# Symmetry aspects of Unconventional Superconductivity

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Badee - Ce-Schieffe Micciche fecdcii



BCS ea fied hê ?  
i e de: 
$$\mathcal{H} = \sum_{\vec{k},s} \xi_{\vec{k}} c^{\dagger}_{\vec{k}s} c_{\vec{k}s} + g \sum_{\vec{k},\vec{k}'} c^{\dagger}_{\vec{k}\uparrow} c^{\dagger}_{-\vec{k}\downarrow} c_{-\vec{k}'\downarrow} c_{\vec{k}'\uparrow}$$
?

dec ig fie aci e b ea f

ea fied : 
$$\rho_{\vec{q}} = \sum_{\vec{k},s} \langle c^{\dagger}_{\vec{k}+\vec{q}s} c_{\vec{k}s} \rangle$$
 a icede i  
 $\vec{S}_{\vec{q}} = \sum_{\vec{k}} \sum_{s,s'} \langle c^{\dagger}_{\vec{k}s} \vec{\sigma}_{ss'} c_{\vec{k}s'} \rangle$  i de i  
 $b_{\vec{k}} = \langle c_{-\vec{k}\downarrow} c_{\vec{k}\uparrow} \rangle$  BCS-"ffdiag a"

BCS ea fied he  
i e de: 
$$\mathcal{H} = \sum_{\substack{\vec{k},s \\ badeeg}} \xi_{\vec{k}} c_{\vec{k}s}^{\dagger} c_{\vec{k}s} + g \sum_{\substack{\vec{k},\vec{k}' \\ \vec{k} \uparrow}} c_{\vec{k} \uparrow}^{\dagger} c_{-\vec{k} \downarrow} c_{-\vec{k}' \downarrow} c_{\vec{k}' \uparrow}^{\dagger}$$
  
aiigieaci  
badeeg  
badeeg  
badeeg  
i gieaci  
 $k_{\vec{k}} = \epsilon_{\vec{k}} - \mu = \frac{\hbar^2}{2m} (\vec{k}^2 - k_F^2)$   
aiigieaci  
 $U(\vec{r} - \vec{r}) = g \delta^{(3)} (\vec{r} - \vec{r}')$   
 $i eaci$   
 $V(\vec{q} = \vec{k} - \vec{k}') = \int d^3 r U(\vec{r}) e^{i\vec{q} \cdot \vec{r}} = g = V_{\vec{k},\vec{k}'}$   
 $c ide caeigbeee e - e eec ai f ie i(i ige)$ 

BCS ea fied he  

$$\mathcal{H}_{mf} = \sum_{\vec{k},s} \xi_{\vec{k}} c^{\dagger}_{\vec{k}s} c_{\vec{k}s} + g \sum_{\vec{k},\vec{k}'} \{b^{*}_{\vec{k}'} c_{-\vec{k}\perp} c_{\vec{k}\uparrow} + b_{\vec{k}'} c^{\dagger}_{\vec{k}\uparrow} c^{\dagger}_{-\vec{k}\perp} - b^{*}_{\vec{k}} b_{\vec{k}'}\}$$

$$\stackrel{?}{=} \sum_{\vec{k},s} \xi_{\vec{k}} c^{\dagger}_{\vec{k}s} c_{\vec{k}s} - \sum_{\vec{k}} \left\{ \Delta^{*} c_{-\vec{k}\perp} c_{\vec{k}\uparrow} + \Delta c^{\dagger}_{\vec{k}\uparrow} c^{\dagger}_{-\vec{k}\downarrow} - \Delta^{*} b_{\vec{k}} \right\}$$
fi d ai aice ae iff  $\frac{\partial}{\partial t} \gamma^{\dagger}_{\vec{k}} = i[\mathcal{H}_{mf}, \gamma^{\dagger}_{\vec{k}}] = E_{\vec{k}} \gamma^{\dagger}_{\vec{k}}$ 
B g b - a f ai  $c_{\vec{k}\uparrow} = u^{*}_{\vec{k}} \gamma_{\vec{k}1} + v_{\vec{k}} \gamma^{\dagger}_{\vec{k}2} |u_{\vec{k}}|^{2} + |v_{\vec{k}}|^{2} = 1$ 

