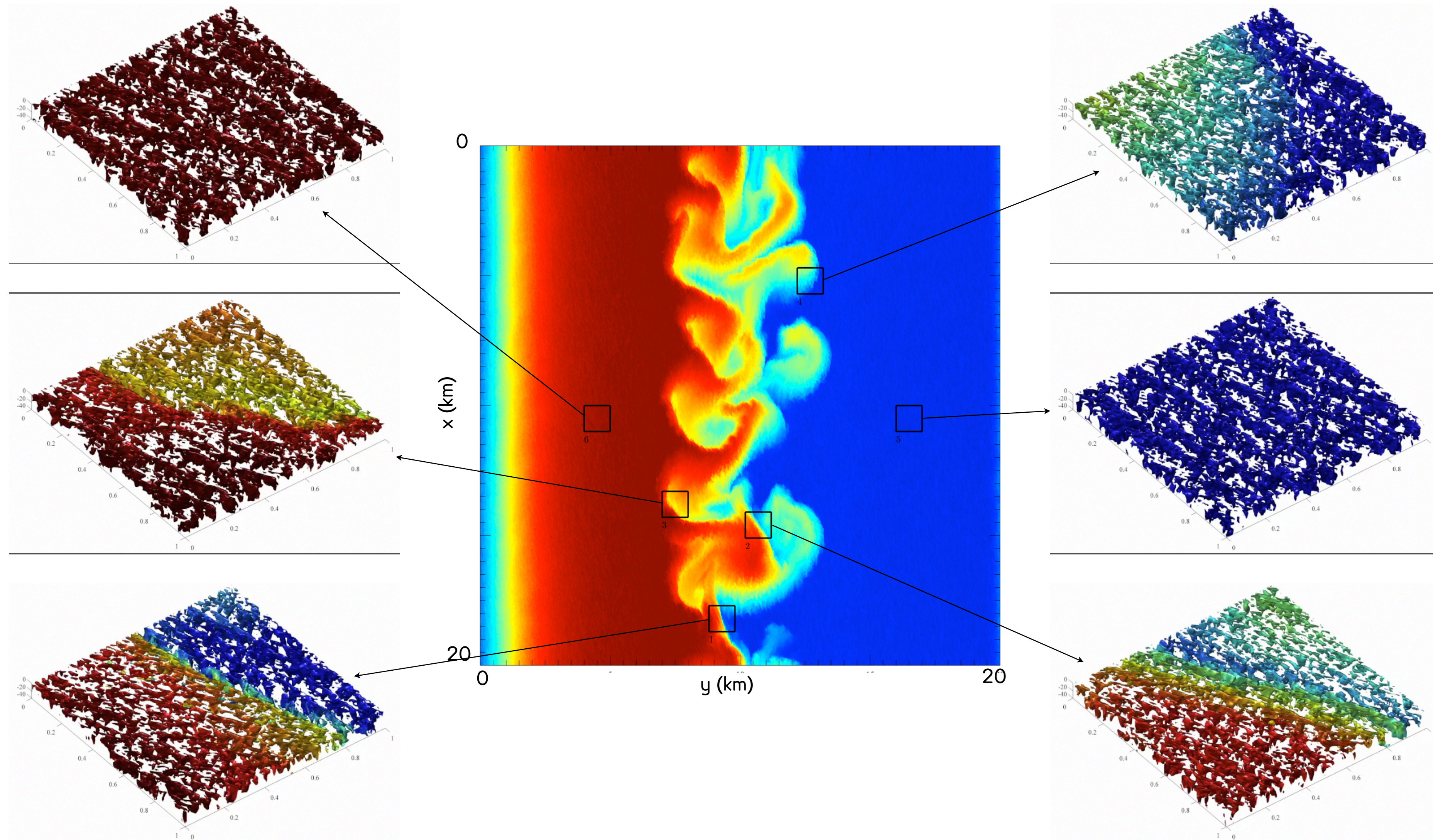
“wavy hydrostatic” if  $\epsilon \gg 1$

$$\frac{\alpha^2}{Re Ri} \left[ w_{,t} + v_j^L w_{,j} + \frac{M_{Ro}}{Ro Ri} w w_{,z} \right] = \boxed{-\pi_{,z} + b - \epsilon v_j^L v_{j,z}^s} + \frac{\alpha^2}{Re Ri} w_{,jj}$$

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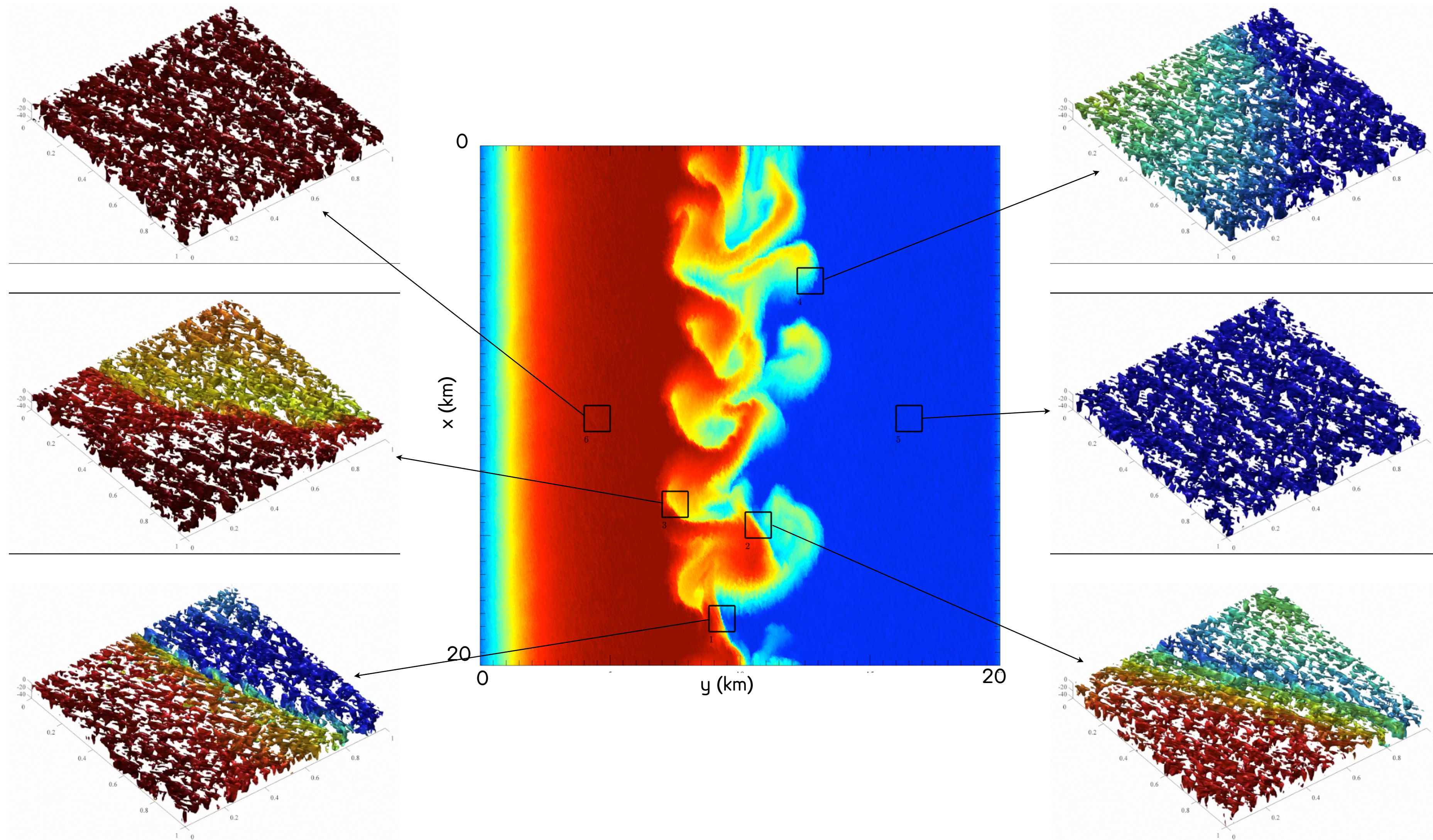


