

Athermal Photofluidization of Glasses or Controlling Viscosity with Light

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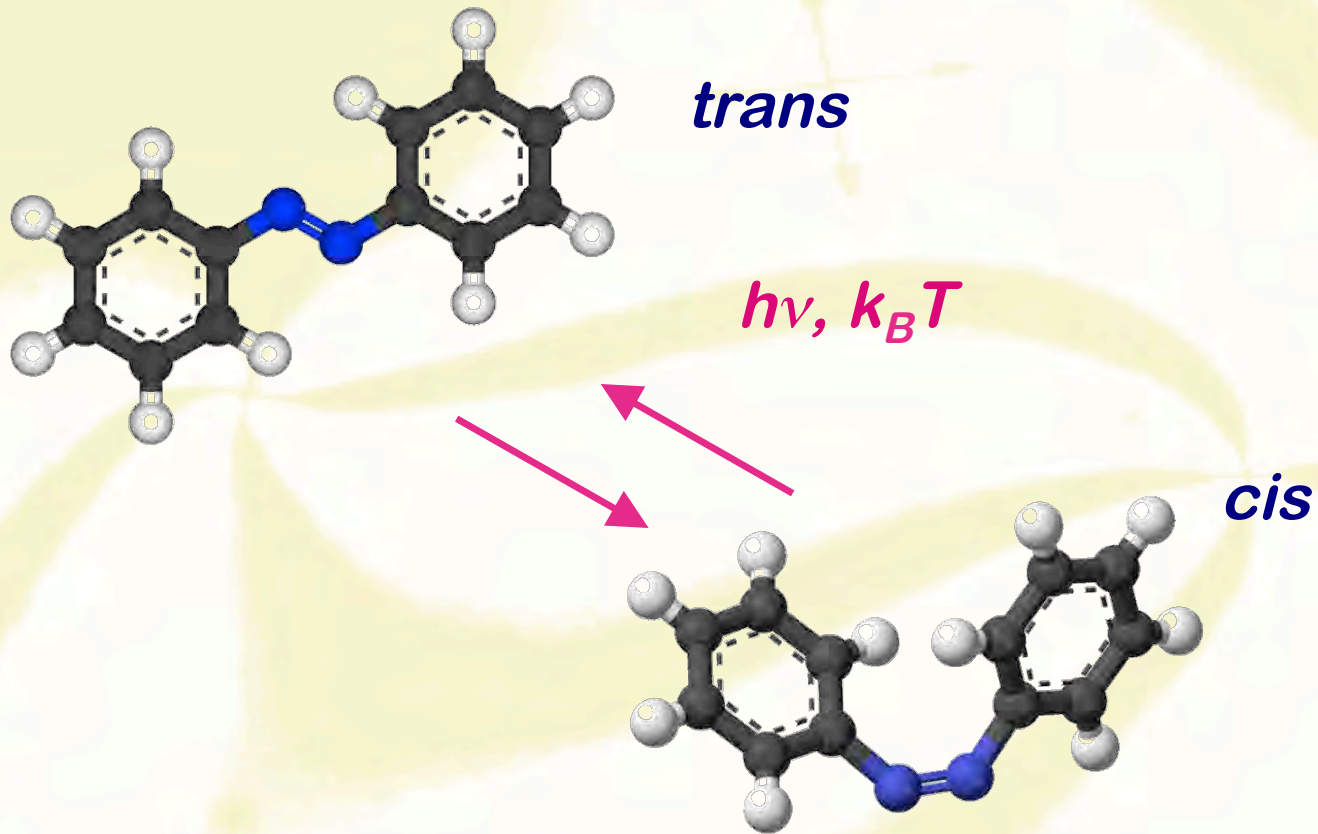
Matt Farrow, Eva Korblova, Dave Walba
Department of Chemistry and Biochemistry
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Department of Physics
Colorado School of Mines

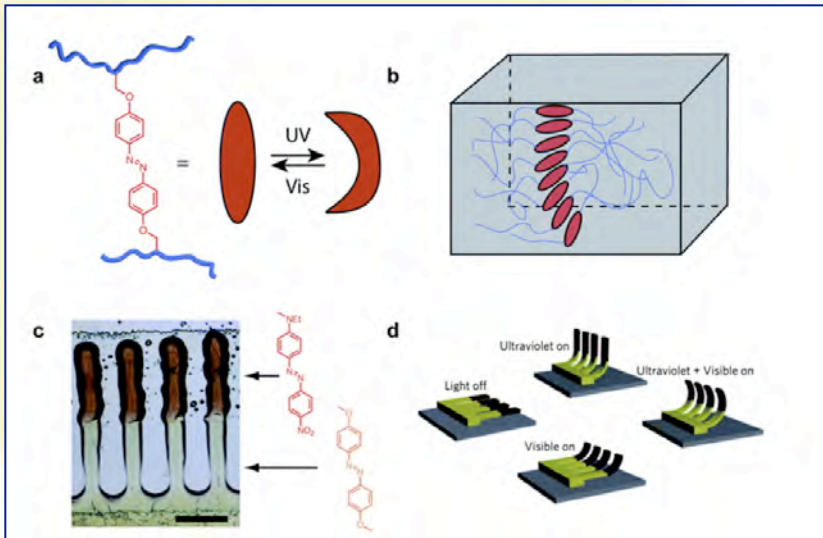
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Nat. Comm. (2013)

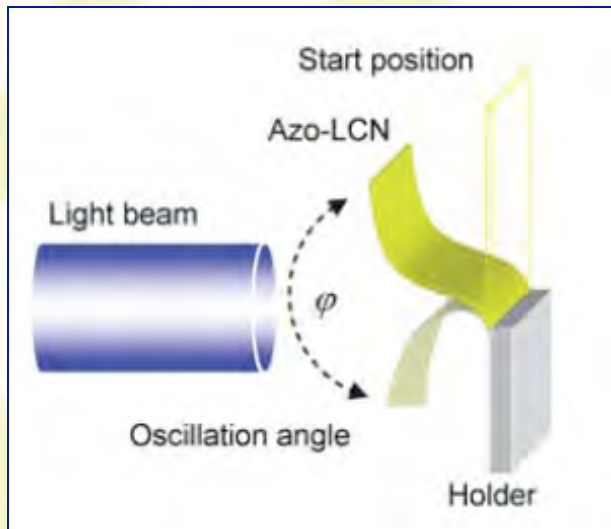
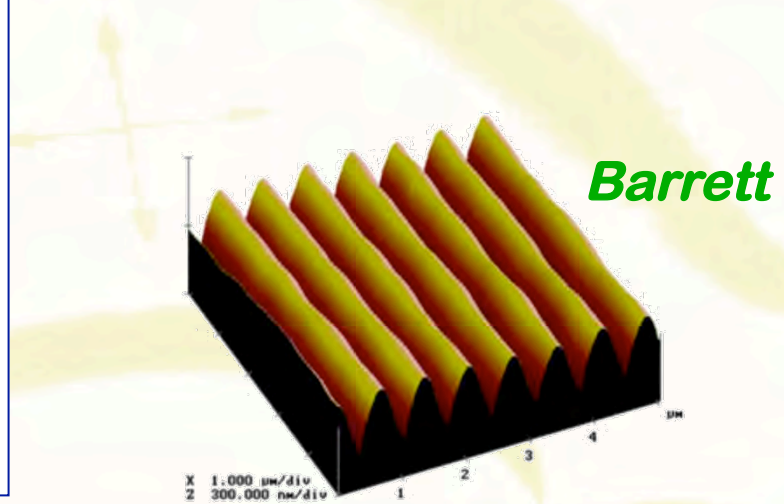
azobenzene



azo – based photo-active glassy polymers



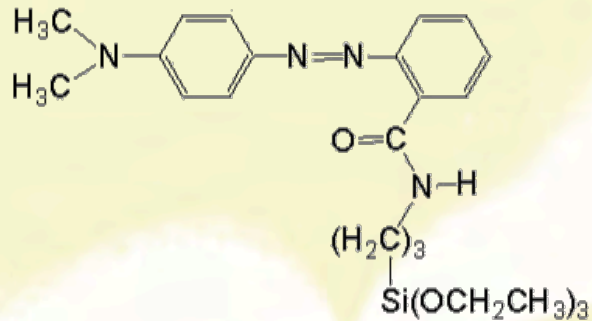
Broer



Wright- Patterson

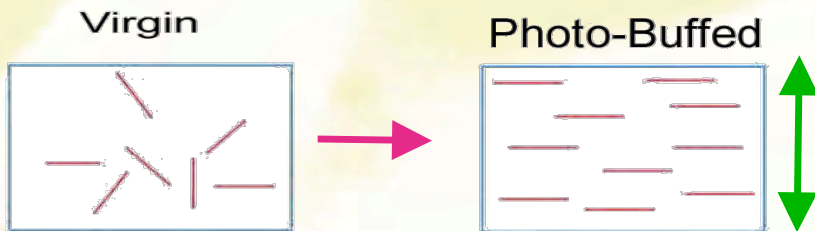
azo-based tethered molecular monolayer (dMR)

- ◆ precursor
(D. Walba)



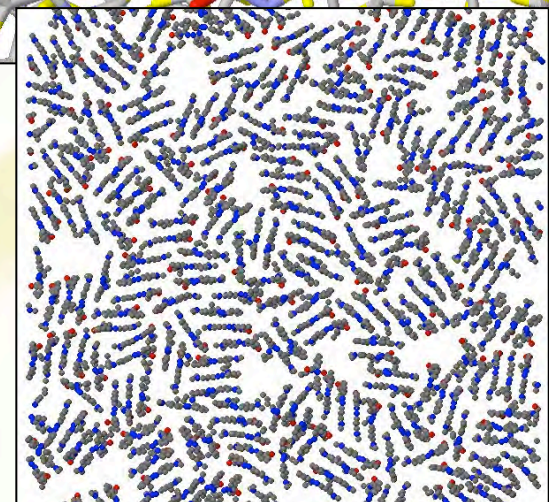
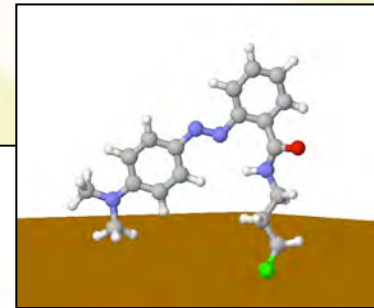
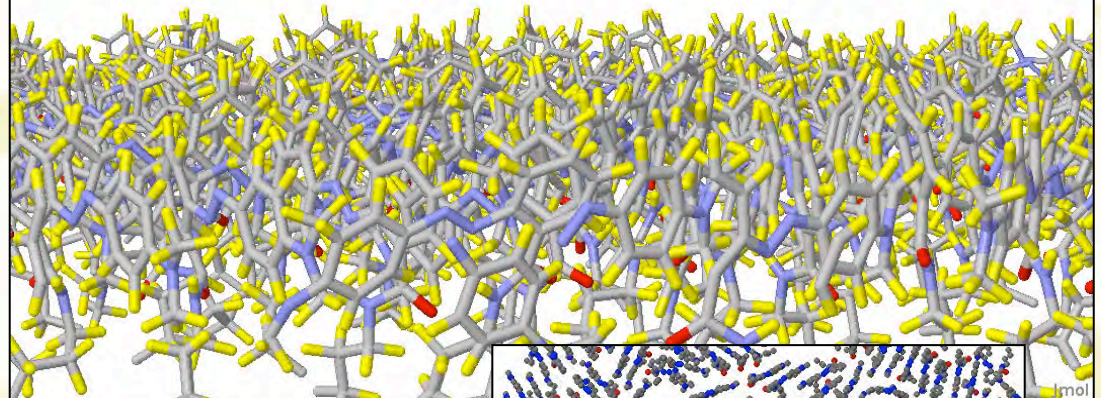
- ◆ covalently bonded to glass surface
- ◆ photo-orientates in-plane

Weigert mechanism:

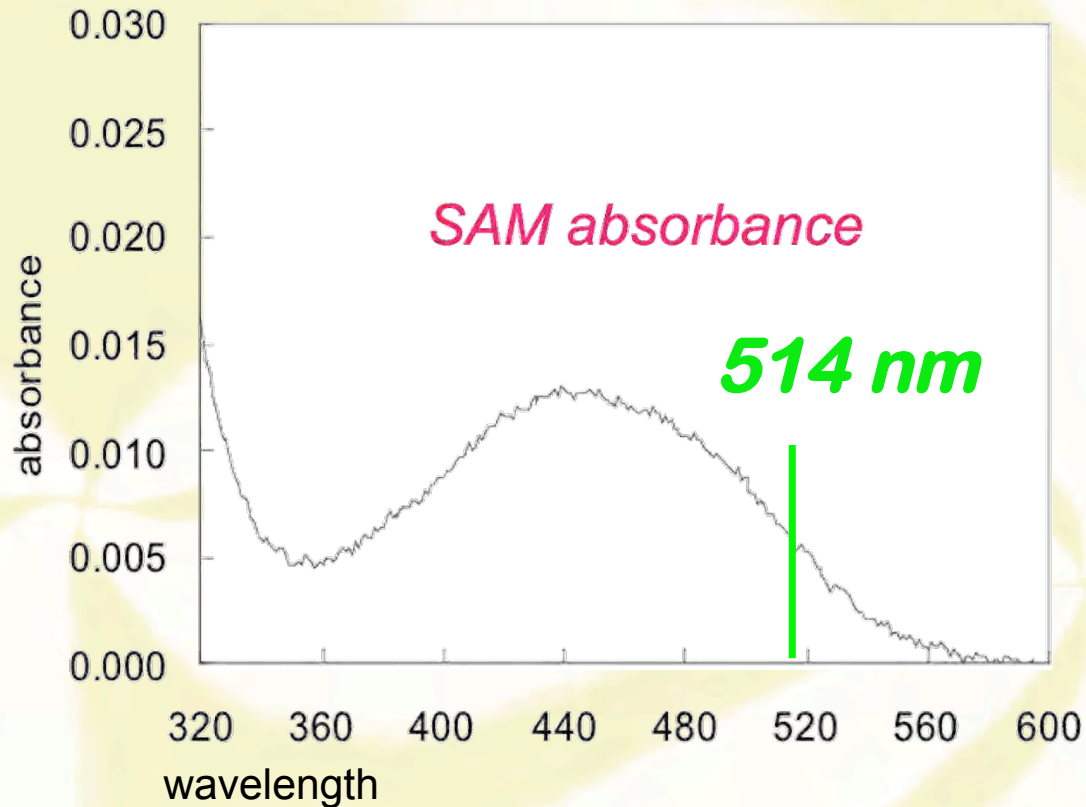


- ◆ photo-orientable 2D XY system

(D. Bedrov
M. Glaser)

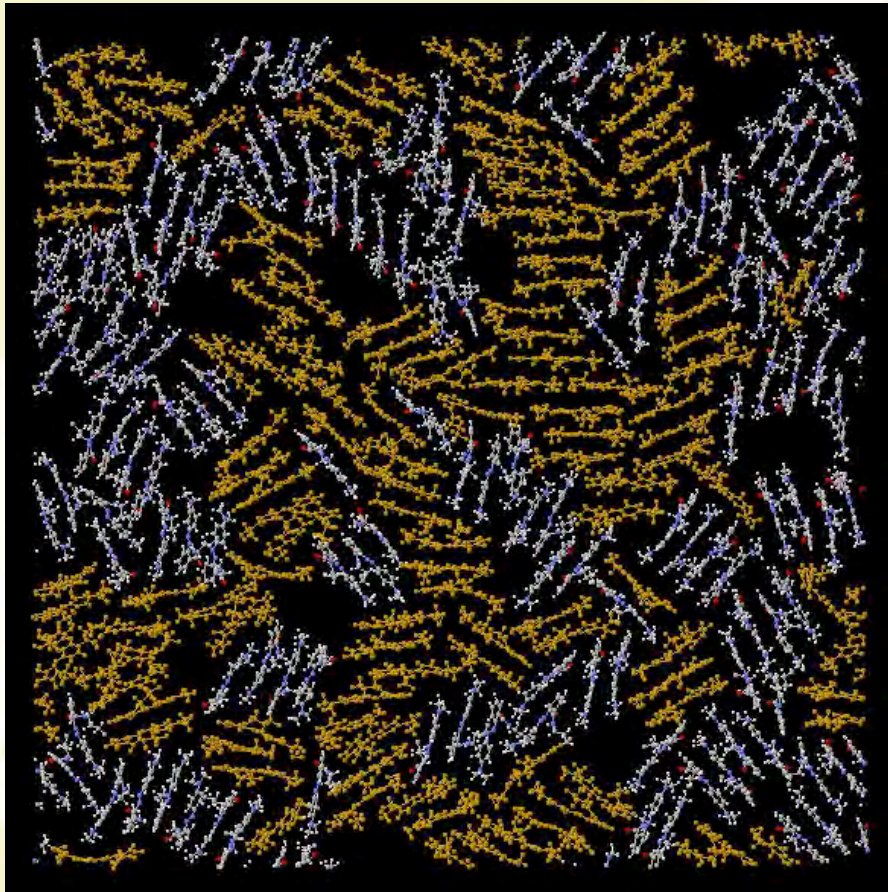


optical absorbance of the dMR monolayer

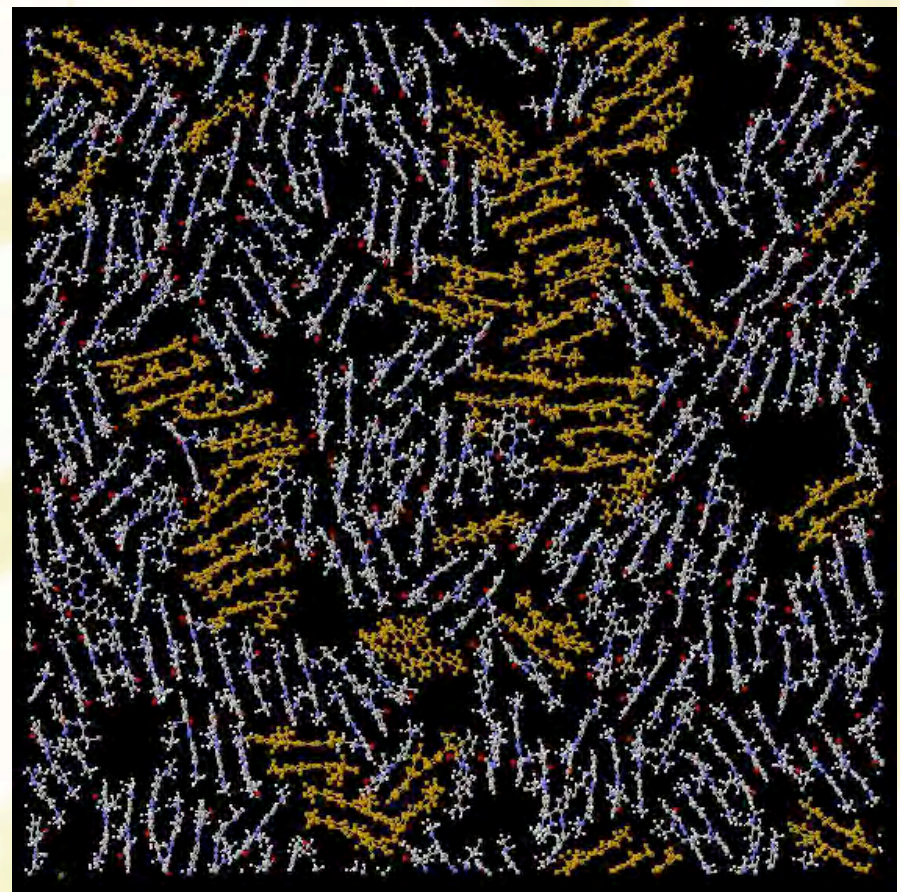


- ◆ *compare to solution (monolayer ~70% coverage)*
- ◆ *a fluence of $20\text{mJ}/\text{cm}^2$ gives 1 photon absorbed/molecule (pa/mol)*