$$\stackrel{\longrightarrow}{=} ai aice eg E_{\vec{k}} = \sqrt{\xi^{2}_{\vec{k}} + \Delta^{2}}$$

$$\stackrel{\longrightarrow}{=} \mathcal{H} = \sum_{\vec{k}} [\xi_{\vec{k}} - E_{\vec{k}} + \Delta b_{\vec{k}}] + \sum_{\vec{k}} E_{\vec{k}} (\gamma^{\dagger}_{\vec{k}1} \gamma_{\vec{k}1} + \gamma^{\dagger}_{\vec{k}2} \gamma_{\vec{k}2})$$











Ze - e e a e Ga a  $1 = -gN(0) \int_{0}^{\epsilon_{c}} \frac{d\xi}{\sqrt{\xi^{2} + \Delta^{2}}} = -gN(0) \sinh^{-1} \frac{\epsilon_{c}}{\Delta}$   $\xrightarrow{?}$   $\Delta \approx 2\epsilon_{c}e^{-1/|g|N(0)} = 1.764k_{B}T_{c}$  ea c i g C de ai e e g a  $E_{cond} = -\frac{1}{2}N(0)|\Delta|^{2}$ difica i f he a i a ice ec  $k_{F}$ 

h e-ie

## Paiigi e aci

eec - h e C bi eaci  
P aiai effec : 
$$\frac{1}{\varepsilon(\vec{q},\omega)} \approx \frac{q^2}{q^2 + k_{TF}^2} + \frac{q^2}{q^2 + k_{TF}^2} \frac{\omega_{\vec{q}}^2}{\omega^2 - \omega_{\vec{q}}^2}$$
ih  $k_{TF}^2 = \frac{6\pi e^2 n_e}{\epsilon_F}$  Th a-Fe i ceeigegh  $\lambda_{TF} = k_{TF}^{-1} \sim 5 - 10 \text{\AA}$   
 $V_{\vec{k},\vec{k}'} = \frac{4\pi e^2}{q^2 \varepsilon(\vec{q},\omega)} = \underbrace{\frac{4\pi e^2}{q^2 + k_{TF}^2}}_{e \quad C \quad b} + \underbrace{\frac{4\pi e^2}{q^2 + k_{TF}^2} \frac{\omega_{\vec{q}}^2}{\omega^2 - \omega_{\vec{q}}^2}}_{eec \quad -h}$ 
 $\vec{q} = \vec{k} - \vec{k}'$ 

$$\omega_q = sq$$



# When Coulomb repulsion is too strong for electron-phonon induced pairing

Adenaatieeways to superconductinity a

#### eaci & C е ai е badeec - h ie aci e h - a ged $(\lambda)$ С В dC e ai a ef ci : H a id C be i? highe -a g a $\psi(\vec{r}, s; \vec{r}', s') = f(|\vec{r} - \vec{r}'|)\chi(s, s')$ е aiig ih $f(r \rightarrow 0) \neq 0$ $l > 0 \longrightarrow f(r \to 0) \propto r^{l}$ ? "cacieaci" effecie eaieaga e **/=0** i a f "c acieaci" $: \Psi_{ss'} \stackrel{?}{\vec{k}} = \langle \hat{c}_{\vec{k}s} \hat{c}_{-\vec{k}s'} \rangle = \Phi(\vec{k}) \chi(s,s')$ e fai fide ica eec S orbital spin aefciaaieic<sup>6</sup> de a iceechage ee ai: i ge еe dd dd ai : i e dd еe $\vec{k} \rightarrow -\vec{k}$ $s \leftrightarrow s'$

### Key symmetries - Anderson's Theorems (1959,1984)

Cooper pairs with total momentum P<sub>tot</sub>=0 form from degenerate quasiparticle states. *How to guarantee degenerate partners*?