# Diverse types of interaction: Stronger Langmuir (small) Turbulence



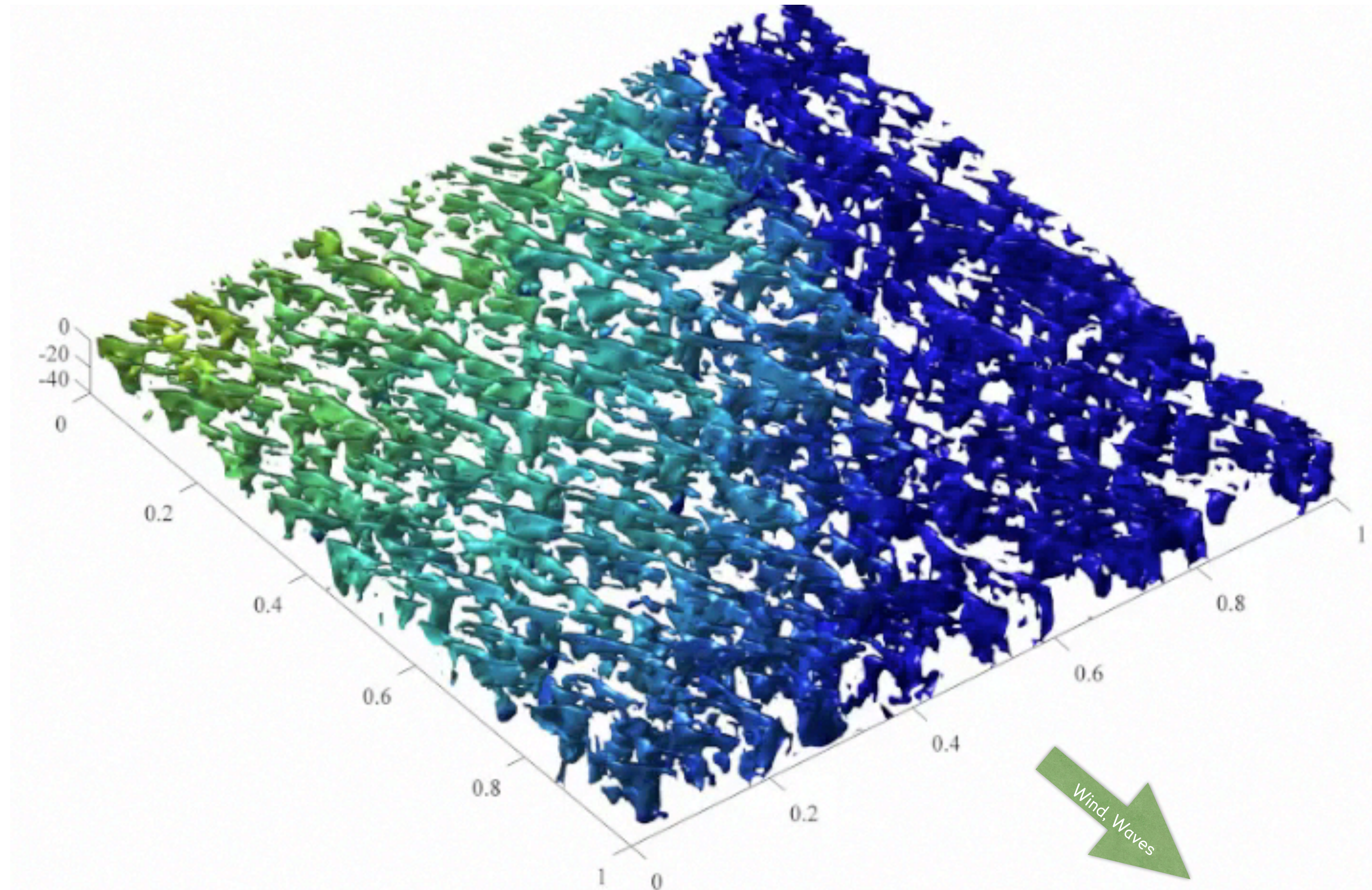
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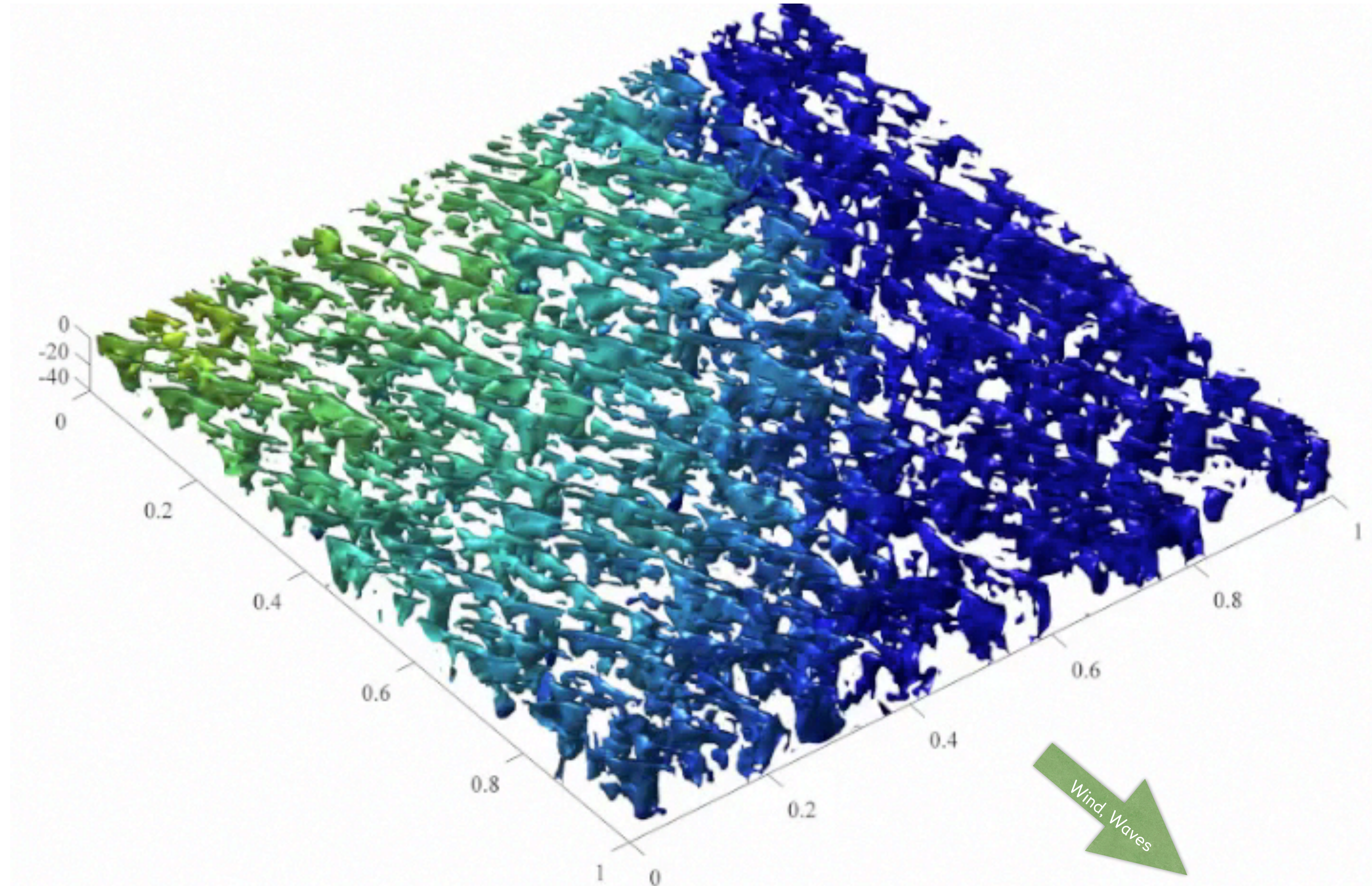
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A front that's not aligned with winds & waves is weaker