Photo-orientation

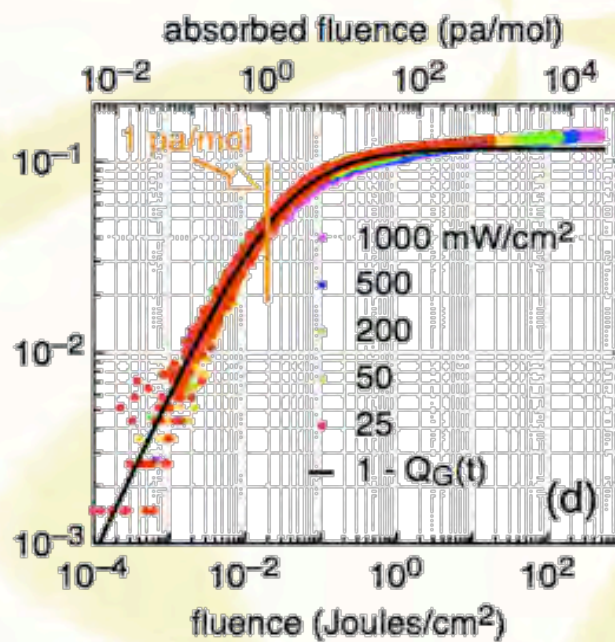
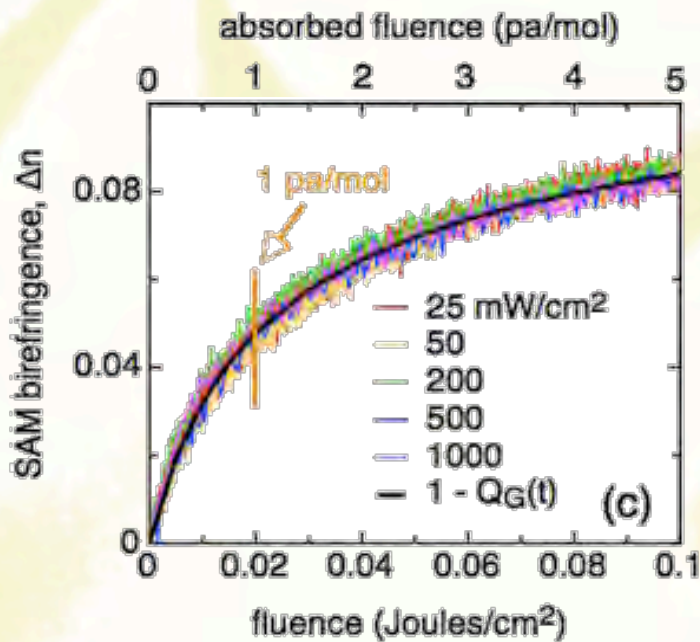
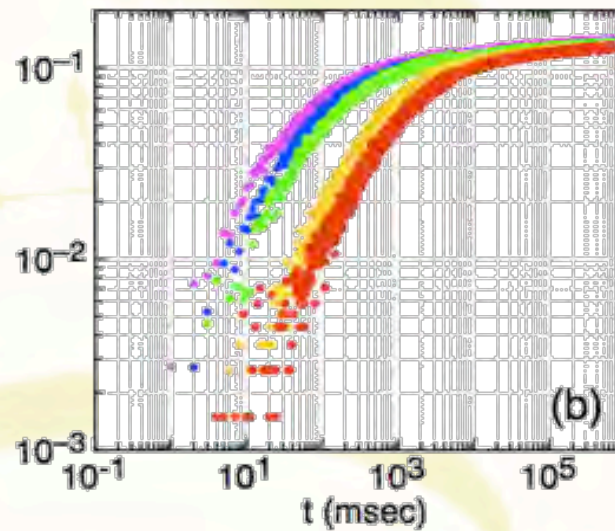
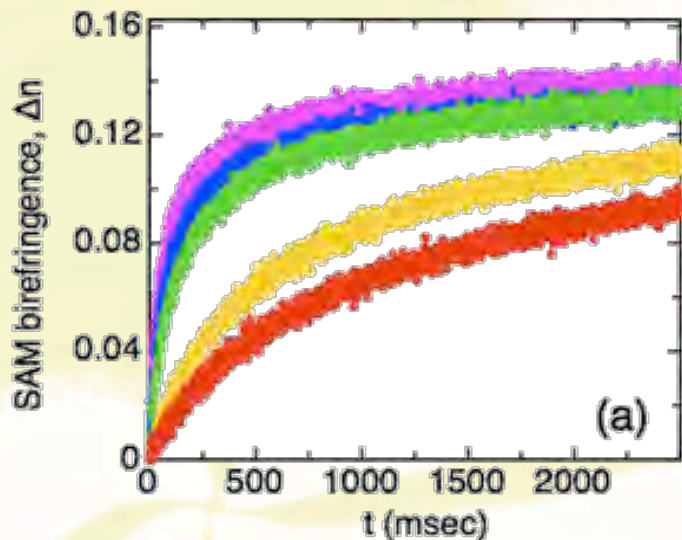


t=0



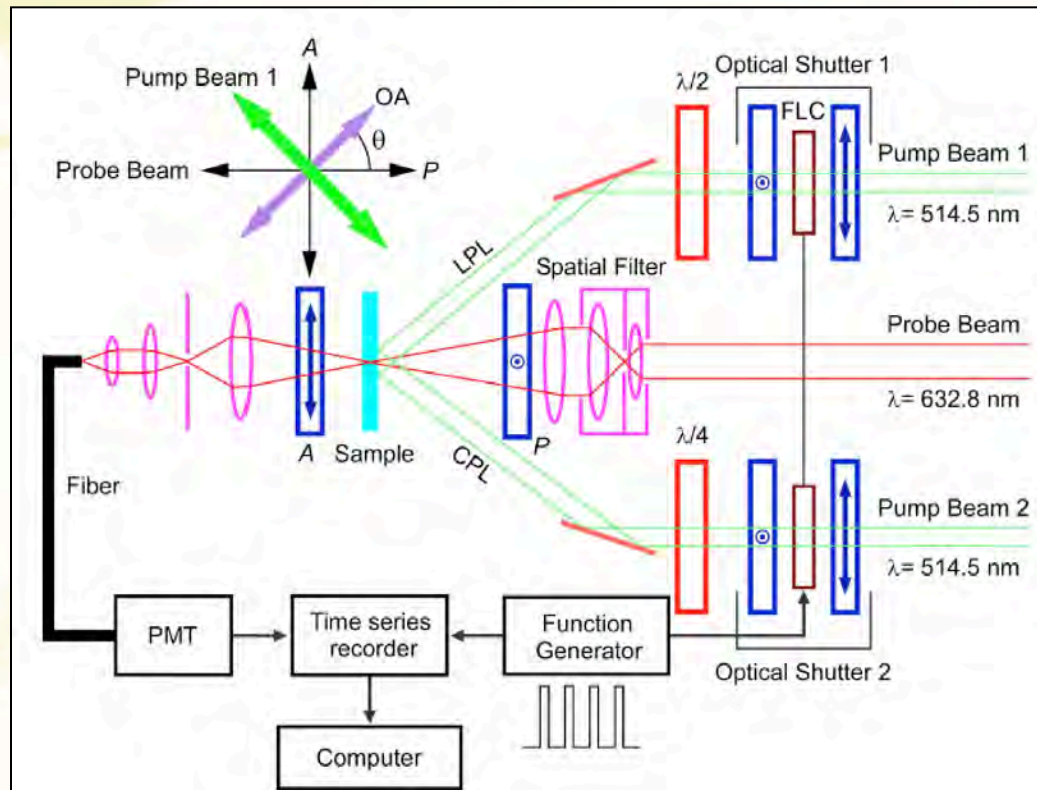
After 9ns

writing dynamics

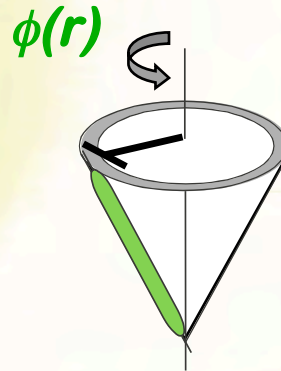
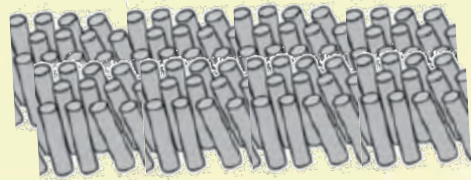


dynamic measurement of in-plane birefringence

- ◆ red probe light (632 nm, focused to 40 μ m)
- ◆ green write light (514 nm focused to 1mm)
- ◆ 100 μ sec time resolution
- ◆ high extinction contrast between crossed polarizers (2.4×10^{-10})

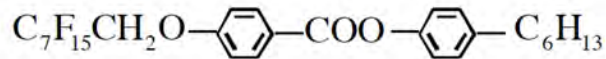


a 2D nematic (smectic C monolayer)



deGennes (1972)

$$\begin{aligned} \langle \Phi_q(0)\Phi_q(\tau) \rangle &= \langle |\Phi_q|^2 \rangle \exp(-Dq^2\tau) \\ &= \frac{k_B T}{Kq^2} \exp(-Dq^2\tau). \end{aligned}$$



$$D = K/\gamma$$

$$\Delta n(t) \sim (t/\tau)^\alpha$$

$$\alpha = T/8\pi K$$

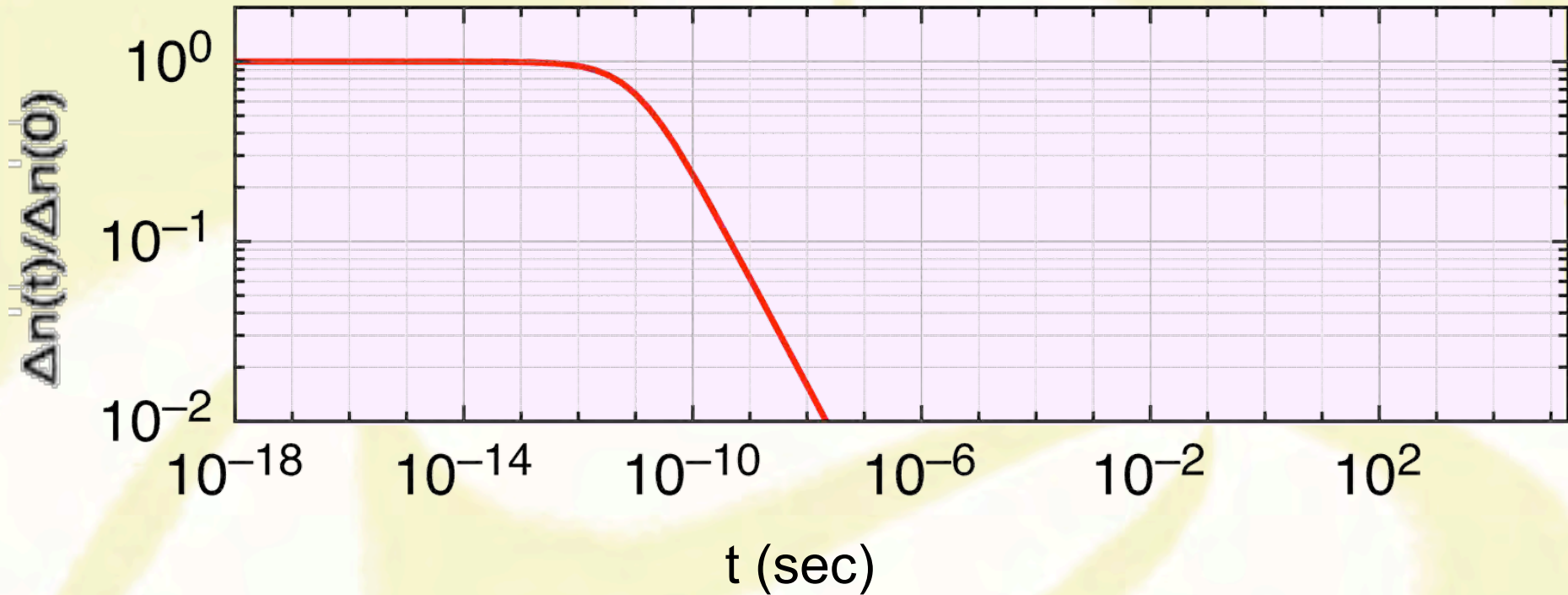
$$\tau = a^2/D = a^2\gamma/K$$

$$\tau \sim 10 \text{ psec}$$

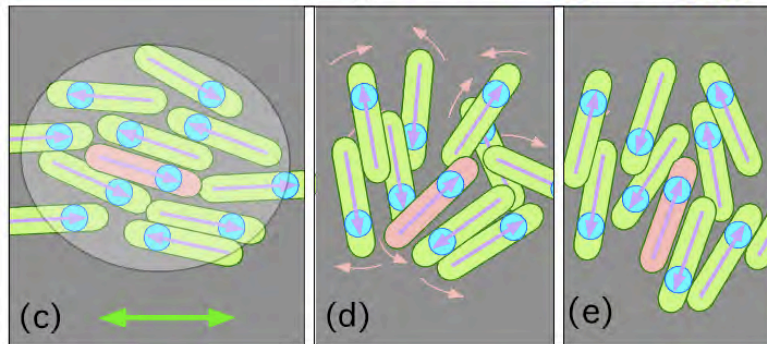
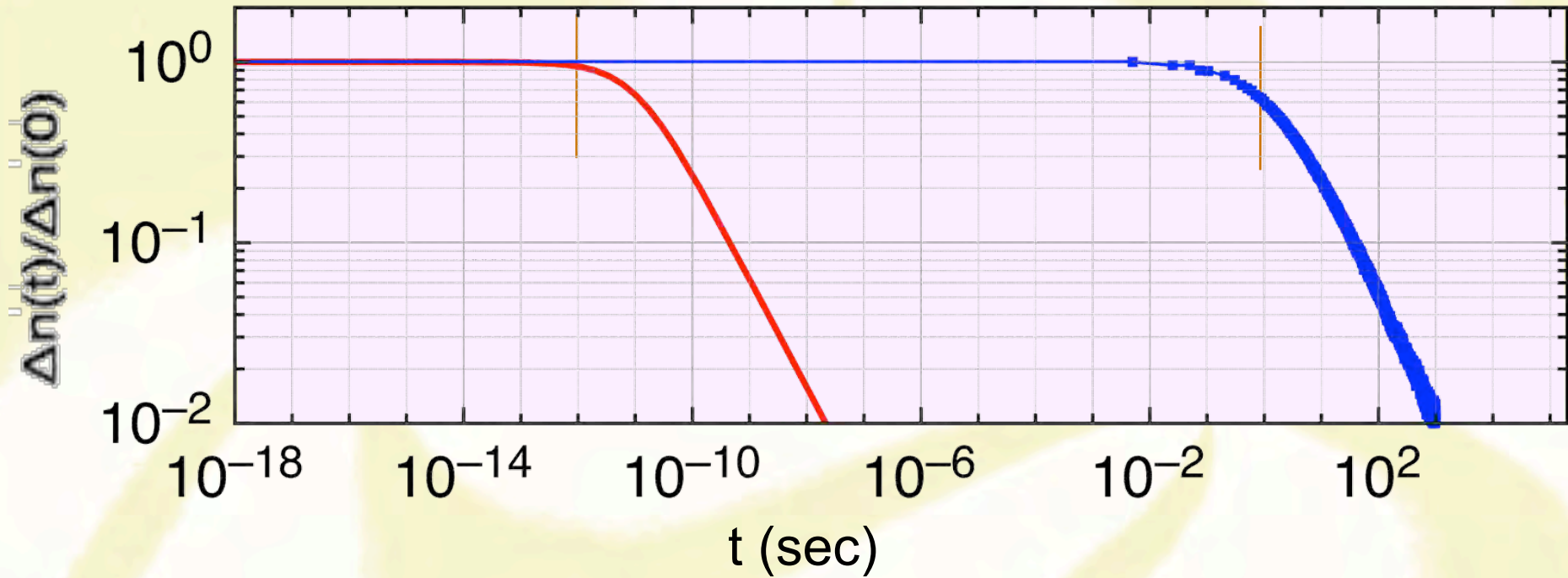
Image_1.png