• Spin singlet pairing: time reversal symmetry  

$$|\vec{k}\uparrow\rangle \longrightarrow \hat{T}|\vec{k}\uparrow\rangle = |-\vec{k}\downarrow\rangle \iff \epsilon_{\vec{k}\uparrow} = \epsilon_{-\vec{k}\downarrow}$$
  
harmful: magnetic impurities, ferromagnetism,  
Zeemann fields (paramagnetic limiting)  
• Spin triplet pairing: time reversal & inversion symmetry

$$|\vec{k}\uparrow\rangle \longrightarrow \begin{cases} \hat{I}|\vec{k}\uparrow\rangle = |-\vec{k}\uparrow\rangle & \epsilon_{\vec{k}\uparrow} &= \epsilon_{-\vec{k}\uparrow} \\ \hat{T}|\vec{k}\uparrow\rangle = |-\vec{k}\downarrow\rangle & \longleftrightarrow & \epsilon_{\vec{k}\downarrow} \\ \hat{T}\hat{I}|\vec{k}\uparrow\rangle = |\vec{k}\downarrow\rangle & \longleftrightarrow & \epsilon_{\vec{k}\downarrow} \end{cases}$$

crystal structure without inversion center



S i f c ai e cha ge echa i  
E cha ge i e aci : 
$$\mathcal{H}_{ex} = \int d^3r d^3r' U \delta(\vec{r} - \vec{r}') \rho_{\uparrow}(\vec{r}) \rho_{\downarrow}(\vec{r}')$$
  
 $\rightarrow$  i-i d ced ca " ag e ic fie d"  $\vec{H}(\vec{r},t) = \frac{I}{\mu_B \hbar} \vec{S}(\vec{r},t)$   
i d ced i ai ai : d a ica i ce ibii  
 $\vec{S}(\vec{r}',t') = \mu_B \int d^3r \, dt \, \chi(\vec{r}' - \vec{r},t'-t) \, \vec{H}(\vec{r},t)$ 

Side i - i de i i e aci :

$$\mathcal{H}_{sf} = -\frac{I^2}{\hbar^2} \int d^3r \ d^3r' \int dt \ dt' \ \chi(\vec{r} - \vec{r}'; t - t') \ \vec{S}(\vec{r}, t) \cdot \vec{S}(\vec{r}', t')$$

$$6$$
i ified if c ai e chage de

,

## A e ai e echa i f C e ai i g

Pail g f e e i e i e a c i : K h & L i ge (1965) c e e e d C b e ia i e a ha g-a g e d c i a ai (ha Fe i e d g e) F i e d e c i a i :  $V(r) \propto r^{13} \cos(2k_F r)$ ai i g i high-a g a e c ha e  $T_c/T_F \sim \exp\{-(2l)^4\}$ e !

Paiigb ageicfcai : Be & Schieffe (1966)



Sifcai echage echai effecie aligieaci:  $? \mathcal{H}'_{sf} = \sum_{\vec{k},\vec{k}'} \sum_{s_1,s_2,s_3,s_4} ? V_{\vec{k},\vec{k}';s_1s_2s_3s_4} c^{\dagger}_{\vec{k},s_1} c^{\dagger}_{-\vec{k},s_2} c_{-\vec{k}',s_5} c_{\vec{k}',s_4}$   $V_{\vec{k},\vec{k}';s_1s_2s_3s_4} = -\frac{I^2}{4} Re\chi(\vec{k}-\vec{k}',\omega=\varepsilon_{\vec{k}}-\varepsilon_{\vec{k}'})\vec{\sigma}_{s_1s_4}\cdot\vec{\sigma}_{s_2s_3}$ d a ica i ce ibii:  $\chi(\vec{q},\omega) = \frac{\chi_0(\vec{q},\omega)}{1-I\chi_0(\vec{q},\omega)} RPA$ f i iceec ga:  $\chi_0(\vec{q},\omega) \approx N(0) \left(1-\frac{\vec{q}'^2}{12k_F^2} + i\frac{\pi}{2}\frac{\omega}{v_F|\vec{q}|}\right)$   $\omega \in \varepsilon$ 