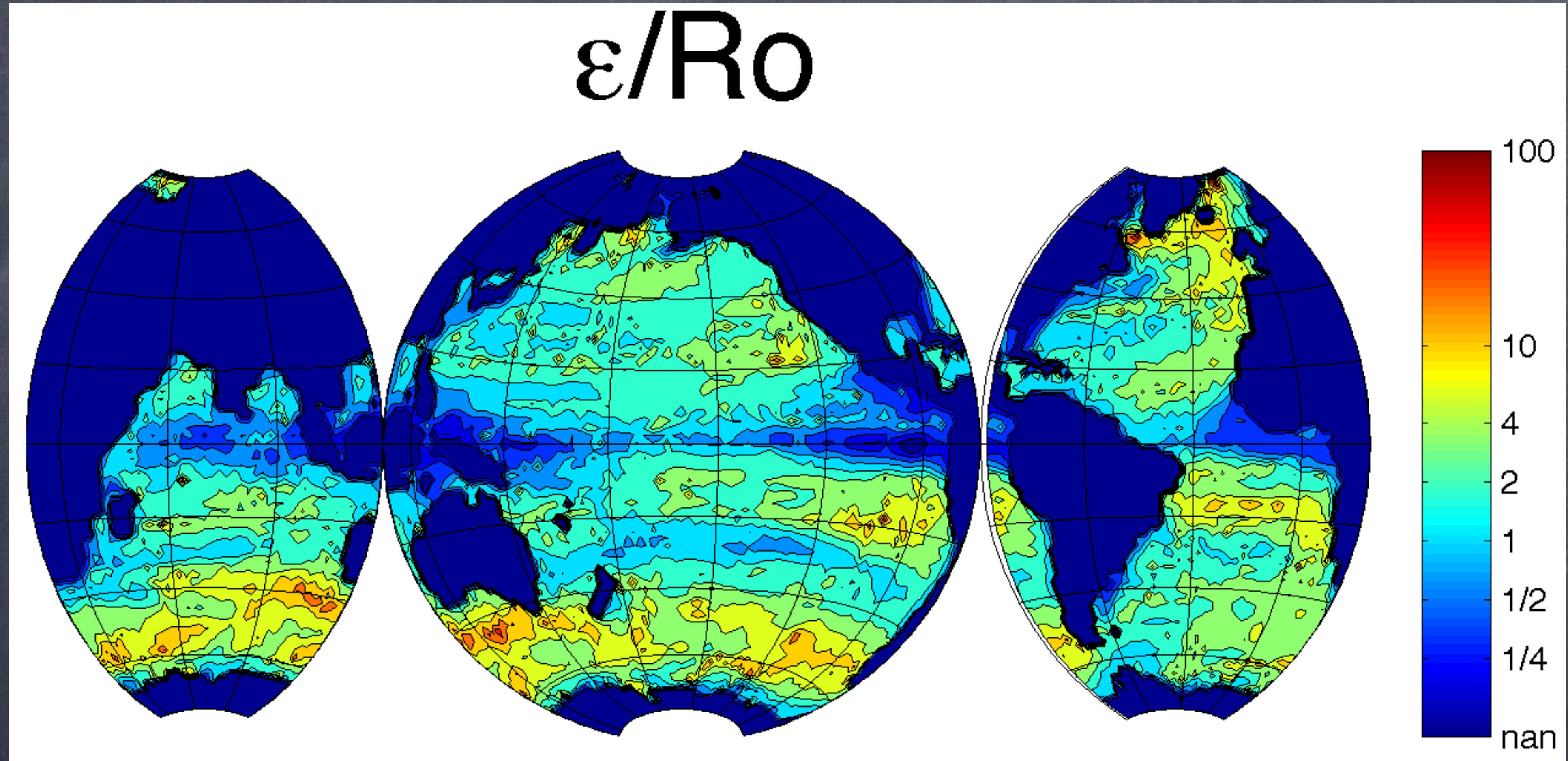
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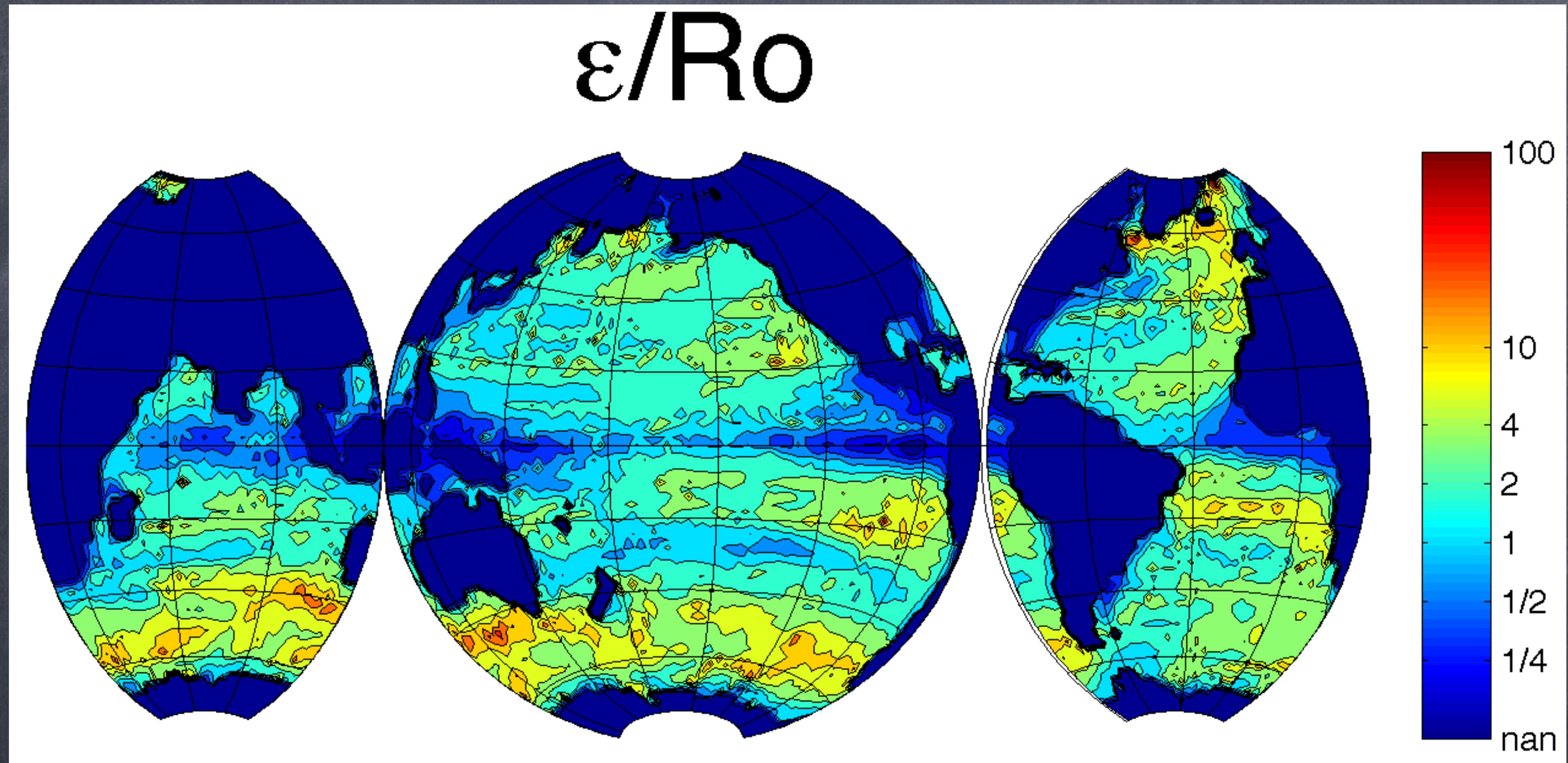
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# Do Stokes force directly affect larger scales?



$$\varepsilon = \frac{V^s H}{f L H_s} \quad Ro = \frac{U}{f L}$$

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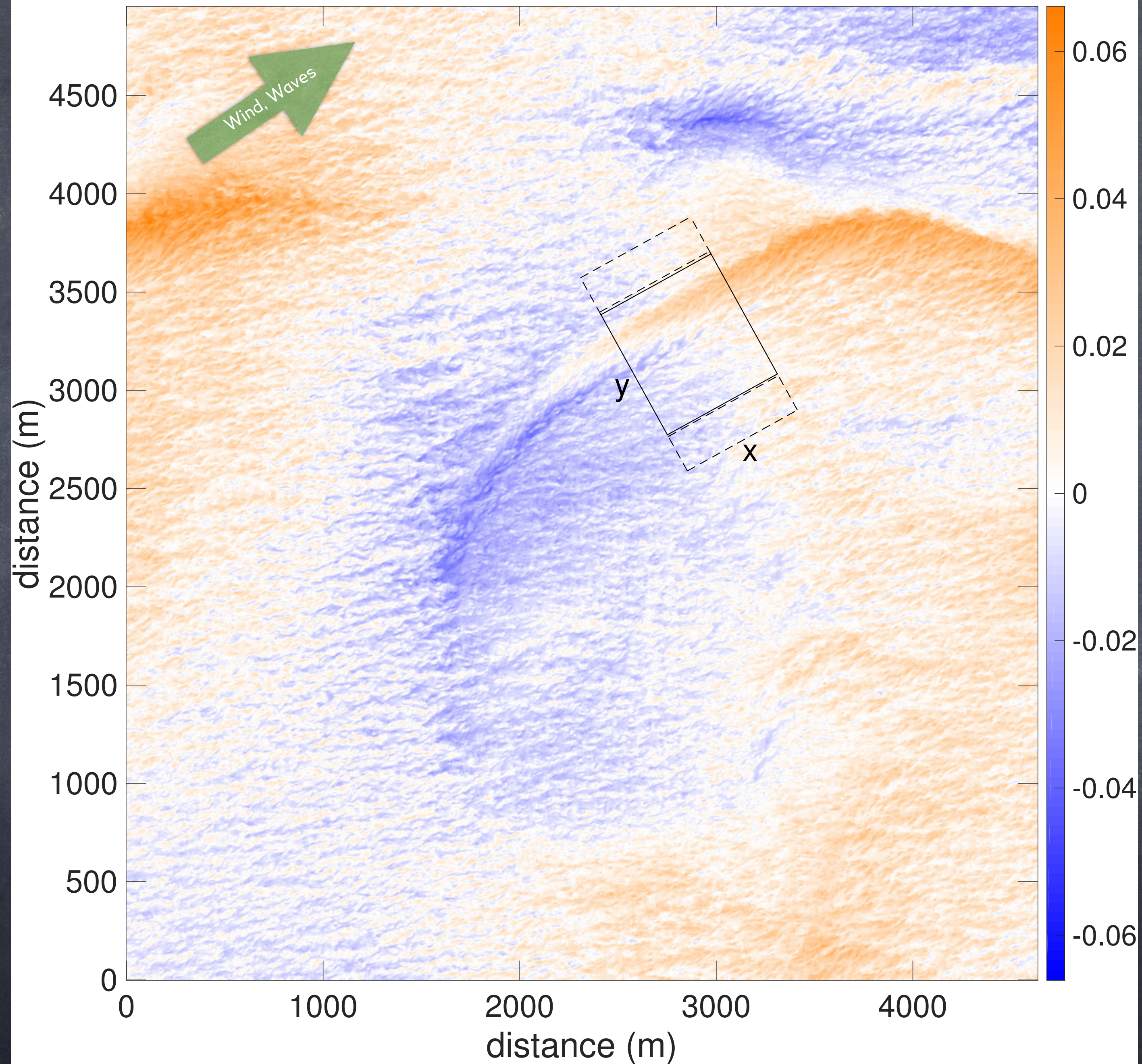
“wavy hydrostatic” if