comparing SmC^* and dMR monolayers

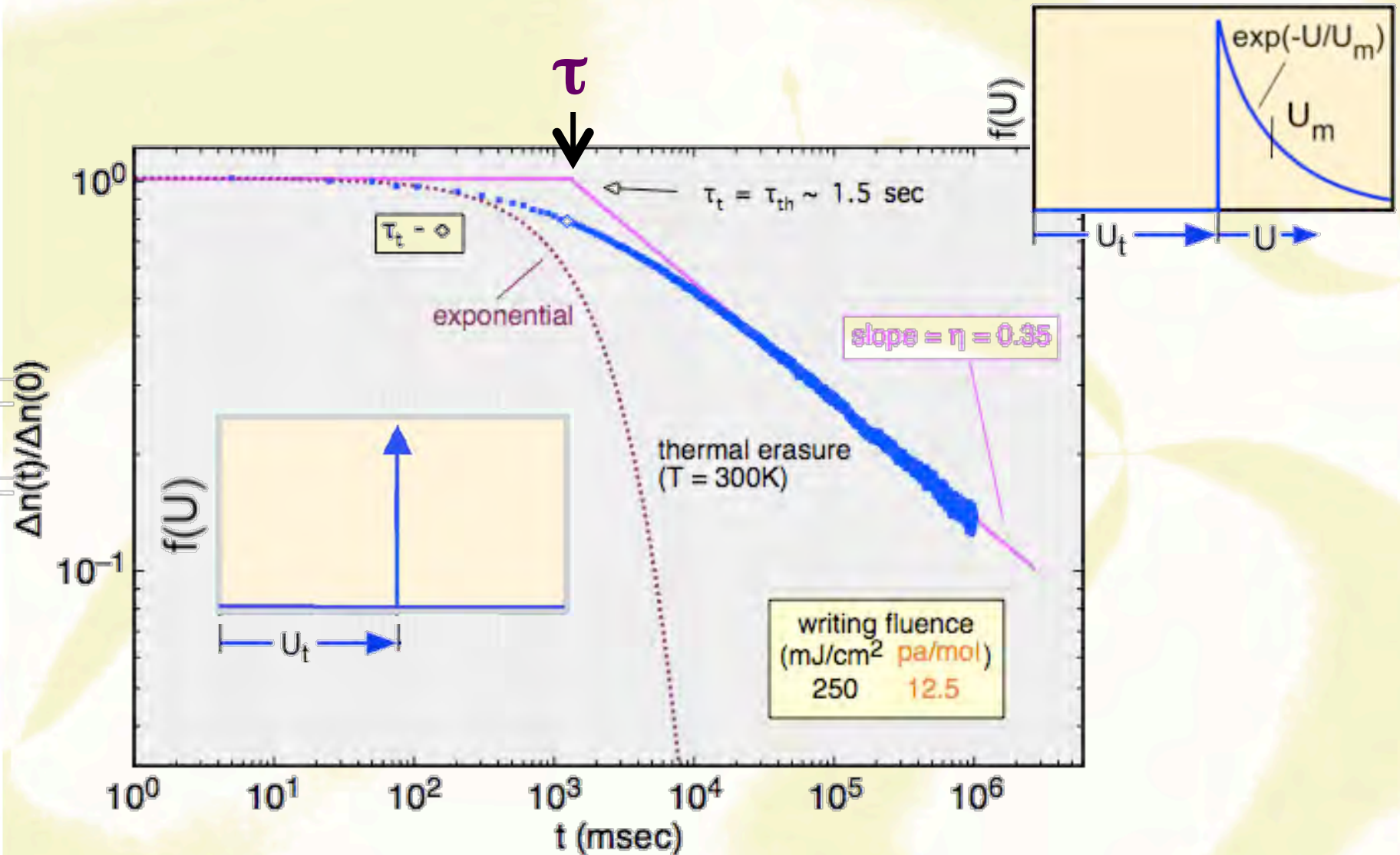


tethered monolayer is a glass



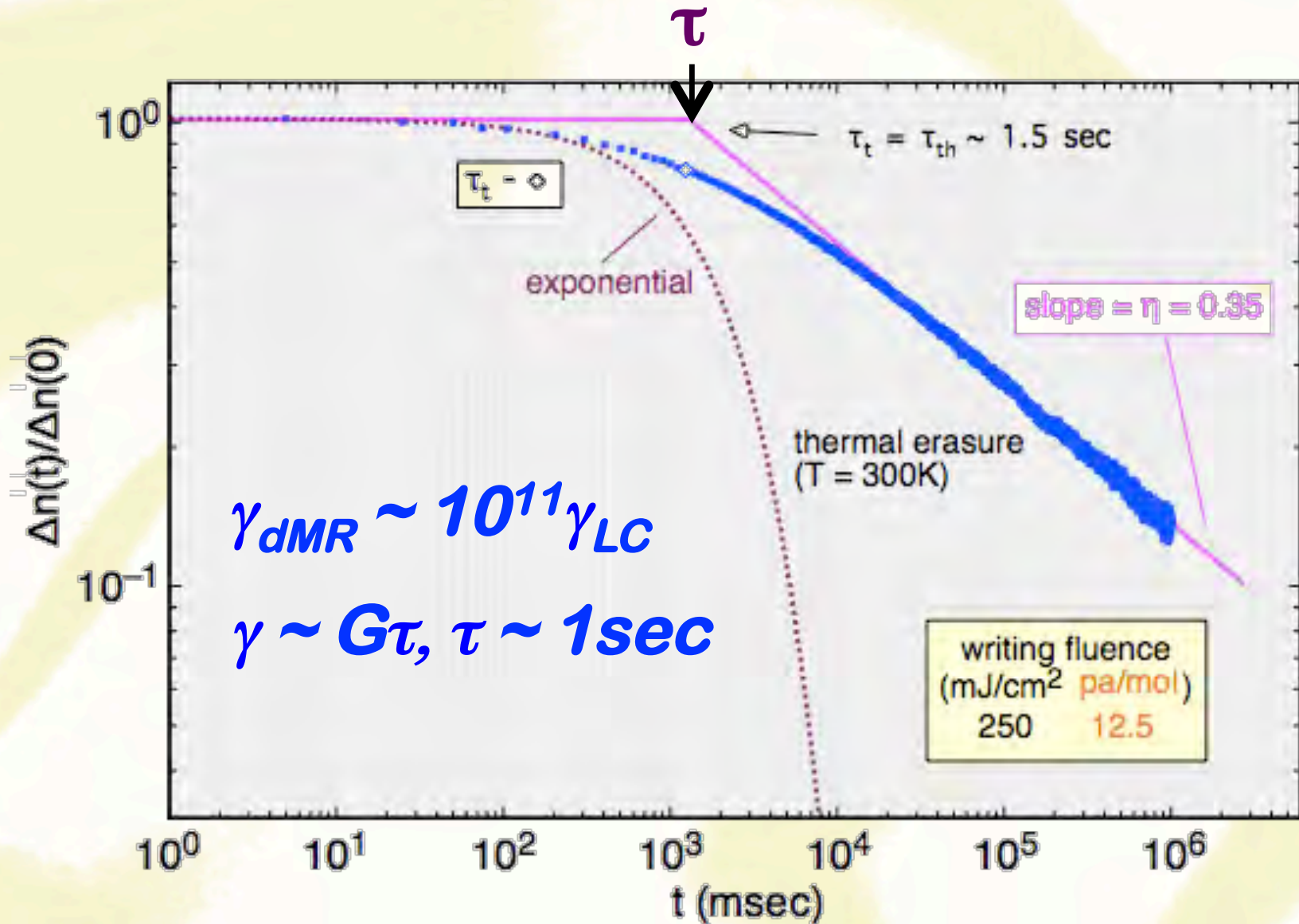
$$\gamma_{DMR} \sim 10^{11} \gamma_{SmC}$$

erasure by $k_B T$ at $T = 300K$



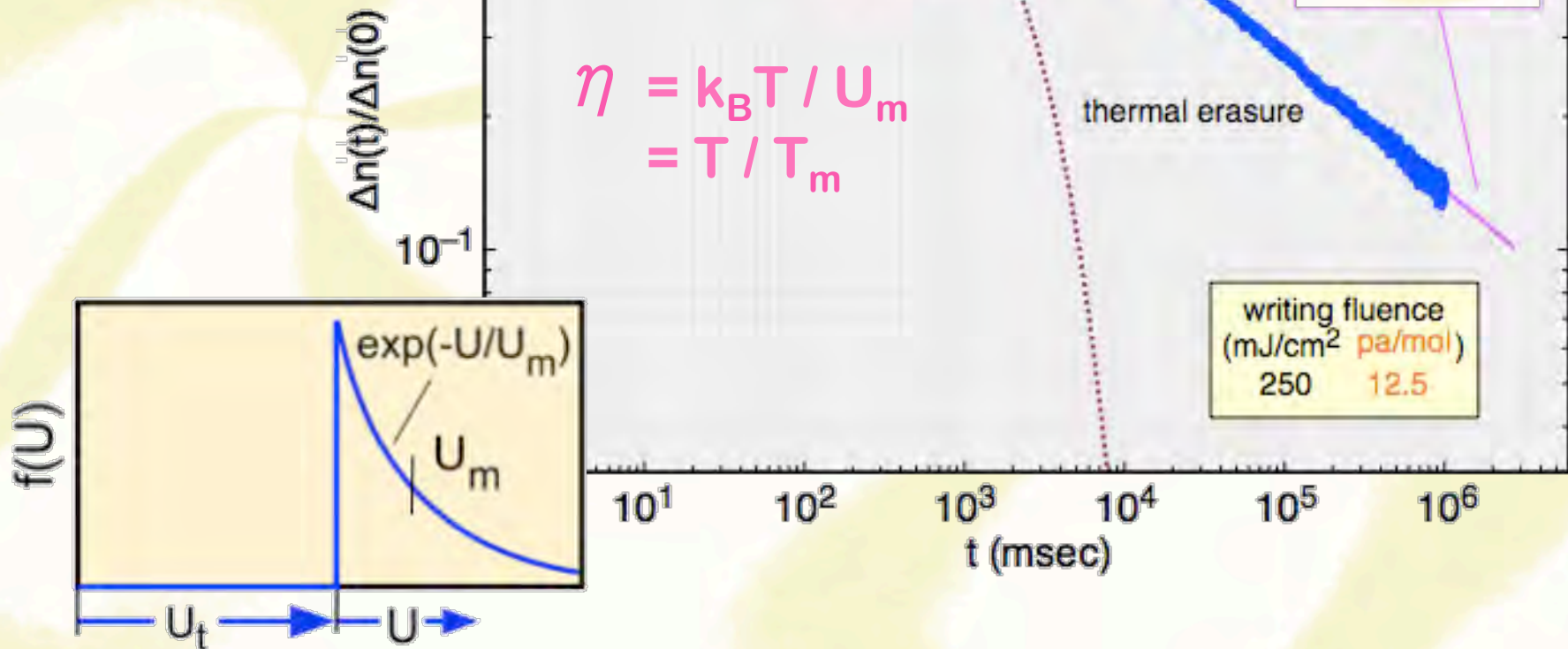
$$U_t \sim k_B T \ln(10^{11}) = 7,500K$$

dMR monolayer is an orientational glass



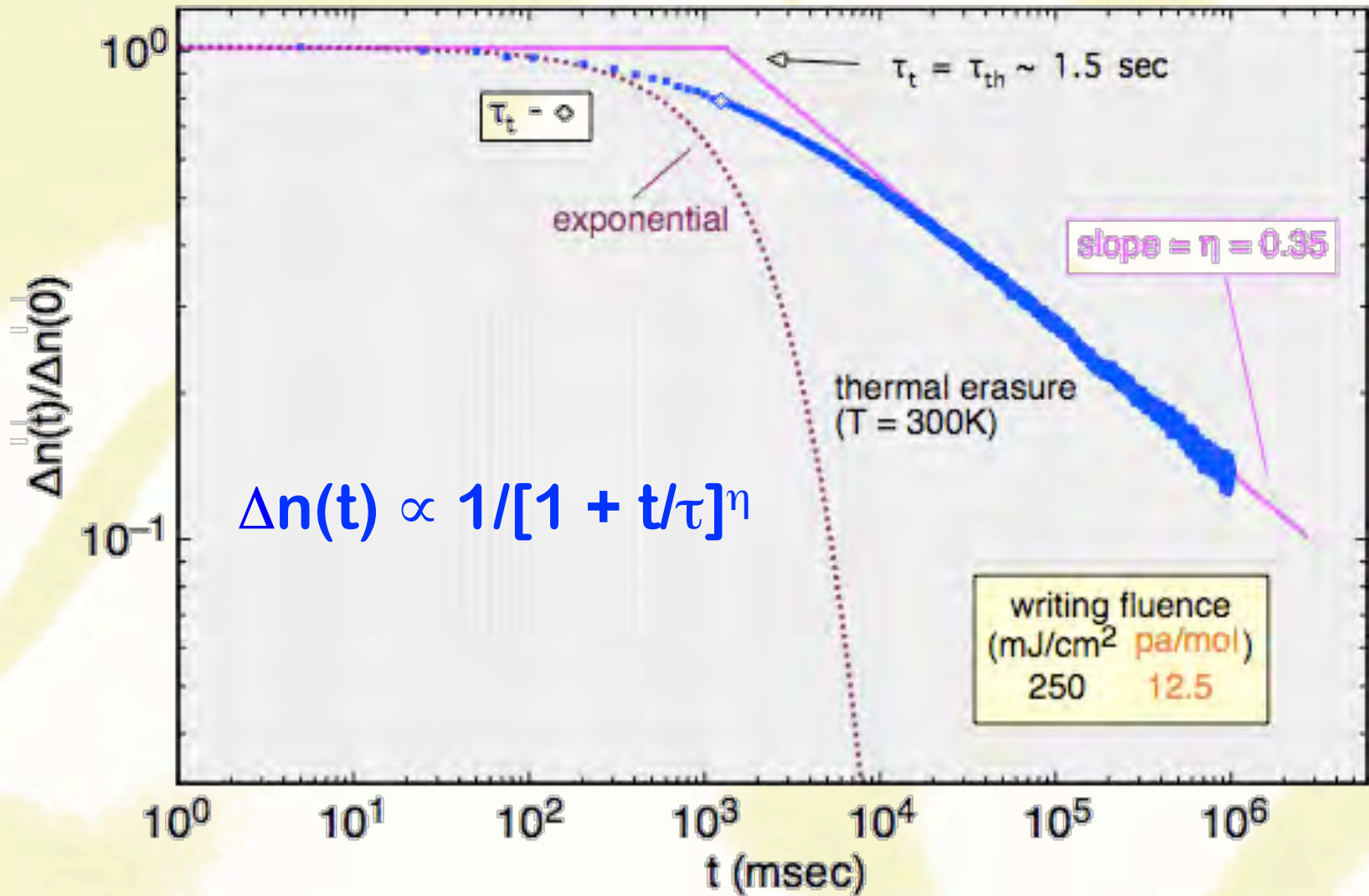
The barrier gap (activation energy)