Sifcai e chage echai  
effecie aiigi e aci:  

$$\mathcal{H}'_{sf} = \sum_{\vec{k},\vec{k}'} \sum_{s_1,s_2,s_3,s_4} V_{\vec{k},\vec{k}';s_1s_2s_3s_4} c^{\dagger}_{\vec{k},s_1} c^{\dagger}_{-\vec{k},s_2} c_{-\vec{k}',s_3} c_{\vec{k}',s_4}$$

$$V_{\vec{k},\vec{k}';s_1s_2s_3s_4} = -\frac{I^2}{4} Re\chi(\vec{k} - \vec{k}', \omega = \varepsilon_{\vec{k}} - \varepsilon_{\vec{k}'}) \vec{\sigma}_{s_1s_4} \cdot \vec{\sigma}_{s_2s_3}$$
C e i cha e:  

$$V_{\vec{k},\vec{k}'} = \frac{3I^2}{4} Re\chi(\vec{k} - \vec{k}', \omega = \varepsilon_{\vec{k}} - \varepsilon_{\vec{k}'}) \qquad S=0 \quad \text{i ige}$$

$$V_{\vec{k},\vec{k}'} = -\frac{I^2}{4} Re\chi(\vec{k} - \vec{k}', \omega = \varepsilon_{\vec{k}} - \varepsilon_{\vec{k}'}) \qquad S=1 \quad \text{i ie}$$
': e ie : a acie

v

S i f c ai e cha ge echa i  
Paiigf i i e 
$$\frac{k_x + ik_y}{\sqrt{2k_F}}$$
  $m = +1$   
a g a c e fga f ci  
 $\Delta_{\vec{k}} = \Delta g_{\vec{k}}^m$   $g_{\vec{k}}^m = \frac{k_z}{k_F}$   $m = 0$   
P ec ed effeci e i e aci :  $\frac{k_x - ik_y}{\sqrt{2k_F}}$   $m = -1$   
 $V(\xi, \xi') = -\frac{I^2}{4} \sum_{i'} g_i \chi(\vec{k} - \vec{k'}, \omega = 0)g_i, \delta(\xi - \xi_i)\delta(\xi' - \xi_i^2) \approx \begin{cases} V_1 & |\xi|, |\xi'| < \epsilon \\ 0 & \text{otherwise} \end{cases}$   
 $V_1 = -\frac{I}{12} \frac{IN(0)}{(1 - IN(0))^2}$   $k_B T_c = \hbar \epsilon_c e^{-1/\lambda_s}$   
 $\epsilon \quad \epsilon \quad \xi \quad cha ace i ice eg : a a ag \quad ec \quad \epsilon_c = \frac{8}{\pi IN(0)} (1 - IN(0))E_F$ 





v

$$\hat{k}_{\vec{k},ss'} = -\sum_{\vec{k}',s_1s_2} V_{\vec{k}',\vec{k};s_1s_2s's} \langle c_{\vec{k}'s_1}^{\dagger} c_{-\vec{k}'s_2}^{\dagger} \rangle \qquad \hat{\Delta}_{\vec{k}} = \begin{pmatrix} \Delta_{\vec{k}\uparrow\uparrow} & \Delta_{\vec{k}\uparrow\downarrow} \\ \Delta_{\vec{k}\downarrow\uparrow} & \Delta_{\vec{k}\downarrow\downarrow} \end{pmatrix}$$

Sef-c i e ga e ai  
B g b a f ai 
$$\longrightarrow$$
 Q ai aice ec  
 $E_{\vec{k}} = \sqrt{\xi_{\vec{k}}^2 + |\Delta_{\vec{k}}|^2}$   $|\Delta_{\vec{k}}|^2 = \frac{1}{2} \operatorname{tr} \left( \hat{\Delta}_{\vec{l}}^{\dagger}_{\vec{k}} \hat{\Delta}_{\vec{k}} \right)$   
N e: ai aice ga i k-de e de  $\hat{\Delta}_{\vec{k}} = \begin{pmatrix} \Delta_{\vec{k}\uparrow\uparrow} & \Delta_{\vec{k}\uparrow\downarrow} \\ \Delta_{\vec{k}\downarrow\uparrow} & \Delta_{\vec{k}\downarrow\downarrow} \end{pmatrix}$   
A i :  
 $\hat{\Delta}_{\vec{k}}^{\dagger} \hat{\Delta}_{\vec{k}} = |\Delta_{\vec{k}}|^2 \hat{\sigma}_0$  i dege e ac : SU(2) e  
i e e e a & i e i e