$$\epsilon \gg 1$$

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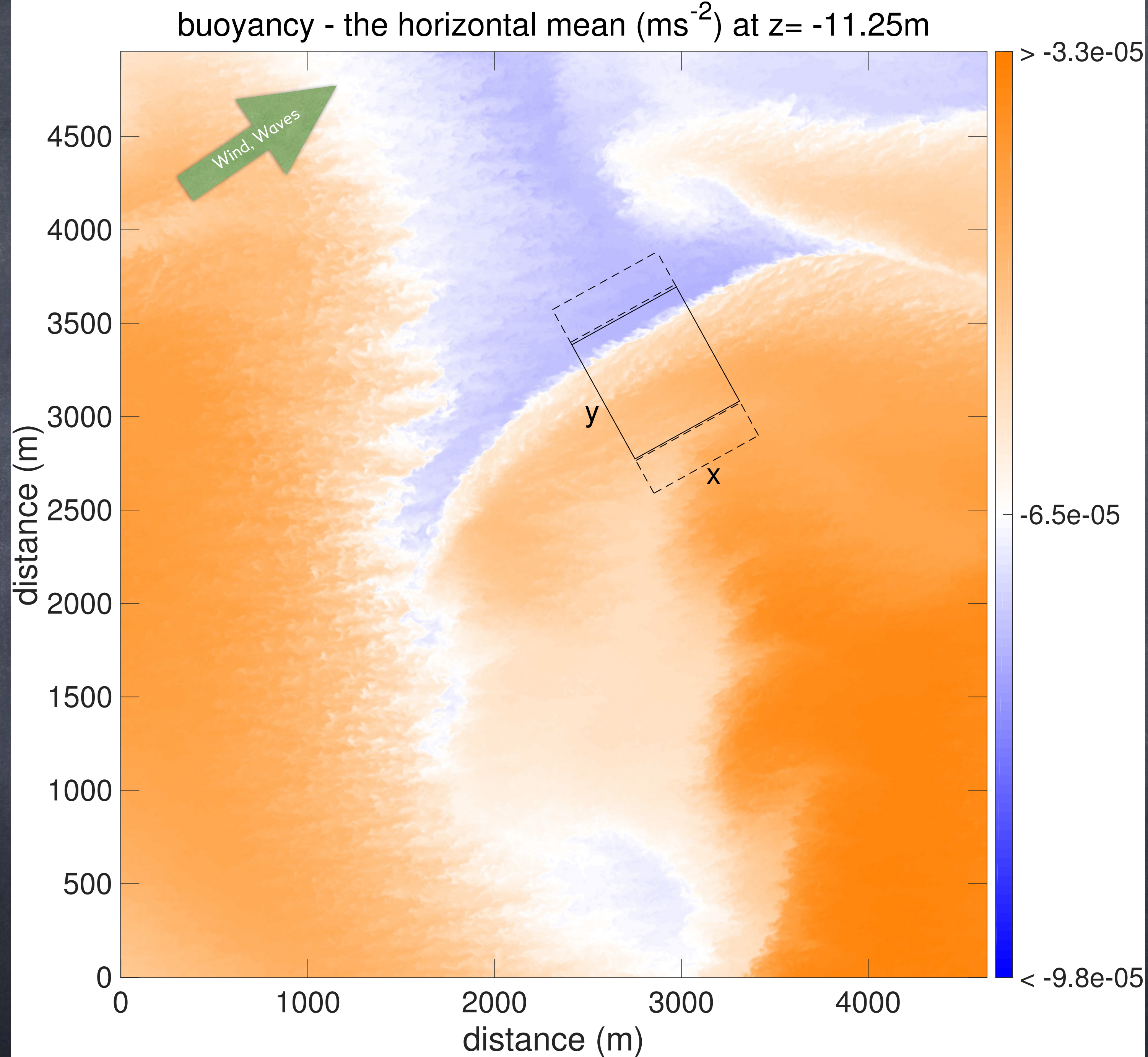
$$Ro = \frac{U}{f L}$$

velocity in the x-direction - the horizontal mean ( $\text{ms}^{-1}$ ) at  $z = -11.25\text{m}$