$$\tau(U) = \tau_r \exp(U / k_B T)$$

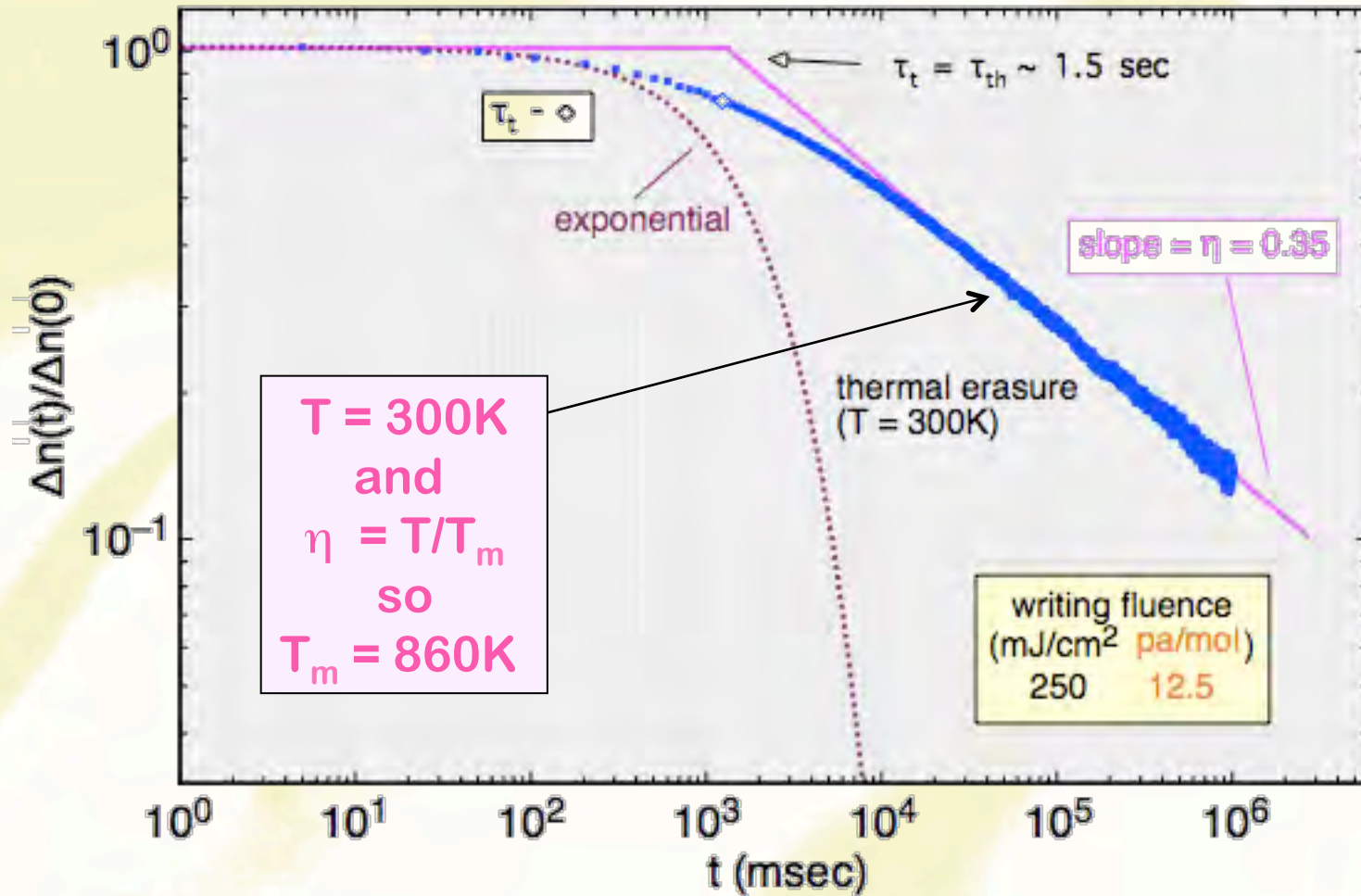


$$U_t \sim k_B T \ln(10^{11}) = 7,500\text{K}$$

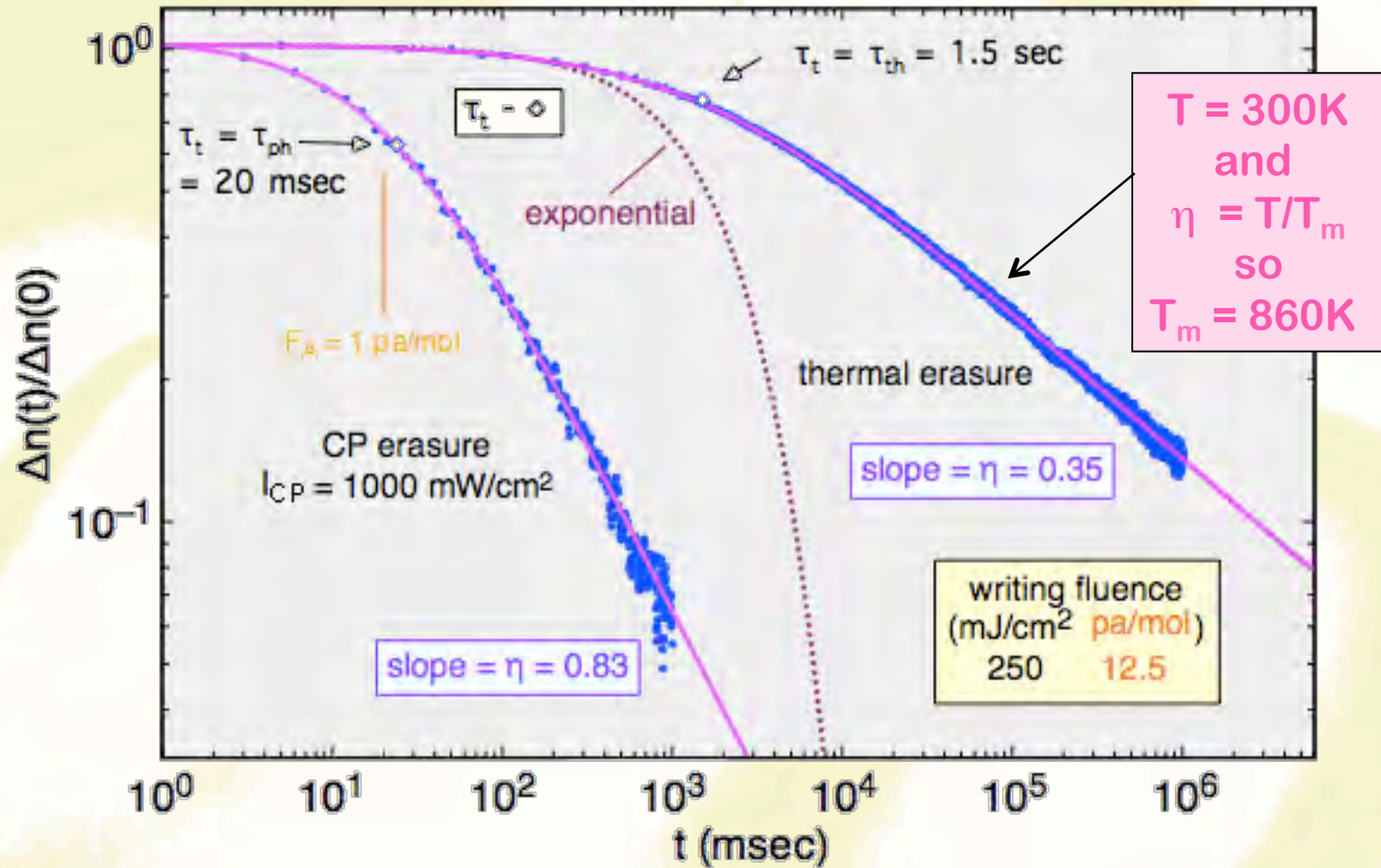
erasure by $k_B T$ at $T = 300K$



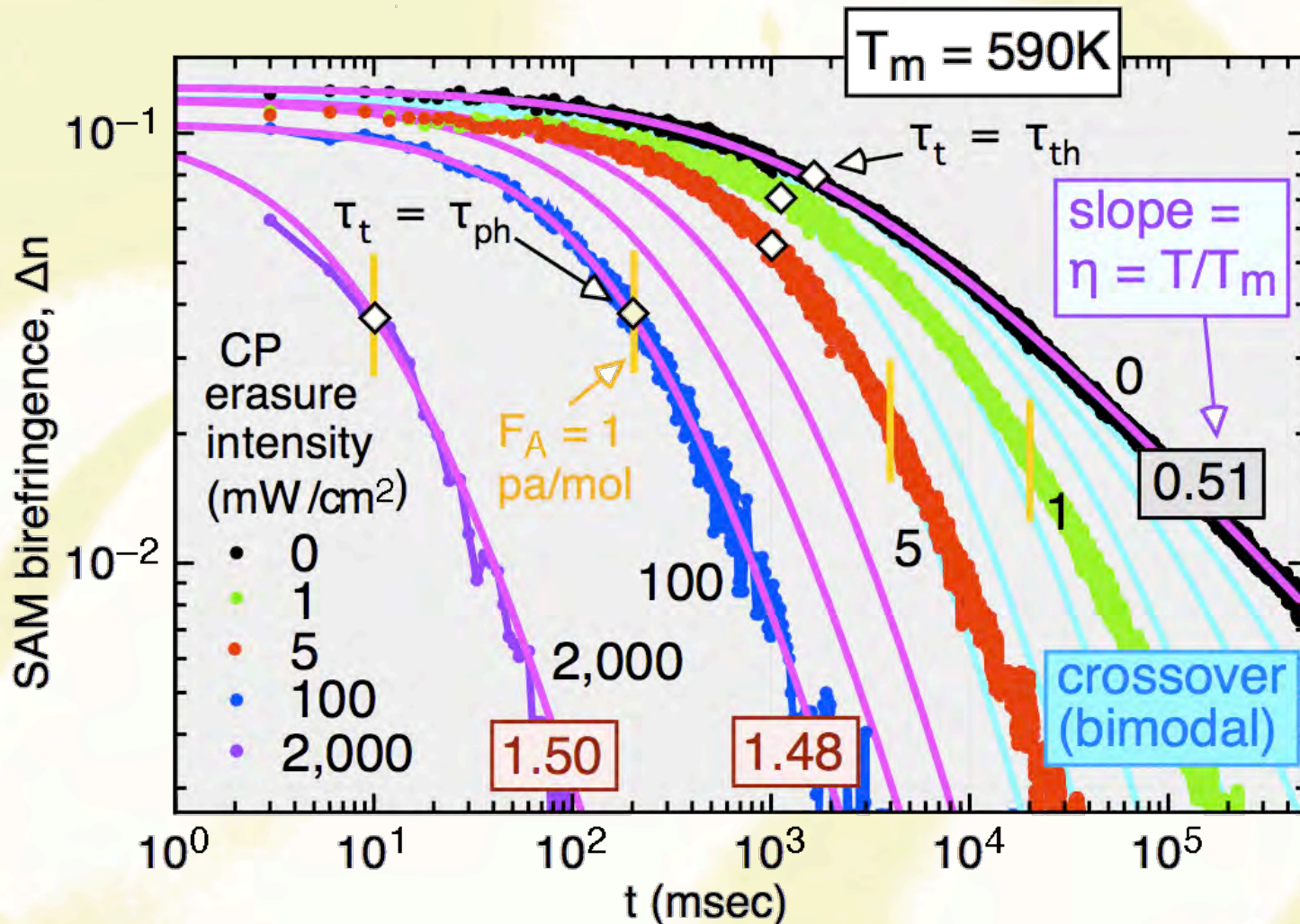
the power law tail



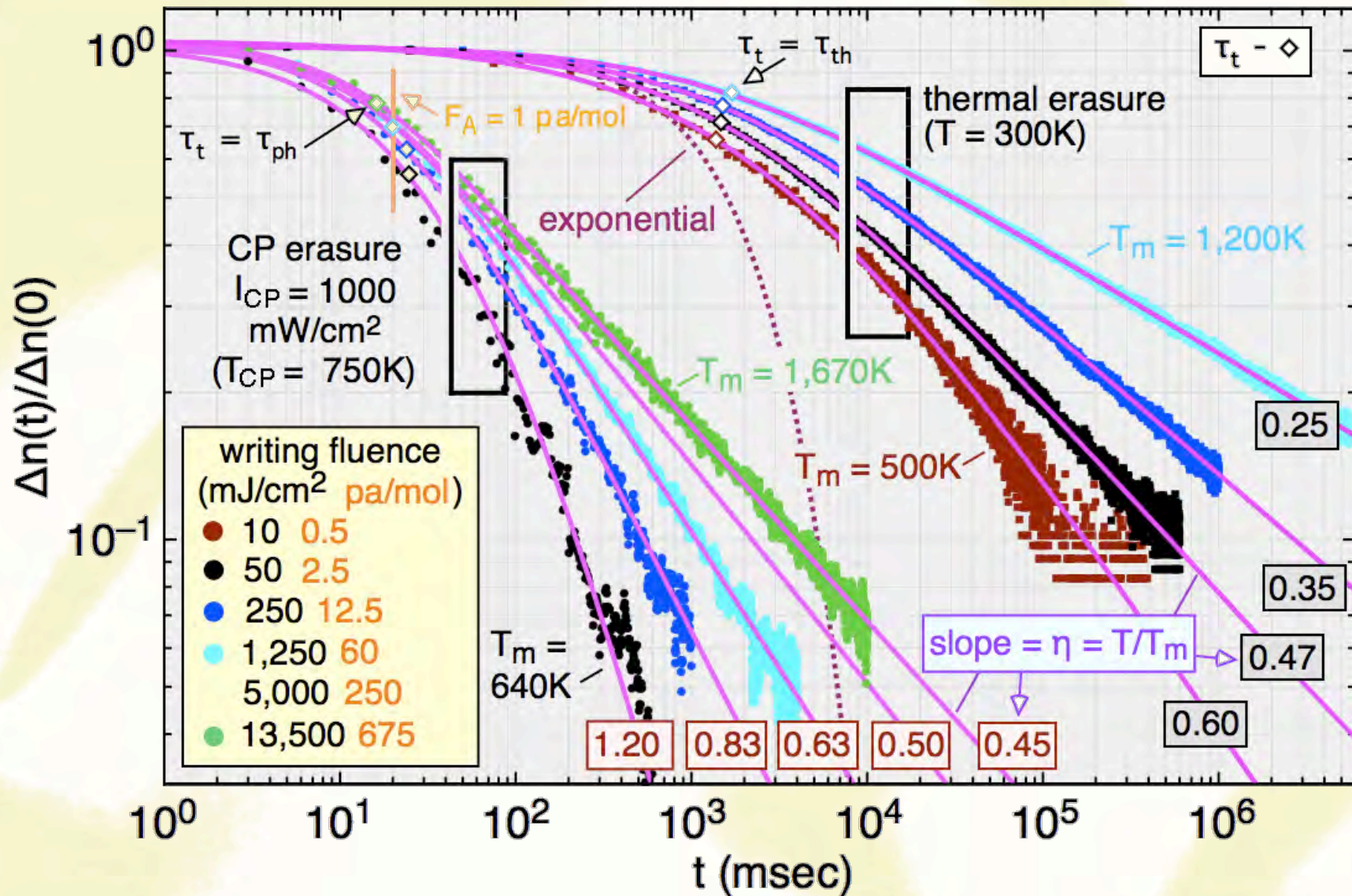
erasure by circularly polarized light



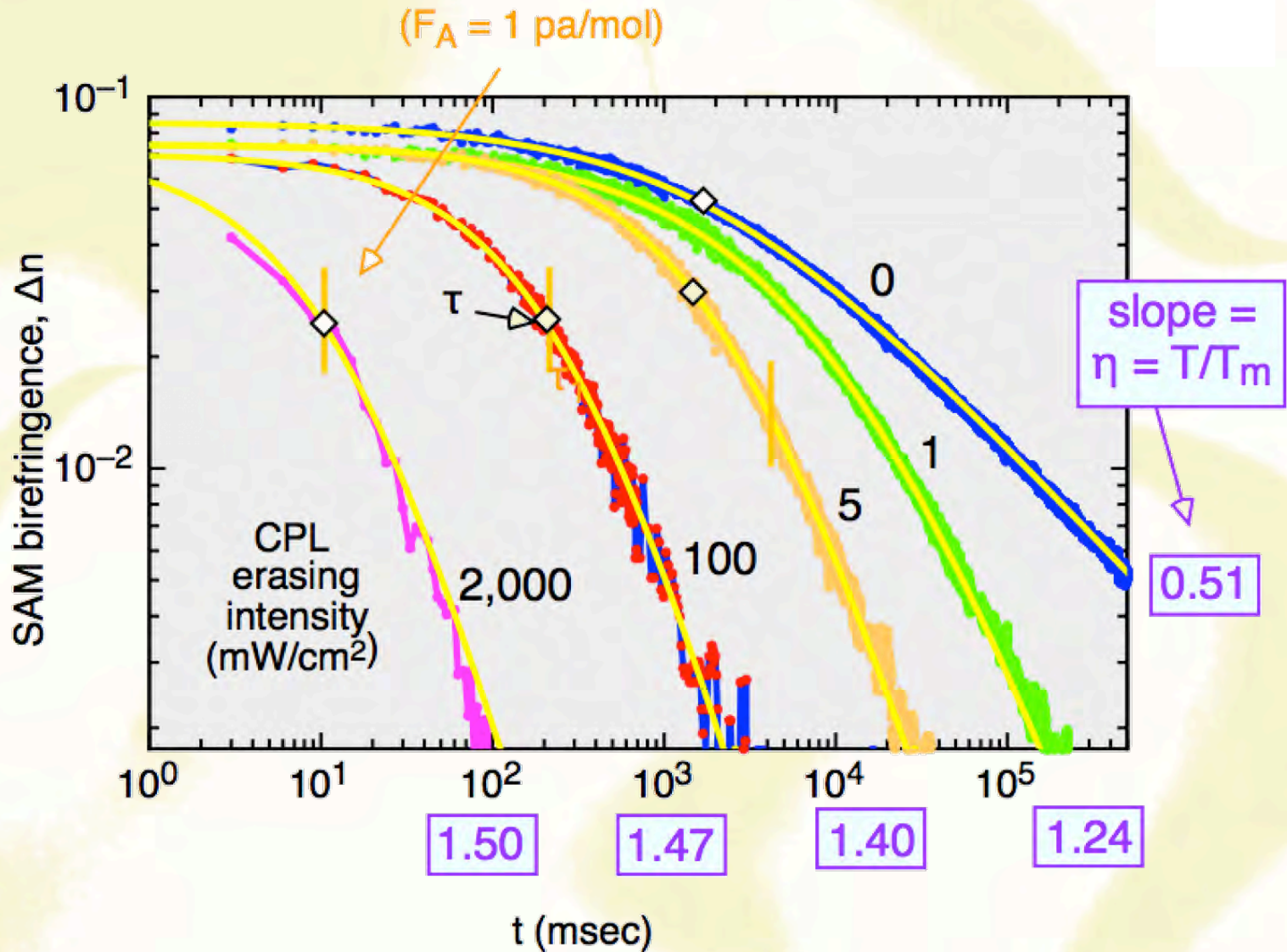
crossover to photo-controlled erasing



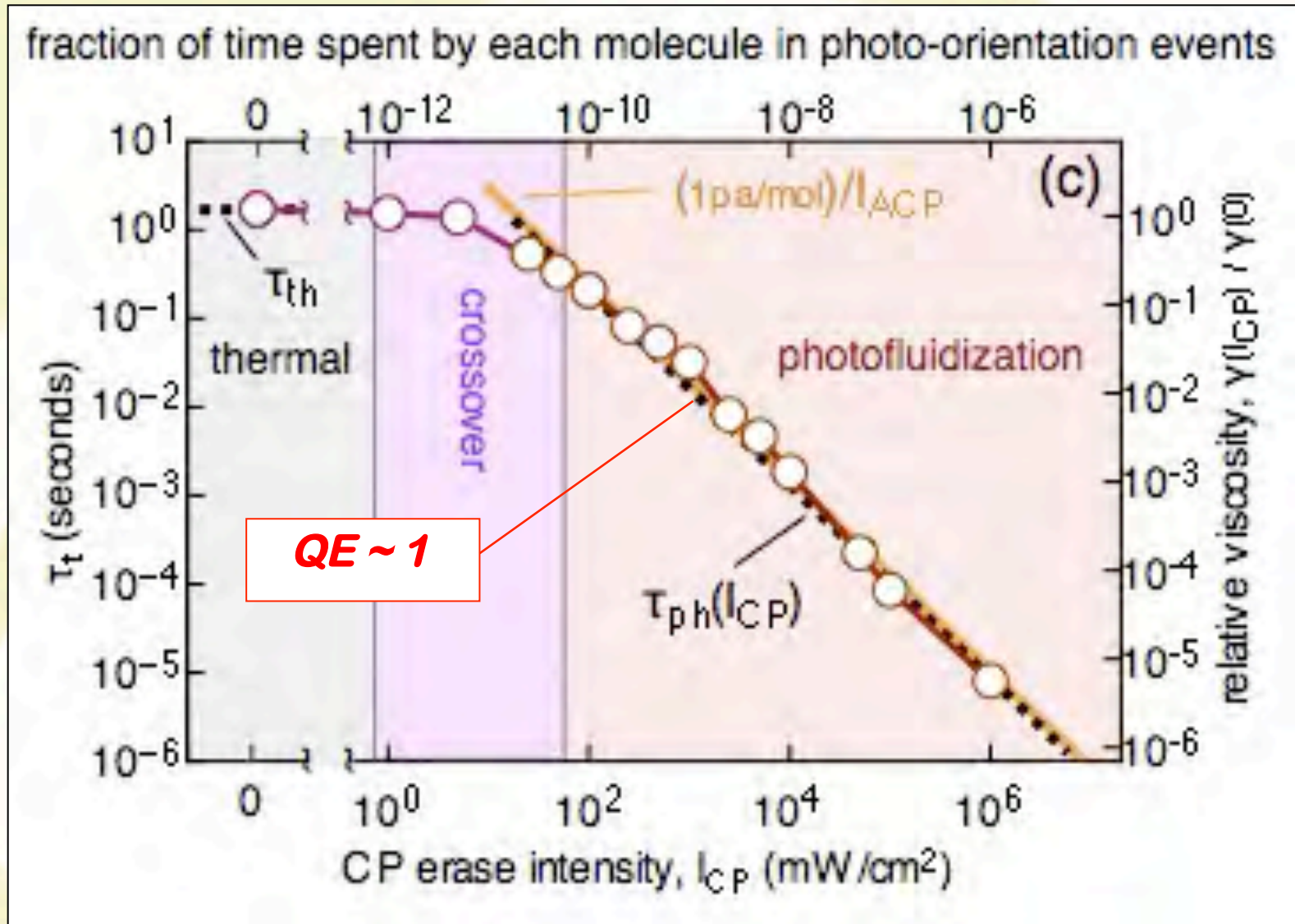
T_m increases with writing fluence



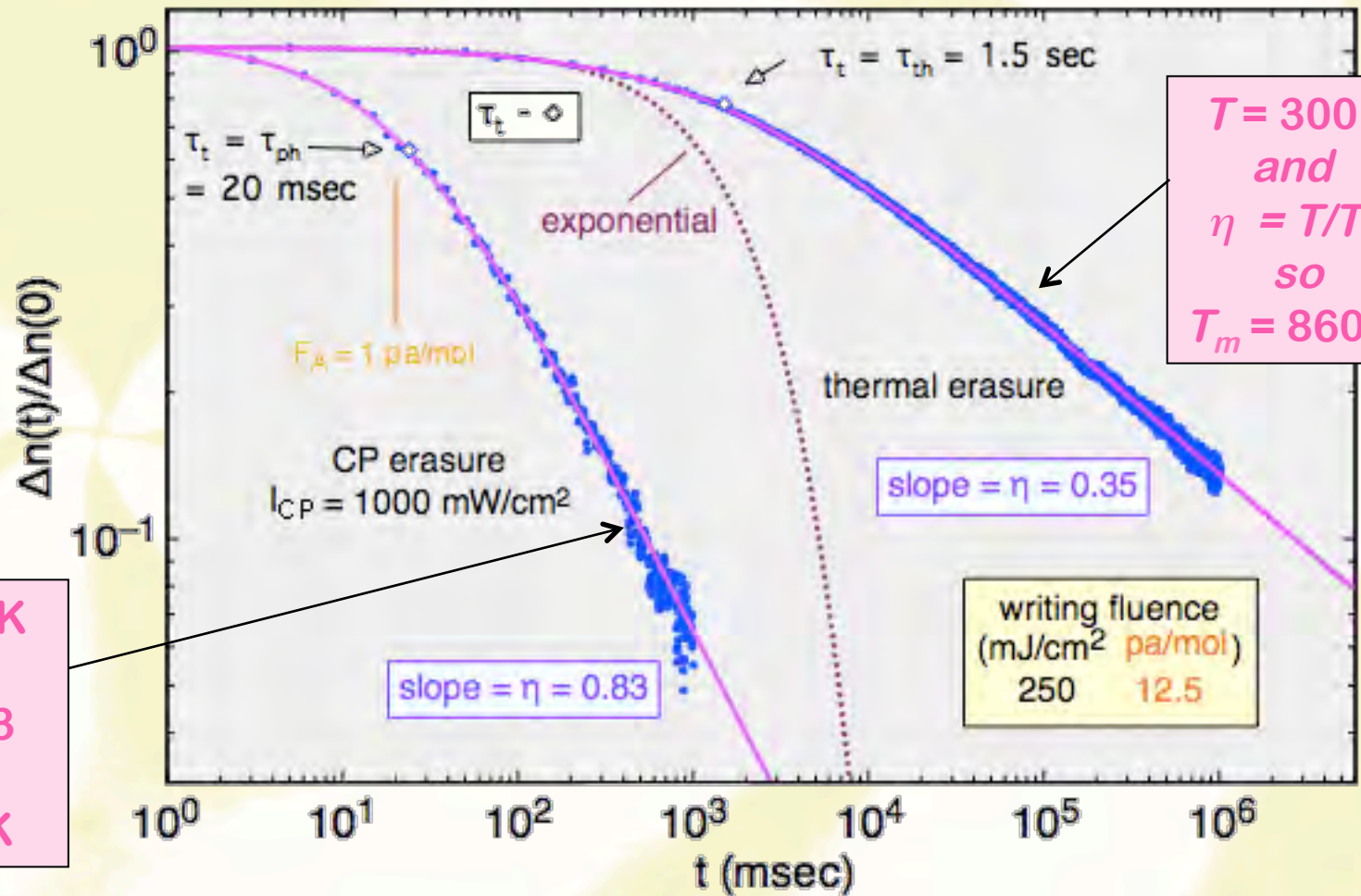
increasing erase intensity



τ occurs at $F = 1$ photon absorbed/molecule