Sefc i e cee ai :  

$$\Delta_{\vec{k},s_1s_2} = -\sum_{\vec{k}',s_3s_4} V_{\vec{k},\vec{k}';s_1s_2s_3s_4} \frac{\Delta_{\vec{k}',s_4s_3}}{2E_{\vec{k}}} \tanh\left(\frac{E_{\vec{k}}}{2k_BT}\right)$$

$$\begin{split} \mathbf{S} \quad \mathbf{c} \quad \mathbf{e} \quad \mathbf{f} \quad \mathbf{he} \quad \mathbf{ga} \quad \mathbf{f} \quad \mathbf{C}_{6}^{i} \quad \mathbf{f} \\ \mathbf{Ga} \quad \mathbf{f} \quad \mathbf{ci} \quad : \quad 2 \quad 2 \quad \mathbf{a} \quad \mathbf{i} \quad \mathbf{i} \quad \mathbf{ace} \qquad \Delta_{\vec{k},ss'} = -\sum_{\vec{k}',s_3s_4} V_{\vec{k},\vec{k}',s} \mathcal{F}_{ss_4}(c_{\vec{k}'s_3}c_{-\vec{k}'s_4}) \\ \hline \left\{ \begin{pmatrix} \mathbf{c}_{-\vec{k},s_1} c_{\vec{k},s_2} \end{pmatrix} = \boldsymbol{\phi}(\vec{k}) \chi_{s_1s_2} \\ \mathbf{bia} \quad \mathbf{i} \end{matrix} \right. \qquad \Delta_{\vec{k},ss'}^* = -\sum_{\vec{k}'s_1s_2} V_{\vec{k}',\vec{k}';s_1s_2s's}(c_{\vec{k}'s_1}^{\dagger}c_{-\vec{k}'s_2}) \\ \mathcal{F}_{\vec{k}'s_1s_2}^{\dagger} = -\sum_{\vec{k}'s_1s_2} V_{\vec{k}',\vec{k}';s_1s_2s's}(c_{\vec{k}'s_1}^{\dagger}c_{-\vec{k}'s_2}) \\ \mathcal{F}_{\vec{k}'s_1s_2}^{\dagger} = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \quad \mathbf{e} \quad \mathbf{e} \quad \mathbf{ai} \quad \mathbf{i} \quad \mathbf{i} \quad \mathbf{g} \\ \mathbf{f}_{\vec{k}'s_1s_2}^{\dagger} = -\phi(-\vec{k}) \quad \Leftrightarrow \quad \chi_{s_1s_2} = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \quad \mathbf{dd} \quad \mathbf{ai} \quad \mathbf{i} \quad \mathbf{i} \quad \mathbf{e} \\ |\downarrow\downarrow\rangle \\ \Delta_{\vec{k},s_1s_2}^{\dagger} = -\Delta_{-\vec{k},s_2s_1}^{\dagger} = \begin{cases} \Delta_{-\vec{k},s_1s_2}^{\dagger} = -\Delta_{\vec{k},s_2s_1}^{\dagger} \quad \mathbf{even} \\ -\Delta_{-\vec{k},s_1s_2}^{\dagger} = \Delta_{\vec{k},s_2s_1}^{\dagger} \quad \mathbf{odd} \end{cases}$$