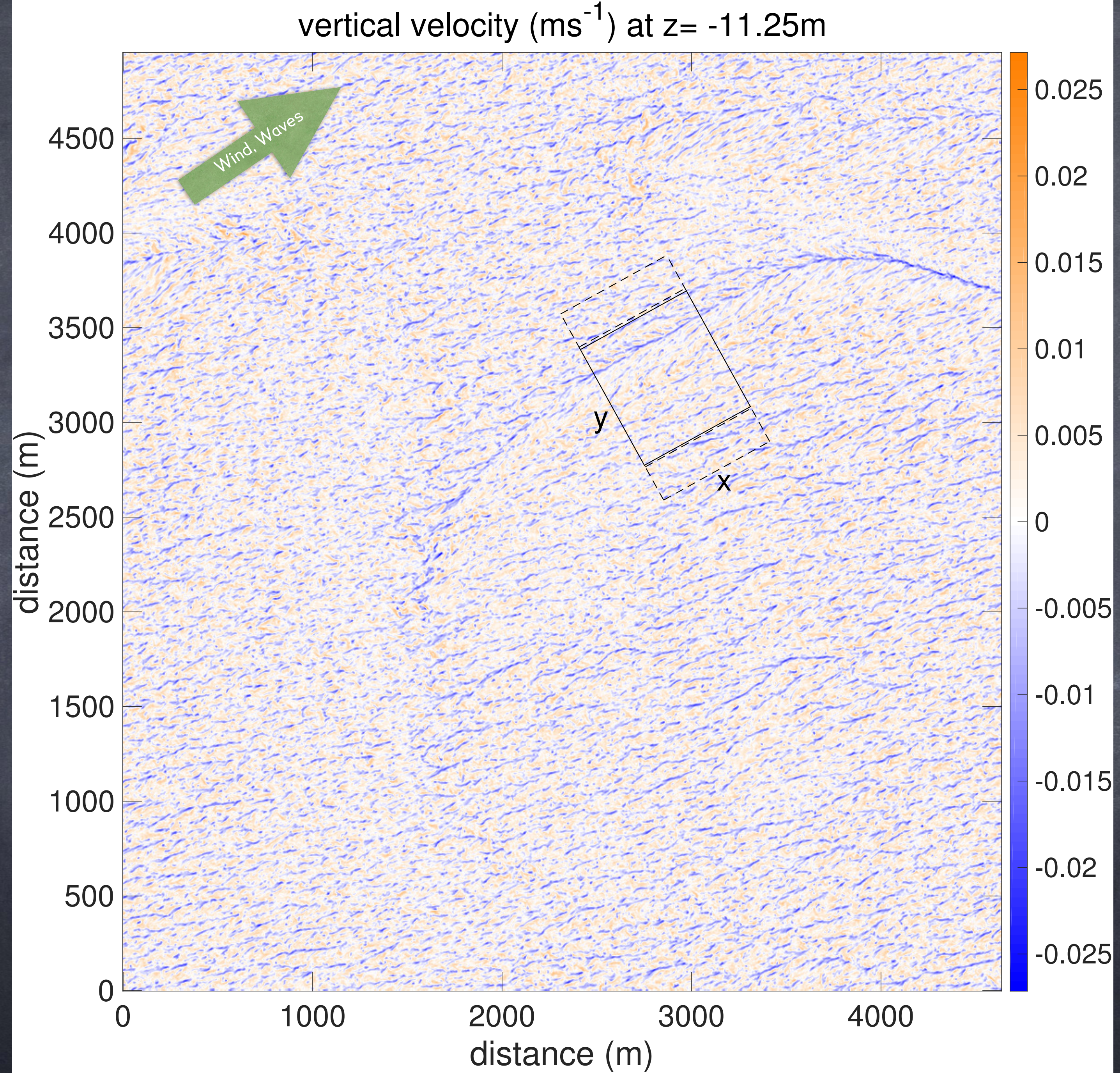
V

T

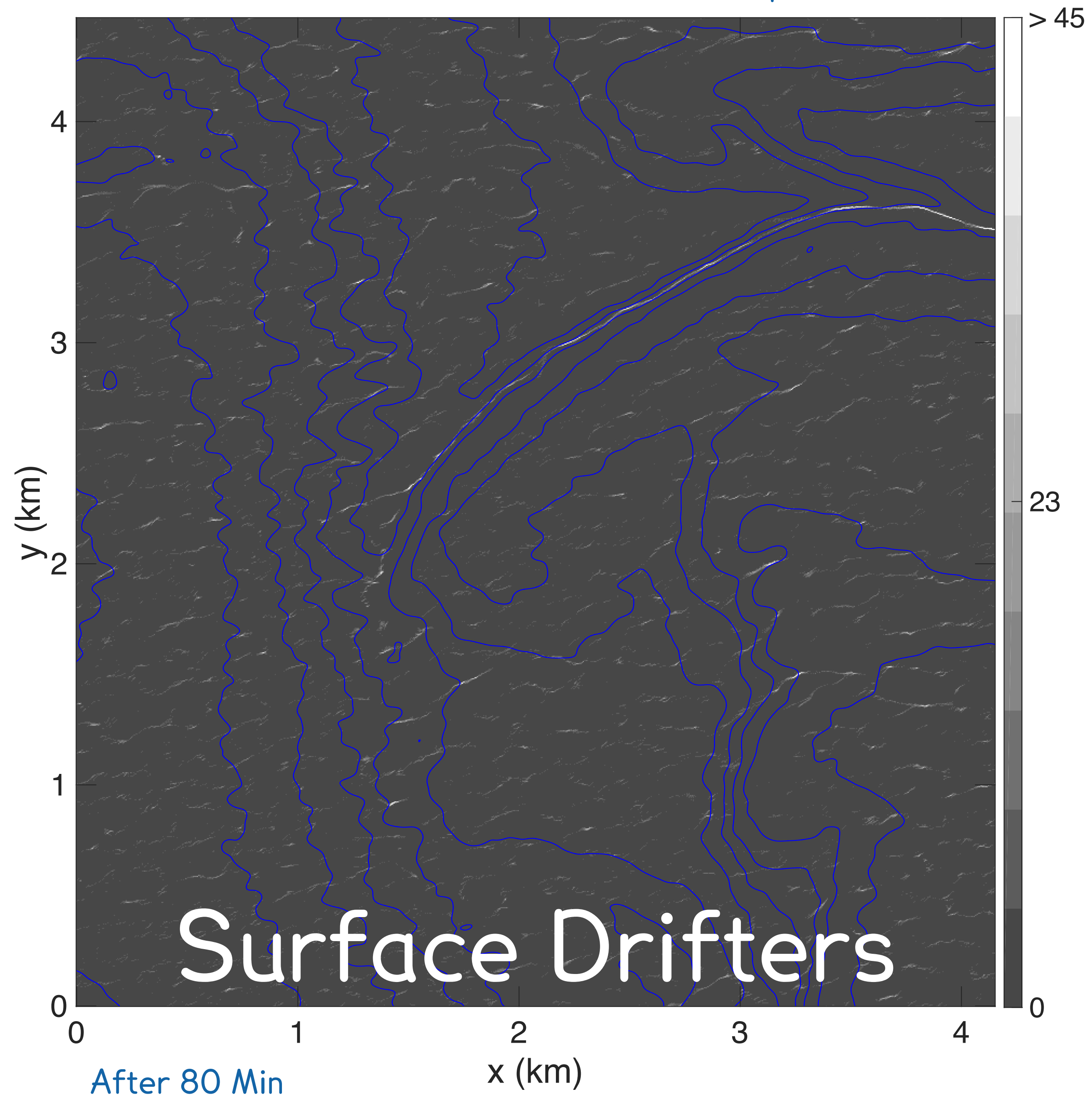


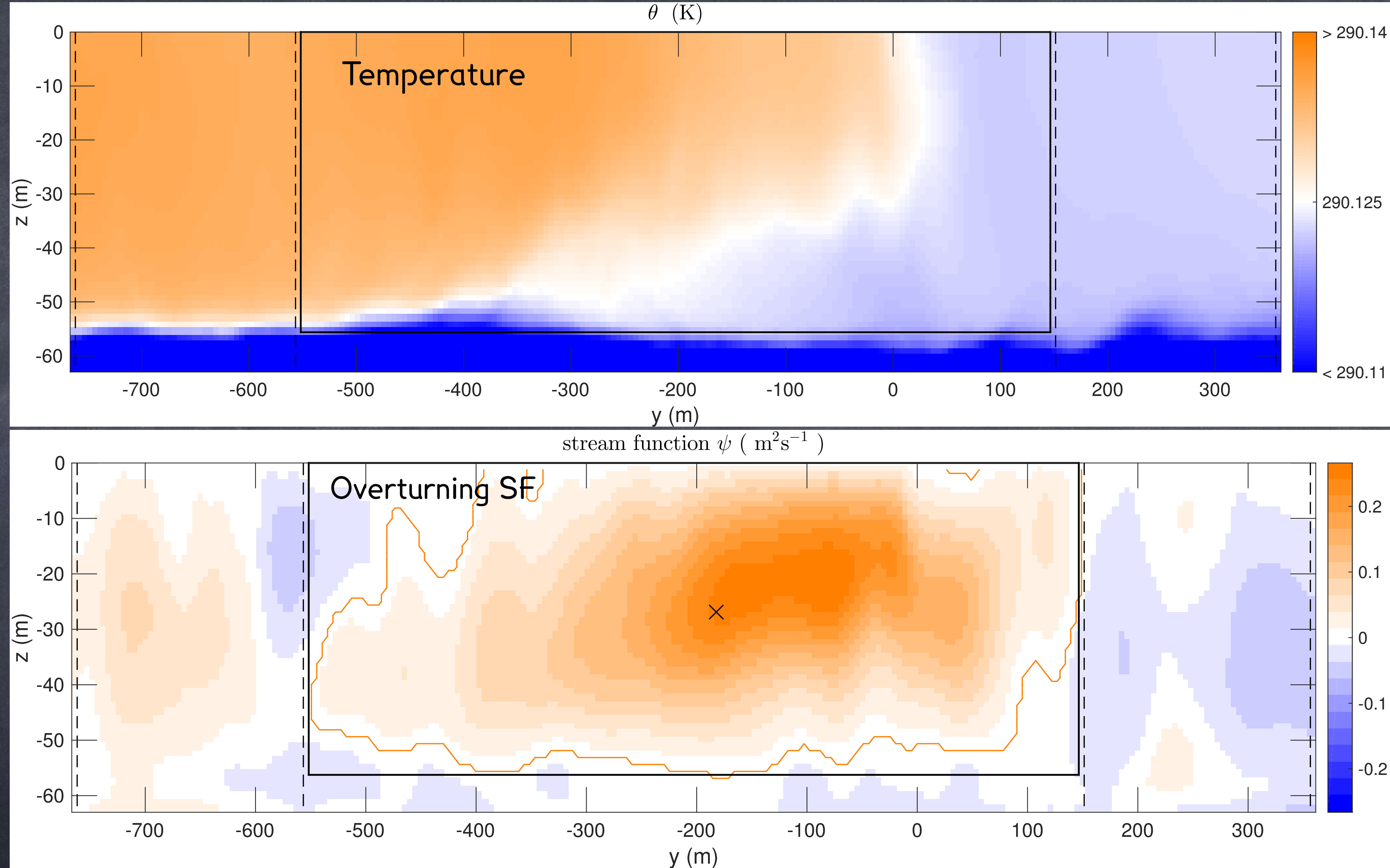


W



Initially every surface node has 1 drifter,  
so there are 851796 drifters in the picture





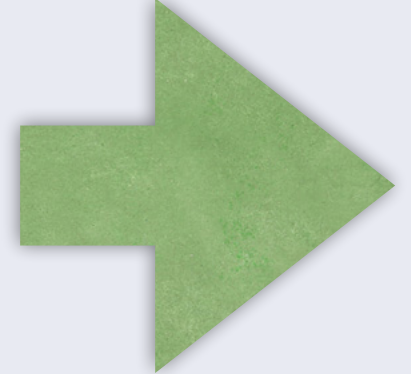
N. Suzuki, BFK, P. E. Hamlington, and L. P. Van Roekel. Surface waves affect frontogenesis. *Journal of Geophysical Research-Oceans*, 121:1-28, 2016.

# Do (wavy hydrostatic) Stokes Forces Matter?

## Yes! At Leading Order (in LES)

**Table 3.** Integrated Budget for Overturning Vorticity<sup>a</sup>

Responsible Force	Relative Value
<i>Relative Tendency of Overturning Circulation along the Cell Boundary</i>	
Net tendency	11 ± 8%
<b>Sources</b>	
Buoyancy anomaly	100%
Stokes shear force anomaly	44 ± 4%
Interaction with $v^H$	44 ± 8%
Frontal anomaly in pressure gradient	6 ± 9%
Nonlinear interaction with $v^B$ :	2 ± 1%
<b>Sinks</b>	
Frontal turbulence anomaly (mostly, imbalance in wavy Ekman relation)	-82 ± 11%
Coriolis on along-front jet	-66 ± 2%
Lagrangian advection of $(v^\psi, w^\psi)$	-36 ± 7%



- N. Suzuki and BFK. Understanding Stokes forces in the wave-averaged equations. *Journal of Geophysical Research-Oceans*, 121:1-18, 2016.
- N. Suzuki, BFK, P. E. Hamlington, and L. P. Van Roekel. Surface waves affect frontogenesis. *Journal of Geophysical Research-Oceans*, 121:1-28, 2016.