erasure by circularly polarized light



$T = 300$ K
and
 $\eta = T/T_m$
so
 $T_m = 860$ K

$T_m = 860$ K
and
 $\eta = 0.83$
so
 $T = 710$ K

every absorbed photon induces a local glass transition

$$U_t = 7,500$$

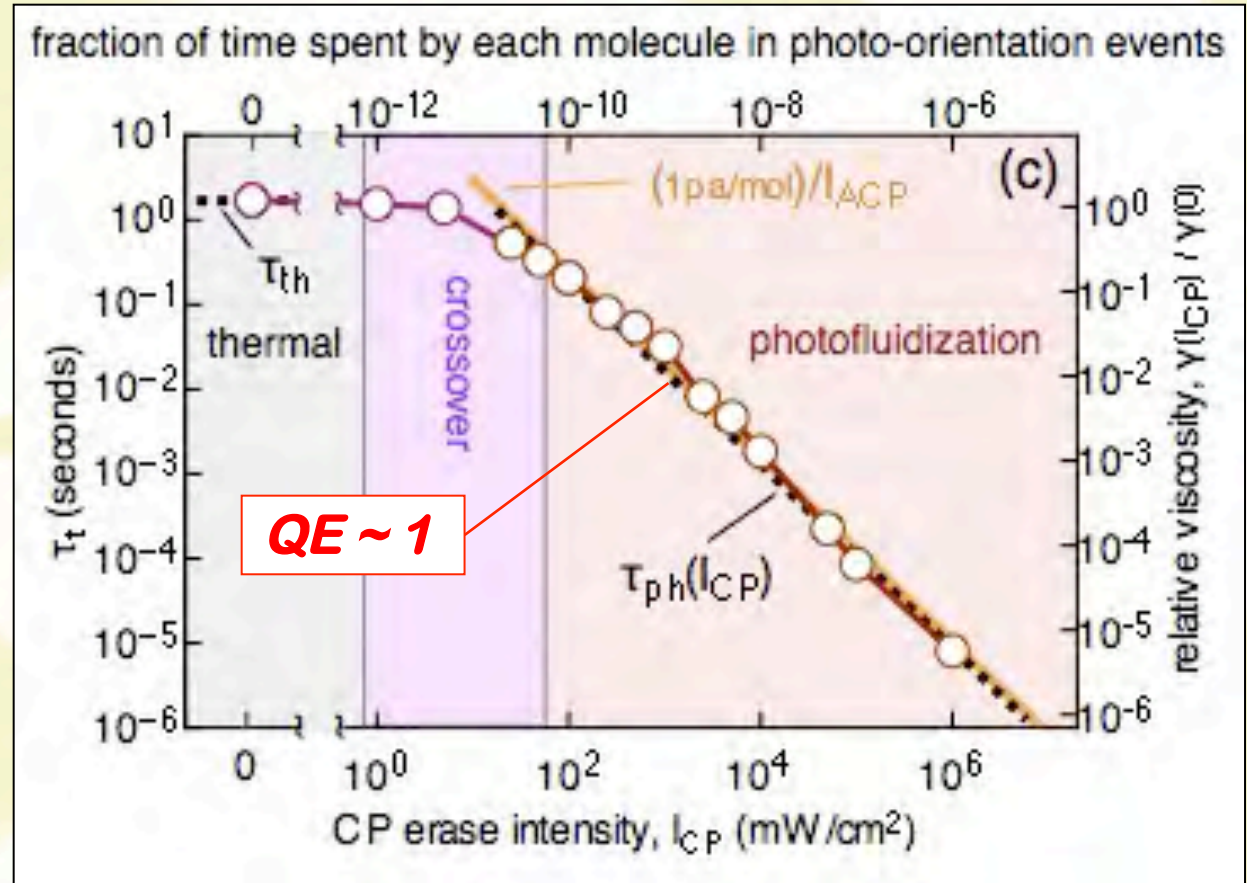
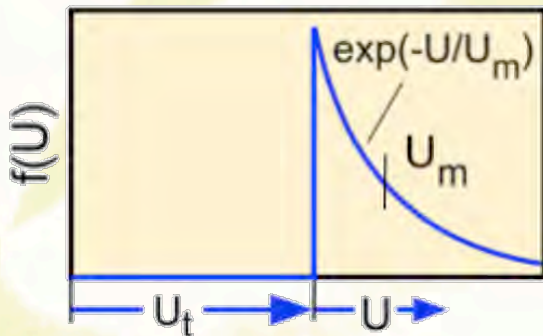
$$T = 800K$$

gives

$$QE \sim \exp(-7,500/800)$$

but we find

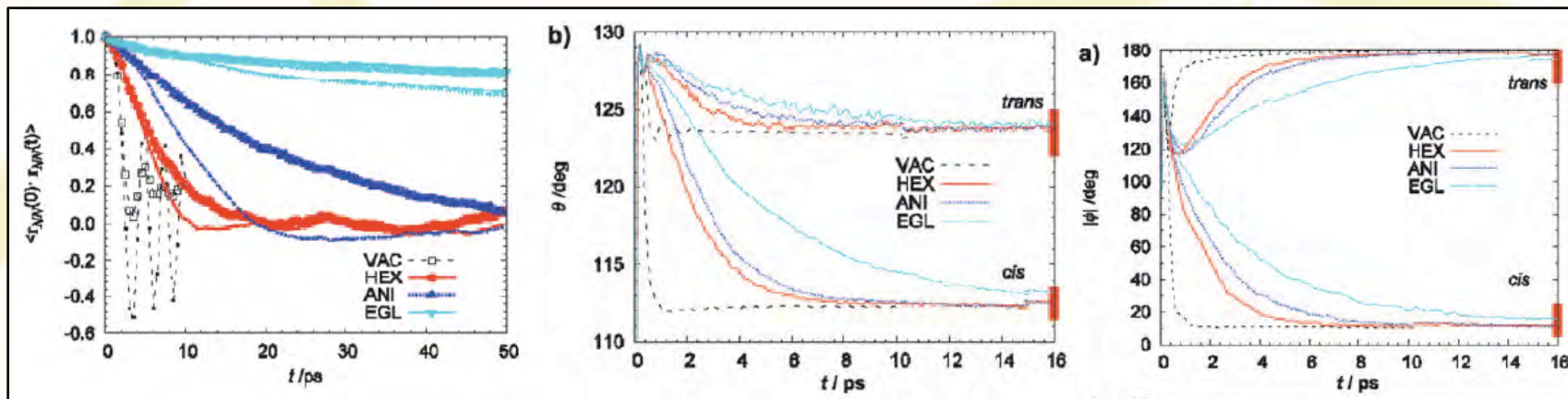
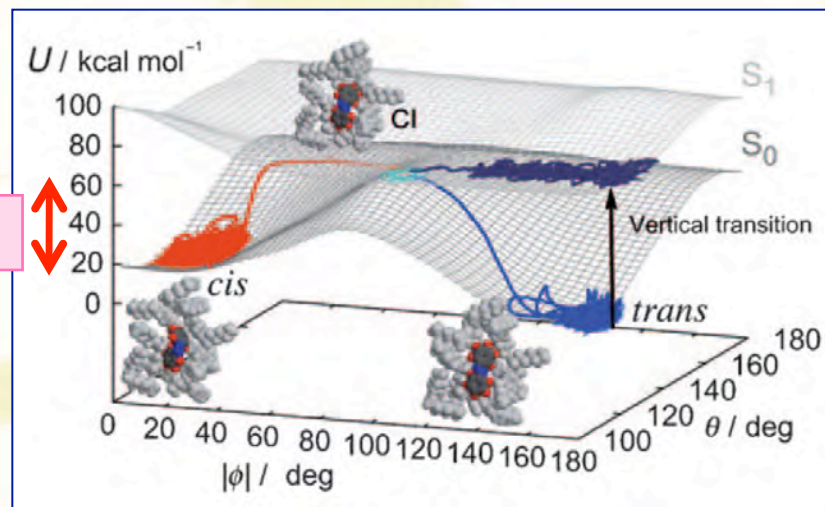
$$QE \sim 1!$$



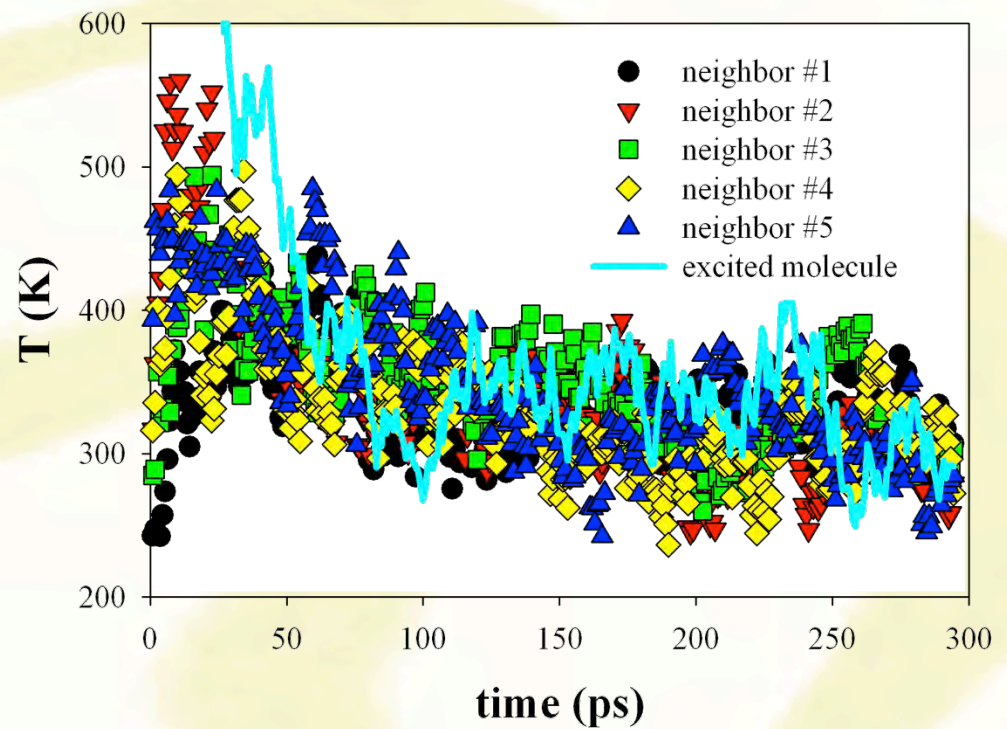
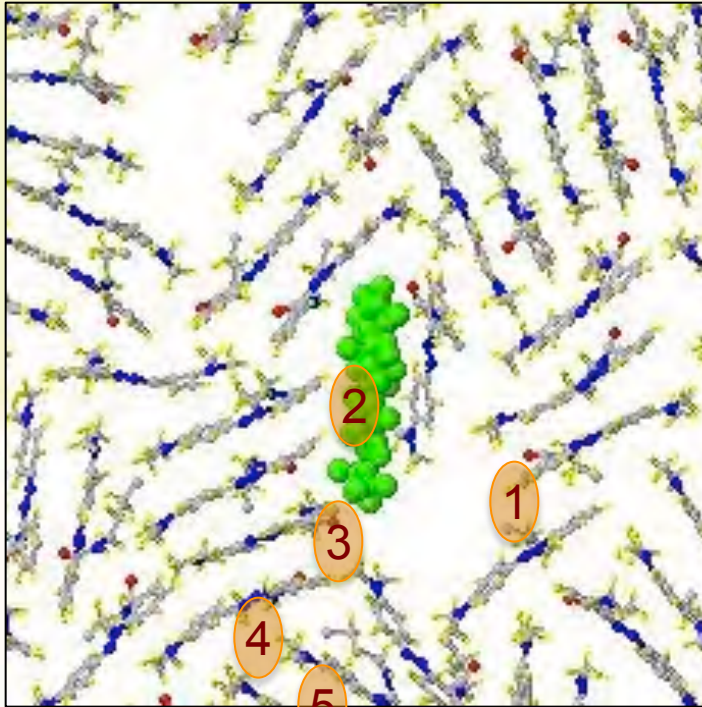
This means that each 800K photo event melts (fluidizes) the local 7,500K barrier, i.e. produces a local glass transition