$S \ c \ e \ f \ he \ ga \ f \ c_{6}^{i}$ Ga f ci : 22 a i i i ace $\Delta_{\vec{k},ss'} = -\sum_{\vec{k}',s_3s_4} V_{\vec{k},\vec{k}';ss's_3s_4} \langle c_{\vec{k}'s_3}c_{-\vec{k}'s_4} \rangle$ $\Delta_{\vec{k},ss'} = -\sum_{\vec{k}'s_1s_2} V_{\vec{k}',\vec{k};s_1s_2s's} \langle c_{\vec{k}'s_1}^{\dagger}c_{-\vec{k}'s_2}^{\dagger} \rangle$
Ee ai ige
$\widehat{\Delta}_{ec{k}} \stackrel{?}{=} \left( egin{array}{cc} \Delta_{ec{k},\uparrow\uparrow} & \Delta_{ec{k},\uparrow\downarrow} \ \Delta_{ec{k},\downarrow\uparrow} & \Delta_{ec{k},\downarrow\downarrow} \end{array}  ight) = \left( egin{array}{cc} 0 & \psi(ec{k}) \ -\psi(ec{k}) & 0 \end{array}  ight) = i \widehat{\sigma}_y \psi(ec{k})$
e e e edb caaf ci $\psi$ $\psi$ e e $E_{\vec{k}} = \sqrt{\xi_{\vec{k}}^2 +  \psi(\vec{k}) ^2}$
Odd ai i e
$\widehat{\Delta}_{\vec{k}} = \begin{pmatrix} -d_x(\vec{k}) + id_y(\vec{k}) & d_z(\vec{k}) \\ d_z(\vec{k}) & d_x(\vec{k}) + id_y(\vec{k}) \end{pmatrix} = i \left( \vec{d}(\vec{k}) \cdot \hat{\vec{\sigma}} \right) \hat{\sigma}_y$
e e e ed b ec f ci $$ dd $E_{\vec{k}} = \sqrt{\xi_{\vec{k}}^2 +  \vec{d}(\vec{k}) ^2}$

S c <sup>?</sup> f he g Gafci: 22 ai	$\begin{array}{cccc} \text{ga f } C & \text{i} \\ & 6 \\ \text{i i ace} \\ & & & \\ \Delta_{\vec{k},ss'} = -\sum_{\vec{k}',s_3s_4} V_{\vec{k},\vec{k}';ss's_3s_4} \langle c_{\vec{k}'s_3}c_{-\vec{k}'s_4} \rangle \\ & & & \\ \Delta_{\vec{k},ss'}^* = -\sum_{\vec{k}'s_1s_2} V_{\vec{k}',\vec{k};s_1s_2s's} \langle c_{\vec{k}'s_1}^\dagger c_{-\vec{k}'s_2}^\dagger \rangle \end{array}$
Ee ai iige	
$\widehat{\Delta}_{\vec{k}} \stackrel{?}{=} \begin{pmatrix} \Delta_{\vec{k},\uparrow\uparrow} \\ \Delta_{\vec{k},\downarrow\uparrow} \end{pmatrix}$ e e e edb caaf	$\begin{split} & \Delta_{\vec{k},\uparrow\downarrow} \\ & \Delta_{\vec{k},\downarrow\downarrow} \end{pmatrix} = \begin{pmatrix} 0 & \psi(\vec{k}) \\ -\psi(\vec{k}) & 0 \end{pmatrix} = i \hat{\sigma}_y \psi(\vec{k}) \\ & \text{ci}  \psi \stackrel{\rightarrow}{\longrightarrow} \psi \stackrel{\rightarrow}{\longrightarrow} e e  E_{\vec{k}} = \sqrt{\xi_{\vec{k}}^2 +  \psi(\vec{k}) ^2} \end{split}$
	, , , , , , , , , , , , , , , , , , , ,
Odd ai i i e	$d_x \leftrightarrow  \downarrow\downarrow\rangle -  \uparrow\uparrow\rangle$
ic figai	$\left. d_y \leftrightarrow -i  ight  \downarrow \downarrow  angle -i  ight  \uparrow \uparrow  angle  ight angle ~~ igstarrow ~~ "ec d \perp ec S ~~ "  ight.$