# Analytic & Numerical Wavy Submesoscale Stability: Symmetric Instabilities

- Hoskins (1974) showed that if a front in thermal wind balance is symmetrically unstable, the PV must be anticyclonic.
- Haney et al (2015) extend Hoskins' analysis to flows in Lagrangian thermal wind balance in the special case of constant Stokes shear.

$$fQ = \underbrace{f^2 N^2 - M^4}_{\text{geostrophic } fQ} - \underbrace{fM^2 U_z^S}_{\text{Stokes-modified } fQ} < 0.$$

Vert. Density Gradient      Horiz. Density Gradient      Anti-Stokes Shear

- In the absence of Stokes drift, this is equivalent to the familiar criteria on Richardson Number, with Stokes drift is distinct.

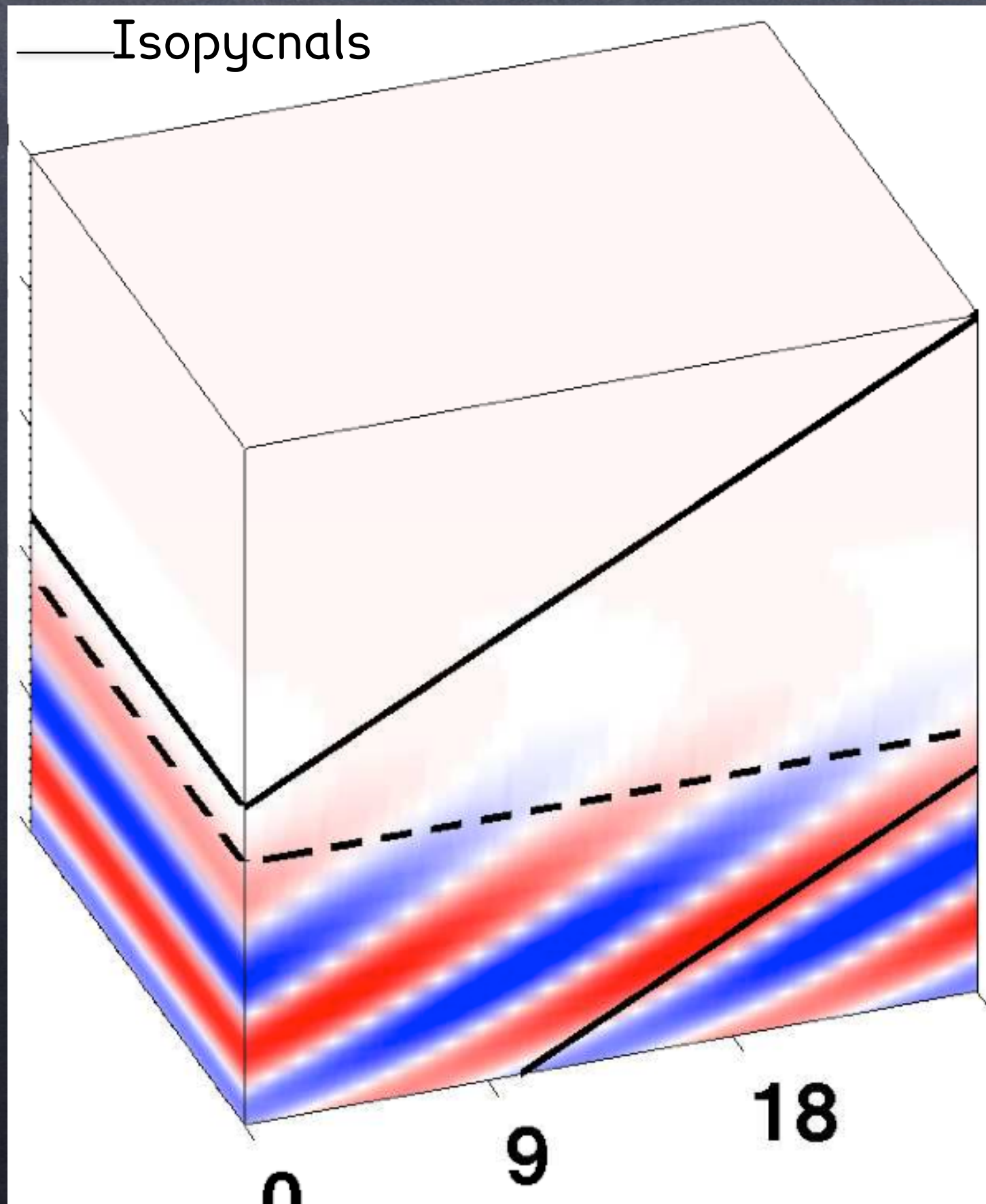
$$\text{Ri} < 1 \Rightarrow \text{SI}$$

Wavy Submesoscale  
Instability Different:  
Symmetric Instability

$$fQ < 0 \Rightarrow \text{SI}$$

$$\text{Ri} = 0.5$$

Wavy Stokes  
Stabilize SI

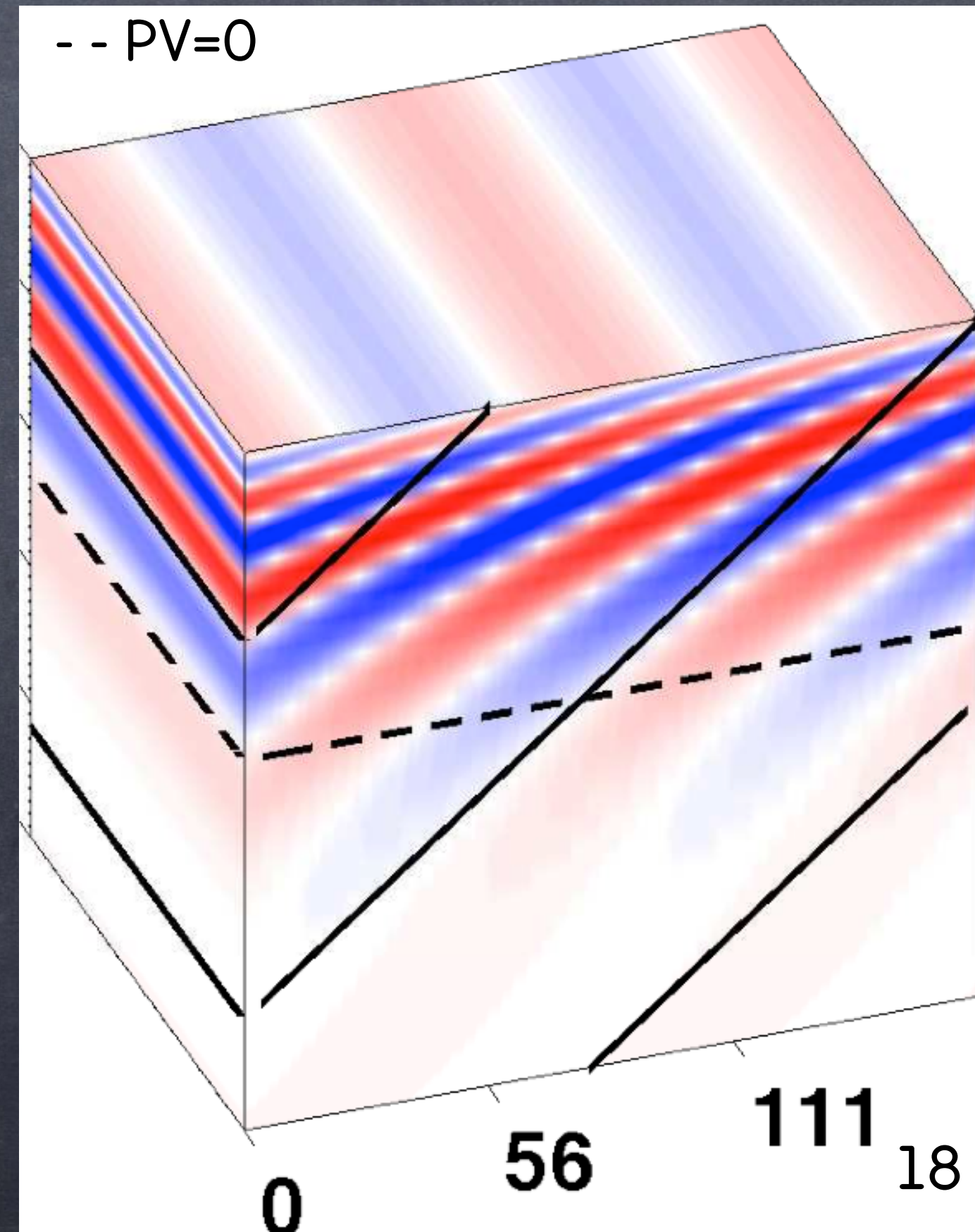


Cross front  
velocity for the  
fastest growing  
mode

S. Haney, BFK, K.  
Julien, and A. Webb.  
Symmetric and  
geostrophic  
instabilities in the  
wave-forced ocean  
mixed layer. JPO  
45:3033-3056,  
2015.