the photo-absorption event

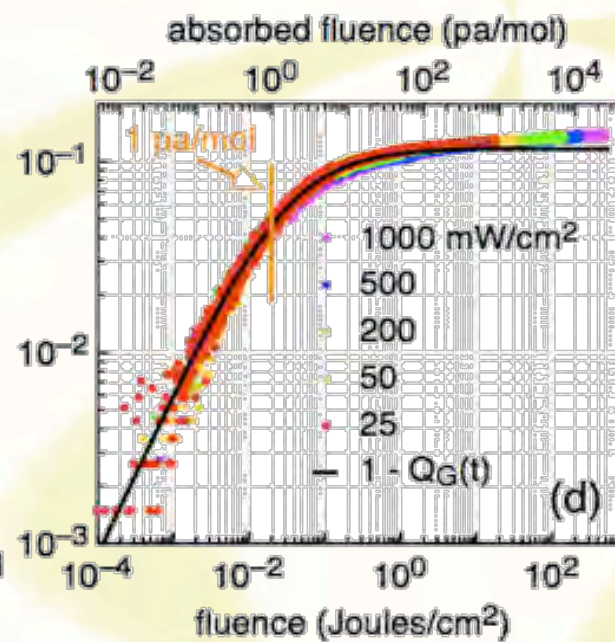
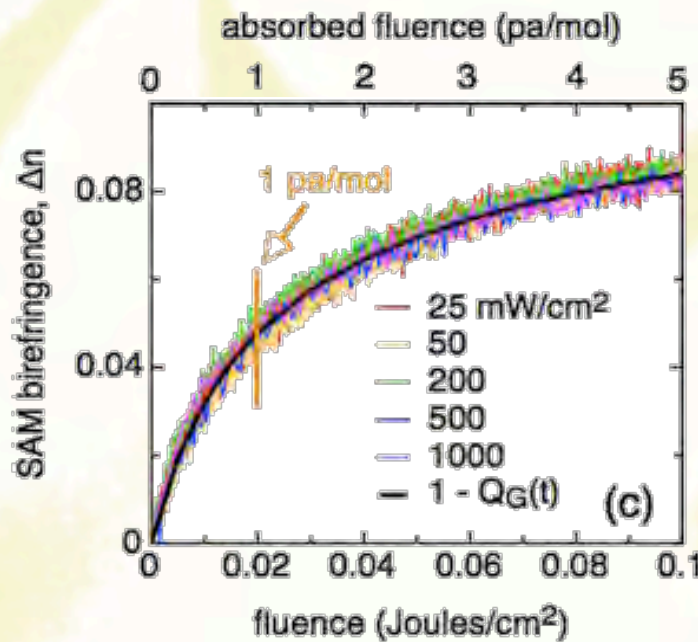
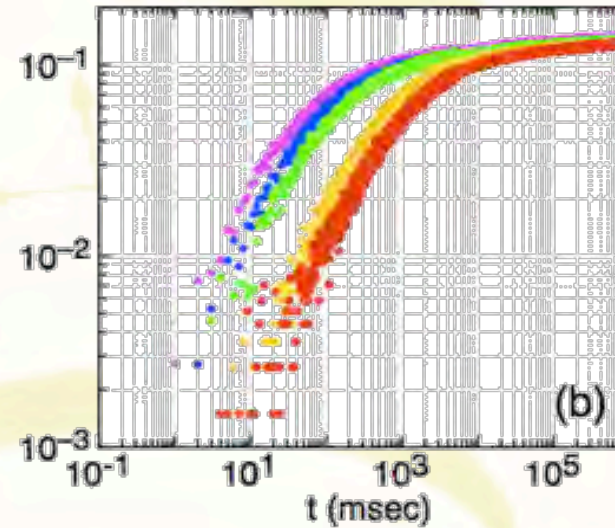
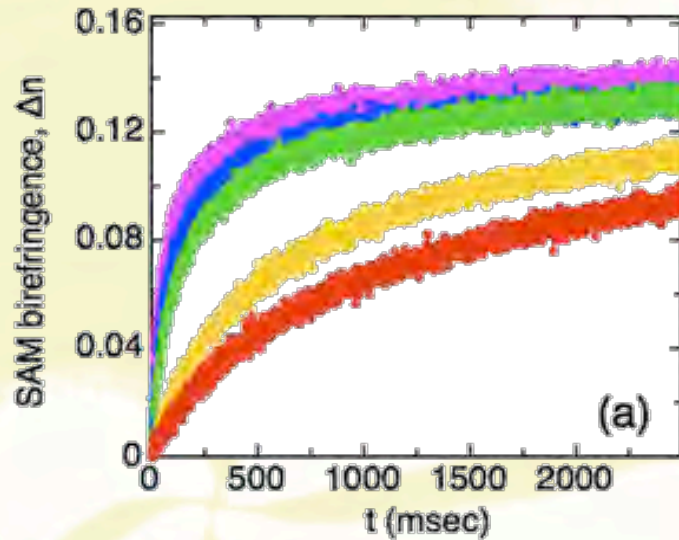
- ◆ $h\nu = 2.4 \text{ eV}$
- ◆ $T = h\nu/k_B = 29,000\text{K}$
- ◆ each molecule:
~100 motional degrees of freedom



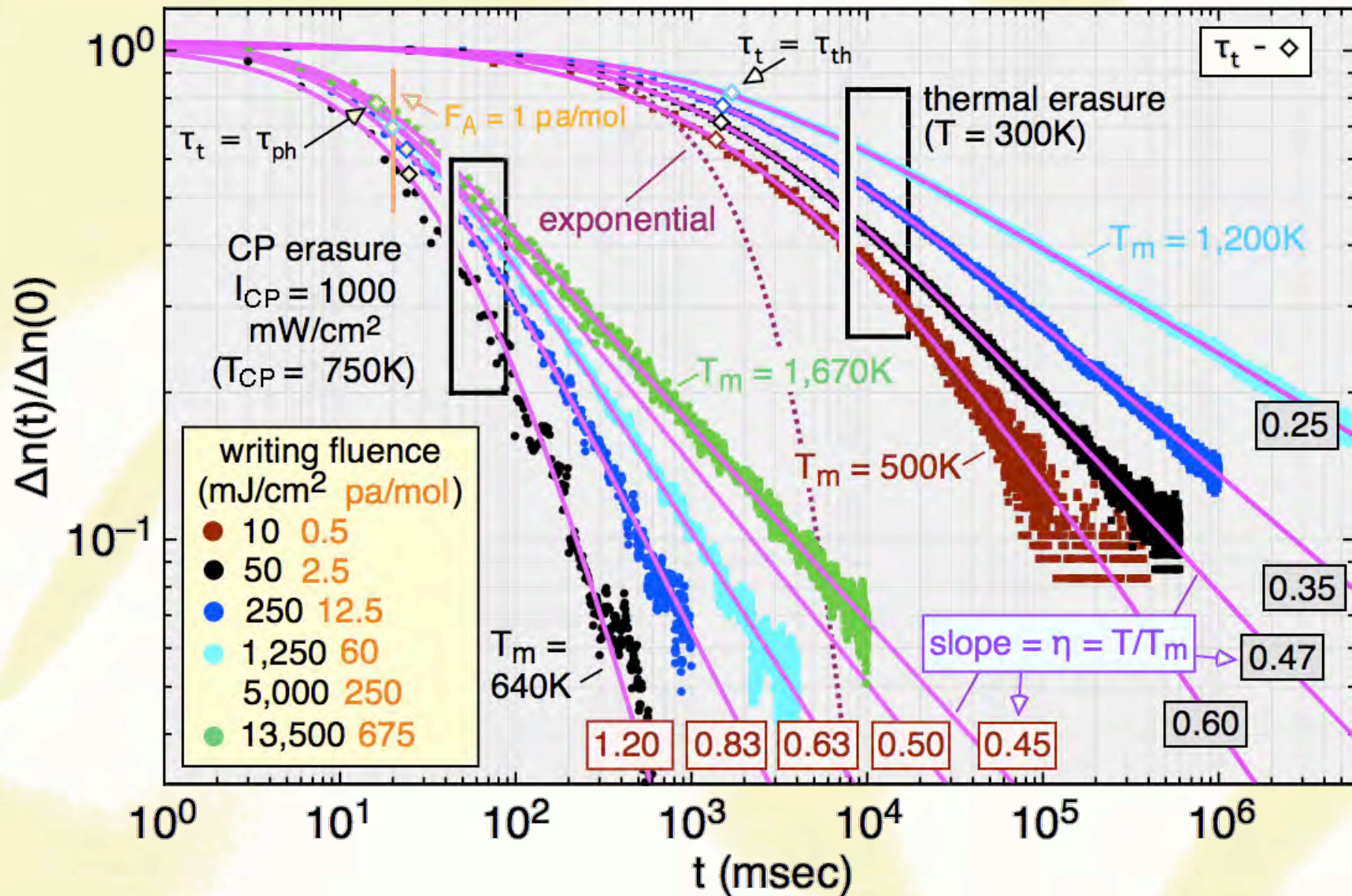
the photo-absorption event



writing dynamics



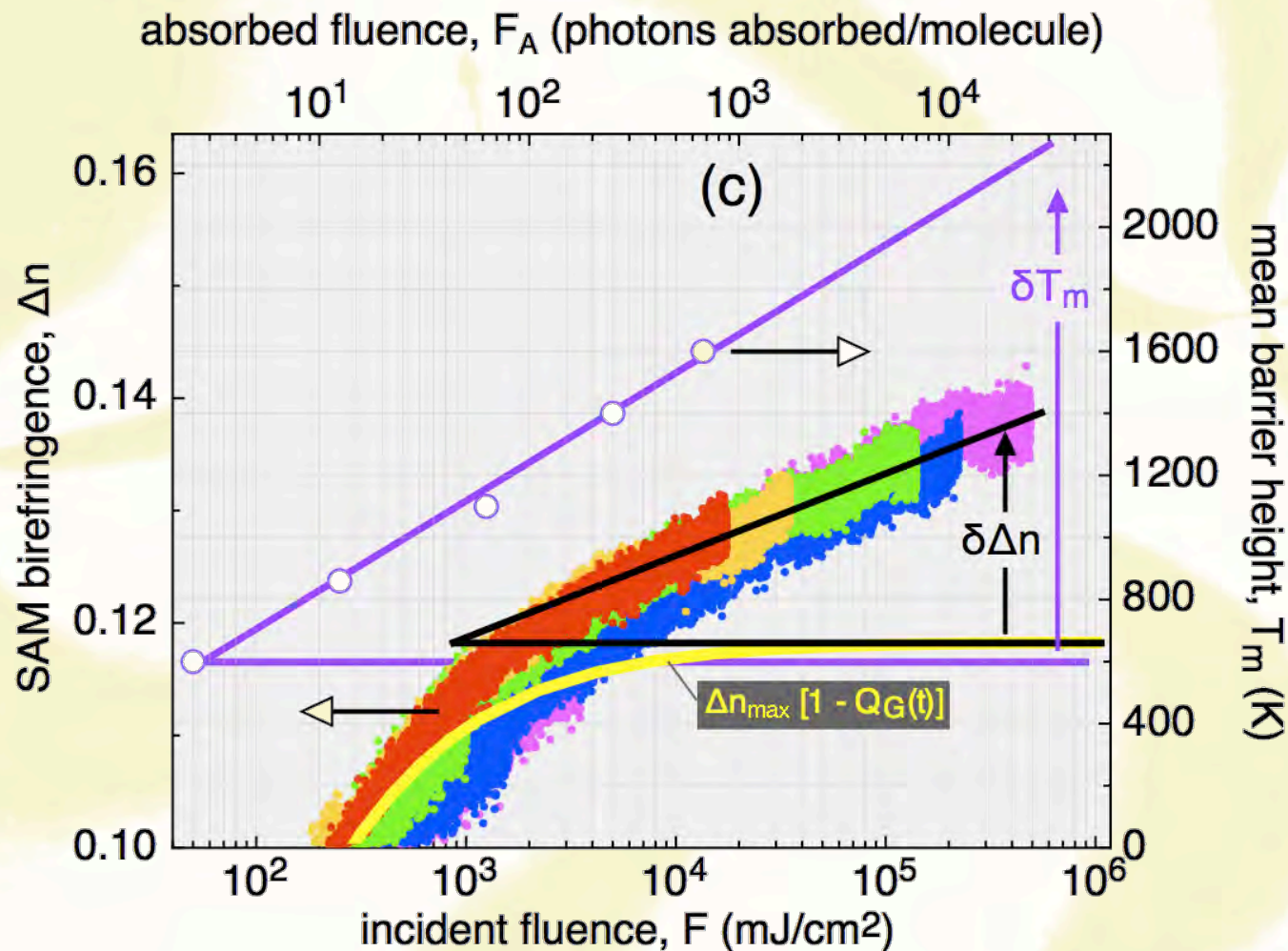
T_m increases with writing fluence



orientational work hardening at large fluence

$$\delta\Delta n \sim \log(t)$$

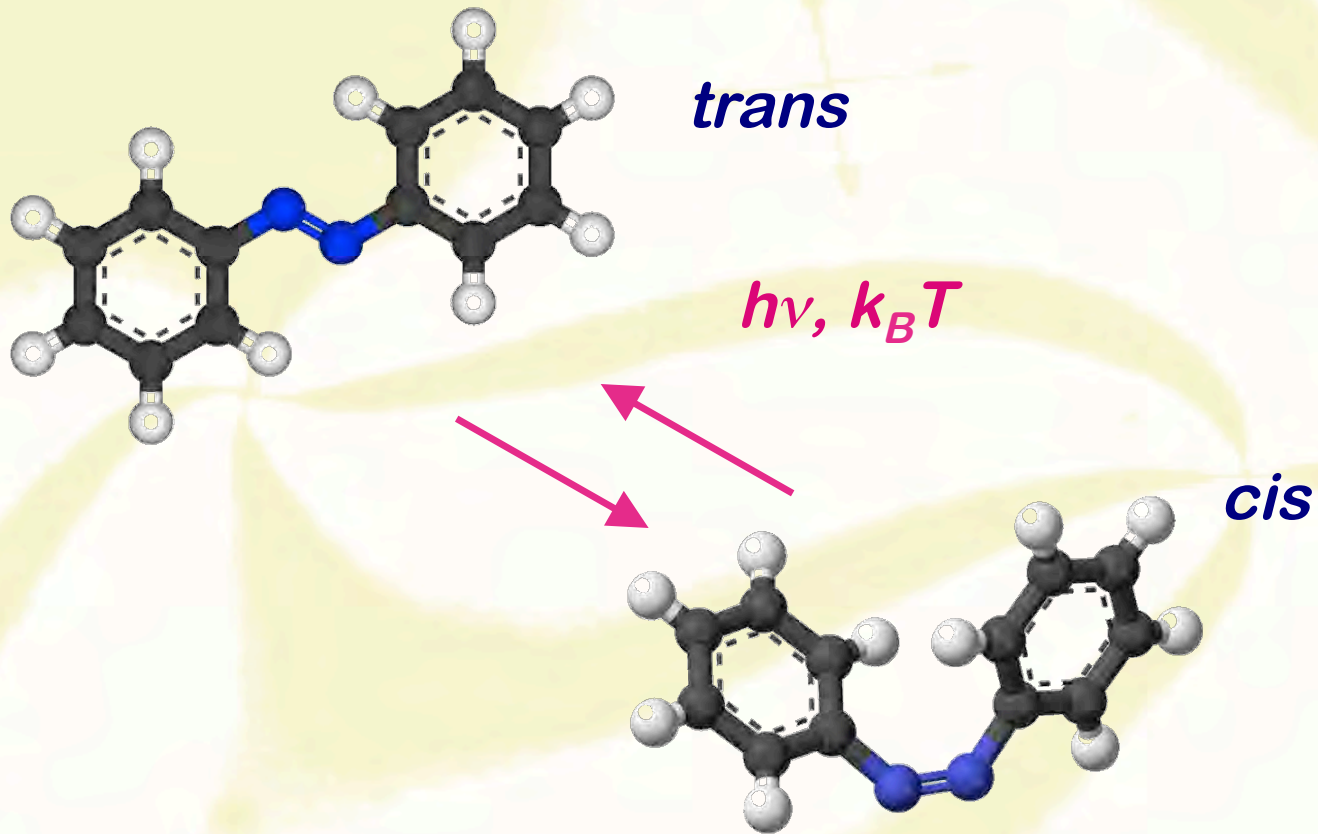
$$\delta T_m \sim \log(t)$$



$$1 \approx P(t) = t\nu_w \exp(-\delta T_m(t)/T)$$

Nabarro Mat. Sci. and Eng. A (2001)

azobenzene



summary

- ◆ *photo-induced relaxation is essentially isothermal*
- ◆ *thermal and photo-induced relaxation events occur by random sequences of discrete local relaxation events.*
- ◆ *isomerizing molecules attack their confining barriers with an effective $T \sim 800\text{K}$*
- ◆ *the $T \sim 800\text{K}$ attacks produce local glass transitions, such that every photon generates a barrier escape event*