Thermodynamic properties and Low.enercy, spectrum

Taii e<sub>2</sub> ea e  $\text{Paiigi eaci}: \quad V_{\vec{k}\,,\vec{k}';s_1s_2s_3s_4} = J^0_{\vec{k}\,,\vec{k}'}\hat{\sigma}^0_{s_1s_4}\hat{\sigma}^0_{s_2s_3} + J_{\vec{k}\,,\vec{k}'}\hat{\sigma}_{s_1s_4}\cdot\,\hat{\sigma}_{s_2s_3}$ de i-de i i - i Sef-c i e ce e a i : dd ai i i e ee ai iiqe  $\psi(\vec{k}) = -\sum_{\vec{k}'} \underbrace{(J_{\vec{k},\vec{k}'}^0 - 3J_{\vec{k},\vec{k}'})}_{= v_{\vec{k},\vec{k}'}^*} \frac{\psi(\vec{k}')}{2E_{\vec{k}'}} \tanh\left(\frac{E_{\vec{k}'}}{2k_BT}\right) \qquad \vec{d}(\vec{k}) = -\sum_{\vec{k}'} \underbrace{(J_{\vec{k},\vec{k}'}^0 + J_{\vec{k},\vec{k}'})}_{= v_{\vec{k},\vec{k}'}^*} \frac{\vec{d}(\vec{k}')}{2E_{\vec{k}'}} \tanh\left(\frac{E_{\vec{k}'}}{2k_BT}\right)$  $\vec{\lambda} \psi(\vec{k}) = -N(0) \langle v^s_{\vec{k},\vec{k}'} \psi(\vec{k}') \rangle_{\vec{k}',FS} -\lambda \vec{d}(\vec{k}) = -N(0) \langle v^t_{\vec{k},\vec{k}'} \vec{d}(\vec{k}') \rangle_{\vec{k}',FS}$ eige a e  $\lambda$   $\implies$   $k_B T_c = 1.14 \epsilon_c e^{-1/\lambda}$ 





$$L = R(0) \frac{E}{\Delta_m} \ln \left| \frac{1 + \frac{E}{\Delta_m}}{1 - \frac{E}{\Delta_m}} \right|$$

$$L = Q(0) \frac{E}{\Delta_m} \ln \left| \frac{1 + \frac{E}{\Delta_m}}{1 - \frac{E}{\Delta_m}} \right|$$

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$$L = Q(0) \frac{E}{\Delta_m} \ln \left| \frac{1 + \frac{E}{\Delta_m}}{1 - \frac{E}{\Delta_m}} \right|$$

ic

$$L \stackrel{1?}{-e} \stackrel{?}{e} \stackrel{e}{a} \stackrel{e}{a} \stackrel{e}{e} \stackrel{ie}{e} \stackrel{field}{=} 6$$
he d a ic i d i aedb he e ci ed a i a ice
$$E_{\vec{k}} = \sqrt{\xi_{\vec{k}}^2 + |\Delta_{\vec{k}}|^2} \quad |\Delta_{\vec{k}}|^2 = \frac{1}{2} \operatorname{tr} \left( \hat{\Delta}_{\vec{k}}^{\dagger} \hat{\Delta}_{\vec{k}} \right) \qquad \Delta_{\vec{k}} = \Delta_{-\vec{k}}$$
e a i :
$$N(E) = \frac{2}{\Omega} \sum_{\vec{k}} \delta(E_{\vec{k}} - E)$$

$$N(E) = N(0) \int \frac{d\Omega_{\vec{k}}}{4\pi} \int d\xi \delta(\sqrt{\xi^2 + |\Delta_m g_{\vec{k}}|^2} - E) = N(0) \int \frac{d\Omega_{\vec{k}}}{4\pi} \frac{E}{\sqrt{E^2 - |\Delta_m g_{\vec{k}}|^2}}$$
A i ic ga f ci :
$$\Delta_{\vec{k}} = \Delta_m \cos\theta$$
i e
$$de$$

$$N(E) = N(0) \frac{E}{\Delta_m} \begin{cases} \frac{\pi}{2} & |E| < \Delta_m \\ \arccos(\frac{\Delta_m}{E}) & \Delta_m \ge |E| \end{cases}$$

### L-e ea e e ie

ea i he aiiedeedig ga g

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ecific hea C(T) ?		
Ld eeai dehλ		
NMR		
heac dcii к		