$$\text{Ri} = 2$$

Wavy Stokes  
Destabilize SI



# Frontal Consequences:

Observing Energy Flux from Global to Dissipative Scales

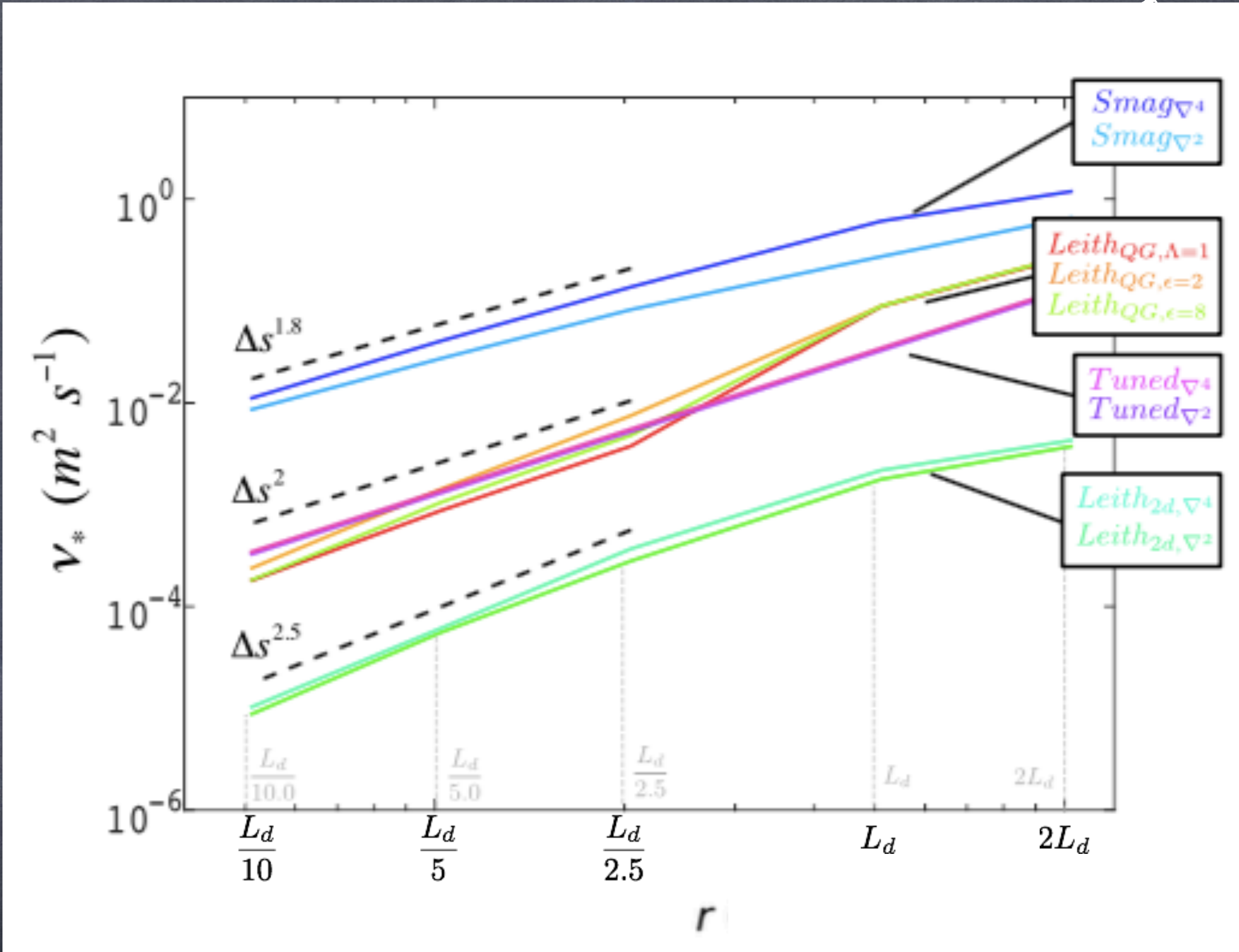
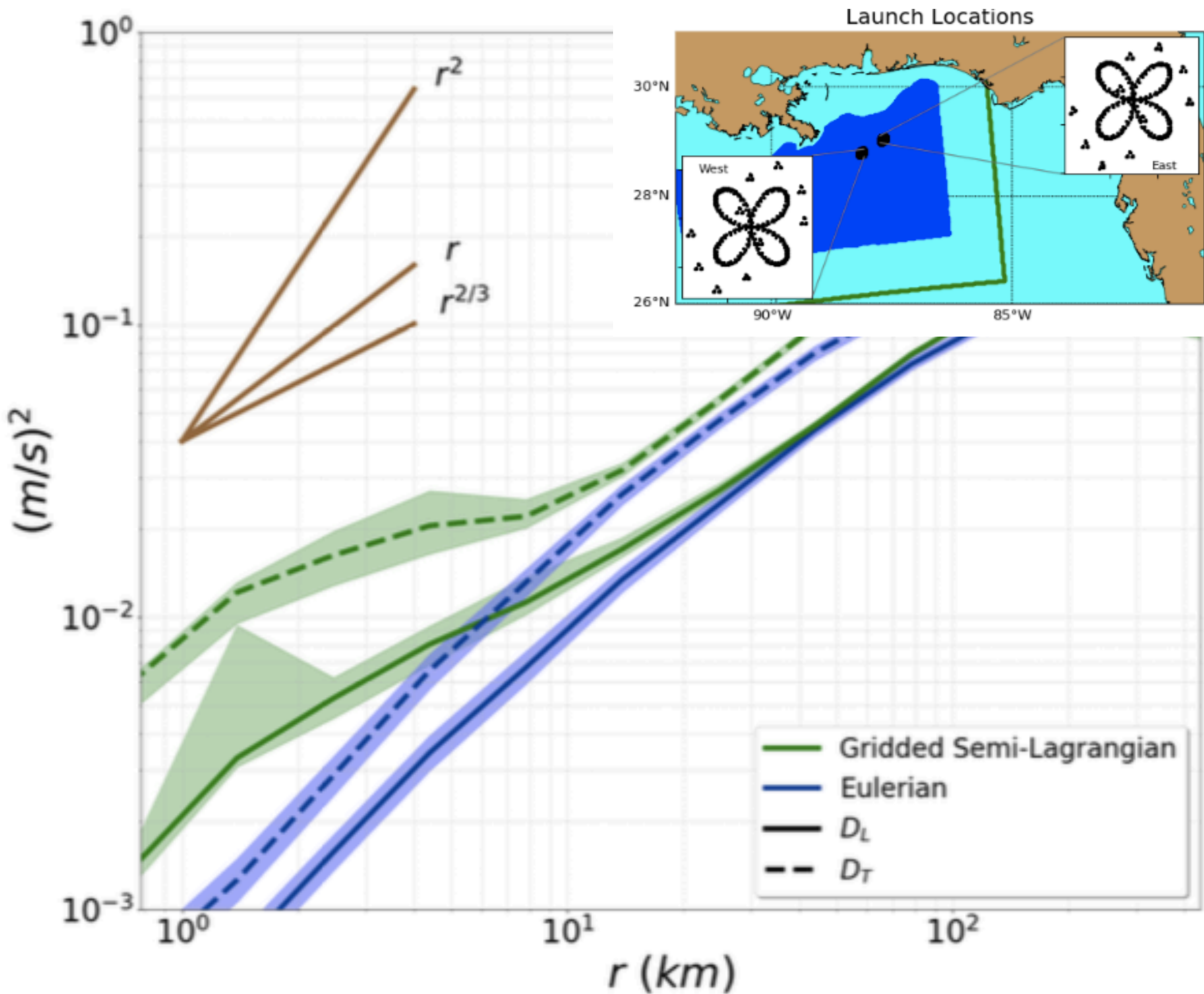
- Energy flows from the global winds & tides to mm-scale dissipation by viscosity
- Fronts are a key concentration effect of energy, aiding in the transfer
- However, the presence of fronts also complicates observations of the energy flow
- Observations by drifters, etc. must handle the strong heterogeneity due to the presence of fronts

N. Suzuki, BFK, P. E. Hamlington, and L. P. Van Roekel. Surface waves affect frontogenesis. *Journal of Geophysical Research-Oceans*, 121:1-28, 2016.

J. C. McWilliams and BFK. Oceanic wave-balanced surface fronts and filaments. *Journal of Fluid Mechanics*, 730:464-490, 2013.