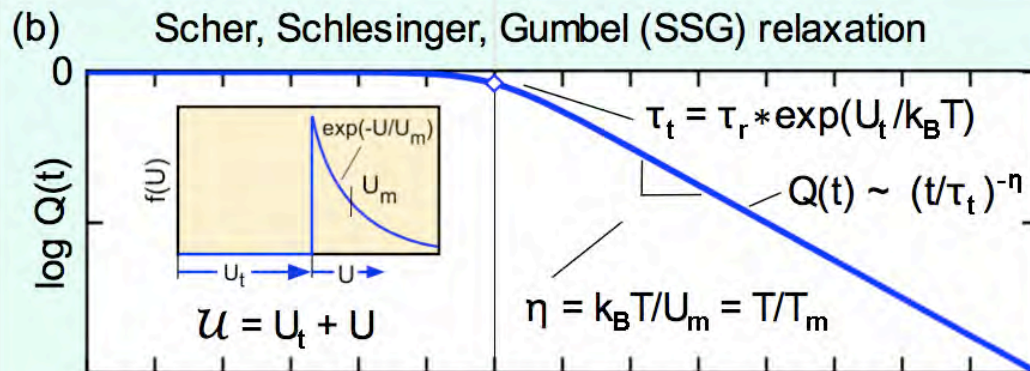
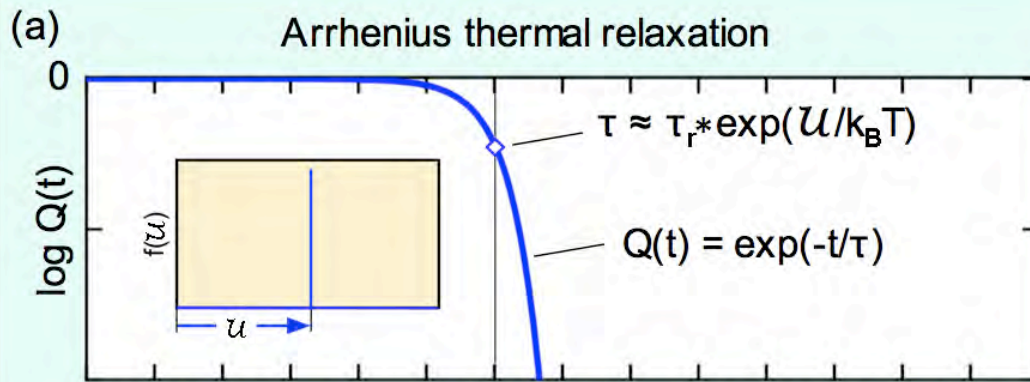
the end



barrier height distributions

$$\tau(U) = \tau_r \exp(U / k_B T)$$

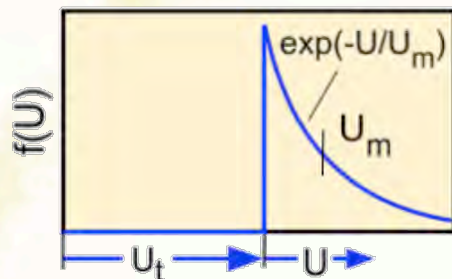
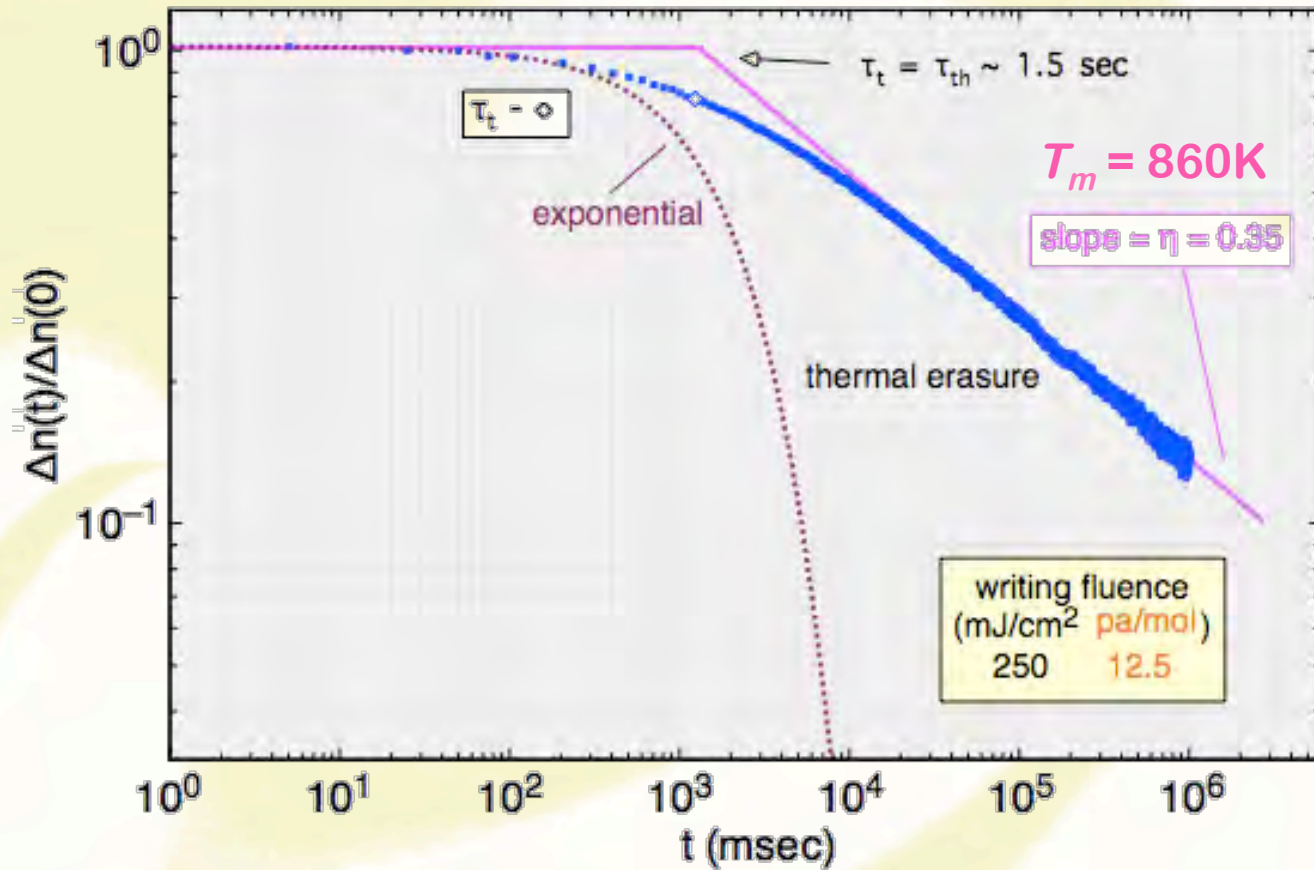
$$Q(t) = A \int_0^{\infty} e^{-(t/\tau)} H(\tau) d\tau = \int_0^{\infty} e^{-(t/\tau(U))} f(U) dU$$



$$\tau = \tau_r \exp(U(t) / k_B T)$$

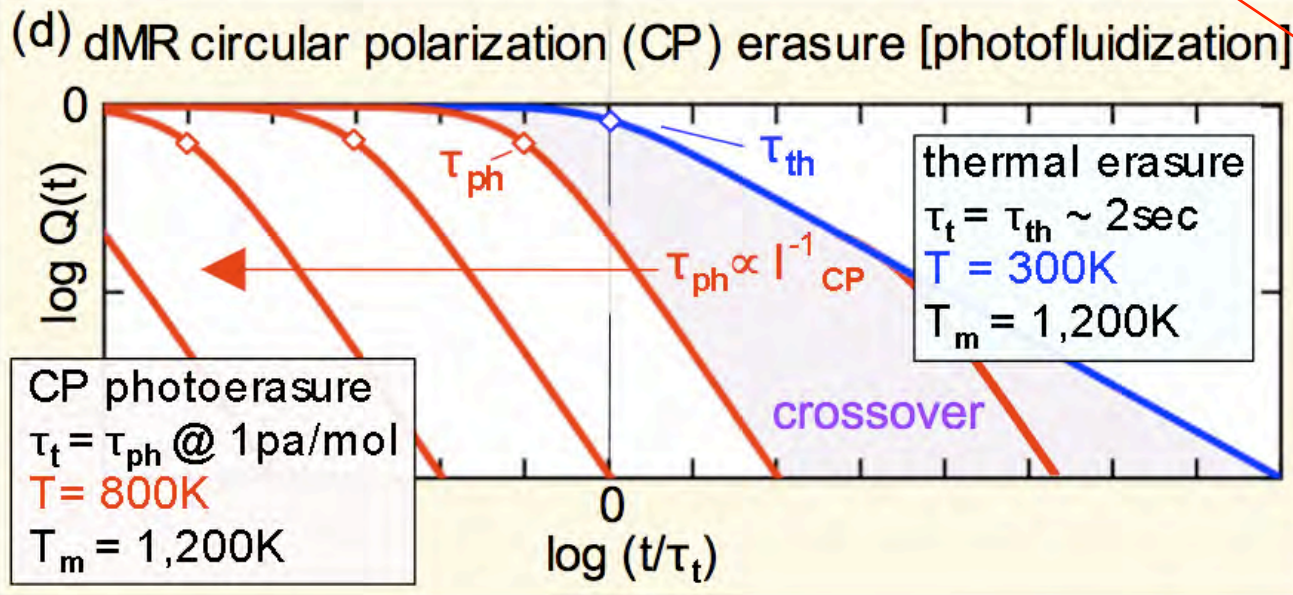
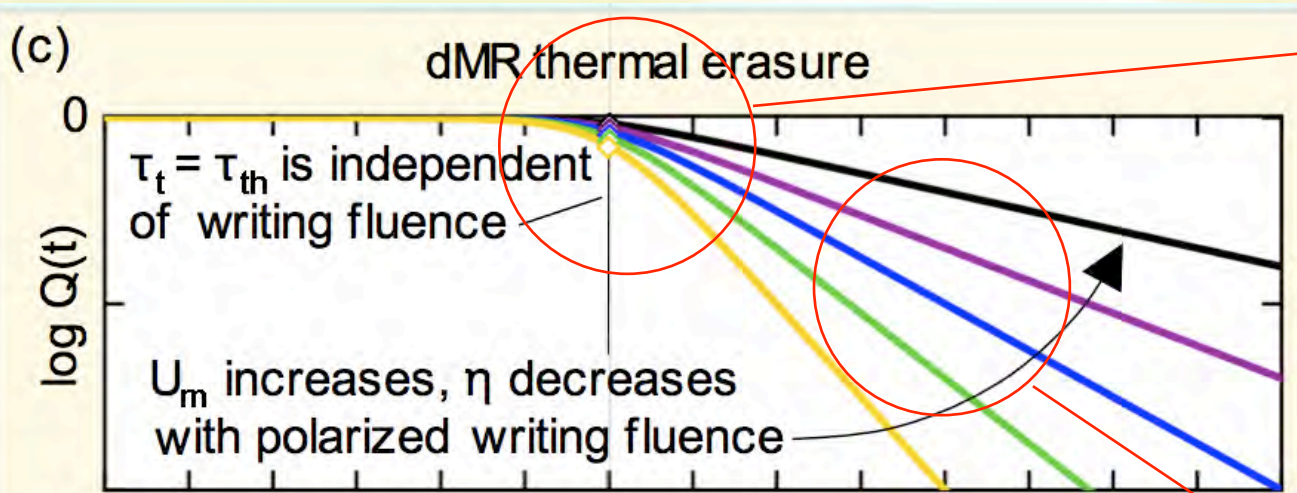
$$U() = k_B T \log(/\tau_r)$$

it's an orientational glass

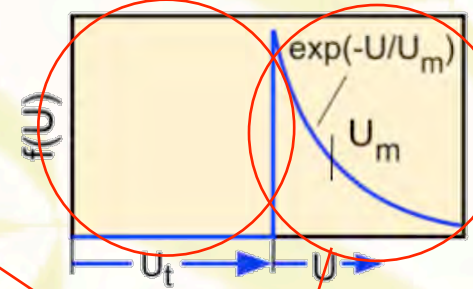


$$U_t \sim k_B T \ln(10^{11}) = 7,500 K$$

dependence on writing and erasing



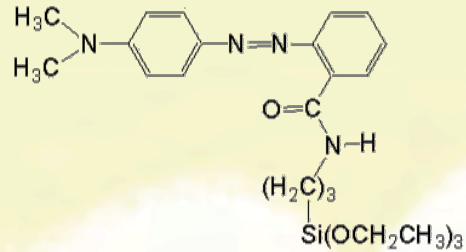
local



collective

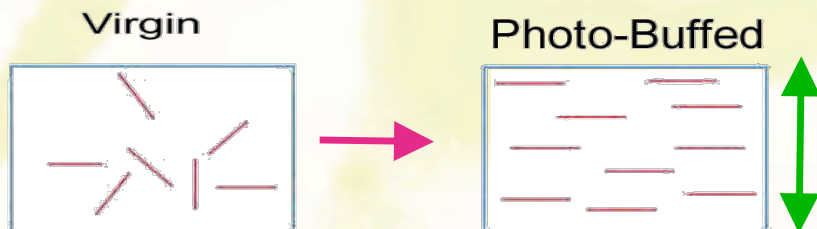
azo-based tethered molecular monolayer

- ◆ precursor
(a derivative of methyl red - dMR)

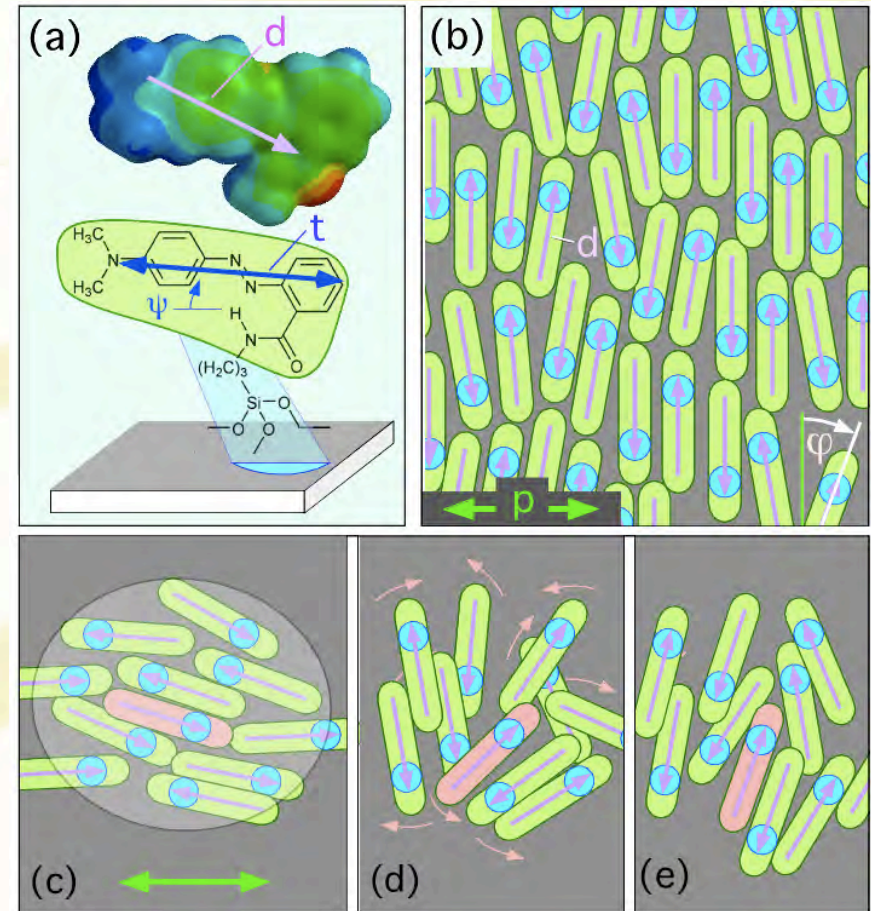


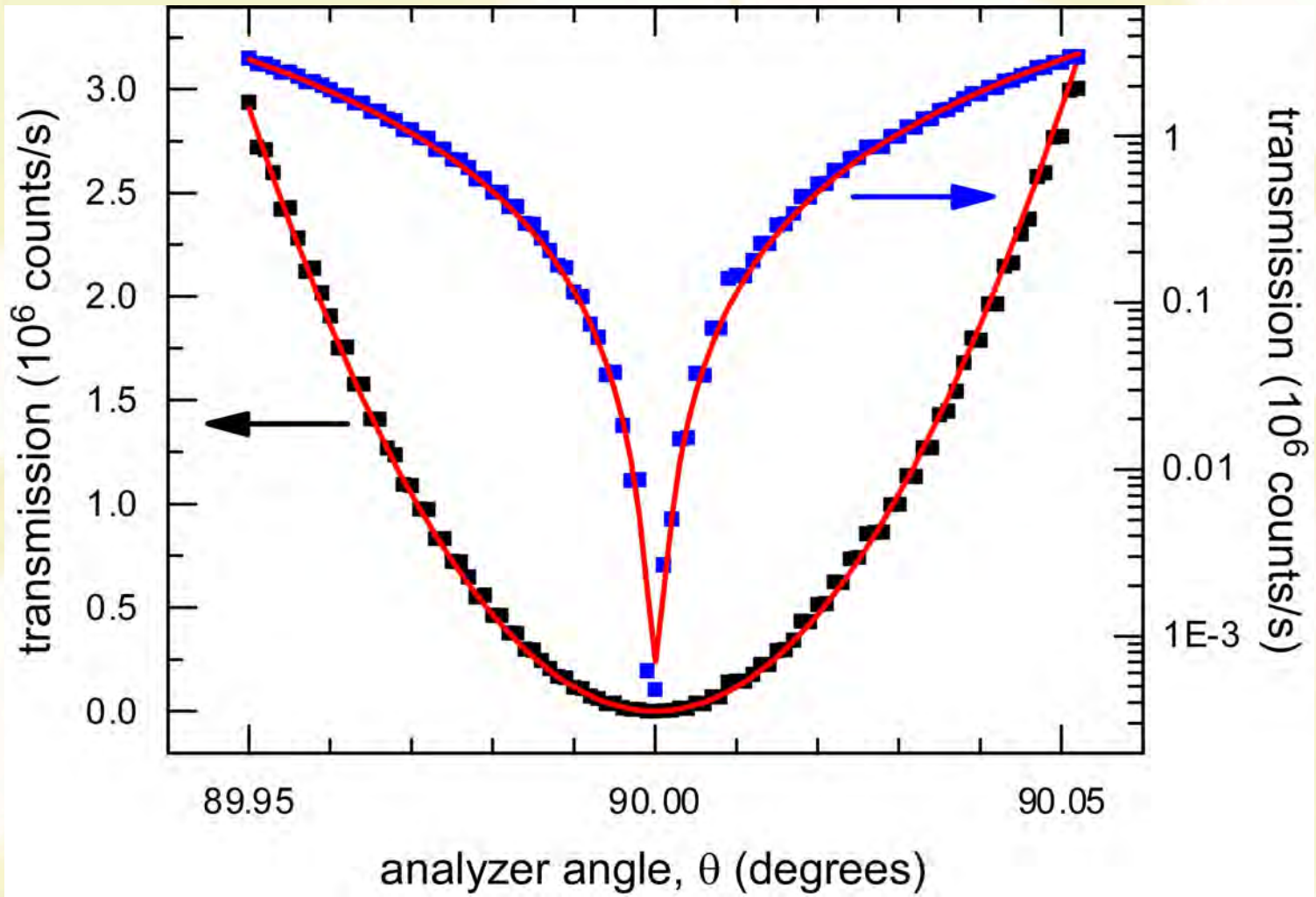
- ◆ covalently bonded to glass surface
- ◆ photo-orientates in-plane