# 'Observed' scale-sensitivity?

Gridded Semi-Lagrangian and Eulerian Second Order Structure Functions



Some Theory/Model combos are inconsistent (e.g., Smagorinsky in a QG regime)

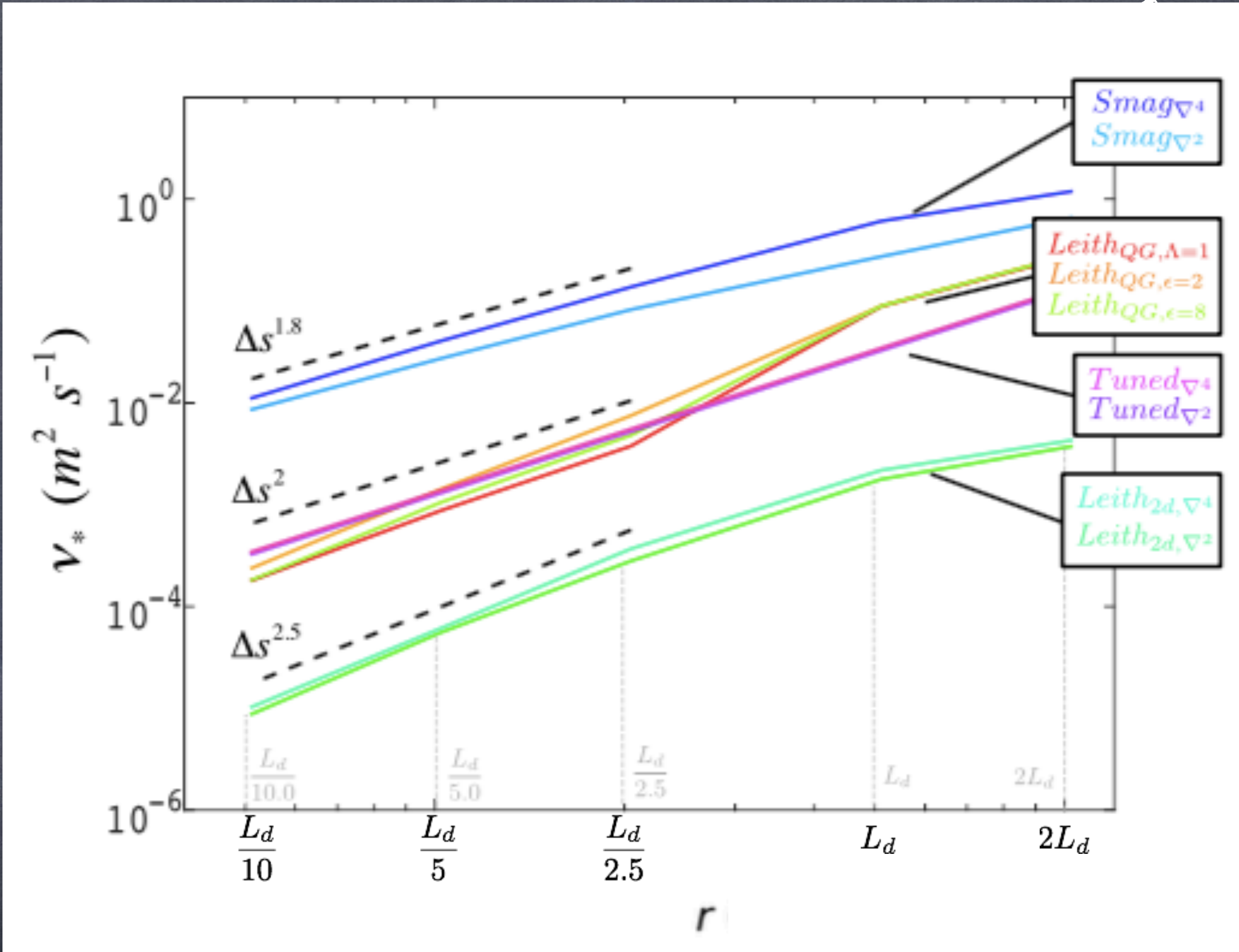
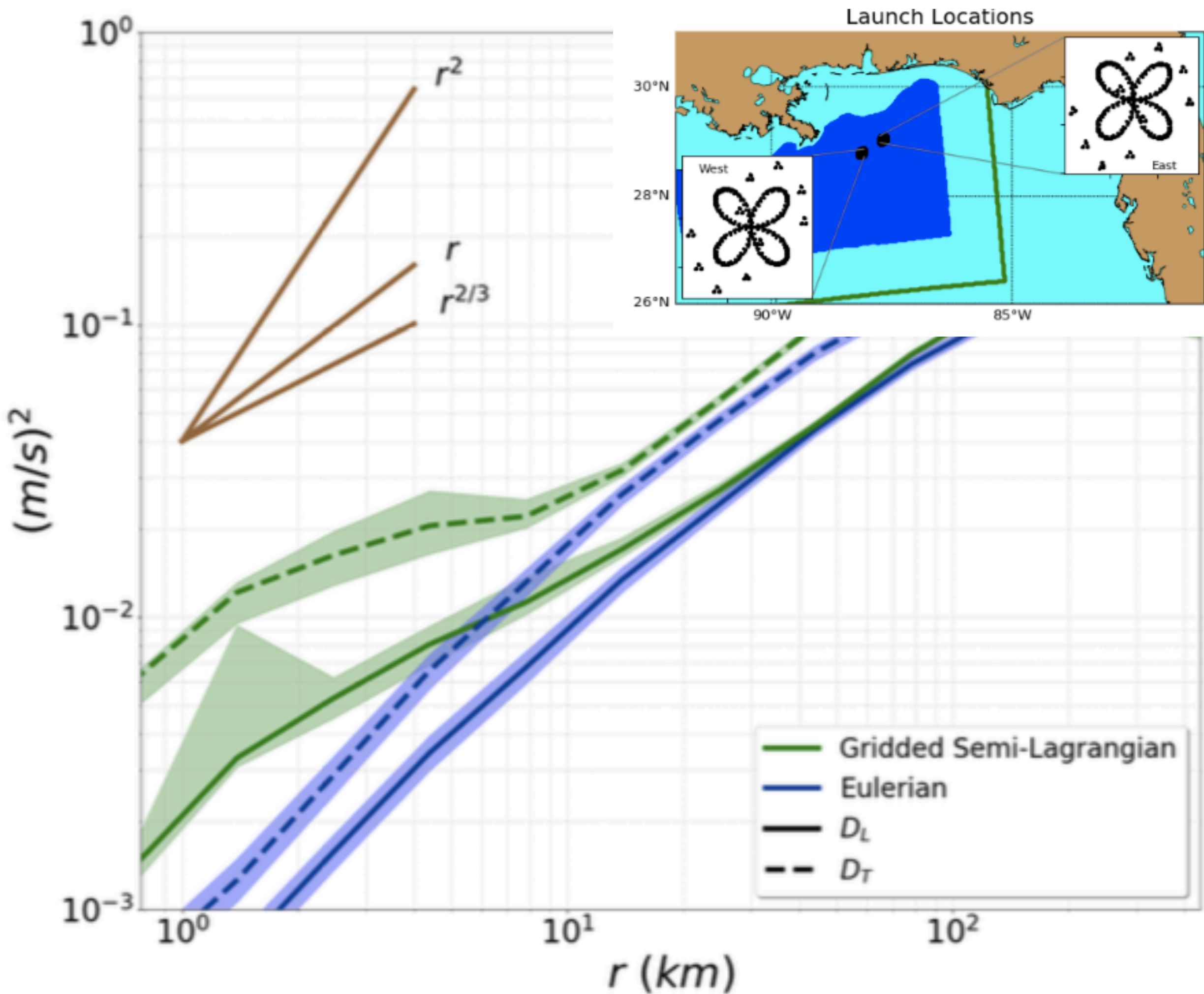
S. D. Bachman, B. Fox-Kemper, and B. Pearson, 2017: A scale-aware subgrid model for quasi-geostrophic turbulence. Journal of Geophysical Research—Oceans, 122:1529–1554.

J. Pearson, B. Fox-Kemper, R. Barkan, J. Choi, A. Bracco, and J. C. McWilliams. Impacts of convergence on Lagrangian statistics in the Gulf of Mexico. Journal of Physical Oceanography, February 2018. Submitted.



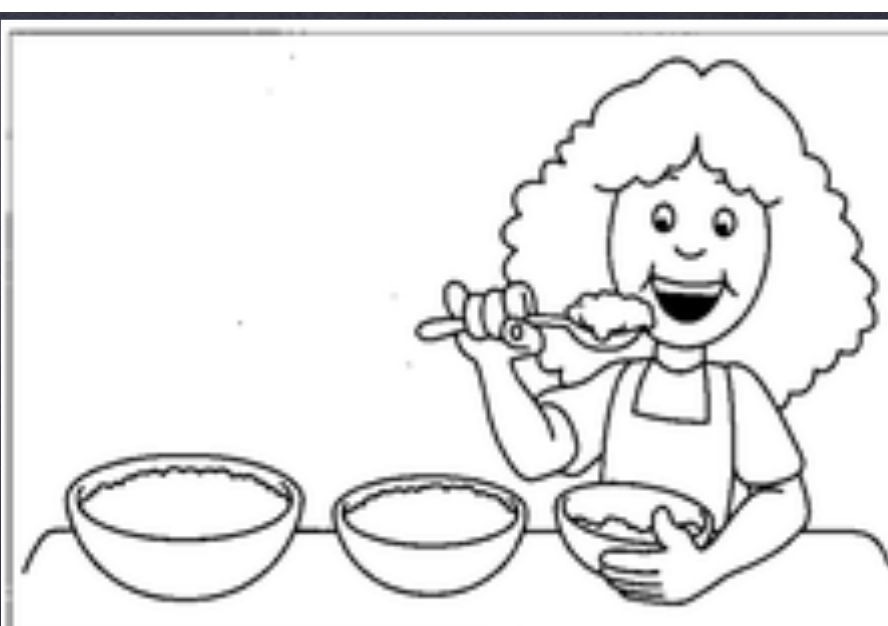
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