Weigert mechanism:



- ◆ photo-orientable 2D XY system





a different model

Michael F. Shlesinger *Ann. Rev. Phys. Chem.* 1988. 39: 269–90

- ◆ *orientational diffusion with a divergent distribution of waiting times:*

$$Q(t) = \int \exp(-t/\tau) H(\tau) d\tau$$

- ◆ $\tau = \tau_0 \exp(\Delta/kT)$, the rate is determined by activated hopping with barrier Δ
- ◆ $f(\Delta) = (kT_0)^{-1} \exp(-\Delta/kT_0)$, the barrier is distributed with a mean width kT_0
- ◆ *With these assumptions Schlesinger shows*

$$H(\tau) \propto (\tau_0/\tau)^\beta$$

with

$$\beta - 1 = T/T_0$$



barrier model results

- ◆ integration gives

$$Q_e(t) = [1 + (t/\tau_o)]^{(1-\beta)}, \text{ with } \beta > 1$$

- ◆ power-law writing and erasing:

- β controls erasing: $\Delta n(t) \sim Q_e(t) = (t/\tau_o)^{(1-\beta)} = (t/\tau_o)^{-T/T_o}$ at large t

- writing: $\Delta n_w \propto (1 - Q_e(t)) \propto (t/\tau_o)$ at small t

- ◆ β depends on net energy that generated the written state through mean barrier height T_o

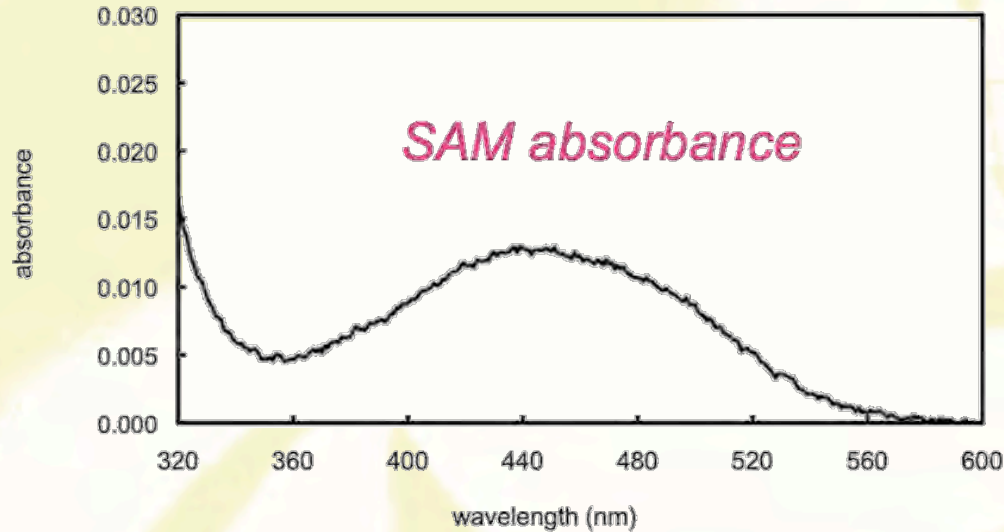
$$\beta - 1 = T/T_o$$

- ◆ T_o increases as $\log(\text{dose})$

$$T_o \propto \delta(t)$$



coverage measured by absorbance



$$\frac{1}{N_s} = \frac{3}{2A \ln 10} \langle \sin^2 \theta \rangle \sigma_{iso}$$

A : absorbance

θ : tilt angle

N_s : number of molecule/area

$$\sigma_{iso} = 1.1 \times 10^{-16} \text{ cm}^2$$

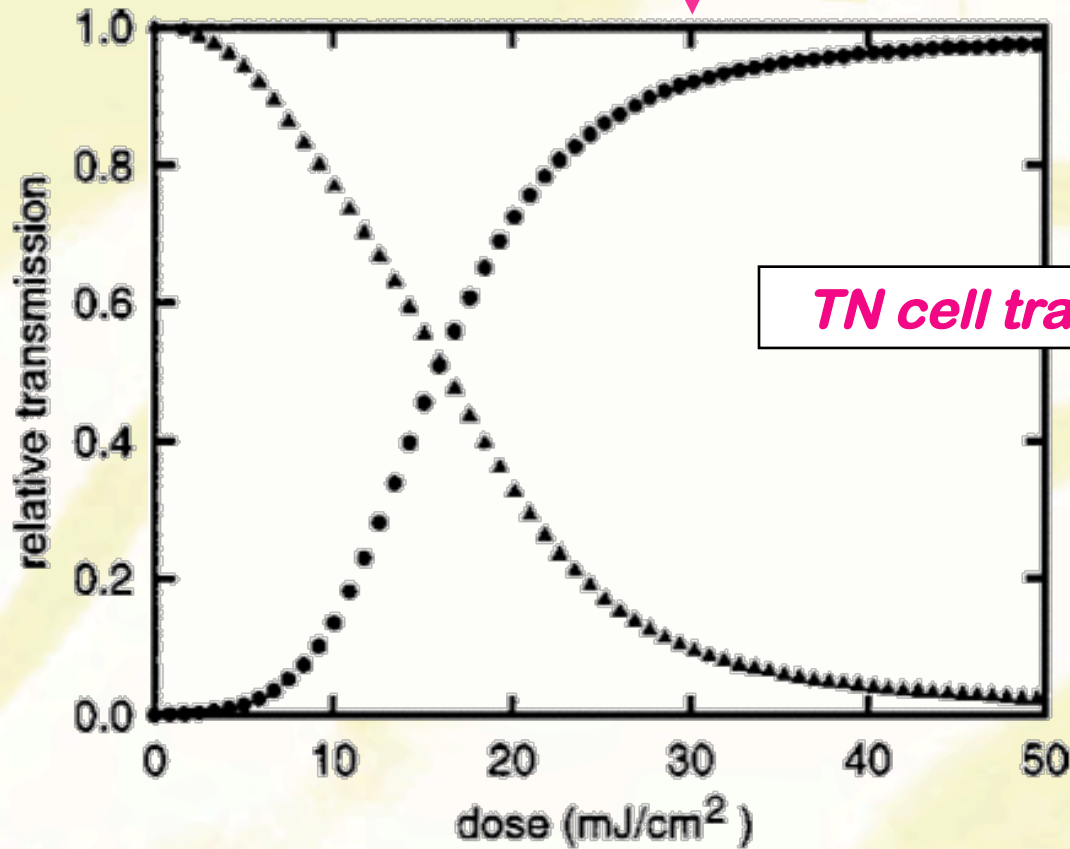
(d-MR in chloroform at 40 nm)

$$\frac{1}{N_s} \approx 1.0 \text{ nm}^2 \rightarrow \text{monolayer } (\sim 70\% \text{ coverage})^*$$

* AFM shows uniform layer morphology

aligns nematics

~2 absorbed 450 nm photons/molecule

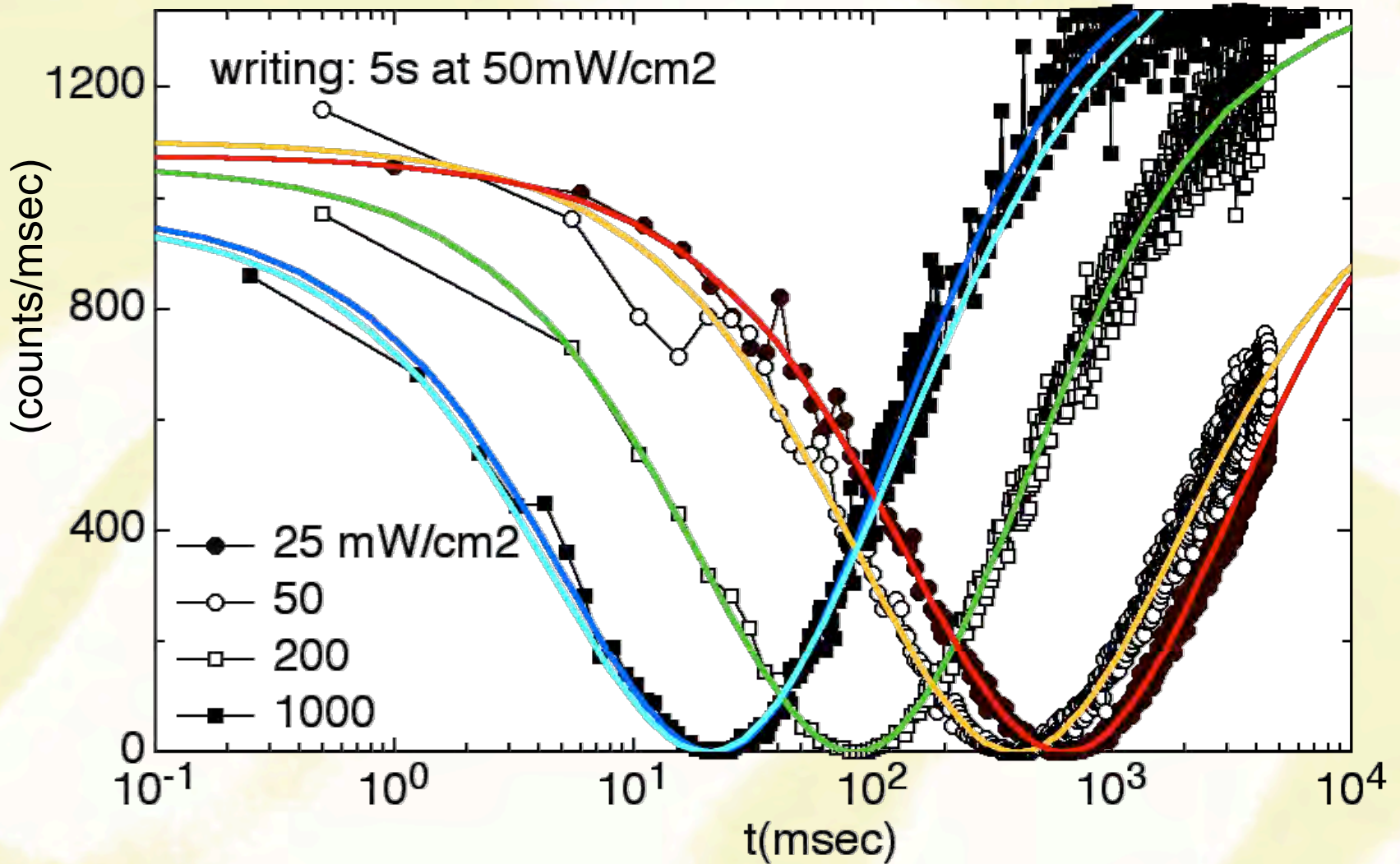


TN cell transmission

Yi & Furtak, APL 2004



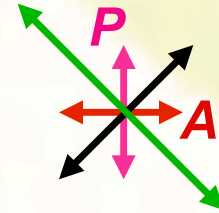
reorienting the writing polarization (+45° to -45°)



model (motivated by data)

◆ $Q(t) = \text{in-plane nematic-like order parameter} \propto \Delta n$

◆ $\text{transmission (thin layer)} = T \propto \Delta n^2 \propto Q^2$

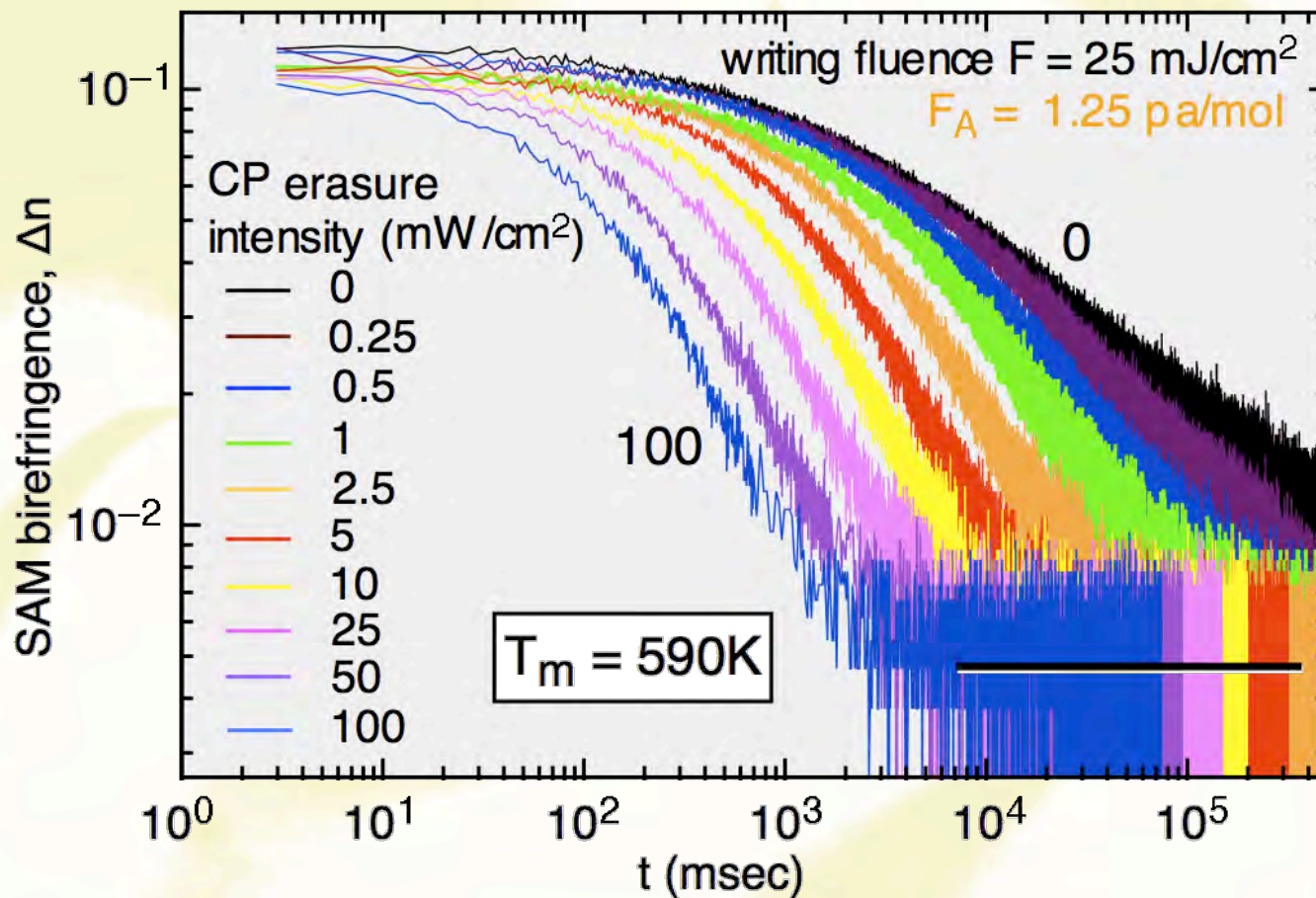


◆ *model: orientational diffusion with a divergent distribution of waiting times:*

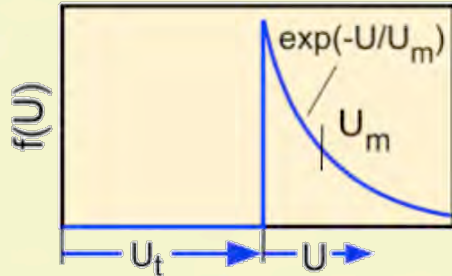
$$\text{relaxation: } Q_e(t) = \int H(\tau) * \exp -(t/\tau)^\alpha d\tau$$

where data indicates: $H(\tau) \sim (\tau_0/\tau)^\beta * \exp -(\tau_0/\tau)^\alpha$

erasing the birefringence with circular polarized light

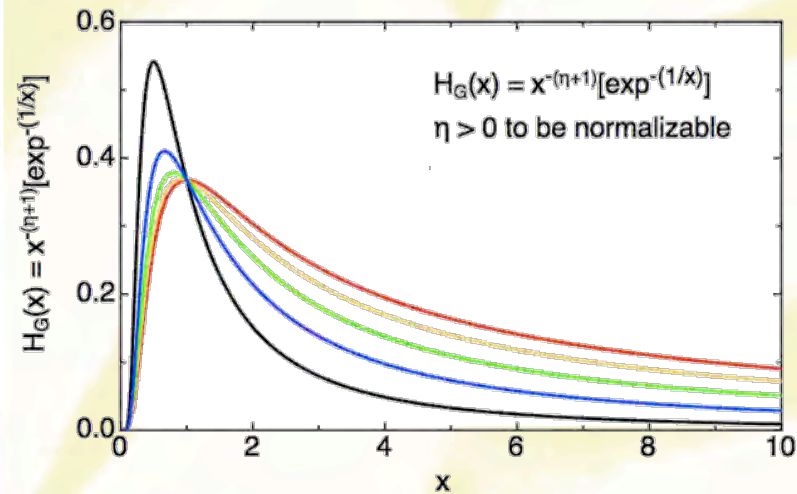


statistics of extreme values of random functions



Scher Schlesinger distribution

Schlesinger, Ann. Rev. Phys. Chem. (1988)



Gumbel distribution

Schlesinger, Ann. Rev. Phys. Chem. (1988)

$$f_G(U) = \left\{ \beta / [U_m \Gamma(1/\beta)] \right\} \left\{ \exp[-U/U_m + \exp(-\beta U/U_m)] \right\}$$

$$H_G(\tau) = [\alpha / \Gamma(h/\beta) \tau_t] [\exp(-\tau_t/\tau)^\alpha] [\tau/\tau_t]^{-(\eta+1)}$$

$$\Delta n(t) \propto Q_G(t) = 1 / [1 + (t/\tau_t)^\alpha]^{\eta/\alpha}$$

dMR monolayer is an orientational glass

